

# Beam Parameters and Physics Reach

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115th ILC@DESY General Project Meeting , DESY, Nov 2014



# Outline

- 1 Machine-design and Delivered beams
- 2 Beam conditions and the Detector
- 3 Physics implications
  - $\tilde{\tau}$  in SPS1' or STCx
  - Higgs
  - Polarisation and Near Degenerate  $\tilde{e}$
- 4 Conclusions

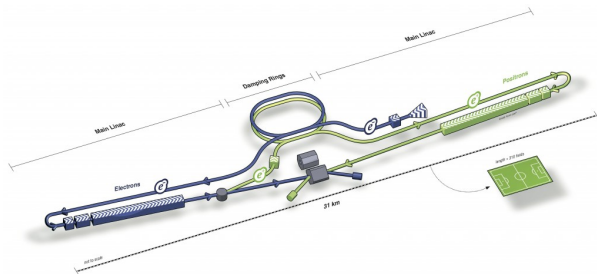
# The ideal linear collider

- Exactly known initial  $e^+ e^-$  state.
- Fully polarised beams.
- As many events that you need at any  $E_{CMS}$  .
- Pure electron/positron beams.
- No background from the machine.
- No  $\gamma\gamma$  background ...

# The real linear collider

- Beam energy has both initial and beam-beam induced **spread**.
- **Low positron polarisation** ( $\sim 30\%$ ), and  $< 100\%$  electron polarisation.
- **Limited** luminosity.
- **Mixed** lepton and photon beams.
- **Huge number** of low energy background particles from the machine.
- $\gamma\gamma$  background exists ...
- There is **ISR** ...

# Elements of the real collider



We need **electrons and positrons**:

- Electron source.
- Positron source.

We need **well defined** beams:

- Damping system.

We need **high energy**:

- Main linac.

We need to get the **beams to the detectors**:

- Beam delivery system.
- Final focus.

# Elements of the real collider: Sources

- Electron source.
  - **Polarised laser** shining on photo-cathodes, specially designed to yield polarised electrons. Collect and pre-accelerate, then send to damping system.
- Positron source
  - High ( $> 150$  GeV) energy electron beam passes an helical undulator acting as a FEL, to produce high intensity, polarisation and energy ( $\sim 10$  MeV) photons. These hit a rotating target to produce  $e^+ e^-$  e-pairs. Positrons are collected, pre-accelerated and sent to damping.
  - Electrons are from the main beam  $\rightarrow$  an additional energy dispersion, due to the synchrotron radiation losses in the undulator.

# Elements of the real collider: Damping system

- From the sources, the **dispersion** both in angle and energy are **way to big**.
- Send beams (now at  $\sim 5$  GeV) to rings where they pass "wigglers" making them **cool off** by synchrotron radiation.
- **Kick out bunches**, one-by-one, every  $\sim 100$  ns to make the bunch train. Bunches are **separated by a few ns**, given by (circumference of damping ring)/(number of bunches).
- All this must be done in the **200 ms between bunch trains**.
- The damping rings are at the centre of the complex: need to **transport the bunches  $\sim 15$  km** to the start of the main linac after damping.

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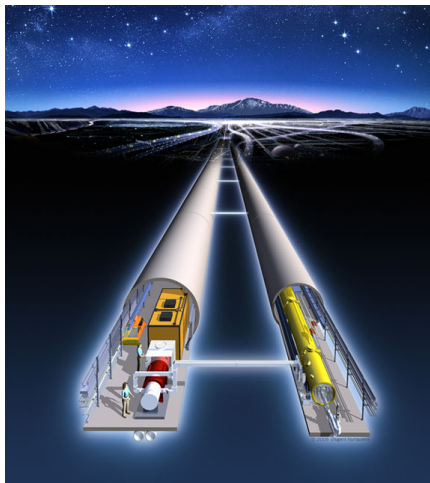
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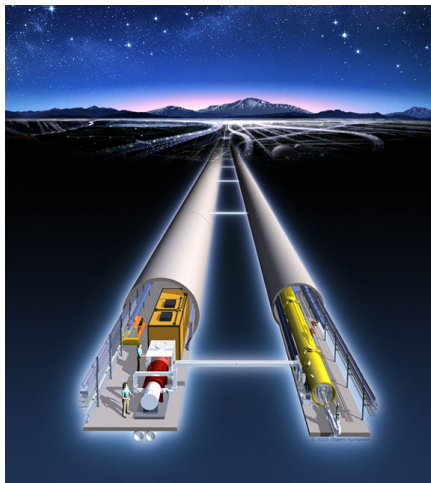
# Elements of the real collider: Main linac

- Super-conducting RF cavities, 31.5 MV/m gradient, 9 cells.
- One RF unit = 3 cryo-units, 2 with 9 cavities, one with 8 + a focusing quadropole.
- 278 of these in the positron linac, 282 in the electron one (more, since energy is lost in the undulator !)  $\Rightarrow$  132 912 cells...
- The linac needs power: Klystrons all along. How many particles one can get/time depends on how many of these one installs.



