Slepton mass measurement in a $\tilde{\tau}$ co-annihilation scenario

Mikael Berggren¹

¹DESY, Hamburg

ILD optimsation and analysis phone meeting, Dec 2014



Outline

This is a status report !!!

Outline



- A bench-mark point: STC4
 - STC4 @ 500 GeV
 - STC4 @ 500 GeV: Globaly
 - STC4 @ 500 GeV: ẽ, μ
 - STC4 @ 500 GeV: τ₁
 - Massive SGV $\gamma\gamma$ production



Suppose SUSY is there and has a rich spectrum of sparticles accessible at the ILC. Then:

• Easy - wrt. things like \tilde{H} only, WIMP only: Lots to see.

• Hard - wrt. things like *H* only, WIMP only: Lots to Disentangle. Specifically:

- When data starts coming in, what is is first light ?
- How do we quickly determine a set of model parameters ?
- What is then the optimal use of beam-time in such a scenario ?
- And in a staged approach ?
- Spectrum in continuum vs. threshold-scans?
- Special points, eg. between τ
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 ₂ thresholds.
- Clean vs. high cross-section.
- And so on ...

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Background from SM:

• Real missing energy + pair of SM-particles = di-boson production, with neutrinos:

- $WW \rightarrow \ell \nu \ell \nu$
- $ZZ \rightarrow f\bar{f}\nu\nu$
- Fake missing energy + pair of SM-particles = $\gamma\gamma$ processes, ISR, single IVB.
 - $e^+e^- \rightarrow e^+e^-\gamma\gamma \rightarrow e^+e^-f\bar{f}$, with both e^+e^- un-detected.
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Observables:

Observable	Gives	lf
Edges (or average and		not too far from
width)	Masses	threshold
Shape of spectrum	Spin	
Angular distributions	Mass, Spin	
Invariant mass distributions		
from full reconstruction	Mass	cascade decays
Angular distributions from		
full reconstruction	Spin, CP,	masses known
Un-polarised Cross-section		
in continuum	Mass, coupling	
Polarised Cross-section	Mass, coupling,	
in continuum	mixing	
Decay product polarisation	Mixing	$\tilde{\tau}$ decays
Threshold-scan	Mass(es), Spin	

Consider $e^+e^- \rightarrow XX$, followed by $X \rightarrow UY$, where Y is a detectable (SM) particle. Then

•
$$E_{Y max(min)} = \frac{E_{Beam}}{2} \left(1 - \left(\frac{M_U}{M_X} \right)^2 \right) \left(1_{(-)}^+ \sqrt{1 - \left(\frac{M_X}{E_{Beam}} \right)^2} \right)$$
, so that
• $M_X = E_{Beam} \sqrt{1 - (\Delta/\Sigma)^2}$

$$(\Delta = E_{Y max} - E_{Y min}; \Sigma = E_{Y max} + E_{Y min})$$

If the spectrum is flat (eg if *X* is a sfermion) between the end-points:

• $\langle E_Y \rangle = (E_{Y max} + E_{Y min})/2$ and $\sigma_{E_Y} = \sqrt{(E_{Y max} - E_{Y min})/12}$, which gives

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$$M_U = E_{Beam} \sqrt{1 - \frac{2 \langle E_Y \rangle}{E_{Beam}}} \sqrt{1 - \left(\frac{6\sigma_{E_Y}^2}{\langle E_Y \rangle}\right)^2}$$

• $M_X = E_{Beam} \sqrt{1 - \left(\frac{12\sigma_{E_Y}^2}{\langle E_Y \rangle}\right)^2}$

SUSY mass determination

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• $M_U = E_{Beam} \sqrt{1 - (\Delta/\Sigma)^2} \sqrt{1 - \Sigma/E_{Beam}}$
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If the spectrum is flat (eg if X is a sfermion) between the end-points:

• $< E_Y >= (E_Y _{max} + E_Y _{min})/2$ and $\sigma_{E_Y} = \sqrt{(E_Y _{max} - E_Y _{min})/12}$, which gives

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Example: STC4

STC4-8

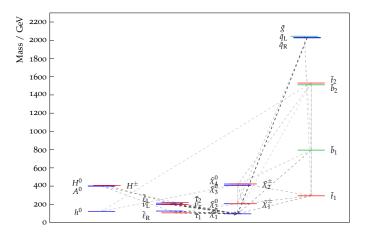
- 11 parameters.
- Separate gluino
- Higgs, un-coloured, and coloured scalar parameters separate

Parameters chosen to deliver all constraints (LHC, LEP, cosmology, low energy).

