## Slepton mass measurement in a $\tilde{\tau}$ co-annihilation scenario

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## Outline

This is a status report !!!
(9) Outline
(2) Studying SUSY in rich models
(3) A bench-mark point: STC4

- STC4 @ 500 GeV
- STC4 @ 500 GeV: Globaly
- STC4 @ $500 \mathrm{GeV}: ~ \tilde{e}, \tilde{\mu}$
- STC4 @ $500 \mathrm{GeV}: \tilde{\tau}_{1}$
- Massive SGV $\gamma \gamma$ production
- Fitting the $\tilde{\tau}_{1}$ end-point

4) Outlook \& Conclusions

## Aim of the study

Suppose SUSY is there and has a rich spectrum of sparticles accessible at the ILC. Then:

- Easy - wrt. things like $\tilde{H}$ only, WIMP only: Lots to see. - Hard - wrt. things like $\tilde{H}$ only, WIMP only: Lots to Disentangle.


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- When data starts coming in, what is is first light?
- How do we quickly determine a set of model parameters?
- What is then the optimal use of beam-time in such a scenario?
- And in a staged approach ?
- Spectrum in continuum vs. threshold-scans?
- Special points, eg. between $\tilde{\tau}_{1} \tilde{\tau}_{2}$ and $\tilde{\tau}_{2} \tilde{\tau}_{2}$ thresholds.
- Clean vs. high cross-section
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## SUSY signatures and backgrounds

## Background from SM:

- Real missing energy + pair of SM-particles = di-boson production, with neutrinos:
- $Z Z \rightarrow f \bar{f} \nu \nu$
- Fake missing energy + pair of SM-particles $=\gamma \gamma$ processes, ISR, single IVB.


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- $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow e^{+} e^{-} \gamma \gamma \rightarrow e^{+} e^{-} f \bar{f}$, with both $e^{+} e^{-}$un-detected.
- $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow f \bar{f} \gamma$, with $\gamma$ un-detected.


## Observables:

| Observable | Gives | If |
| :--- | :--- | :--- |
| Edges (or average and <br> width) | Masses | $\ldots$ not too far from <br> threshold |
| Shape of spectrum | Spin |  |
| Angular distributions | Mass, Spin |  |
| Invariant mass distributions <br> from full reconstruction | Mass | ... cascade decays |
| Angular distributions from <br> full reconstruction | Spin, CP, | $\ldots$ masses known |
| Un-polarised Cross-section <br> in continuum | Mass, coupling |  |

## Observables: Pair-production, two-body decay

Consider $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow X X$, followed by $X \rightarrow U Y$, where Y is a detectable (SM) particle. Then

- $E_{Y \max (\min )}=\frac{E_{\text {Beam }}}{2}\left(1-\left(\frac{M_{u}}{M_{X}}\right)^{2}\right)\left(1_{(-)}^{+} \sqrt{1-\left(\frac{M_{X}}{E_{\text {Beam }}}\right)^{2}}\right)$,
- $M_{X}=E_{\text {Beam }} \sqrt{1-(\Delta / \Sigma)^{2}}$

the spectrum is flat (eg if $X$ is a sfermion) between the end-points:



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- $M_{X}=E_{\text {Beam }} \sqrt{1-(\Delta / \Sigma)^{2}}$
- $M_{U}=E_{\text {Beam }} \sqrt{1-(\Delta / \Sigma)^{2}} \sqrt{1-\Sigma / E_{\text {Beam }}}$
$\left(\Delta=E_{Y \text { max }}-E_{Y \text { min }} ; \Sigma=E_{Y \text { max }}+E_{Y \text { min }}\right)$



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If the spectrum is flat (eg if $X$ is a sfermion) between the end-points:
$\left.0<E_{Y}\right\rangle=\left(E_{Y \text { max }}+E_{Y \min }\right) / 2$ and $\sigma_{E_{Y}}=\sqrt{\left(E_{Y \text { max }}-E_{Y \text { min }}\right) / 12}$,
gives



