The Georgi Algorithms of Jet Clustering

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Howard Georgi, arXiv:1408.1161 SFG, arXiv:1408.3823

Conventional Jet Algorithms [implemented in FastJet]

• Longitudinally invariant k_t algorithm

$$d_{ij}=\min(p_{ti}^2,p_{tj}^2)\Delta R_{ij}^2/R^2$$
 .

• Cambridge/Aachen (C/A) algorithm

$$d_{ij} = \Delta R_{ij}^2/R^2 \,, \qquad d_{iB} = 1 \,.$$

• Anti-k_t algorithm

$$d_{ij} = \min\left(1/p_{ti}^2, 1/p_{tj}^2\right) \Delta R_{ij}^2/R^2, \qquad d_{iB} = 1/p_{ti}^2.$$

• Generalized k_t algorithm

$$d_{ij} = \min(E_i^{2p}, E_j^{2p})(1 - \cos \theta_{ij})/(1 - \cos R),$$

$$d_{iB} = E_i^{2p}.$$

• Durham ($e^+e^- k_t$) algorithm

$$d_{ij} = 2 \times \min(E_i^2, E_j^2)(1 - \cos \theta_{ij}).$$

The Georgi Algorithms [arXiv:1408.1161]

Jet function:

$$J_{\beta}(P_{\alpha}) \equiv E_{\alpha} - eta rac{P_{lpha}^2}{E_{lpha}} = E_{lpha} \left[(1-eta) + eta v_{lpha}^2
ight] \, ,$$

with jet momentum $P_{\alpha} = (E_{\alpha}, \mathbf{P}_{\alpha}) \equiv \sum_{i \in \alpha} p_i$ & velocity $v_{\alpha} \equiv \frac{|\mathbf{P}_{\alpha}|}{E_{\alpha}}$ where α is a set of subjets.

- J_{β} increases when clustering:
 - E_{α} increases due to energy conservation;
 - Jet virtuality (mass) P_{α}^2 doesn't increase that much.
- Not only pair-wisely, but also globally.
- Cone implemented implicitly:

$$J_{\beta}(P_{\alpha}+p_j) = (E_{\alpha}+E_j)\left[(1-\beta) + \beta \frac{|\mathbf{P}_{\alpha}|^2 + 2|\mathbf{P}_{\alpha}||\mathbf{p}_j|\cos\theta + |\mathbf{p}_j|^2}{(E_{\alpha}+E_j)^2}\right]$$

• $1 \rightarrow 2$ Splitting:



• Sudakov Basis:

$$q_i \equiv \alpha_i p + rac{q_i^2 - P_{\perp i}^2}{2\alpha_i} n + P_{\perp i} \,,$$

where p & n are light-like, $p \cdot P_{\perp i} = n \cdot P_{\perp i} = 0 \& p \cdot n = 2$. • Energy fraction:

$$z_i \equiv \frac{\alpha_i}{\alpha_{i-1}} \approx \frac{E_i}{E_{i-1}}$$

• Sudakov Factor for Parton i = q, g:

$$\mathbf{\Delta}_{\mathbf{i}}(\mathbf{t}) = \exp\left[-\int_{t_0}^t \frac{dt'}{t'} \int_{z_{min}(t')}^{z_{max}(t')} dz \sum_j \mathbf{P}_{\mathbf{j}\mathbf{i}}(t', z)\right]$$

.

• Analogy to Decay:

$$\mathbb{P}(\mathbf{t}) = e^{-t\Gamma} = \exp\left(-\int_0^t dt'\Gamma\right) = \exp\left(-\int_0^t dt'\int d\Omega_2 |\mathcal{M}|^2\right)$$

• Survival Probability:

$$\mathbb{P}(\mathbf{t}) = e^{-t \int_0^t dt' \Gamma} \sim \Delta_{\mathbf{i}}(\mathbf{t}).$$

• Decay Width:

$$\begin{split} &\Gamma(t)\equiv\int d\Omega_2|\mathcal{M}|^2\sim\int_{z_{min}(t)}^{z_{max}(t)}dz\sum_j P_{ji}(t,z)\\ \bullet \text{ Energy fraction:} \qquad |\mathcal{M}|^2\sim\sum_j \mathbf{P}_{ji}(t,z). \end{split}$$

• Splitting Functions:

• $q \rightarrow qg$:

$$P_{qq}(z) = C_F \frac{1 + (1 - z)^2}{z}$$

• $g \rightarrow gg$:

$$P_{gg}(z)=C_A\left[rac{z}{1-z}+rac{1-z}{z}+z(1-z)
ight]\,.$$

• g
ightarrow qq:

$$P_{qg}(z) = T_R \left[z^2 + (1-z)^2
ight] \, .$$

- Stopping Conditions:
 - Perturbativity for $\alpha_s[z^2(1-z)^2t] \Rightarrow z^2(1-z)^2t > \Lambda^2$.
 - Larger phase space @ higher scale:

$$rac{1}{2}-\sqrt{rac{1}{4}-rac{\Lambda}{\sqrt{t}}} < z < rac{1}{2}+\sqrt{rac{1}{4}-rac{\Lambda}{\sqrt{t}}}\,.$$

• Tends to emit one soft parton,

$$z
ightarrow 0$$
 .

• Soft parton takes less fraction of energy @ higher scale.

$$\frac{1}{2} - \sqrt{\frac{1}{4} - \frac{\Lambda}{\sqrt{t}}} < z < \frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\Lambda}{\sqrt{t}}}.$$

• Angular ordering



 $\theta_i \approx \frac{t_i}{2\alpha_i^2} = \frac{t_i}{2z_i^2\alpha_{i-1}^2}, \quad \theta_i < \theta_{i-1} \quad \Rightarrow \quad \mathbf{t_i} < (1 - \mathbf{z_{i-1}})^2 \mathbf{t_{i-1}}.$

Consider the clustering of a soft p_j with $P_{\alpha-j} \rightarrow P_{\alpha}$:

$$J_{\beta}(P_{\alpha-j}) \quad \Rightarrow \quad J_{\beta}(P_{\alpha}) \equiv \mathbf{E}_{\alpha} - \beta \frac{\mathbf{P}_{\alpha}^2}{\mathbf{E}_{\alpha}} = E_{\alpha} \left[(1-\beta) + \beta v_{\alpha}^2 \right] \,,$$

• Energy increase is *z*:

$$z = \frac{E_j}{E_\alpha} = \frac{\mathbf{E}_\alpha - E_{\alpha-j}}{E_\alpha}$$

• Jet virtuality (mass) increase is less than *z*: Virtualtiy Reconstruction:

$$P_{\alpha}^{2} = \frac{p_{j}^{2}}{z} + \frac{P_{\alpha-j}^{2}}{1-z} + z(1-z)t \approx \frac{P_{\alpha-j}^{2}}{1-z} + z(1-z)t,$$

$$\frac{1}{E_{\alpha}} \left[\frac{\mathbf{P}_{\alpha}^{2}}{\mathbf{E}_{\alpha}} - \frac{P_{\alpha-j}^{2}}{\mathbf{E}_{\alpha-j}} \right] \approx z(1-z)\frac{t}{E_{\alpha}^{2}}.$$

Generalized Georgi Algorithms

• The original jet function:

$$J_{\beta}(P_{\alpha}) \equiv E_{\alpha} - eta rac{P_{\alpha}^2}{E_{\alpha}} = E_{\alpha} \left[(1-eta) + eta v_{lpha}^2
ight] \, ,$$

- E_{α} increases fast;
- v_{α} decreases a little bit.
- For positive J_{β} ,

$$1 \geq extsf{v}_lpha^2 \geq 1 - rac{1}{eta} \equiv extsf{v}_{ extsf{min}}^2$$
 .

 $\beta \ge 1$ to ensure a positive v_{min}^2 , which is **NOT necessary**. • **Generalized** version:

$$J_{\beta}^{(n)}(P_{\alpha}) \equiv E_{\alpha}^{n} \left[(1-\beta) + \beta v_{\alpha}^{2} \right], \quad \mathbf{n} > \mathbf{0}$$

• $J_{\beta} = J_{\beta}^{(1)}$.

- As long as *n* is not too small, $J_{\beta}^{(n)}$ can increase.
- More d.o.f for kinematics.

Kinematic Properties

• Cone should't shrink, otherwise:

- Not self-consistent.
 - For 2 subjets $\mathbf{p}_{j}(E, |\mathbf{p}|, \theta) \& \mathbf{p}'_{j}(E, |\mathbf{p}|, \theta)$;
 - The 1st subjet p_j is within the cone & can be clustered;
 - Cone shrinks;
 - The 2nd subjet p'_i may be excluded;
 - vice versa
 - The result of jet clustering is procedure-dependent.
- Cannot comply with angular ordering.
- Larger cone for smaller *z*:
 - Soft emission tends to have larger openning angle.

$$heta_i pprox rac{t_i}{2lpha_i^2} = rac{t_i}{2z_i^2 lpha_{i-1}^2}$$

• Cone is bounded from above:

- To avoid mixing up different jets.
- For $e^+e^- \rightarrow jj$, cone < half sphere.