- At E_{CMS} = 500 GeV:
 - All sleptons available.
 - No squarks.
 - Lighter bosinos, up to $\tilde{\chi}^0_3$ (in $e^+e^- \rightarrow \tilde{\chi}^0_1 \tilde{\chi}^0_3$)

(See H. Baer, J. List, arXiv:1307:0782.)

Full STC4 mass-spectrum

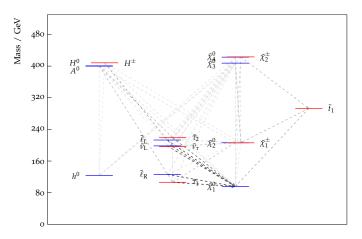


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Zoomed STC4 mass-spectrum



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Image: A matrix

Channels and observables at 250, 350 and 500 GeV

Channel	Threshold	Available at	Can give
$\tilde{ au}_1 \tilde{ au}_1$	212	250	$M_{\tilde{ au}_1}, \tilde{ au}_1$ nature,
			au polarisation
$ ilde{\mu}_{ m R} ilde{\mu}_{ m R}$	252	250+	+ $M_{\tilde{\mu}_{\mathrm{R}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\mu}_{\mathrm{R}}$ nature
$\tilde{e}_R\tilde{e}_R$	252	250+	+ $M_{\tilde{e}_{\rm R}}, M_{\tilde{\chi}_1^0}, \tilde{e}_{\rm R}$ nature
$ ilde{\chi}_1^0 ilde{\chi}_2^{0^{*)}}$	302	350	+ $M_{\tilde{\chi}^0_2}, M_{\tilde{\chi}^0_1}$, nature of $\tilde{\chi}^0_1, \tilde{\chi}^0_2$
$\tilde{ au}_1 \tilde{ au}_2^{*)}$	325	350	+ $M_{\tilde{\tau}_2} \theta_{mix} \tilde{\tau}$
$\tilde{e}_R \tilde{e}_L^{*)}$	339	350	+ $M_{\tilde{e}_L}$, $\tilde{\chi}_1^0$ mixing, \tilde{e}_L nature
$\tilde{\nu}_{\tilde{\tau}}\tilde{\nu}_{\tilde{\tau}}$	392	500	7 % visible BR ($\rightarrow \tilde{\tau}_1 W$)
$\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\pm *)}$	412	500	+ $M_{\tilde{\chi}_1^{\pm}}$, nature of $\tilde{\chi}_1^{\pm}$
$\tilde{e}_L \tilde{e}_L^{*)}$	416	500	+ $M_{\tilde{e}_L}$, $M_{\tilde{\chi}_1^0}$, \tilde{e}_L nature
$ ilde{\mu}_{ m L} ilde{\mu}_{ m L}{}^{*)}$	416	500	+ $M_{\tilde{\mu}_{\mathrm{R}}}, M_{\tilde{\chi}_{1}^{0}}, \tilde{\mu}_{\mathrm{R}}$ nature
$\tilde{\tau}_2 \tilde{\tau}_2^{*)}$	438	500	+ $M_{\tilde{\tau}_2}, M_{\tilde{\chi}_1^0}, \tilde{\tau}_2$ nature, $\theta_{mix \ \tilde{\tau}}$
$ ilde{\chi}_1^0 ilde{\chi}_3^{0^{*)}}$	503	500+	+ $M_{\tilde{\chi}_3^0}, M_{\tilde{\chi}_1^0}$, nature of $\tilde{\chi}_1^0, \tilde{\chi}_3^0$

- *): Cascade decays.
- + invisible $\tilde{\chi}_{1}^{0}\tilde{\chi}_{1}^{0}$, $\tilde{\nu}_{\tilde{e}}\tilde{\mu}\tilde{\nu}_{\tilde{e}}\tilde{\mu}$.

