

# Application of Matrix Element Method

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# recent development: matrix element method (approach the true likelihood of each event)

- @ ILC: well defined initial states, precisely measured final states, MEM perfectly fits for precision measurement
- MEM tools for full simulation released with latest ilcsoft-v01-17-06: Physsim-v00-01
- basic verification done
- principle demonstrated in first application for  $e^+e^- \rightarrow eeH$  via ZZ-fusion
- main focus now is to apply for other major Higgs measurement (today for recoil mass)

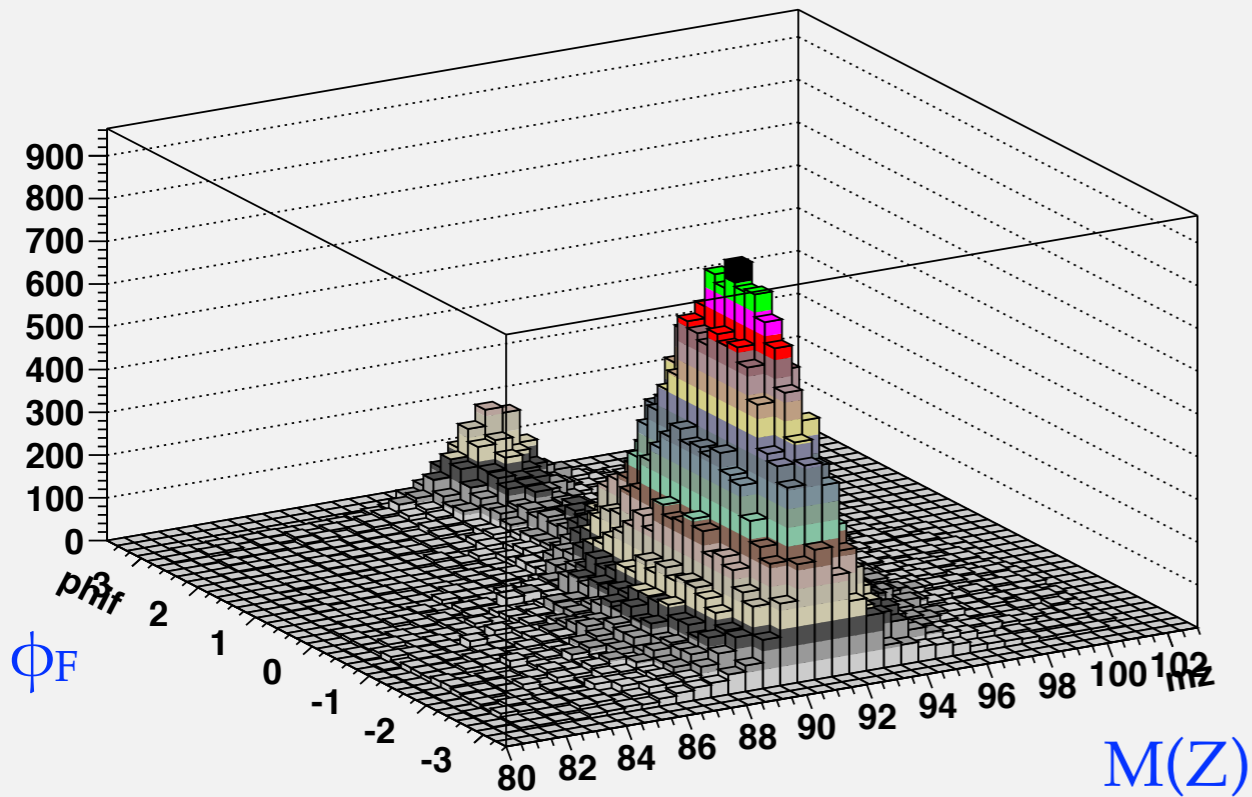
svn co <https://svnsrv.desy.de/basic/physsim/Physsim/trunk>

# reminder: verification of calculated matrix element (ZHH events)

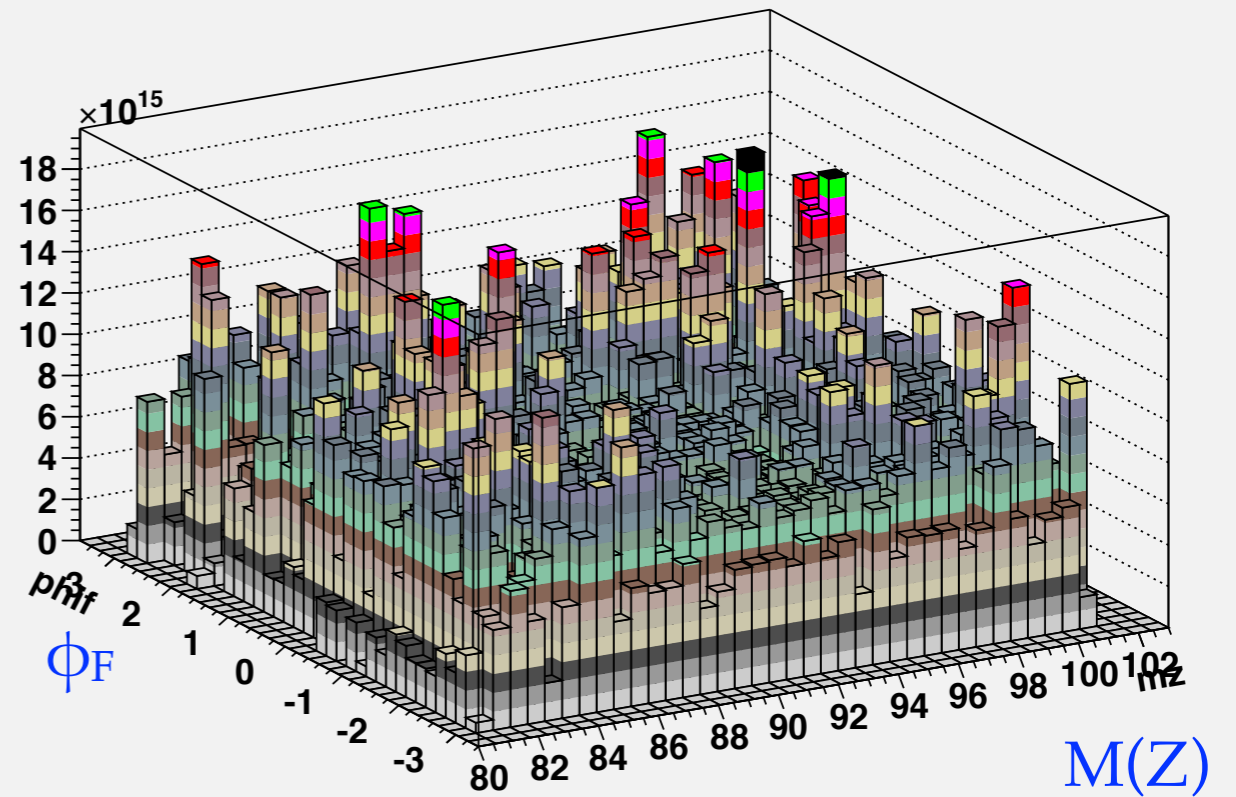
original events

weighted by  $\frac{1}{|ME|^2}$

phif:mz {(abs(mz-91)<10&&abs(phif)<3&&1./sigmall<10.E15)}

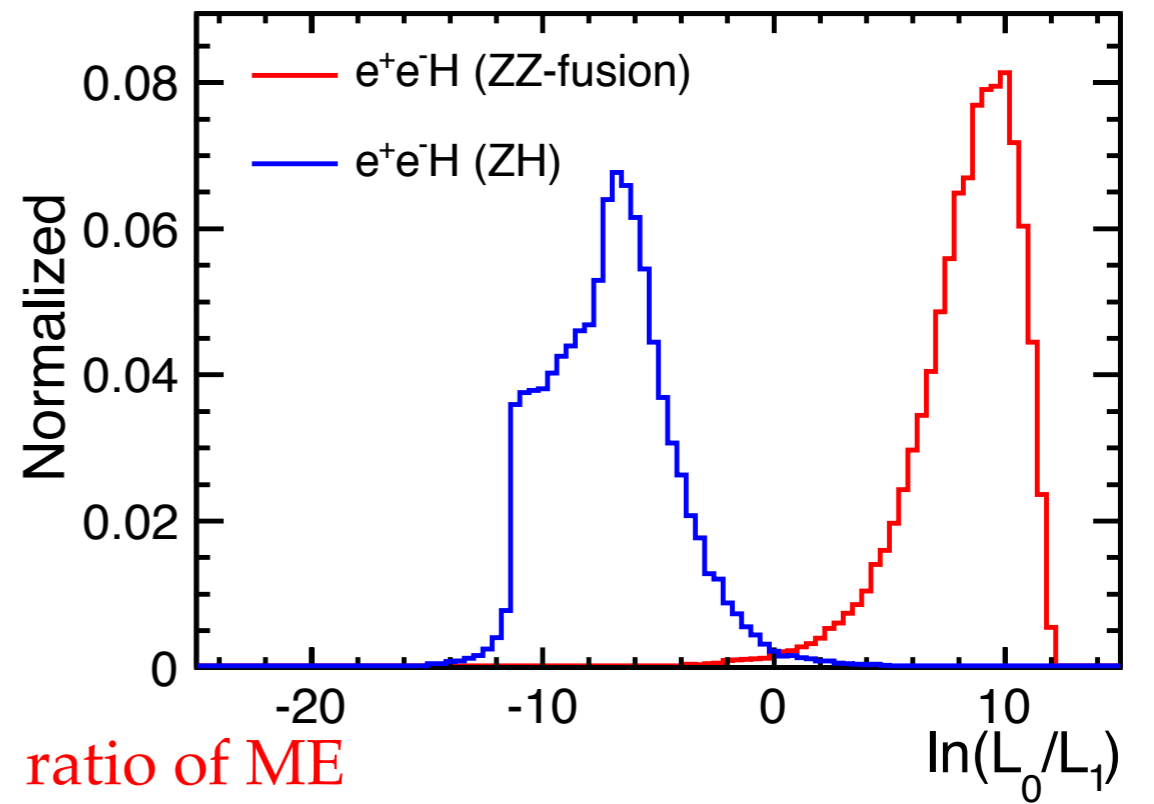
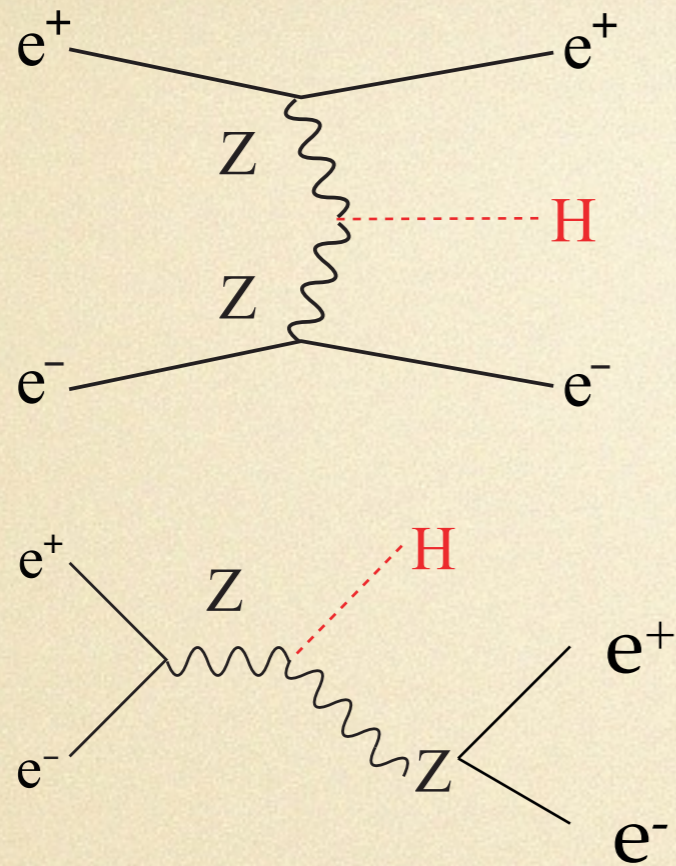