Property 1: Cone shouldn't shrink



• Inclusion cone:

$$J^{(n)}_eta(P_lpha)> \max\left\{J^{(n)}_eta(P_lpha-p_j),J^{(n)}_eta(p_j)
ight\}\,.$$

• Exclusion cone:

$$J^{(n)}_eta(P_lpha)> \max\left\{J^{(n)}_eta(P_lpha+p_j'),J^{(n)}_eta(p_j')
ight\}\,.$$

Property 1: Cone shouldn't shrink

When **expanded**, the conditions on $\theta_{in}^{(n)}(-) \& \theta_{ex}^{(n)}(+)$, $\bar{\beta} + \beta v_{\alpha}^2 > z^n \left[\bar{\beta} + \beta v_j^2\right], \qquad \bar{\beta} \equiv 1 - \beta, z_{\pm} \equiv 1 \pm z.$ $\bar{\beta} + \beta v_{\alpha}^2 > z_{\mp}^n \left[\bar{\beta} + \frac{\beta}{z_{\mp}^2} (v_{\alpha}^2 + z^2 v_j^2 \mp 2z \cos \theta v_{\alpha} v_j)\right],$

• The 1st is phase space broadening:

$$v_{lpha}^2 - v_{min}^2 > z^n \left(v_j^2 - v_{min}^2\right) \,.$$

• The **2nd**:

$$\cos \theta_{in}^{(n)} \equiv \frac{\left(1 - z_{-}^{2-n}\right) \left(v_{\alpha}^{2} - v_{min}^{2}\right) + z^{2} \left(v_{j}^{2} - v_{min}^{2}\right) + 2z v_{min}^{2}}{2z v_{\alpha} v_{j}},$$

$$\cos \theta_{ex}^{(n)} \equiv \frac{\left(z_{+}^{2-n} - 1\right) \left(v_{\alpha}^{2} - v_{min}^{2}\right) - z^{2} \left(v_{j}^{2} - v_{min}^{2}\right) + 2z v_{min}^{2}}{2z v_{\alpha} v_{j}}.$$

Property 1: Cone shouldn't shrink

• Distance between inclusion & exclusion cones:

$$\cos\theta_{in}^{(n)} - \cos\theta_{ex}^{(n)} = \frac{2 - z_{-}^{2-n} - z_{+}^{2-n}}{2zv_{\alpha}v_{j}}(v_{\alpha}^{2} - v_{min}^{2}) + \frac{z}{v_{\alpha}v_{j}}(v_{j}^{2} - v_{min}^{2})$$

To ensure self-consistency, $n \leq 2$.

• Expansion for $z \rightarrow 0$,

$$\cos \theta_{in}^{(n)} - \cos \theta_{ex}^{(n)} \approx -\frac{(2-n)(1-n)(v_{\alpha}^2 - v_{min}^2) + 2(v_j^2 - v_{min}^2)}{2v_{\alpha}v_j}z$$

Together with phase space broadening,

$$\left[\frac{1}{z^n} - \frac{(2-n)(1-n)}{2}\right](v_{\alpha}^2 - v_{min}^2) \ge 0\,,$$

which is always true for $n \ge 0$.

Property 2: Larger cone for smaller z

$$\cos \theta_{in}^{(n)} \equiv \frac{\left(1 - z_{-}^{2-n}\right) \left(v_{\alpha}^{2} - v_{min}^{2}\right) + z^{2} \left(v_{j}^{2} - v_{min}^{2}\right) + 2z v_{min}^{2}}{2z v_{\alpha} v_{j}},$$

$$\cos \theta_{ex}^{(n)} \equiv \frac{\left(z_{+}^{2-n} - 1\right) \left(v_{\alpha}^{2} - v_{min}^{2}\right) - z^{2} \left(v_{j}^{2} - v_{min}^{2}\right) + 2z v_{min}^{2}}{2z v_{\alpha} v_{j}}.$$

• Mirror Symmetry:

$$\cos \theta_{ex}^{(n)}(z) = \cos \theta_{in}^{(n)}(-z).$$

• **Property 1 = Property 2**:

$$\cos heta_{ex}^{(n)} - \cos heta_{in}^{(n)} > 0 \quad \Rightarrow \quad heta_{in}^{(n)}(z) ext{ increases with } z$$