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- For $\tilde{\tau}_1$: $E_{\tau,min} = 2.3 \text{ GeV}$, $E_{\tau,max} = 45.5 \text{ GeV}$: $\gamma\gamma - background \Leftrightarrow pairs - background$.
- For $\tilde{\tau}_2$: : $E_{\tau,min} = 52.4 \text{ GeV}, E_{\tau,max} = 150.0 \text{ GeV}$: $WW \rightarrow l\nu l\nu - background \Leftrightarrow Polarisation.$
- For ẽ_Ror μ̃_R: :E_{l,min} = 7.3 GeV, E_{l,max} = 99.2 GeV: Neither γγ nor WW → lνlν background severe.
- For pol=(1,-1): σ(ẽ_Rẽ_R) = 1.3 pb !
- $\tilde{\tau}$ NLSP $\rightarrow \tau$:s in most SUSY decays \rightarrow SUSY is background to SUSY.
- For pol=(-1,1): $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$ and $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-)$ = several hundred fb and BR(X $\rightarrow \tilde{\tau}$) > 70 %. For pol=(1,-1): $\sigma(\tilde{\chi}_2^0 \tilde{\chi}_2^0)$ and $\sigma(\tilde{\chi}_1^+ \tilde{\chi}_1^-) \approx 0$.

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STC4 @ 500 GeV

Strategy:

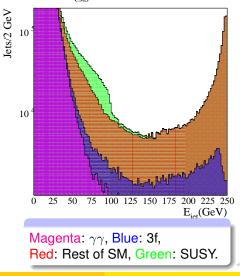
- Global preselection to reduce SM, while efficiency for all signals stays above \sim 90 %.
- The further select for all sleptons ($\tilde{e}_R, \tilde{e}_L, \tilde{\mu}_R, \tilde{\mu}_L, \tilde{\tau}_1$).
- Next step: specific selections for \tilde{e}_R and $\tilde{\mu}_R$, for \tilde{e}_L and $\tilde{\mu}_L$, and for $\tilde{\tau}_1$.
- Last step: add particle id to separate ẽ and μ̃, special cuts for τ̃₁.
- Check results both for RL and LR beam-polarisation.

STC4 global

After a few very general cuts:

- Missing energy > 100
- Less than 10 charged tracks
- $|\cos \theta_{Ptot}| < 0.95$
- Exactly two τ-jets
- Visible mass < 300 GeV
- θ_{acop} between 0.15 and 3.1

E_{CMS}=500 GeV, Pol=+0.8,-0.3



• Selections for $\tilde{\mu}$ and \tilde{e} :

- Correct charge.
- P_T wrt. beam and one ℓ wrt the other.
- Tag and probe, ie. accept one jet if the other is "in the box".

• Further selections for R:

- Cuts on polar angle and angle between leptons.
- E_{jet}, beam-pol 80%,-30%...
- Further selections for L (LR):
 - $q_{jet} \cos \theta_{jet}$
 - $M_{vis} \neq M_Z$

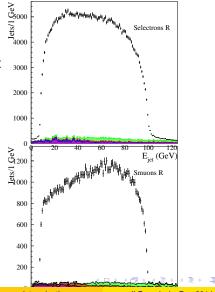
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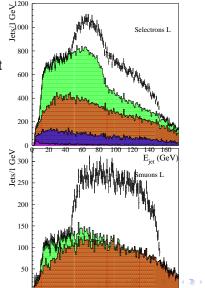
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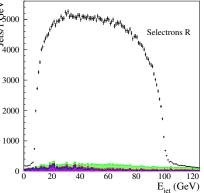


Masses from $\tilde{e}, \tilde{\mu}$ in the continuum

- In R[E_{min}, E_{max}], the MVB exists and is min(max)(E_ℓ) (!)
- In presence of background this won't work.
- Try to mitigate the effect of extreme cases:
 - Exclude highest/lowest x%, and/or
 - Subdivide in sub-samples and average.
- Also calculate masses using mean and s.d. of entire spectrum and compare.
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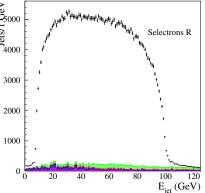
Mikael Berggren (DESY)

ILD analysis. Dec 2014 15 / 21

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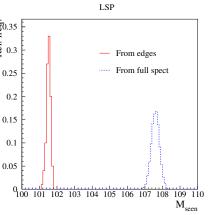