phif:mz {1./sigmall\*(abs(mz-91)<10&&abs(phif)<3&&1./sigmall<10.E15)}

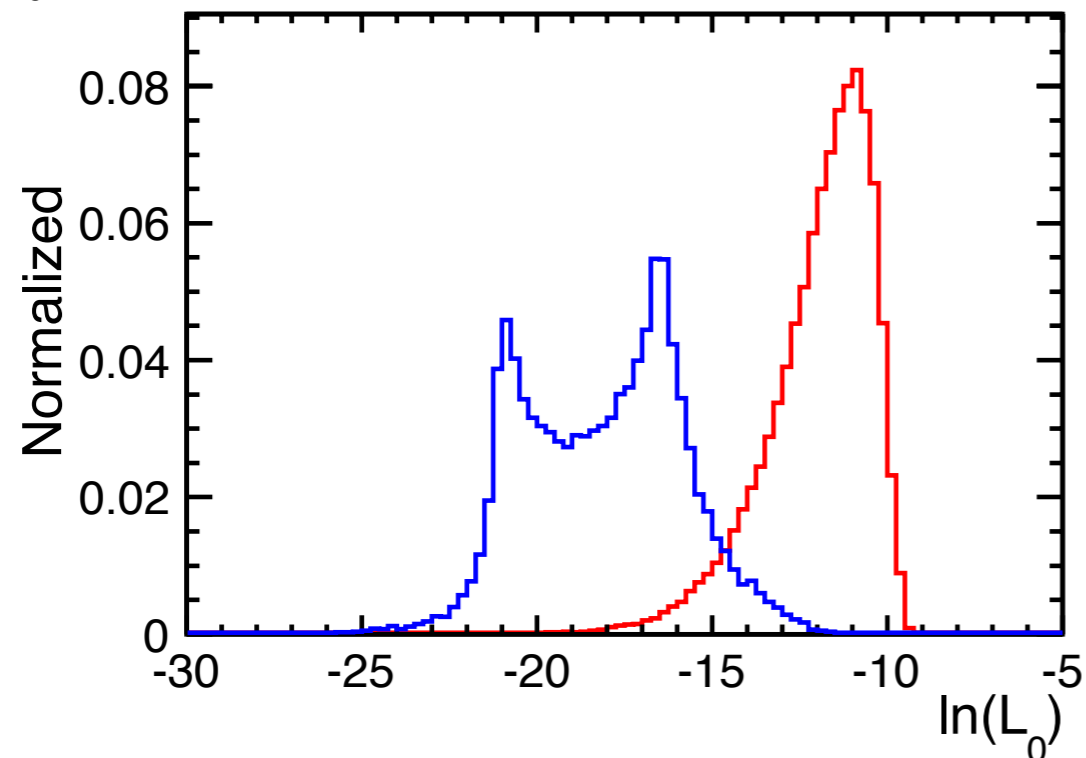


$\phi_F$ : azimuthal angle defined in Z rest frame of fermion from  $Z \rightarrow ff$

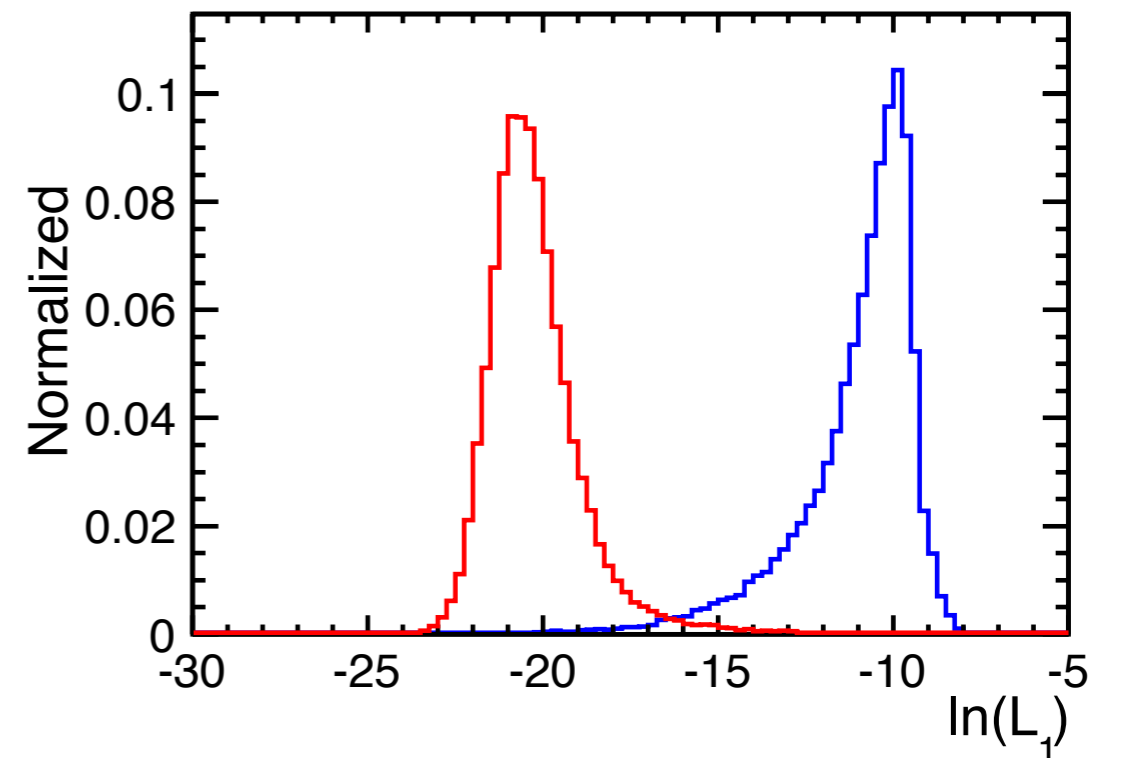
# first application: $e^+e^- \rightarrow e^+e^-H$