Property 3: Cone is bounded from above

- The largest opening angle is associated with the first emission & usually the smallest z.
- In the expanded form,

$$\begin{aligned} \cos \theta_{ex}^{(n)} &\approx \frac{1}{2v_{\alpha}v_{j}} \left[(2-n)(v_{\alpha}^{2}-v_{min}^{2})+2v_{min}^{2} \right] \\ &+ \frac{1}{2v_{\alpha}v_{j}} \left[\frac{(2-n)(1-n)}{2}(v_{\alpha}^{2}-v_{min}^{2})-(v_{j}^{2}-v_{min}^{2}) \right] z \end{aligned}$$

There is an **upper limit** if $1 \le n \le 2$,

$$\begin{aligned} \cos \theta_{ex}^{(n)} &\gtrsim & \frac{1}{2v_{\alpha}v_{j}} \left[(2-n)(v_{\alpha}^{2}-v_{min}^{2})+2v_{min}^{2} \right] \\ &+ & \frac{1}{2v_{\alpha}v_{j}} \left[\frac{(2-n)(1-n)}{2}(v_{\alpha}^{2}-v_{min}^{2})-(v_{j}^{2}-v_{min}^{2}) \right] \\ && \left(\frac{v_{\alpha}^{2}-v_{min}^{2}}{v_{j}^{2}-v_{min}^{2}} \right)^{\frac{1}{n}}. \end{aligned}$$

Property 3: Cone is bounded from above

• There is an upper limit if $1 \le n \le 2$,

$$egin{aligned} \cos heta_{ex}^{(n)} &\gtrsim & rac{1}{2 v_lpha v_j} \left[(2-n) (v_lpha^2 - v_{min}^2) + 2 v_{min}^2
ight] \ &+ & rac{1}{2 v_lpha v_j} \left[rac{(2-n)(1-n)}{2} (v_lpha^2 - v_{min}^2) - (v_j^2 - v_{min}^2)
ight] \ && \left(rac{v_lpha^2 - v_{min}^2}{v_j^2 - v_{min}^2}
ight)^{rac{1}{n}}. \end{aligned}$$

• In the relativistic limit ${\bf v}_{\alpha}\approx {\bf v}_{j}\approx 1$,

$$2\sin^2\left(\frac{1}{2} heta_{ex}^{(n)}
ight)\lesssim \frac{n(5-n)}{4}\frac{1}{eta}, \quad \Rightarrow \quad oldsymbol{eta}>rac{4}{n(5-n)}$$

β & n v.s. phase space boundary/maximal cone size.
n = 1 (original), β > 1, v²_{min} > 0.
1 < n ≤ 2 (generalized), v²_{min} < 0.

Lorentz Invariance

$$J^{(n)}_{eta}(P_{lpha})\equiv E^n_{lpha}\left(1-etarac{P^2_{lpha}}{E^2_{lpha}}
ight)\,.$$

- Jet virtuality P_{α}^2 is invariant.
- Change only comes from jet energy,

$$E_{\alpha} \rightarrow \gamma_B E_{\alpha}$$

• The effect of boost on the dimensionless part can be compensated by,

$$\beta \to \gamma_B^2 \beta$$

jet by jet.

• Jet function should be redefined,

$$J^{(n)}_{\beta} o \gamma^{-n}_B J^{(n)}_{\gamma^2_B \beta}$$

to retain the jet function value.

• Jet sequence not affected!

Summary

$$J^{(n)}_{eta}(P_{lpha})\equiv E^n_{lpha}\left[(1-eta)+eta v^2_{lpha}
ight]\,.$$

- J_{β} increases when clustering:
 - E_{α} increases due to energy conservation;
 - Jet virtuality P_{α}^2 doesn't increase that much.
- Not only pair-wise, but also can be defined globally.
- Cone implemented implicitly:

$$J_{\beta}^{(n)}(P_{\alpha}+p_{j}) = (E_{\alpha}+E_{j})^{n} \left[\bar{\beta}+\beta \frac{|\mathbf{P}_{\alpha}|^{2}+2|\mathbf{P}_{\alpha}||\mathbf{p}_{j}|\cos\theta+|\mathbf{p}_{j}|^{2}}{(E_{\alpha}+E_{j})^{2}}\right]$$

- Kinematic Properties:
 - Cone shouldn't shrink;
 - Larger cone for smaller *z*;
 - Cone is bounded from above.
- Parameter space: $1 \le n \le 2$, $\beta > 4/n(5-n)$.
- Lorentz invariance.

Thank You!