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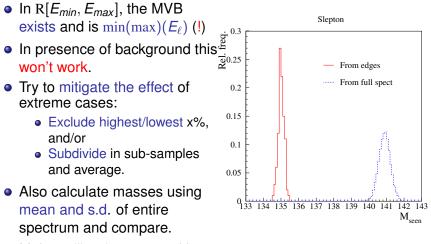
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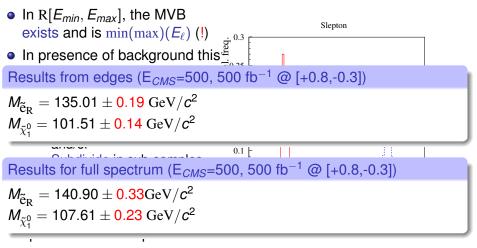
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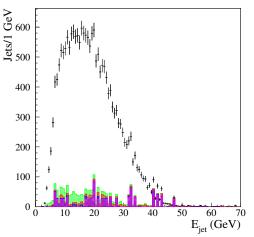
Mikael Berggren (DESY)

STC4 sleptons @ 500 GeV: $\tilde{\tau}_1$

Selections for $\tilde{\tau}_1$:

- Correct charge.
- P_T wrt. beam and one τ wrt the other.
- $M_{jet} < M_{ au}$
- *E_{vis}* < 120 GeV,*M_{vis}* ∈ [20, 87] GeV.
- Cuts on polar angle and angle between leptons.
- Little energy below 30 deg, or not in τ-jet.
- At least one *τ*-jet should be hadronic.
- Anti- $\gamma\gamma$ likelihood.

Mikael Berggren (DESY)



(B)

• Note the few $\gamma\gamma$ events just at the end-point !

- Don't want to do "dirty tricks" to fit the end-point \Rightarrow need more stat.
- But I've already used all existing generated events, and that only represents 20 fb⁻¹, but is nevertheless 580 GB in 1134 stdheps \Rightarrow
- Generate on-the-fly inside SGV.
- Callable Whizard is already an option in SGV, but:
 - Most control in Whizard steering-file, not in SGV.
 - Can't set unique seed to event-generation per job in SGV.
- So, I extended the interface, so that
 - Generator seed,
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- Using the existing meta-data files for aa_2f, easy to script a massive production.
- Specify wanted integrated lumi, maximum number of events/job, and the set of meta-data files to parse.
- On the German NAF:
 - 1615 jobs of 0.5 MEvents
 - Total generated: 0.8 GEvents.
 - Wall-clock time first started to last completed: 3 hours (with typically 200-300 jobs running at the same time).
 - Written to analysis ntuple: 83 MEvents, size 330 GB (compare: would have needed TB:s of stdheps!)
- Some notes:
 - The cross-section of the channels was corrected wrt. the DBD numbers (to take care of not only the ratio of number of γ:s to electrons, but also the different beam-profiles.
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Mikael Berggren (DESY)

SUSY mass determination

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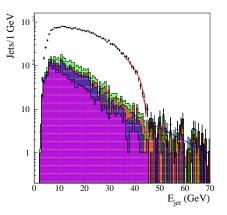
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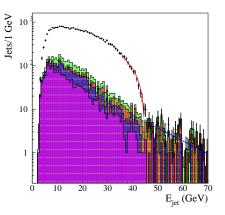
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 Fit exponential and
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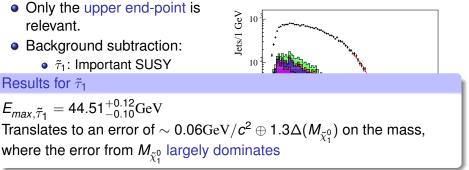


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Results for $\tilde{\tau}_1$

 $E_{max,\tilde{\tau}_1} = 44.51^{+0.12}_{-0.10} \text{GeV}$



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Reminder: SPS1a' results (Phys.Rev.D82:055016,2010)

The previous $\tilde{\tau}$ study in the very similar model SPS1a' gave:

Results for $\tilde{\tau}_1$	Results from cross-section for $\tilde{\tau}_1$
$egin{aligned} M_{\widetilde au_1} = \ 107.73^{+0.03}_{-0.05} { m GeV}/c^2 \oplus 1.3\Delta(M_{\widetilde\chi^0_1}) \end{aligned}$	$\Delta(N_{signal})/N_{signal} = 3.1\% ightarrow \Delta(M_{\widetilde{ au}_1}) = 3.2 { m GeV}/c^2$
The error from $M_{\tilde{\chi}_1^0}$ largely	
dominates	Results from cross-section for $\tilde{\tau}_2$
Results for $\tilde{\tau}_2$	$\Delta(N_{signal})/N_{signal} = 4.2\% \rightarrow$
_	$\Delta(M_{ ilde{ au}_2}) = 3.6 { m GeV}/c^2$
$M_{\widetilde{\tau}_2} = 183^{+11}_{-5} \mathrm{GeV}/c^2 \oplus 18\Delta(M_{\widetilde{\chi}^0_1})$	End-point + Cross-section
The error from the endpoint	$ ightarrow \Delta(M_{ ilde{\chi}_1^0}) = 1.7 { m GeV}/c^2$
largely dominates	
	Also: $ au$ polarisation in $ ilde{ au}_1$ decays

$$\Delta(\mathcal{P}_{\tau})/\mathcal{P}_{\tau}$$
 = 9 %.