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If the spectrum is flat (eg if $X$ is a sfermion) between the end-points:
- $\left\langle E_{Y}\right\rangle=\left(E_{Y \text { max }}+E_{Y \text { min }}\right) / 2$ and $\sigma_{E_{Y}}=\sqrt{\left(E_{Y \text { max }}-E_{Y \text { min }}\right) / 12}$, which gives
- $M_{U}=E_{\text {Beam }} \sqrt{1-\frac{2\left\langle E_{Y}\right\rangle}{E_{\text {Beam }}}} \sqrt{1-\left(\frac{6 \sigma_{E_{Y}}^{2}}{\left\langle E_{Y}\right\rangle}\right)^{2}}$
- $M_{X}=E_{B e a m} \sqrt{1-\left(\frac{12 \sigma_{E_{Y}}^{2}}{\left.<E_{Y}\right\rangle}\right)^{2}}$


## Example: STC4

## STC4-8

- 11 parameters.
- Separate gluino
- Higgs, un-coloured, and coloured scalar parameters separate

Parameters chosen to deliver all constraints (LHC, LEP, cosmology, low energy).
At $E_{C M S}=500 \mathrm{GeV}$ :

- All sleptons available.
- No squarks.
- Lighter bosinos, up to $\tilde{\chi}_{3}^{0}$ (in $\left.\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \tilde{\chi}_{1}^{0} \tilde{\chi}_{3}^{0}\right)$
(See H. Baer, J. List, arXiv:1307:0782.)


## Full STC4 mass-spectrum



## Zoomed STC4 mass-spectrum



## Channels and observables at 250, 350 and 500 GeV

| Channel | Threshold | Available at | Can give |
| :---: | :---: | :---: | :---: |
| $\tilde{\tau}_{1} \tilde{\tau}_{1}$ | 212 | 250 | $M_{\tilde{\tau}_{1}}, \tilde{\tau}_{1}$ nature, <br> $\tau$ polarisation |
| $\tilde{\mu}_{\mathrm{R}} \tilde{\mu}_{\mathrm{R}}$ | 252 | 250+ | $+M_{\tilde{\mu}_{\mathrm{R}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\mu}_{\mathrm{R}}$ nature |
| $\tilde{\mathrm{e}}_{\mathrm{R}} \tilde{\mathrm{e}}_{\mathrm{R}}$ | 252 | 250+ | $+M_{\tilde{\mathrm{e}}_{\mathrm{R}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\mathrm{e}}_{\mathrm{R}}$ nature |
| $\tilde{\chi}_{1}^{0} \tilde{\chi}_{2}^{0 *)}$ | 302 | 350 | $+M_{\tilde{\chi}_{2}^{0}}, M_{\tilde{\chi}_{1}^{0}}$, nature of $\tilde{\chi}_{1}^{0}, \tilde{\chi}_{2}^{0}$ |
| $\tilde{\tau}_{1} \tilde{\tau}_{2}{ }^{*}$ | 325 | 350 | $+M_{\tilde{\tau}_{2}} \theta_{\text {mix }} \tilde{\tau}$ |
| $\tilde{\mathrm{e}}_{\mathrm{R}} \tilde{\mathrm{e}}_{\mathrm{L}}{ }^{*)}$ | 339 | 350 | $+M_{\tilde{e}_{L}}, \tilde{\chi}_{1}^{0}$ mixing, $\tilde{e}_{\text {L }}$ nature |
| $\tilde{\nu}_{\sim}^{\sim} \tilde{\nu}_{\tilde{\tau}}$ | 392 | 500 | $7 \%$ visible BR $\left(\rightarrow \tilde{\tau}_{1} W\right)$ |
| $\tilde{\chi}_{1}^{ \pm} \tilde{\chi}_{1}^{ \pm *)}$ | 412 | 500 | $+M_{\tilde{\chi}_{1}^{ \pm}}$, nature of $\tilde{\chi}_{1}^{ \pm}$ |
| $\tilde{\mathrm{e}}_{\mathrm{L}} \tilde{\mathrm{e}}_{\mathrm{L}}{ }^{*)}$ | 416 | 500 | $+M_{\tilde{\chi}_{\mathrm{L}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{e}_{\mathrm{L}}$ nature |
| $\tilde{\mu}_{\mathrm{L}} \tilde{\mu}_{\text {L }}{ }^{*)}$ | 416 | 500 | $+M_{\tilde{\mu}_{\mathrm{R}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\mu}_{\mathrm{R}}$ nature |
| $\tilde{\tau}_{2} \tilde{\tau}_{2}{ }^{*)}$ | 438 | 500 | $+M_{\tilde{\tau}_{2}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\tau}_{2}$ nature, $\theta_{\text {mix }} \tilde{\tau}$ |
| $\tilde{\chi}_{1}^{0} \tilde{\chi}_{3}^{0 *)}$ | 503 | 500+ | $+M_{\tilde{\chi}_{3}^{0}}, M_{\tilde{\chi}_{1}^{0}}$, nature of $\tilde{\chi}_{1}^{0}, \tilde{\chi}_{3}^{0}$ |

*): Cascade decays.