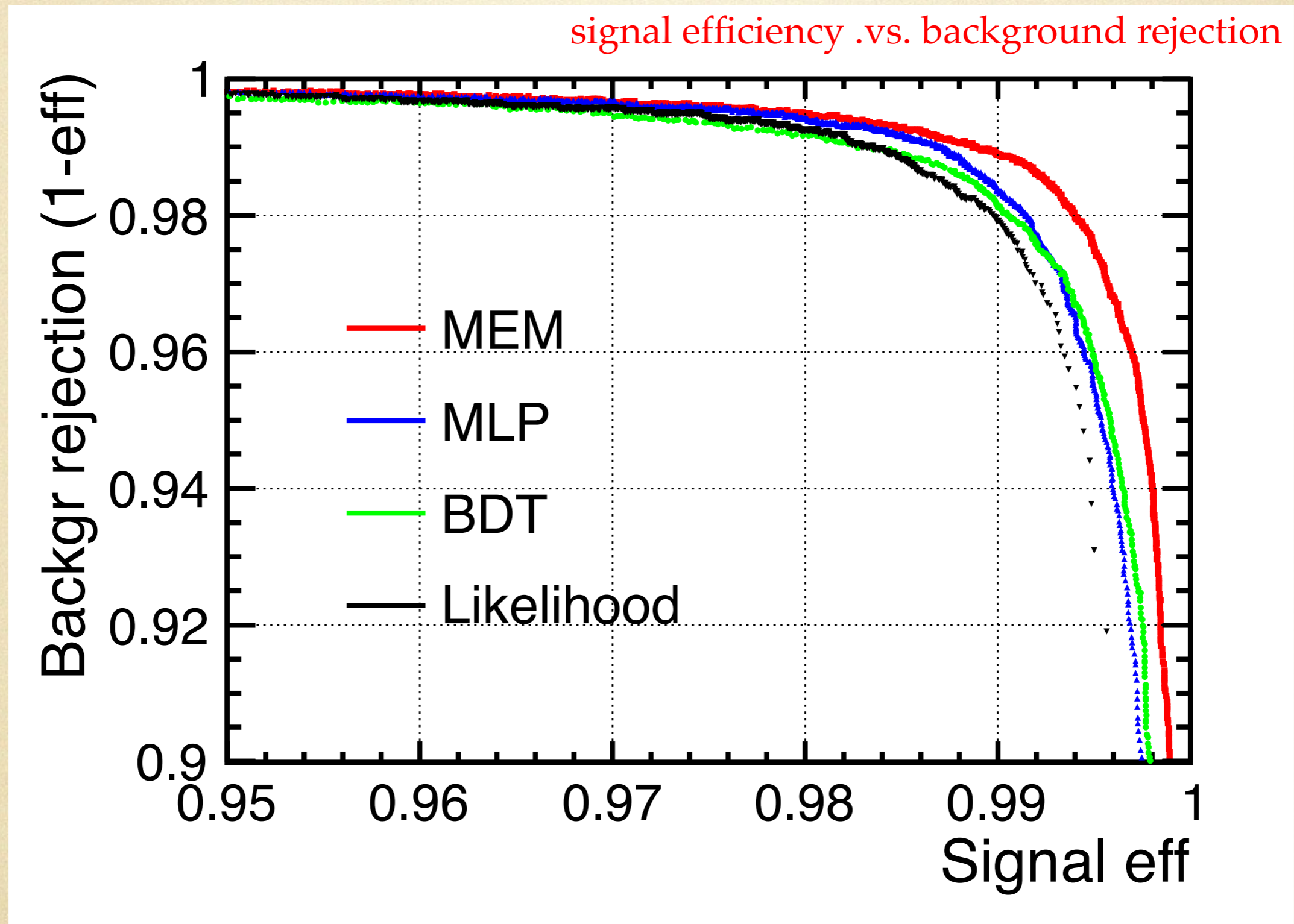
$L_0$ : ME as from ZZ-fusion



$L_1$ : ME as from ZH



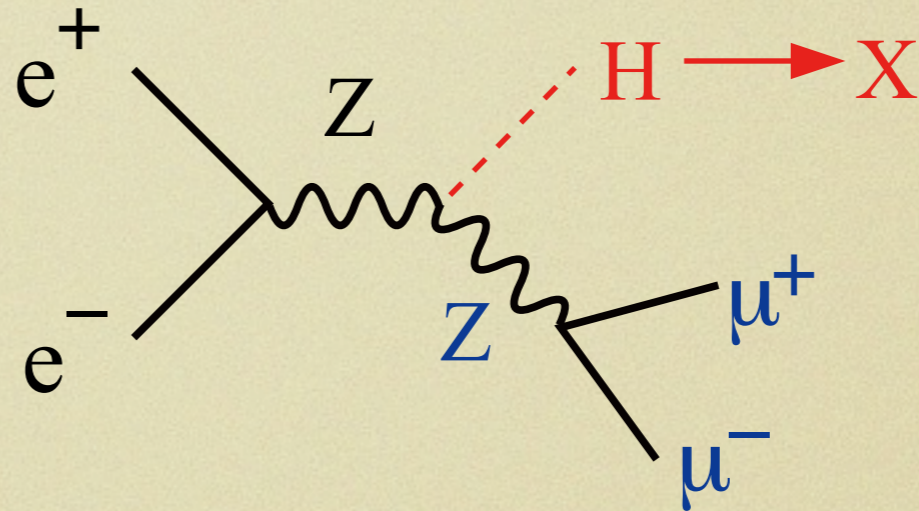
principle demonstrated: physicist .vs. statistician



signal: eeH via ZZ-fusion; background: eeH via ZH

apply for more critical measurement

recoil mass analysis

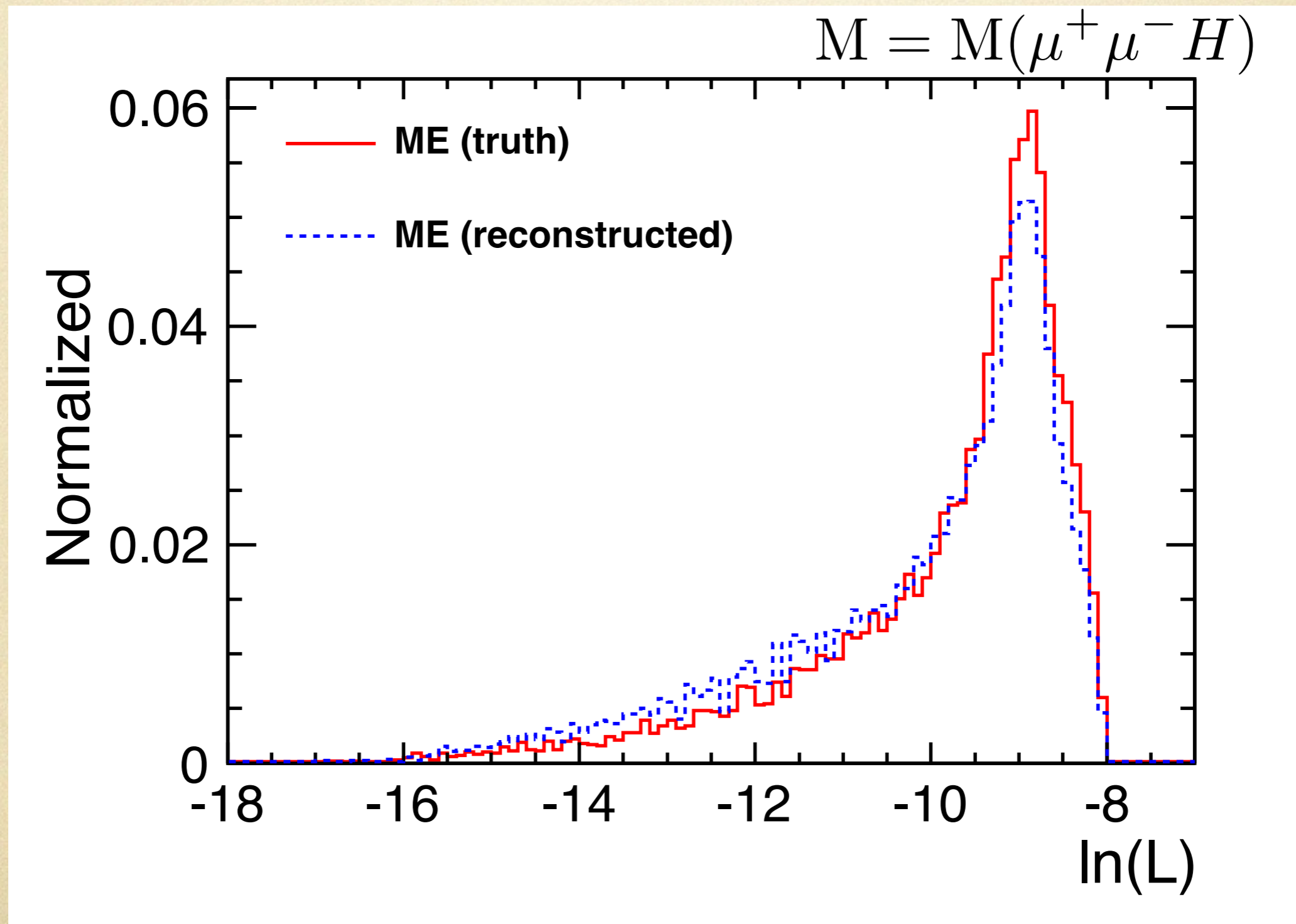


- ✓ momentum of muon are well measured
- ✓ since ECM is well defined, momentum of Higgs can be well determined as well
- ✓ therefore matrix element can be precisely calculated without assuming any Higgs decay!