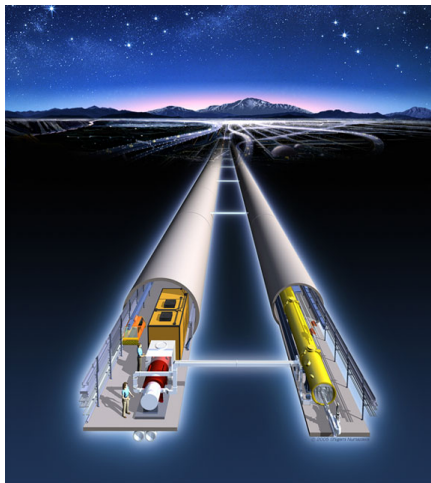
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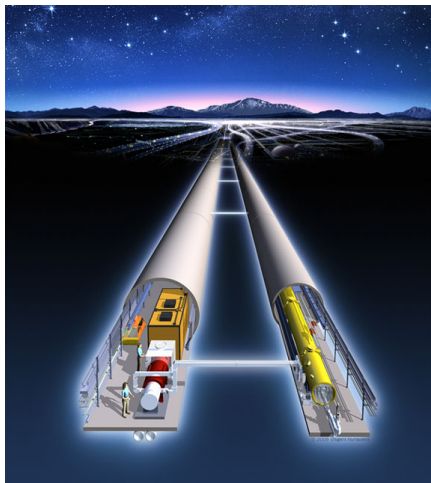
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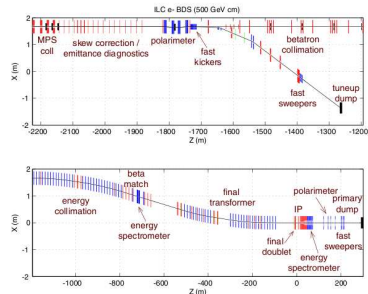
# Elements of the real collider: BDS and final focus

## BDS:

- Last 2 km.
- Monitor and measure beam.
- Clean up beam-halo etc.
- Protect detectors.
- Anything the beam hits here will give secondaries ( $E_{beam}$  is up to 500 GeV!), that might hit the detectors.

## Final focus:

- Last 20 m.
- Focuses beams to few 100 nm horizontally, and few nm vertically.



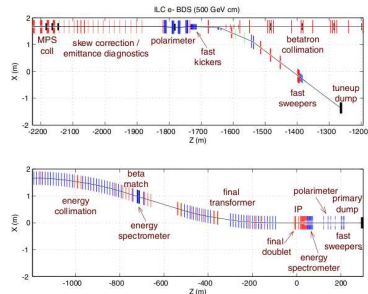
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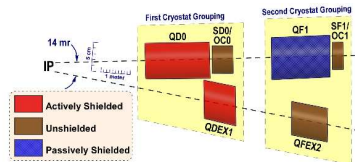




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# Elements of the real collider: Beam-strahlung

Due to the very **strongly focused beams**, the fields (both E and B) has a large bending power on the other beam. Consequences:

- Primary beam is focused by the other beam.
- Strong bending  $\rightarrow$  much synchrotron radiation. Widens the distribution of the primary  $e^\pm$  energy.
- Photons
  - ... get Compton-backscattered  $\rightarrow$  photon component of beam, long tail to lower energies for the  $e^\pm$ .
  - ... interact with photons (synchrotron ones, or virtual ones) in the other beam  $\rightarrow e^\pm$ -pairs.
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The wrong-sign  $e^\pm$ :s gets a maximum kick if they are at the outer edge of the beam. The kick is independent of the (longitudinal) momentum of the particle.

$p_T$  and  $\theta$  anti-correlates, and accumulate at the edge.

To study the effect, also draw the detector in these coordinates:

Place it at the  $p_T$ - $\theta$  corresponding to the  $p_T$  and  $\theta$  a particle should have to turn back at the radius and  $z$  of the detector.

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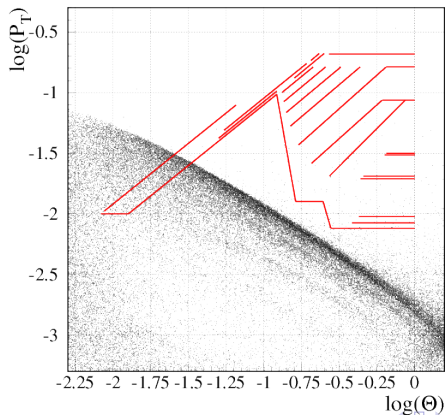
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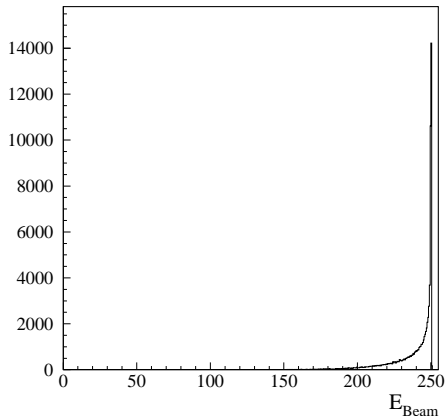
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Example: Pairs in ILD, RDR nominal parameters. Generated with GuineaPig. **124000 particles/BX.**