+ invisible $\tilde{\chi}_{1}^{0} \tilde{\chi}_{1}^{0}, \tilde{\nu}_{\tilde{\mathrm{e}}} \tilde{\nu}^{\tilde{\nu}_{\mathrm{e}}^{2}} . \tilde{u}$.


## Features of STC4 @ 500 GeV

- The $\tilde{\tau}_{1}$ is the NLSP.
- For $\tilde{\tau}_{1}: E_{\tau, \text { min }}=2.3 \mathrm{GeV}, E_{\tau, \max }=45.5 \mathrm{GeV}$ :
$\gamma \gamma-$ background $\Leftrightarrow$ pairs - background.
- For $\tilde{\tau}_{2}:: E_{\tau, \min }=52.4 \mathrm{GeV}, E_{\tau, \max }=150.0 \mathrm{GeV}$ :
$W W \rightarrow I \nu I \nu-$ background $\Leftrightarrow$ Polarisation.
- For $\tilde{e}_{\mathrm{R}}$ or $\tilde{\mu}_{\mathrm{R}}:: E_{I, \min }=7.3 \mathrm{GeV}, E_{I, \max }=99.2 \mathrm{GeV}$ : Neither $\gamma \gamma$ nor $W W \rightarrow \mid \nu I_{\nu}$ background severe.
- For pol=(1,-1): $\sigma\left(\tilde{e}_{R} \tilde{e}_{R}\right)=1.3 \mathrm{pb}$ !
- $\tilde{\tau}$ NLSP $\rightarrow \tau:$ in most SUSY decays $\rightarrow$ SUSY is background to
- For pol=(-1,1): $\sigma\left(\tilde{\chi}_{2}^{0} \tilde{\chi}_{2}^{0}\right)$ and $\sigma\left(\tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}\right)=$several hundred fb and


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- For pol=(-1,1): $\sigma\left(\tilde{\chi}_{2}^{0} \tilde{\chi}_{2}^{0}\right)$ and $\sigma\left(\tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}\right)=$several hundred fb and $\mathrm{BR}(\mathrm{X} \rightarrow \tilde{\tau})>70 \%$. For pol=(1,-1): $\sigma\left(\tilde{\chi}_{2}^{0} \tilde{\chi}_{2}^{0}\right)$ and $\sigma\left(\tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}\right) \approx 0$.


## STC4 @ 500 GeV

Strategy:

- Global preselection to reduce SM, while efficiency for all signals stays above ~ $90 \%$.
- The further select for all sleptons ( $\tilde{e}_{R}, \tilde{e}_{\mathrm{L}}, \tilde{\mu}_{\mathrm{R}}, \tilde{\mu}_{\mathrm{L}}, \tilde{\tau}_{1}$ ).
- Next step: specific selections for $\tilde{e}_{\mathrm{R}}$ and $\tilde{\mu}_{\mathrm{R}}$, for $\tilde{\mathrm{e}}_{\mathrm{L}}$ and $\tilde{\mu}_{\mathrm{L}}$, and for $\tilde{\tau}_{1}$.
- Last step: add particle id to separate ẽ and $\tilde{\mu}$, special cuts for $\tilde{\tau}_{1}$.
- Check results both for RL and LR beam-polarisation.


## STC4 global

After a few very general cuts:

- Missing energy > 100
- Less than 10 charged tracks
- | $\cos \theta_{\text {Ptot }} \mid<0.95$
- Exactly two $\tau$-jets
- Visible mass < 300 GeV
- $\theta_{\text {acop }}$ between 0.15 and 3.1

$$
\mathrm{E}_{\mathrm{CMS}}=500 \mathrm{GeV}, \text { Pol }=+0.8,-0.3
$$



Magenta: $\gamma \gamma$, Blue: 3f, Red: Rest of SM, Green: SUSY.