$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H}$$

# matrix element: reconstructed .vs. truth

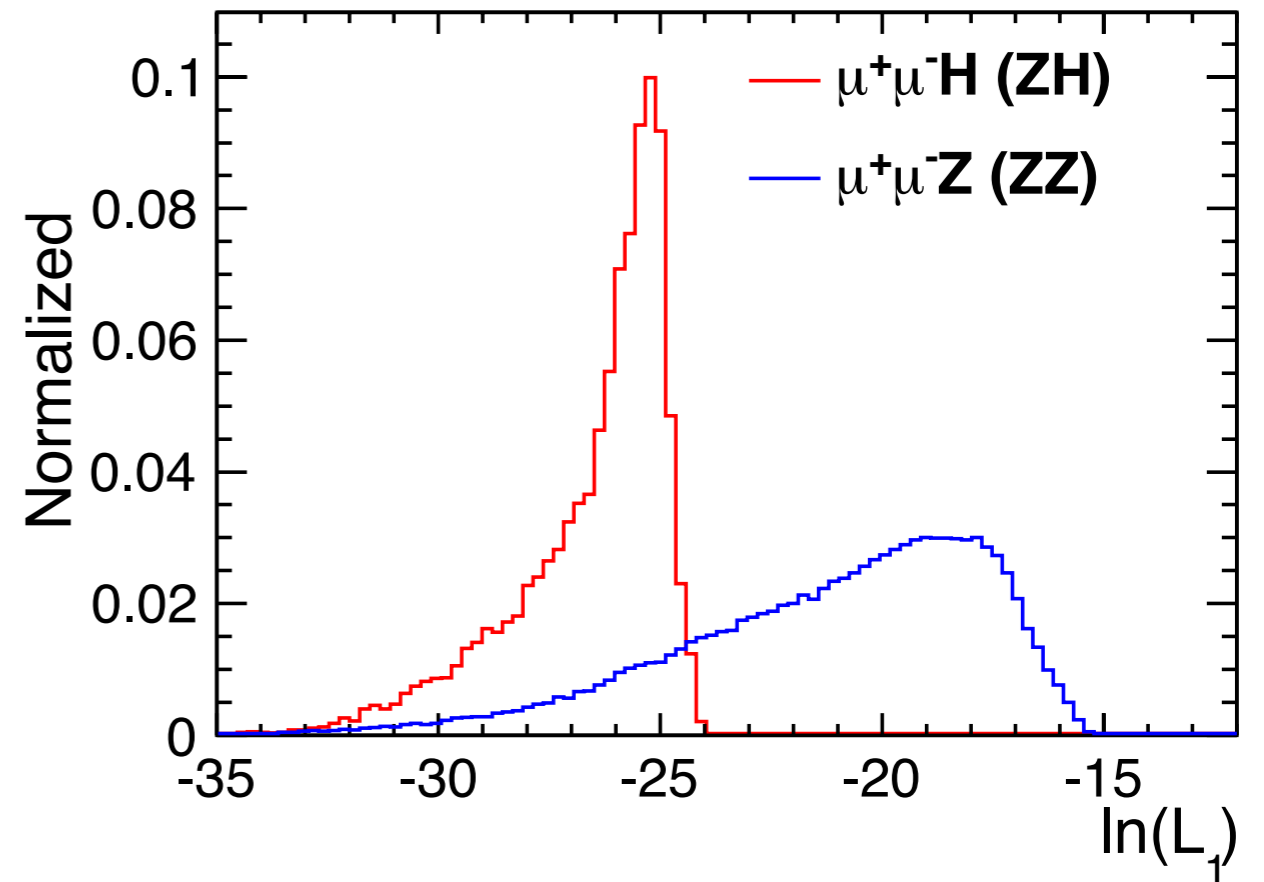
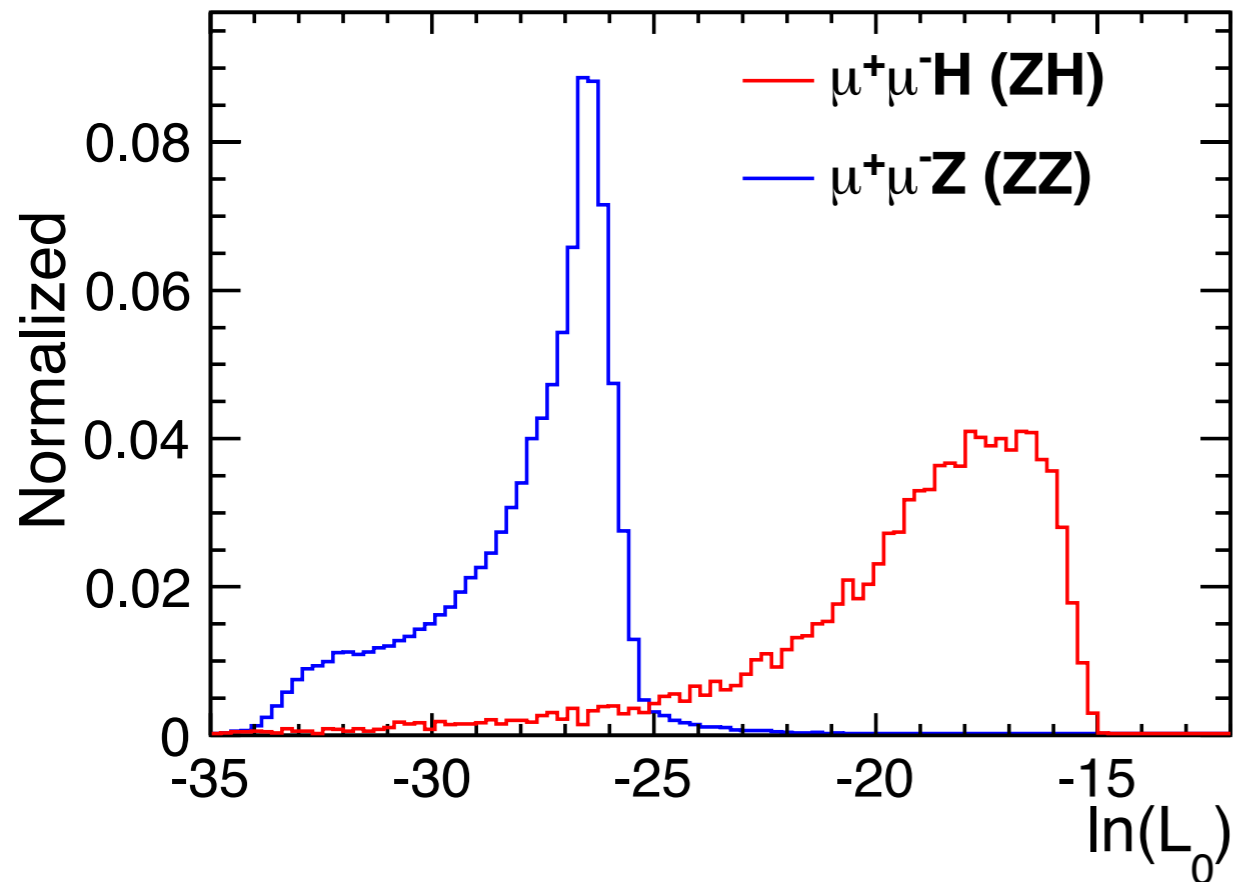
(step-by-step test: samples are generated w/o ISR and beam strahlung)



matrix element: signal .vs. background (ZZ)

$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H}$$

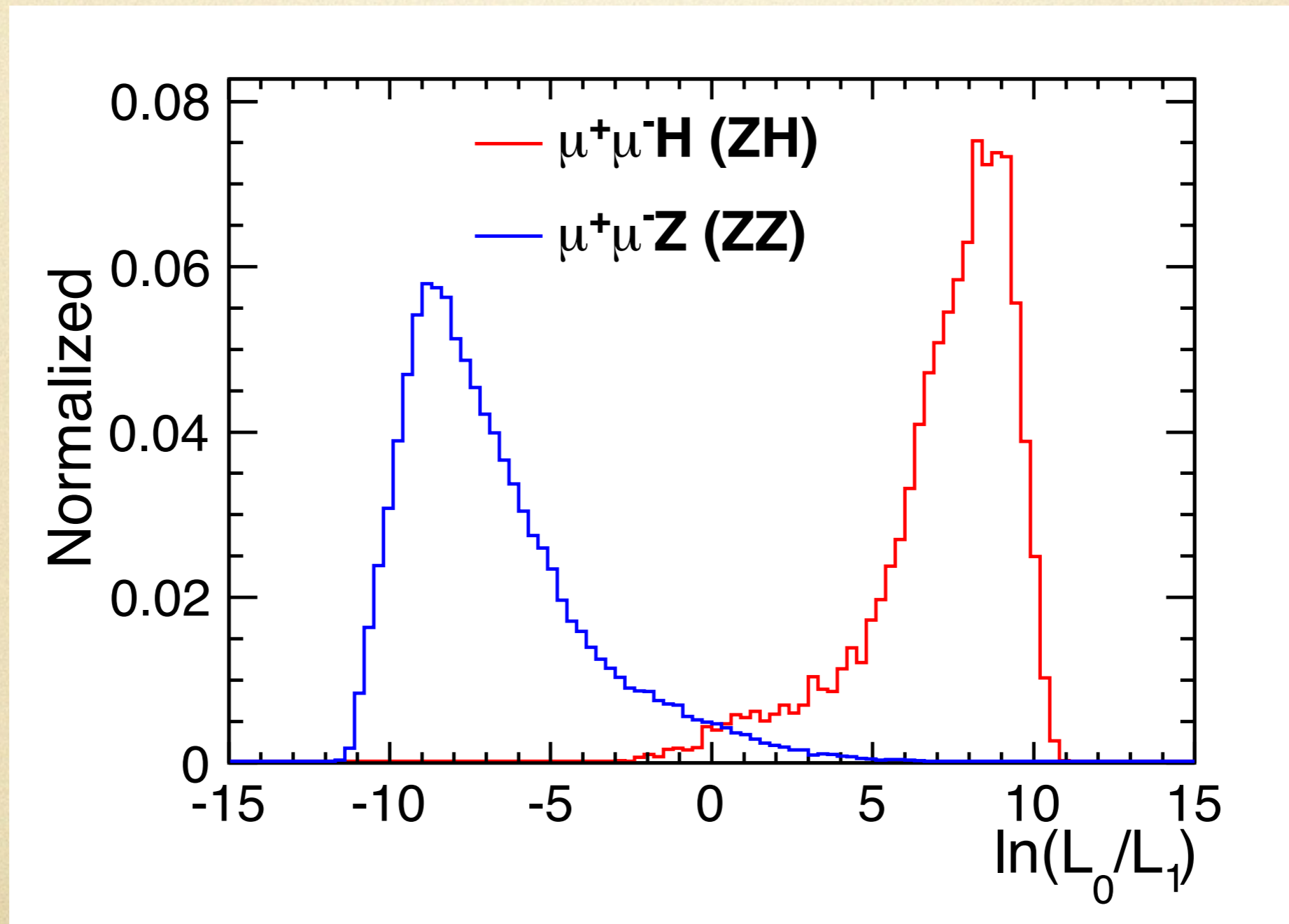
$$M = M(\mu^+ \mu^- Z) \frac{1}{p_Z^2 - m_Z^2 + im_Z \Gamma_Z}$$





matrix element: signal .vs. background (ZZ)