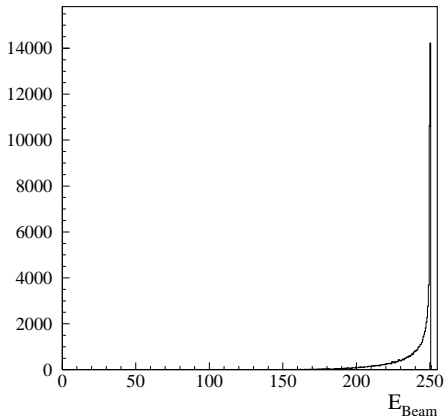
# Elements of the real collider: Beam-spectrum

- On the large scale:
  - Beam-strahlung gives a long tail to lower energies
- Zoom to the peak
  - Noticeable spread in energy.
  - Can be comparable to detector resolution.
  - Note that  $E_{e^-}$  is wider: the undulator.



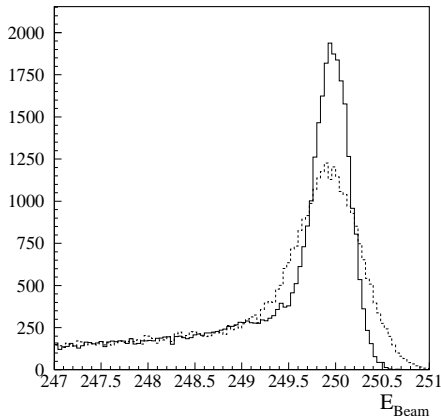
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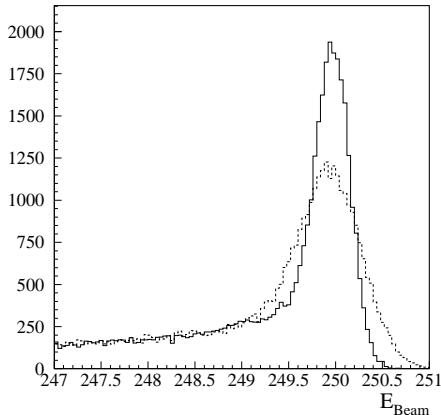
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# Elements of the real collider: Luminosity

Luminosity(L) = Density of particles that pass each other per time-unit.  
 Number of interactions/time=L×cross-section.

- $L = N^2/(t \times A)$
- $N^2/t = (\text{particles in bunch})^2 \times (\text{number of bunches in train}) \times (\text{number of trains per second (= "rep rate")}) = n^2 N_{bunch} f_{rep}$
- RF-power ( $P_{RF}$ )= $E_{cm}(nN_{bunch}f_{rep}) \times \eta$  ( $\eta$ = efficiency of the transfer from the RF-system → beam)

So:  $L \propto P_{RF} n / A E_{cm}$

- $A = \text{cross-section of beams at IP} \propto \sigma_x \times \sigma_y$
- $\sigma \propto \sqrt{\epsilon\beta} = \sqrt{\epsilon_{norm}\beta/\gamma}$ .  $\epsilon$ =emittance,  $\epsilon_{norm}$ =normalised emittance=what the damping system achieves.  $\beta$ =focusing-power of the final focus system.

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- Relative energy-loss due to beam-strahlung:

$$\delta_{BS} \propto (E_{cm}/\sigma_z) \times (n^2/(\sigma_x + \sigma_y)^2)$$

- To reduce beam-strahlung: keep  $\sigma_x + \sigma_y$  big.
- To get high L :  $\sigma_x \times \sigma_y$  small.

Need a flat beam!

$$\sigma_y \ll \sigma_x \rightarrow \delta_{BS} \propto (E_{cm}/\sigma_z) \times (n^2/\sigma_x^2)$$

$$n/\sigma_x \propto \sqrt{\delta_{BS}\sigma_z/E_{cm}}$$

So:

$$L \propto P_{RF} \times \sqrt{\delta_{BS}\sigma_z}/(\sigma_y E_{cm}^{3/2})$$

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Need a flat beam!

$$\sigma_y \ll \sigma_x \rightarrow \delta_{BS} \propto (E_{cm}/\sigma_z) \times (n^2/\sigma_x^2)$$

$$n/\sigma_x \propto \sqrt{\delta_{BS}\sigma_z/E_{cm}}$$

So:

$$L \propto P_{RF} \times \sqrt{\delta_{BS}\sigma_z}/(\sigma_y E_{cm}^{3/2})$$

or

$$L \propto (nN_{bunch}f_{rep}) \times \sqrt{\delta_{BS}\sigma_z}/(\sigma_y E_{cm}^{1/2})$$

or

$$L \propto (n^2 N_{bunch} f_{rep})/(\sigma_x \sigma_y) \propto (n^2 N_{bunch} f_{rep} E_{cm})/(\epsilon_{norm} \beta)$$



# Elements of the real collider: Luminosity

- Relative energy-loss due to beam-strahlung:  
 $\delta_{BS} \propto (E_{cm}/\sigma_z) \times (n^2/(\sigma_x + \sigma_y)^2)$
- To reduce beam-strahlung: keep  $\sigma_x + \sigma_y$  big.
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# Beam conditions and the Detector

## RDR, SB2009 and TDR

During the SB2009 exercise, a number of studies were done comparing physics and detector implications of two different beam-parameter sets.