Study best method to analyse spectra, eg

- Optimal statistic for clean signals.
- Specific reconstruction methods for e, μ , and τ .
- Make a coherent SGV study of all channels, at all *E_{CMS}* stages.
 - Also channels not studied in SPS1a
 - Exploit more complex decay cascades.
- Status:
 - All signals generated.
 - All Background exists at 500, but $\gamma\gamma$ is missing at 250 & 350.
 - At 500, good selections are at hand for the sleptons. In particular, $\tilde{\tau}_1$ compares well with SPS1a' analysis.
 - Need to further study the parameter extraction for L-sleptons (SUSY background).
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Thank You !

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Backup



BACKUP SLIDES

- So, there are two SUSY parameters, and two independent observables in the spectrum.
- Any pair of observables can be chosen, edges, average, standard deviation, width, ...
- Which choice is the best depends on the situation.
- Just a bit of algebra to extract the two SUSY masses.
- Note that if *E_{beam}* >> *M_X*, there is just one observable (low edge becomes 0, width becomes average/2), so one should not operate too far above threshold !
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- Note that there are two decays in each event: two measurements per event.
- Also note that there are not enough measurements to make a constrained fit, even assuming that the two SUSY particles in the two decays are the same: (2 × 4 unknown components of 4-momentum (=8)) (total E and p conservation (=4) + 2

Observables: Pair-production, two-body decay

However:

- If the masses are known from other measurements, there are enough constraints.
- Then the events can be completely reconstructed ...
- ... and the angular distributions both in production and decay can be measured.
- From this the spins can be determined, which is essential to determine that what we are seeing is SUSY.

Furthermore:

- Looking at more complicated decays, such as cascade decays, there are enough constraints if some (but not all) masses are known.
- Allows to reconstruct eg. the slepton mass in [˜]_{χ2} → ℓℓ → ℓℓ[˜]_{χ1} if chargino and LSP masses are known.

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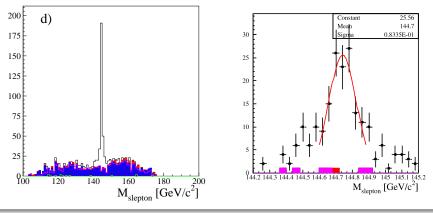
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- Order-of-magnitude better mass resolution.

Backup

Observables: Pair-production, two-body decay

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- The cross-section in e⁺e⁻ →XX close to threshold depends both on coupling, spin and kinematics (= β).
- The distribution of the angle between the two SM-particles depends on β, in a complicated, but calculable way.
- The cross-section is different for L and R SUSY particles.
- So checking how much the cross-section changes when switching beam-polarisations measures mixing.
- Measure the helicity of the SM particle → properties of the particles in the decay, ie. in addition to the produced X, also the invisible U. In one case this is possible: In τ̃ → τχ̃₁⁰ → Xν_τχ̃₁⁰.

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Extracting the $\tilde{\tau}$ properties

See Phys.Rev.D82:055016,2010

Use polarisation (0.8,-0.22) to reduce bosino background.

From decay kinematics:

- $M_{\tilde{\tau}}$ from $M_{\tilde{\chi}_{\tau}^{0}}$ and end-point of spectrum = $E_{\tau,max}$.
- Other end-point hidden in γγ background:Must get M_{χ̃1} from other sources. (μ̃, ẽ, ...)

From cross-section:

•
$$\sigma_{\tilde{\tau}} = A(\theta_{\tilde{\tau}}, \mathcal{P}_{beam}) \times \beta^3/s$$
, so
• $M_{\tilde{\tau}} = E_{beam} \sqrt{1 - (\sigma s/A)^{2/3}}$: no $M_{\tilde{\chi}_1^0}$!