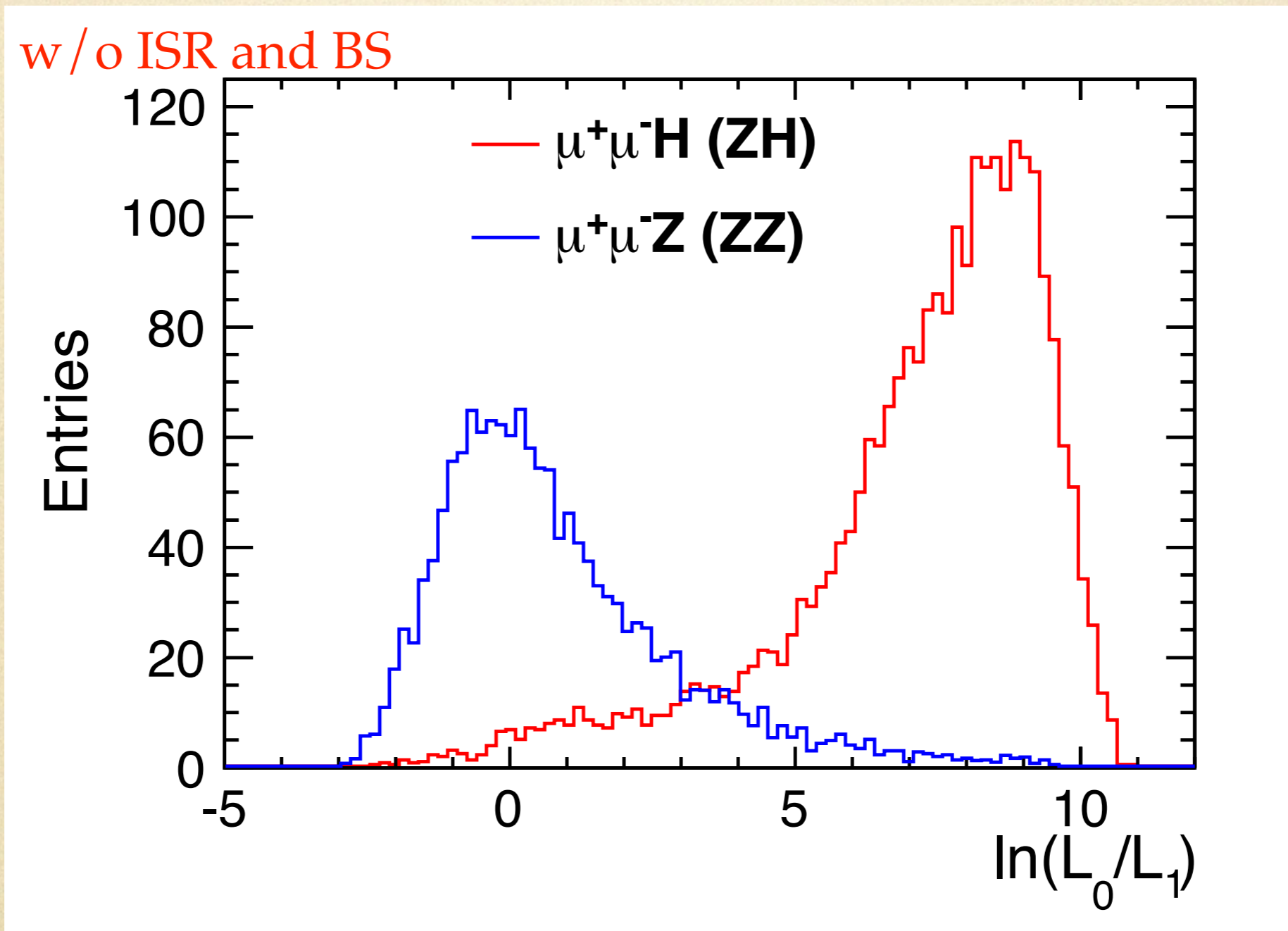
$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H} \quad / \quad M = M(\mu^+ \mu^- Z) \frac{1}{p_Z^2 - m_Z^2 + im_Z \Gamma_Z}$$



# matrix element: signal .vs. background (ZZ)

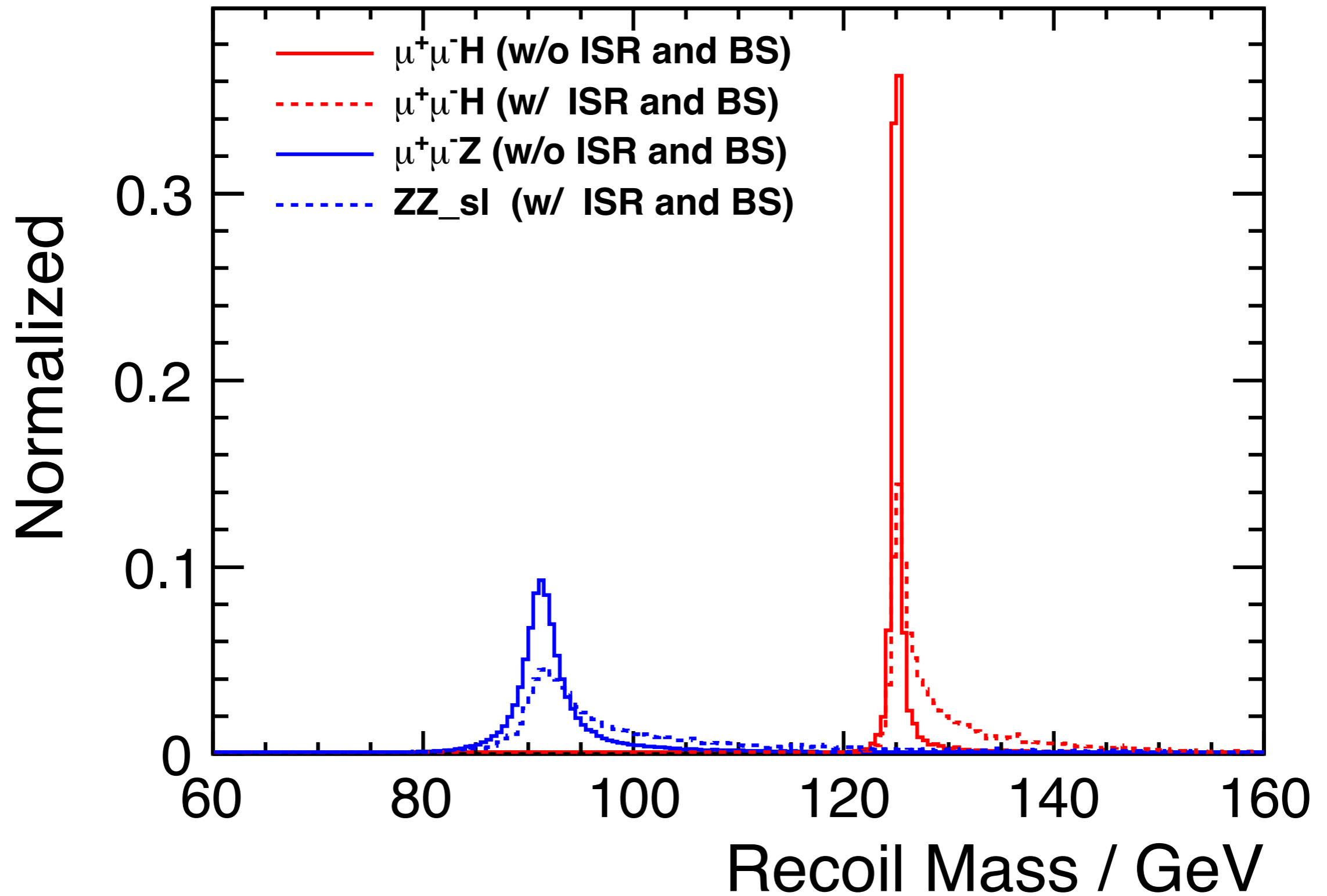
(normalized to expected number of events w/ 250 fb<sup>-1</sup>)

$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H} \quad / \quad M = M(\mu^+ \mu^- Z) \frac{1}{p_Z^2 - m_Z^2 + im_Z \Gamma_Z}$$