- The RDR parameters.
- The SB2009 parameters:
  - Half the number of bunches.
  - Recuperate the luminosity by more aggressive focusing.
  - In addition: modified positron-source, giving low luminosity below  $E_{cms} = 300$  GeV

The TDR largely follows the SB2009, except that the loss of luminosity at low  $E_{cms}$  has been mitigated.

Here I will present some of the observations from the SB2009 exercise.

# RDR and SB2009

- Twice as much beam-strahlung:
  - more overlaid tracks (real or fake)
  - Twice as much energy in BeamCal

## At 500 GeV

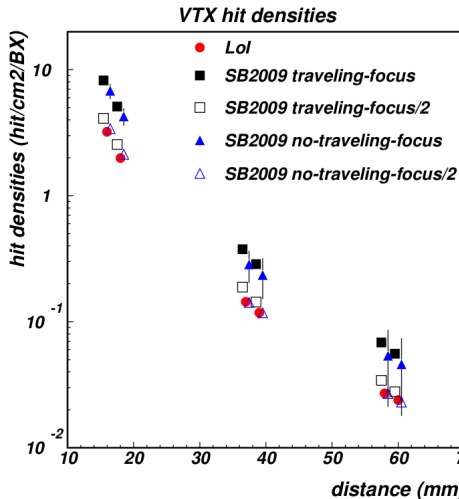
- Total luminosity unchanged RDR  $\rightarrow$  SB2009 w TF, but reduced by %25 for SB2009 w/o TF.
- $P(e^+)$  goes from 33 % to 22 %.
- Incoming energy-spread grows from 0.16 to 0.21
- Luminosity within 1 % of nominal reduced from 0.83 to 0.72.

## At 250 GeV

- Lumi reduced by a factor three.

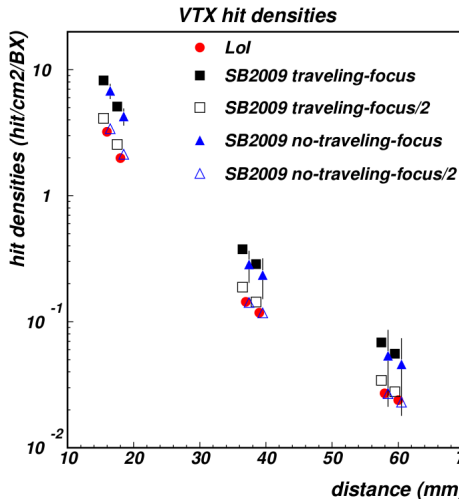
# Beam-strahlung: Hits in Vertex detector

- Full simulation (Mokka), with crossing-angle and anti-DID field.
- The ILD VTX integrates over a certain time-window  $\rightarrow$  Half as many BX:es overlaid in SB2009 wrt. RDR.
- $\Rightarrow$  The net effect is small.
- The SiD VTX time-stamps every BX  $\rightarrow$  Twice as many background hits with SB2009.



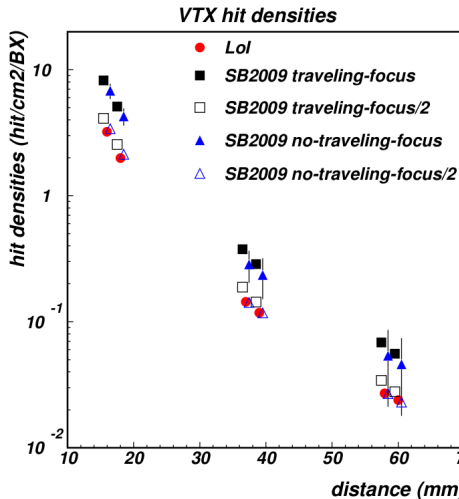
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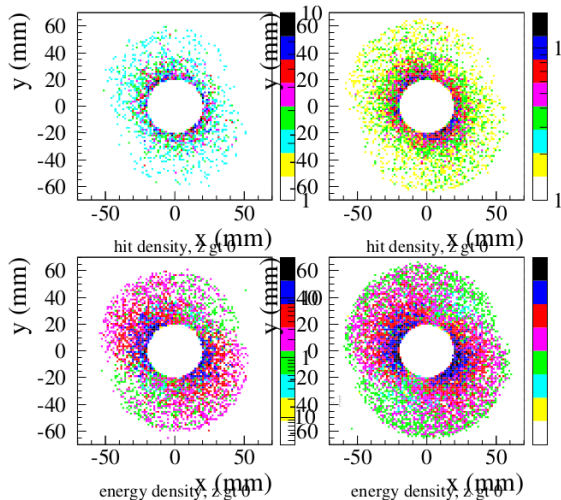
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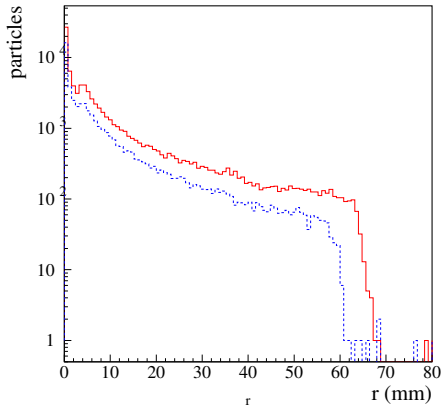
# Beam-strahlung: BeamCal

- Only GP, but with crossing-angle and anti-DID.
- Both hit-densities (top) and energy-density (bottom) matters.
- The issue: can one still see a  $\approx 250$  GeV electron from a  $\gamma\gamma$  process over the pairs-background?