## STC4 sleptons @ 500 GeV:ẽ, $\tilde{\mu}$

- Selections for $\tilde{\mu}$ and ẽ:
- Correct charge.
- $P_{T}$ wrt. beam and one $\ell$ wrt the other.
- Tag and probe, ie. accept one jet if the other is "in the box".
- Further selections for R:
- $\mathrm{E}_{j e t}$, beam-pol 80\%,-30\%...
- Further selections for L (LR):


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- $M_{\text {vis }} \neq M_{Z}$


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## Masses from $\tilde{e}, \tilde{\mu}$ in the continuum

- In R[E $\left.E_{\min }, E_{\max }\right]$, the MVB
exists and is $\min (\max )\left(E_{\ell}\right)(!)$
- In presence of background this
won't work.
- Try to mitiaate the effect of extreme cases:
- Also calculate masses using
mean and s.d. of entire
spectrum and compare.
- Make calibration curve with


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LSP
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Slepton


- Make calibration curve with ToyMC.


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- In R $\left[E_{\min }, E_{\max }\right]$, the MVB exists and is $\min (\max )\left(E_{\ell}\right)(!)$
- In presence of background this $\square$
$M_{\tilde{\mathrm{e}}_{\mathrm{R}}}=135.01 \pm 0.19 \mathrm{GeV} / c^{2}$
$M_{\tilde{\chi}_{1}^{0}}=101.51 \pm 0.14 \mathrm{GeV} / c^{2}$
-……
0.1 -

Results for full spectrum ( $\left.E_{C M S}=500,500 \mathrm{fb}^{-1} @[+0.8,-0.3]\right)$
$M_{\tilde{e}_{R}}=140.90 \pm 0.33 \mathrm{GeV} / c^{2}$
$M_{\tilde{\chi}_{1}^{0}}=107.61 \pm 0.23 \mathrm{GeV} / c^{2}$

- Make calibration curve with ToyMC.


## STC4 sleptons @ $500 \mathrm{GeV}: \tilde{\tau}_{1}$

Selections for $\tilde{\tau}_{1}$ :

- Correct charge.
- $\mathrm{P}_{T}$ wrt. beam and one $\tau$ wrt the other.
- $M_{j e t}<M_{\tau}$
- $E_{\text {vis }}<120 \mathrm{GeV}, M_{\text {vis }} \in[20,87]$ GeV .
- Cuts on polar angle and angle between leptons.
- Little energy below 30 deg, or not in $\tau$-jet.
- At least one $\tau$-jet should be hadronic.

- Anti- $\gamma \gamma$ likelihood.


## Massive SGV $\gamma \gamma$ production

- Note the few $\gamma \gamma$ events just at the end-point!
- Don't want to do "dirty tricks" to fit the end-point $\Rightarrow$ need more stat.
- But l've already used all existing generated events, and that only represents $20 \mathrm{fb}^{-1}$, but is nevertheless 580 GB in 1134 stdheps $\Rightarrow$
- Generate on-the-fly inside SGV.
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- So, I extended the interface, so that
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- On the German NAF:
- 1615 jobs of 0.5 MEvents
- Total generated: 0.8 GEvents.
- Wall-clock time first started to last completed: 3 hours (with typically 200-300 jobs running at the same time).
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- Some notes:
- The cross-section of the channels was corrected wrt. the DBD numbers (to take care of not only the ratio of number of $\gamma$ :s to electrons, but also the different beam-profiles.
- In the DBD-production, an artificial $\mathrm{p}_{T}$-kick was applied to the events, which was not done now.


## Fitting the $\tilde{\tau}_{1}$ end-point

- Only the upper end-point is relevant.
- Background subtraction:
- $\tilde{\tau}_{1}$ : Important SUSY
background,but region
above 45 GeV is signal free.
Fit exponential and
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Results for $\tilde{\tau}_{1}$
$E_{\max , \tilde{\tau}_{1}}=44.51_{-0.10}^{+0.12} \mathrm{GeV}$
Translates to an error of $\sim 0.06 \mathrm{GeV} / c^{2} \oplus 1.3 \Delta\left(M_{\tilde{\chi}_{1}^{0}}\right)$ on the mass, where the error from $M_{\tilde{\chi}_{1}^{0}}$ largely dominates


## Reminder: SPS1a' results (Phys. Rev.D82:055016,20010)

The previous $\tilde{\tau}$ study in the very similar model SPS1a' gave:

## Results for $\tilde{\tau}_{1}$

$M_{\tilde{\tau}_{1}}=$
$107.73_{-0.05}^{+0.03} \mathrm{GeV} / \mathrm{c}^{2} \oplus 1.3 \Delta\left(M_{\tilde{\chi}_{1}^{0}}\right)$
The error from $M_{\tilde{\chi}_{1}^{0}}$ largely dominates

$$
\begin{aligned}
& \text { Results for } \tilde{\tau}_{2} \\
& M_{\tilde{\tau}_{2}}=183_{-5}^{+11} \mathrm{GeV} / \mathrm{c}^{2} \oplus 18 \Delta\left(M_{\tilde{\chi}_{1}}\right) \\
& \text { The error from the endpoint } \\
& \text { largely dominates }
\end{aligned}
$$

Results from cross-section for $\tilde{\tau}_{1}$

$$
\Delta\left(N_{\text {signal }}\right) / N_{\text {signal }}=3.1 \% \rightarrow
$$

$$
\Delta\left(M_{\tilde{\tau}_{1}}\right)=3.2 \mathrm{GeV} / c^{2}
$$

Results from cross-section for $\tilde{\tau}_{2}$

$$
\begin{aligned}
& \Delta\left(N_{\text {signal }}\right) / N_{\text {signal }}=4.2 \% \rightarrow \\
& \Delta\left(M_{\tilde{\tau}_{2}}\right)=3.6 \mathrm{GeV} / \mathrm{c}^{2}
\end{aligned}
$$

End-point + Cross-section

$$
\rightarrow \Delta\left(M_{\tilde{\chi}_{1}}\right)=1.7 \mathrm{GeV} / c^{2}
$$

## Also: $\tau$ polarisation in $\tilde{\tau}_{1}$ decays

$\Delta\left(\mathcal{P}_{\tau}\right) / \mathcal{P}_{\tau}=9 \%$.

## Outlook \& Conclusions

- Study best method to analyse spectra, eg
- Optimal statistic for clean signals.
- Specific reconstruction methods for $e, \mu$, and $\tau$. - Make a coherent SGV study of all channels, at all $E_{C M S}$ stages.
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- Exploit more complex decay cascades.
- Status:
- All signals generated.
- All Background exists at 500 , but $\gamma \gamma$ is missing at 250 \& 350 .
- At 500, good selections are at hand for the sleptons. In particular, $\tilde{\tau}_{1}$ compares well with SPS1a' analysis.
- Need to further study the parameter extraction for L-sleptons (SUSY background).
- Need the same for bosinos.


## Thank You!

## BACKUP

## BACKUP SLIDES

## Observables: Pair-production, two-body decay (less text)

- So, there are two SUSY parameters, and two independent observables in the spectrum.
- Any pair of observables can be chosen, edges, average, standard deviation, width, ...
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- Just a bit of algebra to extract the two SUSY masses.
becomes 0, width becomes average/2), so one should not operate
- Note that there are two decays in each event: two measurements per event.
Also note that there are not enough measurements to make a
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## Observables: Pair-production, two-body decay

However:

- If the masses are known from other measurements, there are enough constraints.
- Then the events can be completely reconstructed ...
- ... and the angular distributions both in production and decay can be measured.
- From this the spins can be determined, which is essential to determine that what we are seeing is SUSY.
Furthermore:


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## Furthermore:

- Looking at more complicated decays, such as cascade decays, there are enough constraints if some (but not all) masses are known.
- Allows to reconstruct eg. the slepton mass in $\tilde{\chi}_{2}^{0} \rightarrow \tilde{\ell} \ell \rightarrow \ell \tilde{\chi}_{1}^{0}$ if chargino and LSP masses are known.
- Order-of-magnitude better mass resolution.


## Observables: Pair-production, two-body decay

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## Observables

But this is not all!

- The cross-section in $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow X X$ close to threshold depends both on coupling, spin and kinematics $(=\beta)$.
- The distribution of the angle between the two SM-particles depends on $\beta$, in a complicated, but calculable way.
- The cross-section is different for $L$ and R SUSY particles.
- So checking how much the cross-section changes when switching beam-polarisations measures mixing.
- Measure the helicity of the SM particle -> properties of the particles in the decay, ie. in addition to the produced $X$, also the invisible $U$. In one case this is possible: $\operatorname{In} \tilde{\tau} \rightarrow \tau \tilde{\chi}_{1}^{0} \rightarrow X \nu_{\tau} \tilde{\chi}_{1}^{0}$.