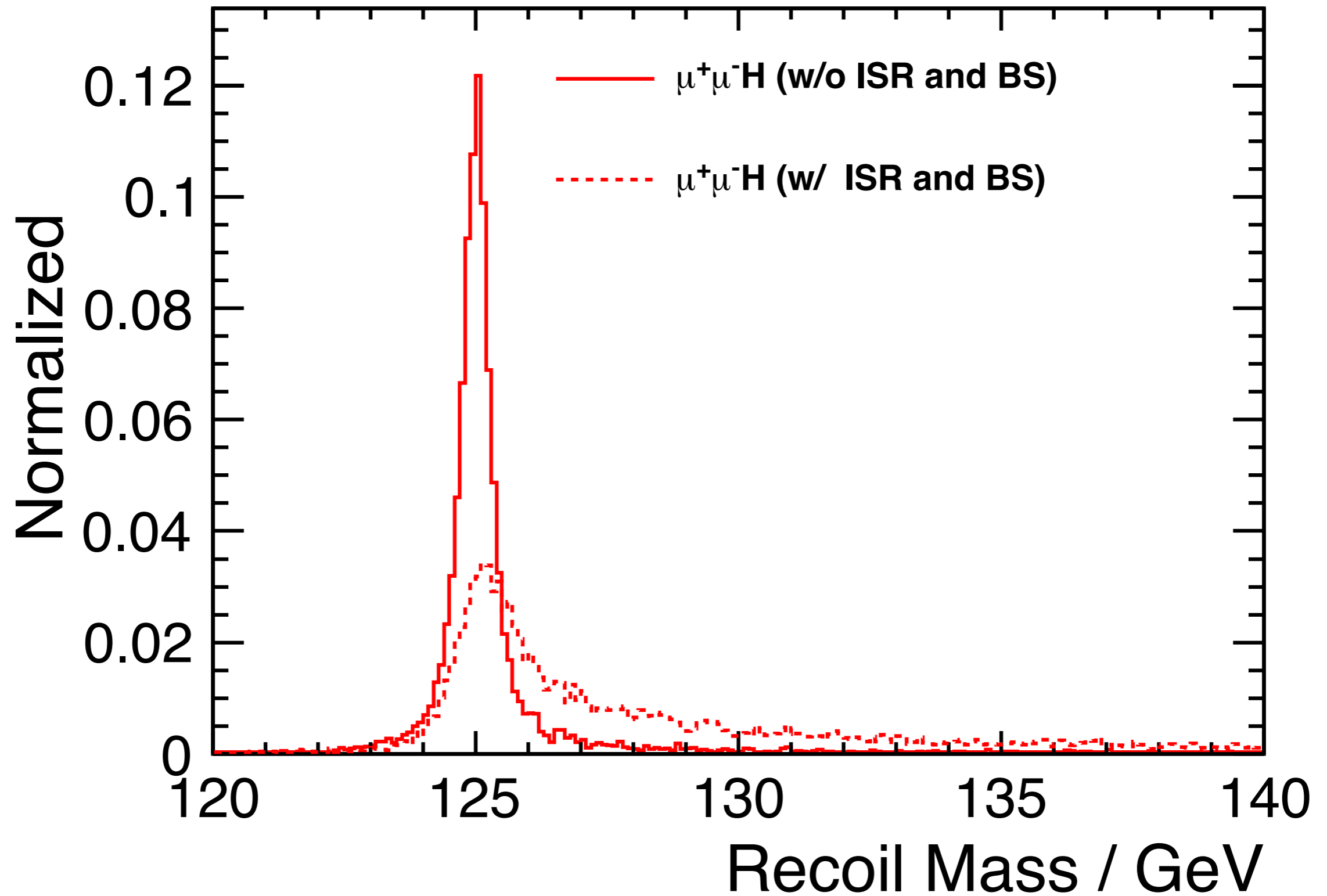


only cut on recoil mass > 110 GeV applied, already very well separated

including everything: ISR + beam spectrum



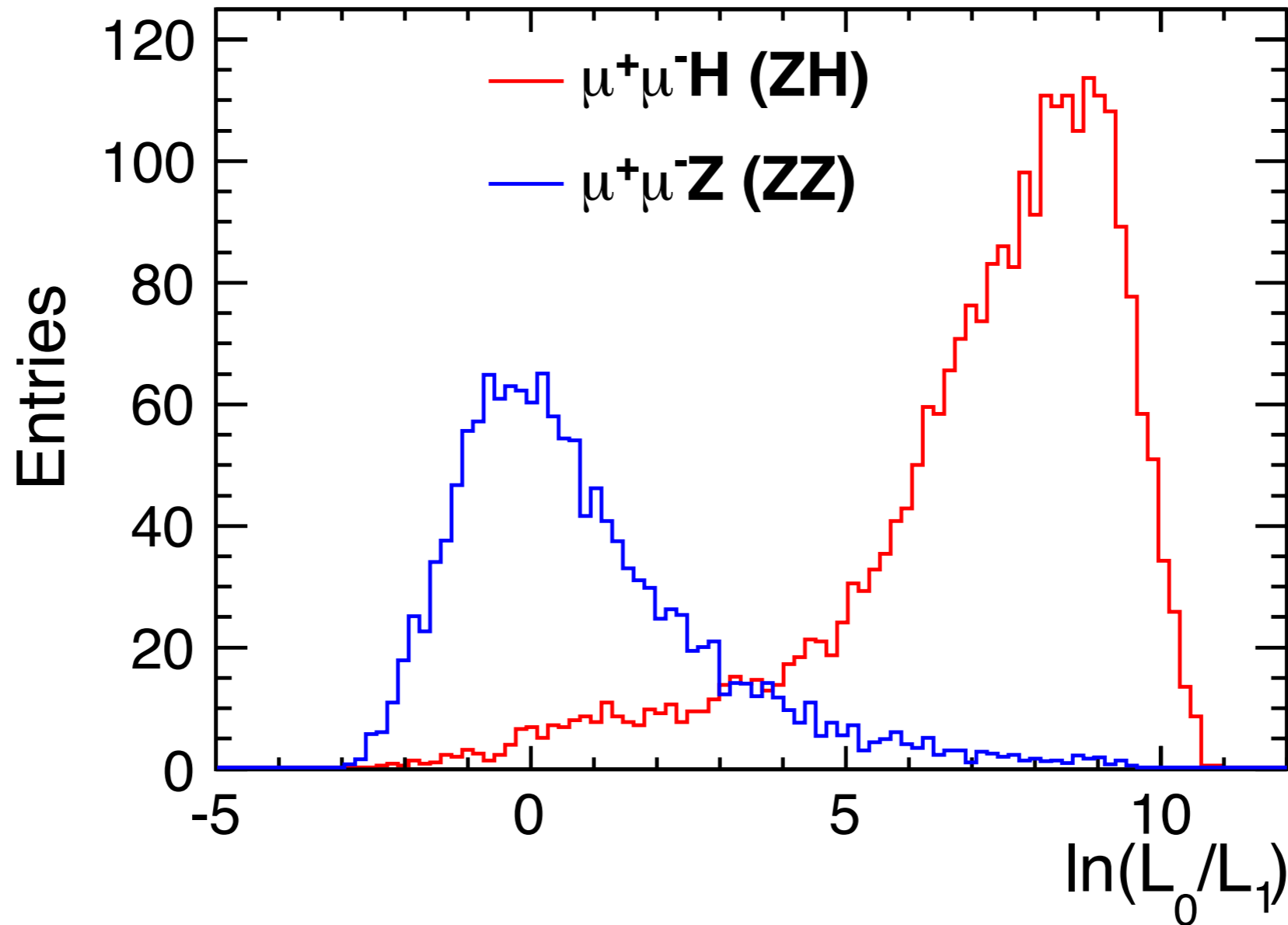
including everything: ISR + beam spectrum



# matrix element: signal .vs. background (ZZ)

$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H} \quad / \quad M = M(\mu^+ \mu^- Z) \frac{1}{p_Z^2 - m_Z^2 + im_Z \Gamma_Z}$$

w/o ISR and BS

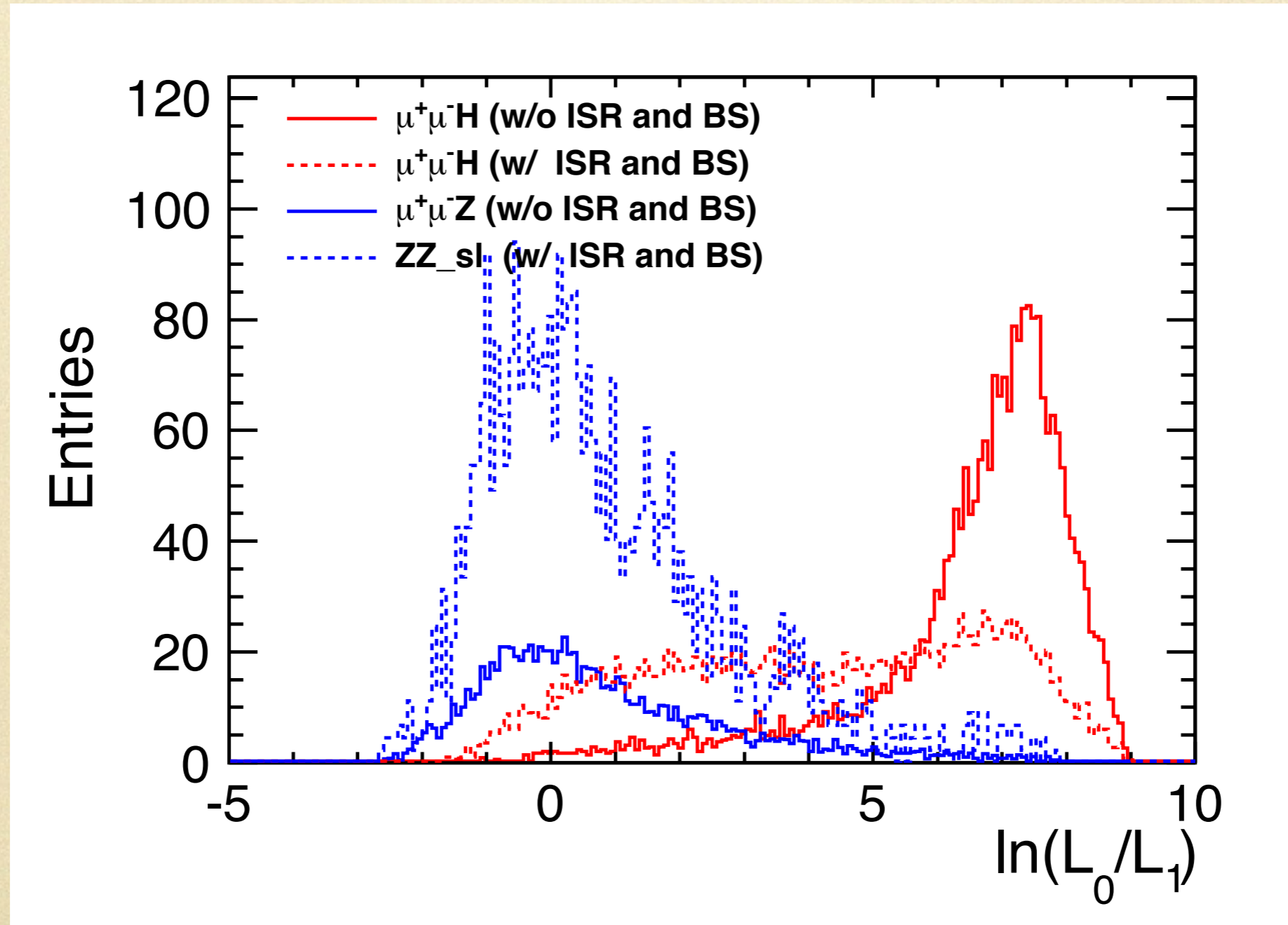


only cut on recoil mass  $> 110$  GeV applied, already very well separated

# matrix element: signal .vs. background (ZZ)

(including everything; normalized to expected number of events w/ 250 fb<sup>-1</sup>)

$$M = M(\mu^+ \mu^- H) \frac{1}{p_H^2 - m_H^2 + im_H \Gamma_H} \quad / \quad M = M(\mu^+ \mu^- Z) \frac{1}{p_Z^2 - m_Z^2 + im_Z \Gamma_Z}$$



on ISR + BS significantly degraded the separation power against ZZ

## proposal to remove ISR

- ☑ ISR enters detector: identification (see Tomita-san's study, eff  $\sim 90\%$ )
- ☐ what if ISR goes to beam pipe? (dominant)


we can resolve it!

$$|P_z(\gamma)| = |P(\gamma)|$$

$$P_t(H) = P_t(Z)$$

$$P_z(H) = -(P_z(Z) + P_z(\gamma))$$

$$E(Z) + \sqrt{P_t^2(H) + P_z^2(H) + m^2(H) + |P(\gamma)|} = \sqrt{s}$$


$$P(\gamma) = \frac{s - 2\sqrt{s}E(Z) + m^2(H) - m^2(Z)}{2(P(Z) \pm (\sqrt{s} - E(Z)))}$$