# Beam-strahlung: BeamCal

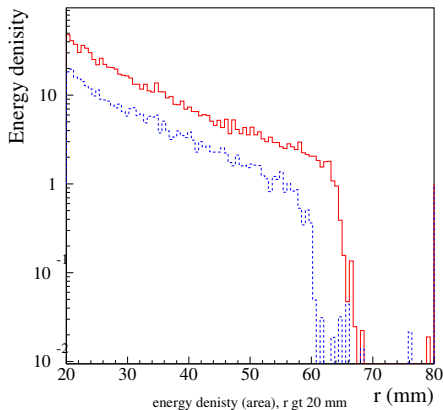
- Distribution of particle energy for  $r > 20$  mm.
- Total energy in BeamCal per BX: 24 TeV for SB2009TF, 10 TeV for RDR nom.
- Number of particles per BX 11500 for SB2009TF, 5400 for RDR nom.
- Energy density vs Radius. SB2009TF has about twice at any given radius, and extends 5 mm further.
- All the relevant numbers double





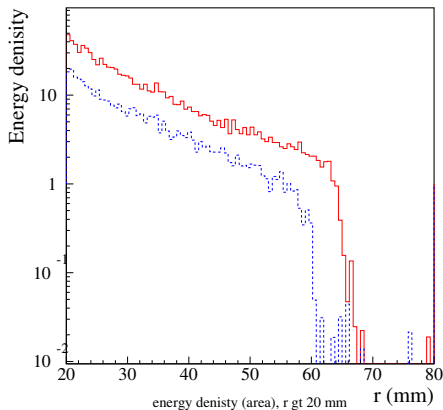
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# SB2009 and physics: $\tilde{\tau}$ in SPS1' or STCx

Small mass difference  $\tilde{\tau}$ -LSP ( $\sim 10$  GeV)  $\Rightarrow$  soft  $\tau$ :s.

Potential effects:

- Decrease of  $P(e^+)$ : More background, less-signal for  $\tilde{\tau}_1$
- Incoming energy-spread grows: end-point blurred.
- Luminosity within 1 % of nominal reduced: lower signal.
- Twice as much beam-strahlung:
  - more overlaid tracks (real or fake): Destroys  $\tau$  topology.
  - Twice as much energy in BeamCal: More  $\gamma\gamma$
- Higher probability for a  $\gamma\gamma$  event in the same BX as the physics event.
- (Total luminosity decrease for SB2009 w/o TF.)

SB2009 and physics:  $\tilde{\tau}$  - finally selected events

	Events for end-point analysis					
case	$\tilde{\tau}_1$			$\tilde{\tau}_2$		
	SM	SUSY	signal	SM	SUSY	signal
RDR	317	998	10466	1518	241	1983
SB09(TF)	814	956	8410	1346	223	1555
SB09(noTF)	611	717	6308	1009	167	1166
	Events for cross-section analysis					
	$\tilde{\tau}_1$			$\tilde{\tau}_2$		
	SM	SUSY	signal	SM	SUSY	signal
RDR	17.6	47.7	2377	1362	33.7	1775
SB09(TF)	17.6	45.7	1784	1194	32.4	1366
SB09(noTF)	13.2	34.3	1337	895	24.3	1025

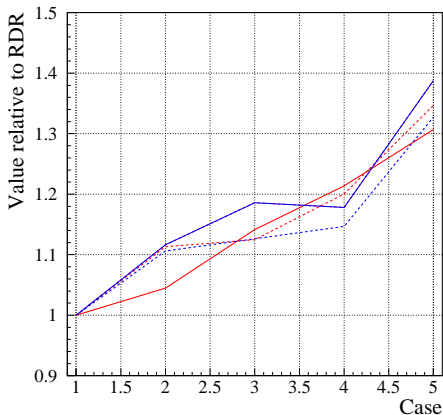
SB2009 and physics:  $\tilde{\tau}$  - SUSY model parameters

## Errors on end-point (GeV)

case	#	$\tilde{\tau}_1$	$\tilde{\tau}_2$
RDR	1	0.129	1.83
+SB bck	2	0.144	2.02
+SB ppol	3	0.153	2.06
+SB spect	4	0.152	2.10
+SB noTF	5	0.179	2.42

## Errors on cross-section (%)

case	#	$\tilde{\tau}_1$	$\tilde{\tau}_2$
RDR	1	2.90	4.24
+SB bck	2	3.03	4.72
+SB ppol	3	3.31	4.77
+SB spect	4	3.52	5.09
+SB noTF	5	3.79	5.71



Red: cross-section, Blue:  
end-point, Solid :  $\tilde{\tau}_1$  , Dashed:  $\tilde{\tau}_2$ .