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## Extracting the $\tilde{\tau}$ properties

See Phys.Rev.D82:055016,2010
Use polarisation (0.8,-0.22) to reduce bosino background.
From decay kinematics:

- $M_{\tilde{\tau}}$ from $M_{\tilde{\chi}_{1}^{0}}$ and end-point of spectrum $=E_{\tau, \text { max }}$.
- Other end-point hidden in $\gamma \gamma$ background:Must get $M_{\tilde{\chi}_{1}^{0}}$ from other sources. ( $\tilde{\mu}, \tilde{e}, \ldots)$
From cross-section:
- $\sigma_{\tilde{\tau}}=A\left(\theta_{\tilde{\tau}}, \mathcal{P}_{\text {beam }}\right) \times \beta^{3} / s$, so
- $M_{\tilde{\tau}}=E_{\text {beam }} \sqrt{1-(\sigma s / A)^{2 / 3}}$ : no $M_{\tilde{\chi}_{1}^{0}}$ !

From decay spectra:

- $\mathcal{P}_{\tau}$ from exclusive decay-mode(s): handle on mixing angles $\theta_{\tilde{\tau}}$ and $\theta_{\tilde{\chi}_{1}^{0}}$


## Topology selection

Take over SPS1a' $\tilde{\tau}$ analysis principle थ properties:

- Only two particles (possibly $\tau: s: s)$ in the final state.
- Large missing energy and momentum.
- High Acolinearity, with little correlation to the energy of the $\tau$ decay-products.
- Central production.
- No forward-backward asymmetry.
- Exactly two jets.
- $N_{c h}<10$
- Vanishing total charge.
- Charge of each jet $= \pm 1$
- $M_{j e t}<25 \mathrm{GeV} / \mathrm{C}^{2}$
- $E_{\text {vis }}$ significantly less than $E_{C M S}$.
- Mmiss significantly less than $M_{C M S}$
- No particle with momentum


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Select this by:

- Exactly two jets.
- $N_{c h}<10$
- Vanishing total charge.
- Charge of each jet $= \pm 1$,
- $M_{j e t}<2.5 \mathrm{GeV} / \mathrm{c}^{2}$,
- Evis significantly less than $\mathrm{E}_{\text {CMS }}$.
- $M_{\text {miss }}$ significantly less than $M_{\text {CMS }}$.
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+ anti $\gamma \gamma$ cuts.

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## $\tilde{\tau}_{1}$ and $\tilde{\tau}_{2} f u r t h e r ~ s e l e c t i o n s ~$

- $\tilde{\tau}_{1}$ :
- $\left(E_{j e t 1}+E_{j e t 2}\right) \sin \theta_{a c o p}<30$ GeV .
- $\tilde{\tau}_{2}$ :
- Other side jet not e or $\mu$
- Most energetic jet not e or $\mu$
- Cut on Signal-SM LR of $\mathrm{f}\left(q_{j e t 1} \cos \theta_{j e t 1}, q_{j e t 2} \cos \theta_{j e t 2}\right)$


## Efficiency 15 (22) \%



## $\tilde{\tau}_{1}$ and $\tilde{\tau}_{2}$ further selections

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## $\tilde{\mu}$ channels

Use "normal" polarisation (-0.8,0.22).

- $\tilde{\mu}_{\mathrm{L}} \tilde{\mu}_{\mathrm{L}} \rightarrow \mu \mu \tilde{\chi}_{1}^{0} \tilde{\chi}_{1}^{0}$
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## Selections

- $\theta_{\text {missingp }} \in[0.1 \pi ; 0.9 \pi]$
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- $M_{\mu \mu} \notin[80.100] \mathrm{GeV}$ and $>30$ $\mathrm{GeV} / \mathrm{c}^{2}$
Masses from edges. Beam-energy spread dominates error.




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$\Delta\left(M_{\tilde{\chi}_{1}^{0}}\right)=920 \mathrm{MeV} / c^{2}$
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## $\tilde{\mu}_{\mathrm{R}}$ threshold scan

From these spectra, we can estimate $M_{\tilde{\mathrm{e}}_{\mathrm{R}}}, M_{\tilde{\mu}_{\mathrm{R}}}$ and $M_{\tilde{\chi}_{1}^{0}}$ to $<$ 1 GeV .

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