## summary and next step

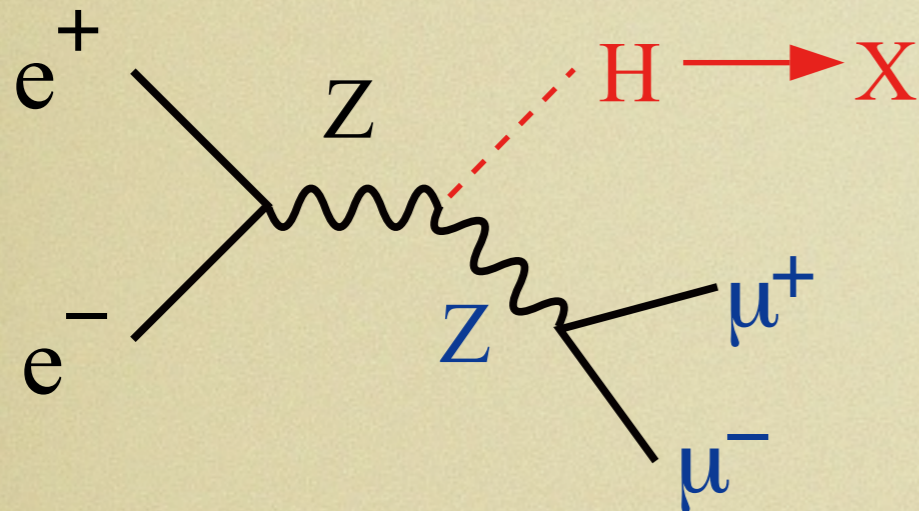
- ME tools have been developed and verified by first look.
- ME indeed can be reconstructed very precisely at the ILC.
- selection will be much simplified using MEM.
- to use Higgs propagator, in principle we need know Higgs mass first, which can be obtained by fitting recoil mass (note: to get mass more cuts on decay part can be applied); to improve  $\sigma(ZH)$  measurement would then be main purpose of using MEM.
- complete analysis is needed to incorporate MEM for recoil mass study, propose to be done by Watanuki-san, (including CP study)
- going to have a look at next application for Higgs self-coupling



back up

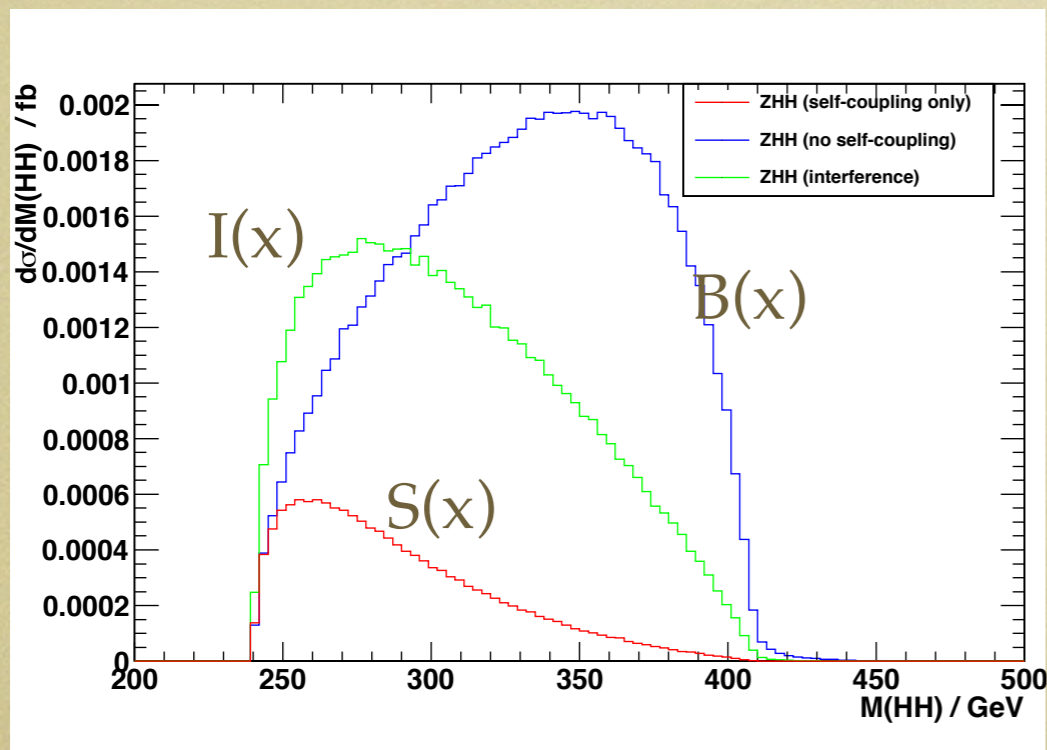
more potential applications:

## recoil mass analysis



- momentum of muon are well measured
- since ecm is well defined, momentum of Higgs can be well measured as well
- hence matrix element can be well calculated without assuming any Higgs decay!

## Higgs self-coupling



$$\frac{d\sigma}{dx} = B(x) + \lambda I(x) + \lambda^2 S(x)$$

$$w_0(x) = c \cdot \frac{I(x) + 2S(x)}{\sigma(x)}$$

what if we use ME as weighting?

# What is Matrix Element (Amplitude)

(squared  $\sim$  differential cross section)

## Cross Section Formula

$$e^+ e^- \rightarrow X_1 + \cdots + X_f + \cdots + X_n$$

$\begin{array}{ccc} \vdots & & \vdots \\ (p^+, s^+) & & (p^-, s^-) \\ \vdots & & \vdots \\ & & (p_f, s_f) \end{array}$

$$d\sigma = \frac{1}{2s\beta_e} \sum_{s^+, s^-, s_f} w_{s^+} w_{s^-} |\mathcal{T}_{fi}|^2 d\Phi_n$$

$\begin{array}{c} \vdots \\ \text{spin weight for } e^- \\ \vdots \\ \text{spin weight for } e^+ \end{array}$

$$w_{s=\pm} = \frac{1 \pm P_s}{2} \quad \left( -1 \leq P_s = \frac{N_+ - N_-}{N_+ + N_-} \leq +1 \right)$$

$$\mathcal{T}_{fi} = \langle p_f, s_f | \hat{T} | p^+, s^+; p^-, s^- \rangle$$

(technically,  $|\text{ME}|^2$  is the weight of each phase space point)

# tool to calculate ME – HELAS (one of many others)

original fortran version by H. Murayama, etc.; c++ version by K. Fujii

## Helicity Amplitudes: HELAS

### External Lines



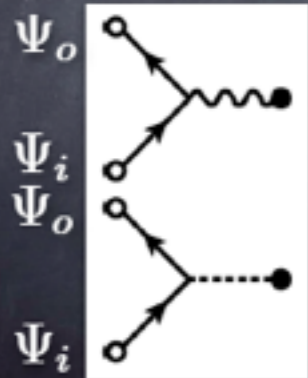
$\Psi_i$

$\Psi_o$

4-momentum  
helicity  
particle  
spinor  
mass  
anti-particle

$IXXXXX(p, m, \lambda, \pm 1, \Psi_i)$   
 $OXXXXX(p, m, \lambda, \pm 1, \Psi_o)$

### Currents



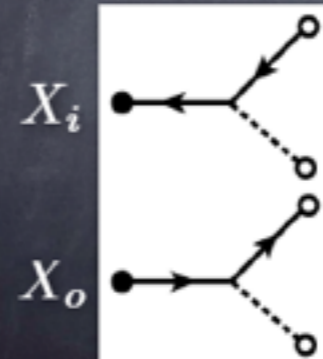
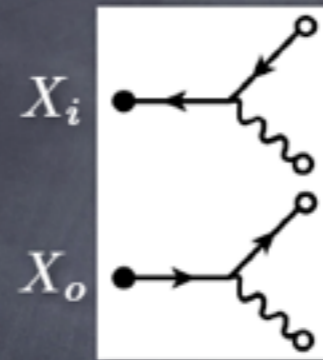
$V$

$S$

incoming spinor  
outgoing spinor  
width  
mass  
wave fun.  
 $G_V(1)$ : left  
 $G_V(2)$ : right

$JIOXXX(\Psi_i, \Psi_o, G_V, m_V, \Gamma_V, V)$   
 $HIOXXX(\Psi_i, \Psi_o, G_S, m_S, \Gamma_S, S)$

### Virtual Fermions



$\Psi_i$

$V$

$\Psi_o$

$V$

$\Psi_i$

$S$

$\Psi_o$

$S$

incoming spinor  
vector  
mass  
width  
incoming virtual spinor  
 $G_V(1)$ : left  
 $G_V(2)$ : right

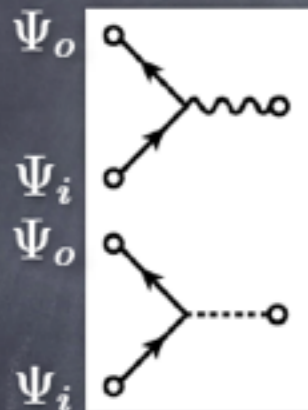
$FVIXXX(\Psi_i, V, G_V, m_X, \Gamma_X, X_i)$   
 $FVOXXX(\Psi_o, V, G_V, m_X, \Gamma_X, X_o)$

outgoing spinor  
outgoing virtual spinor

scalar

$FSIXXX(\Psi_i, S, G_S, m_X, \Gamma_X, X_i)$   
 $FSOXXX(\Psi_o, S, G_S, m_X, \Gamma_X, X_o)$

### Vertices



$V$

$S$

$V_1$

$V_2$

$S$

incoming spinor  
outgoing spinor  
vector  
amplitude  
 $G_V(1)$ : left  
 $G_V(2)$ : right