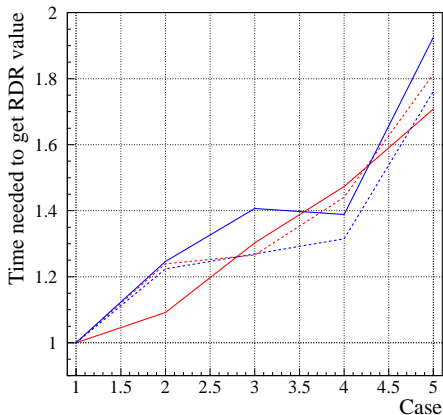
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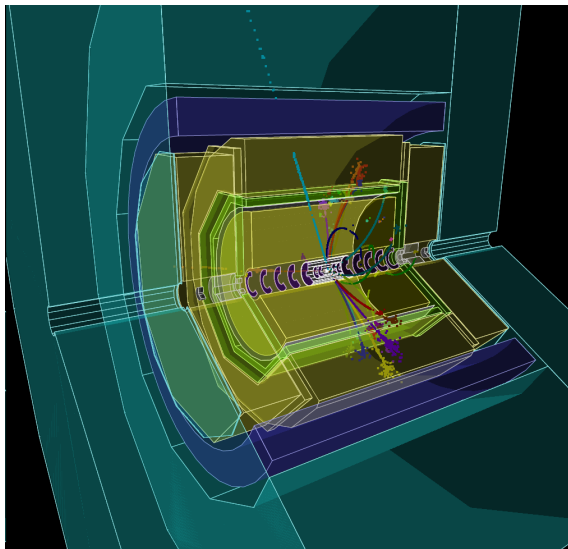
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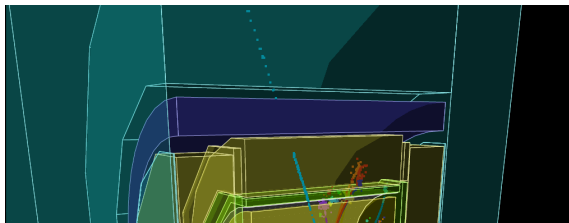


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# SB2009 and physics: SM Higgs at 120 GeV



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## Topic: Model independent Higgs mass

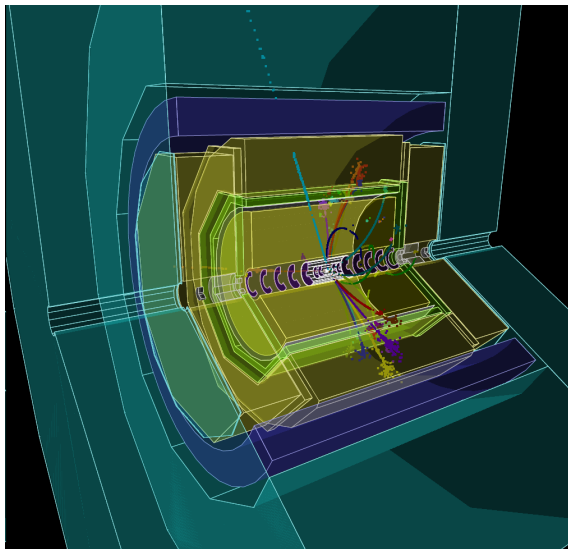
Recoil mass measurement:

- Only reconstruct the  $Z \rightarrow \mu\mu$
- Using E & p conservation the Higgs mass can be measured from the recoil **independent of the decay mode**





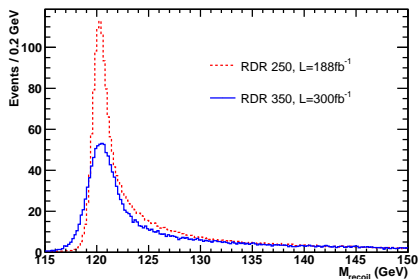
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Compare recoil-mass peak obtainable with the same running-time at 250 or 350 GeV, for...

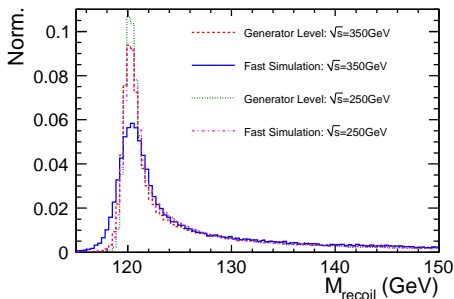
- RDR
- Peak is broader at 350 GeV due to detector-resolution. Higher momentum gives higher error !



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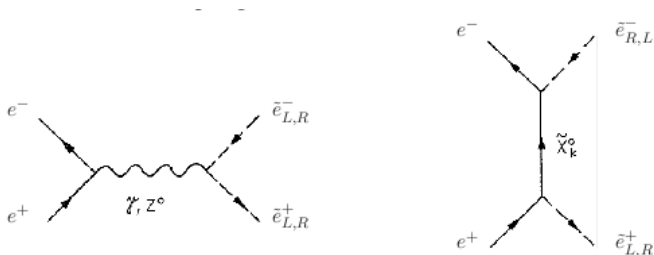
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# SB2009 and physics: Polarisation and Near Degenerate $\tilde{e}$

Super-symmetry associates scalars to chiral (anti)fermions

$$e_{L,R}^- \leftrightarrow \tilde{e}_{L,R}^- \quad \text{and} \quad e_{L,R}^+ \leftrightarrow \tilde{e}_{R,L}^+ \quad (1)$$



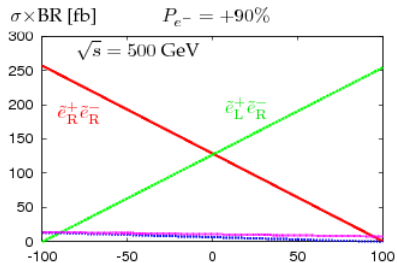
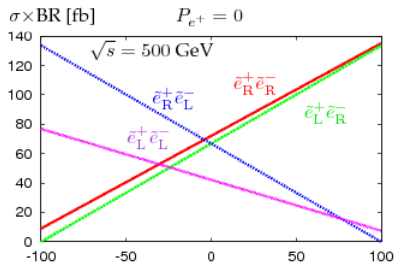
What if  $M_{\tilde{e}_L} \approx M_{\tilde{e}_R}$ , so that thresholds can't separate  $e^+e^- \rightarrow \tilde{e}_L\tilde{e}_L, \tilde{e}_R\tilde{e}_R$  and  $\tilde{e}_R\tilde{e}_L$ ?

# SB2009 and physics: Polarisation and Near Degenerate $\tilde{e}$

Model: SPS1a' like, but:

$M_{\tilde{e}_L} = 200$  GeV and  $M_{\tilde{e}_R} = 195$  GeV. Both decay 100 % to  $\tilde{\chi}_1^0 e$ .

Even with  $P_{e^-} \geq +90\%$ , one can't disentangle the pairs  $\tilde{e}_L^+ \tilde{e}_R^-$  and  $\tilde{e}_R^+ \tilde{e}_R^-$ : Ratio of the cross sections  $\approx$  constant.



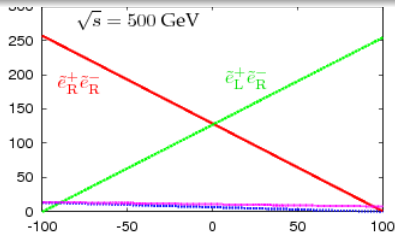
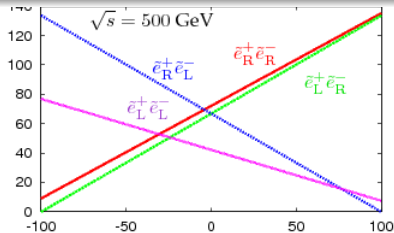
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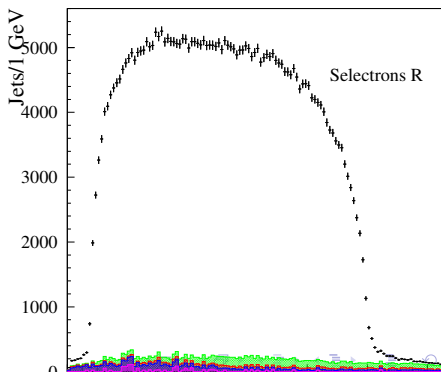
**Polarised positrons a must !**



# SB2009 and physics: Polarisation and Near Degenerate $\tilde{e}$

Background and efficiency from Full-sim SPS1a' sample, kinematics from Whizard simulation of the model.

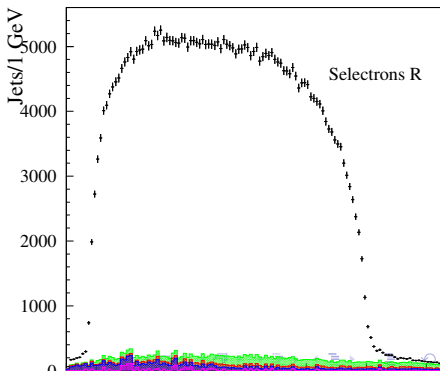
- The  $\tilde{e}$  signal was extracted from the **same sample** as was used for the  $\tilde{\tau}$  study, using the **same cuts** except
  - Demand exactly two well identified electrons.
  - Reverse the  $\tilde{\tau}$  anti-SUSY background cut
  - Some cuts could be loosened
- Almost **background-free** !



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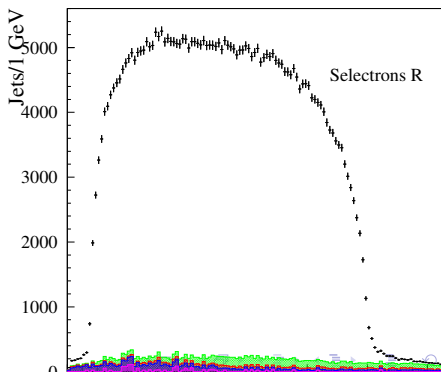




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# SB2009 and physics: Polarisation and Near Degenerate $\tilde{e}$

The handle:

Opposite polarisation beams produces  $\tilde{e}$ :s in **both** s- and t-channel.  
 Same polarisation produces  $\tilde{e}$ :s in t-channel **only**  $\Rightarrow$

## Modification of $\Theta$ distribution with changed positron polarisation

However, the effect is small since t-channel always dominates !  $\tilde{e}$ :s are heavy (and are scalars)  $\Rightarrow$  t- and s- channel kinematic distributions of the electrons are not very different.