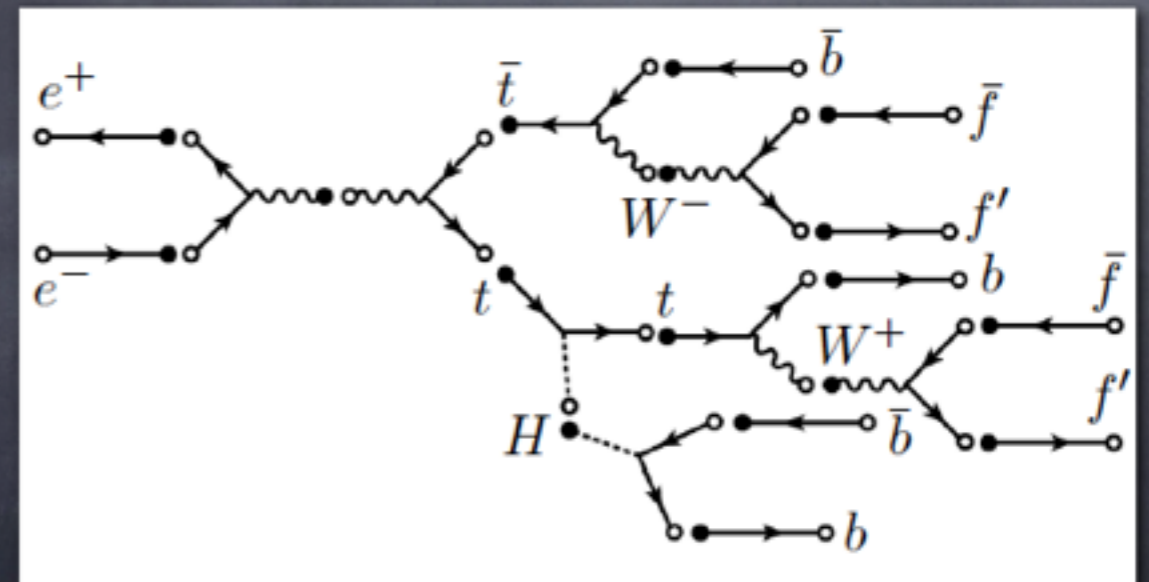
$IOVXXX(\Psi_i, \Psi_o, V, G_V, A)$   
 $IOSXXX(\Psi_i, \Psi_o, S, G_S, A)$

scalar

$VVSXXX(V_1, V_2, S, G_{VVS}, A)$

### Composition of Full Amplitude

$$e^+e^- \rightarrow t\bar{t}H$$



Note: there are some other diagrams  
See physim/top/TTHStudy

Note: there are some more subroutines in HELAS (see manual)

# development of ME tools for physics analysis: status

(for full detector simulation, within Marlin framework)

- basic idea: provide some libraries to calculate ME for given processes, which can be directly used in analyses by just passing four momentum.
- core libraries implemented: **HELLib** (C++ HELAS, helicity amplitude subroutines for feynman diagrams).
- **LCME** (linear collider matrix element libs), so far major Higgs production implemented: LCMEZH, LCMENNH, LCMEEEH, LCMEZHH, LCMENHH, LCMEEEZ.
- verified by using MC truth information.
- Physsim v1.0 released, available now on svn.
- typical way to use it: check out; compile; include **libPhyssim.so** in your \$MARLIN\_DLL; using namespace lcme; follow example marlin processor (included in the package).

svn co <https://svnsrv.desy.de/basic/physsim/Physsim/trunk>

## to be implemented: Detector Effect

- the four momenta we measured have resolution  
—> we need **detector transfer function** (jet-energy resolution, momentum resolution, etc.) and integrate all possible truth four momenta.
- some four momenta can not be measured (missing neutrinos) —> integrate all possible truth four momenta.

$$L(\mathbf{p}_i^{\text{vis}} | \mathbf{a}) = \frac{1}{\sigma_{\mathbf{a}}} \left[ \prod_{j \in \text{inv.}} \int \frac{d^3 p_j}{(2\pi)^3 2E_j} \right] \left[ \prod_{k \in \text{vis.}} \int \frac{d^3 p_k}{(2\pi)^3 2E_k} W_i(\mathbf{p}_i^{\text{vis}} | p_k, \mathbf{a}) \right] |M(p_j, p_k; \mathbf{a})|^2$$

# example core code to calculate matrix element

```
//-----  
// Double Higgs Production Amplitude  
//-----  
HELFermion em(fK[0], kM_e, fHelInitial[0], +1, kIsIncoming);  
HELFermion ep(fK[1], kM_e, fHelInitial[1], -1, kIsOutgoing);  
HELScalar h1(fP[0]);  
HELScalar h2(fP[1]);  
HELFermion f(fP[2], fM[2], fHelFinal[2], +1, kIsOutgoing);  
HELFermion fb(fP[3], fM[3], fHelFinal[3], -1, kIsIncoming);
```

```
Double_t v = 2.*kM_w/kGw;  
Double_t ghhh = -TMath::Power(fMass,2)/v*3.;  
Double_t gzzh = kGz*kM_z;  
Double_t gzzhh = kGz*kGz/2.;
```

```
HELVector zf(fb, f, glz, grz, kM_z, gamz);  
HELVector zs(em, ep, glze, grze, kM_z, gamz);
```

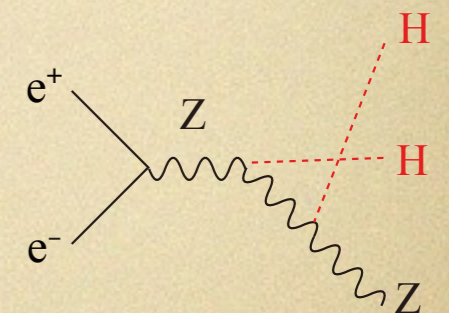
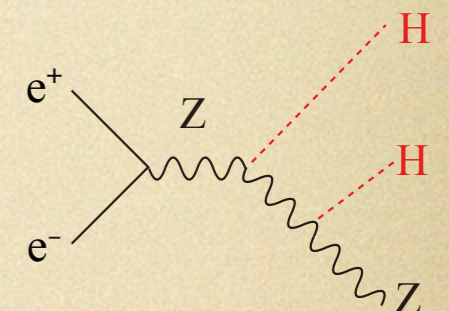
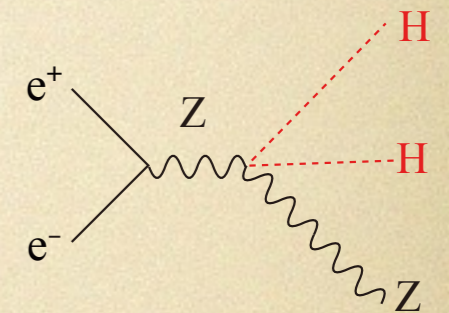
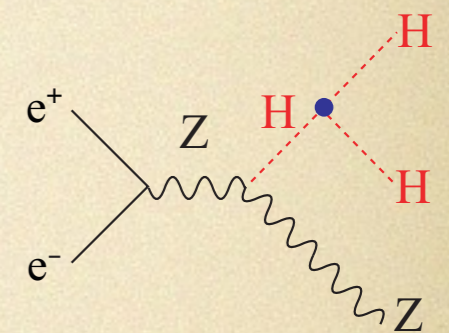
```
HELScalar hh(h1, h2, ghhh, fMass, 0.);  
HELVertex amp1(zs, zf, hh, gzzh); // HHH self-coupling
```

```
HELVertex amp2(zs, zf, h1, h2, gzzhh); // ZZHH 4-point
```

```
HELVector vz1(zf, h1, gzzh, kM_z, gamz);  
HELVertex amp3(zs, vz1, h2, gzzh); // double H-strahlung
```

```
HELVector vz2(zf, h2, gzzh, kM_z, gamz);  
HELVertex amp4(zs, vz2, h1, gzzh); // double H-strahlung
```

```
Complex_t amp = amp1 + amp2 + amp3 + amp4;
```



# example code in your marlin processor

```
// initialize LCMEZHH with Higgs mass of 125 GeV and beam
polarisations P(e-,e+) = (0.,0.)
_zhh = new LCMEZHH("LCMEZHH", "ZHH", 125., 0., 0.);
// set mode of Z decay
_zhh->SetZDecayMode(5);

// -----
// calculate the matrix element
// -----
// put four-momenta of final states to an array
TLorentzVector vLortzMC[4] = {lortzLep1MC, lortzLep2MC, lortzH1MC, lortzH2MC};
// pass kinematics to ME object
_zhh->SetMomentumFinal(vLortzMC);
// matrix element can be given for each combination of initial and final helicities
Int_t vHelLL[2] = {-1,-1};
Int_t vHelLR[2] = {-1,1};
Int_t vHelRL[2] = {1,-1};
Int_t vHelRR[2] = {1,1};
Double_t dSigmaLL = _zhh->GetMatrixElement2(vHelLL);
Double_t dSigmaLR = _zhh->GetMatrixElement2(vHelLR);
Double_t dSigmaRL = _zhh->GetMatrixElement2(vHelRL);
Double_t dSigmaRR = _zhh->GetMatrixElement2(vHelRR);
// if no combination of helicities specified, final combinations are summed
// and initial combinations are weighted by beam polarisations
Double_t dSigma = _zhh->GetMatrixElement2();
// that's all need to do to get matrix element for each event
// -----
```



# verification

$$L(\mathbf{p}_i^{\text{vis}} | \mathbf{a}) = \frac{1}{\sigma_{\mathbf{a}}} \left[ \prod_{j \in \text{inv.}} \int \frac{d^3 p_j}{(2\pi)^3 2E_j} \right] \left[ \prod_{k \in \text{vis.}} \int \frac{d^3 p_k}{(2\pi)^3 2E_k} W_i(\mathbf{p}_i^{\text{vis}} | p_k, \mathbf{a}) \right] |M(p_j, p_k; \mathbf{a})|^2$$

- stdhep events are generated without any ISR and BS, helicity combinations of both initial and final states can be controlled.
- detector transfer function ( $W(p_i | p_k, \mathbf{a})$ ) becomes a delta function, and no invisible variables. (test with MC information).
- ME calculated for each event in this way should be as same as the ME used in event generation (as event weights).
- to verify the calculated ME, if each event is weighted by  $(1. / |ME|^2)$ , all variables should be uniformly distributed.