Need to reconstruct the  $\tilde{e}$  direction:

- 8 Unknown  $\tilde{\chi}_1^0$  momentum components
- Assume  $M_{\tilde{e}}$  and  $M_{\tilde{\chi}_1^0}$  known  $\rightarrow$
- 8 constraints (E and p conservation, 4 mass-relations)

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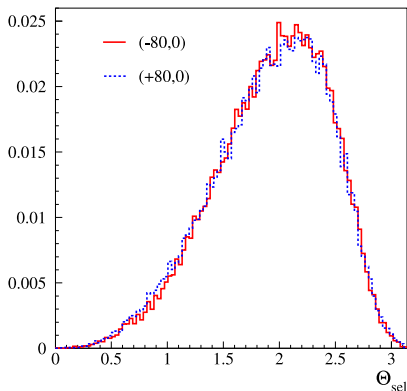
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Analyse assuming  $100 \text{ fb}^{-1}$  for each of the polarisations configurations.

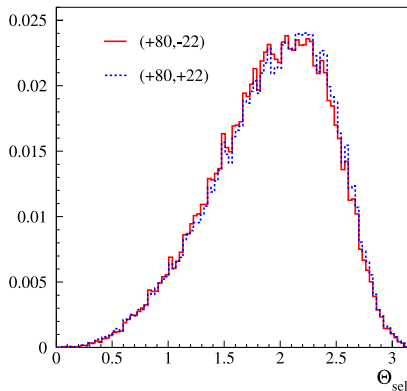
- For  $P(e^-) = \pm 80\%$   $P(e^+) = 0$  and then ..
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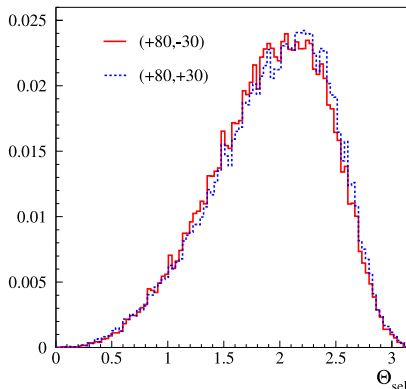
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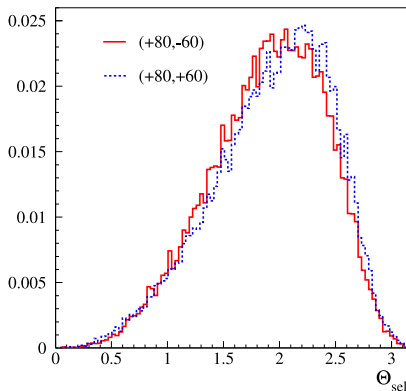
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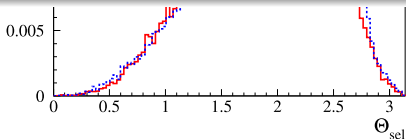


# SB2009 and physics: Polarisation and Near Degenerate $\tilde{\epsilon}$

Analyse assuming  $100 \text{ fb}^{-1}$  for each of the polarisations configurations.

- For  $P(e^-) = +80\%$   $P(e^+) = 0$

$ P(e^+) $ (%)	significance of shift( $\sigma$ )	Title of paper
22	2.4	"Limit on ..."
30	3.5	"Evidence for ..."
60	6.6	"Observation of ..."



# Conclusions...

## Machine-parameters:

- Depending on what is built in to the machine ( $P_{RF}$ ,  $f_{rep}$ ,  $N_{bunch}$ ,  $\delta_{BS}...$ ), luminosity scales differently with  $E_{cm}$ , but it's never constant.
- Different machine setups give different gives different luminosity scaling, different polarisation scaling, different energy within 1 % to nominal, different spread in  $E_{beam}$ .

## Lessons from $\tilde{\tau}$ :s:

- For “fragile” signals, beam-background influences signal directly.
- For any ‘low  $\Delta(M)$  ( $< 10$  GeV) signal, beam-background should be taken into consideration when estimating  $\gamma\gamma$  background.
- RDR  $\rightarrow$  SB2009: 15-20 % degradation (end-point and cross-section,  $\tilde{\tau}_1$  and  $\tilde{\tau}_2$ ).
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# Conclusions

## Lessons from SM Higgs:

- Results will **not scale** with cross-section if  $E_{cm}$  changes: Detector resolution depends on energy, while beam energy spread doesn't.

## Lessons Degenerate $\tilde{e}$ :

- There are physics that **can't be done** without positron polarisation.
- In the studied case, going from 22 % to 60 % positron polarisation  $\Leftrightarrow$  **7.5 times more luminosity !**

# Conclusions

Lessons from SM Higgs:

- Results will **not scale** with cross-section if  $E_{cm}$  changes: Detector resolution depends on energy, while beam energy spread doesn't.

Lessons Degenerate  $\tilde{e}$ :

- There are physics that **can't be done** without positron polarisation.
- In the studied case, going from 22 % to 60 % positron polarisation  $\Leftrightarrow$  **7.5 times** more luminosity !

# Thank You !