From decay spectra:

• \mathcal{P}_{τ} from exclusive decay-mode(s): handle on mixing angles $\theta_{\tilde{\tau}}$ and $\theta_{\tilde{\chi}_{1}^{0}}$

Backup a

$ilde{ au}$ channels

Topology selection

Take over SPS1a' $\tilde{\tau}$ analysis principle

 $\tilde{\ell}$ properties:

- Only two particles (possibly *τ*:s:s) in the final state.
- Large missing energy and momentum.
- High Acolinearity, with little correlation to the energy of the τ decay-products.
- Central production.
- No forward-backward asymmetry.
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- Exactly two jets.
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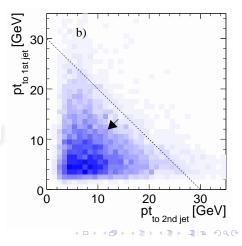
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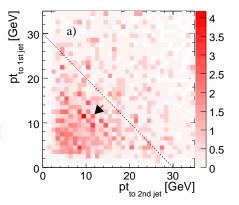
• $(E_{jet1} + E_{jet2}) \sin \theta_{acop} < 30$ GeV.

- Other side jet not e or μ
- Most energetic jet not e or μ
- Cut on Signal-SM LR of f(*q_{iet1}* cos θ_{iet1}, *q_{iet2}* cos θ_{iet2}



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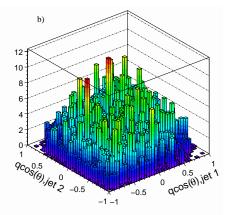
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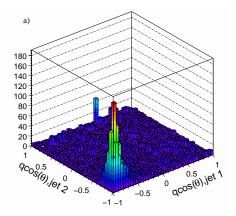
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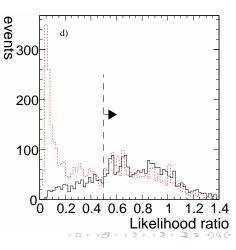
 $\tilde{\tau}$ channels

$\tilde{\tau}_1$ and $\tilde{\tau}_2$ further selections

• $(E_{jet1} + E_{jet2}) \sin \theta_{acop} < 30$ GeV.

τ₂:

- Other side jet not ${\it e}$ or μ
- Most energetic jet not e or μ
- Cut on Signal-SM LR of f(q_{jet1} cos θ_{jet1}, q_{jet2} cos θ_{jet2})



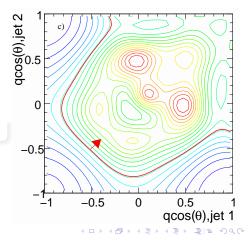
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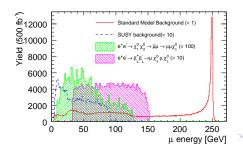
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$\tilde{\mu}$ channels

Use "normal" polarisation (-0.8,0.22).

- $\tilde{\mu}_{L}\tilde{\mu}_{L} \rightarrow \mu\mu\tilde{\chi}_{1}^{0}\tilde{\chi}_{1}^{0}$ • $\tilde{\chi}_{1}^{0}\tilde{\chi}_{2}^{0} \rightarrow \mu\tilde{\mu}_{R}\tilde{\chi}_{1}^{0} \rightarrow \mu\mu\tilde{\chi}_{1}^{0}\tilde{\chi}_{1}^{0}$
- Momentum of *µ*:s
- E_{miss}

• $M_{\mu\mu}$



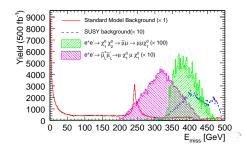
Channels with μ :s

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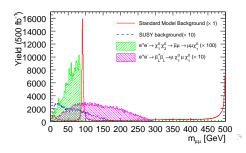


Channels with μ :s

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- $\tilde{\chi}_1^0 \tilde{\chi}_2^0 \to \mu \tilde{\mu}_R \tilde{\chi}_1^0 \to \mu \mu \tilde{\chi}_1^0 \tilde{\chi}_1^0$
- Momentum of *µ*:s
- E_{miss}
- $M_{\mu\mu}$



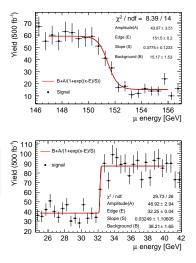
$\tilde{\mu}_{\mathrm{L}}\tilde{\mu}_{\mathrm{L}}$

Selections

- $\theta_{missingp} \in [0.1\pi; 0.9\pi]$
- $E_{miss} \in [200, 430] \text{GeV}$
- $M_{\mu\mu} \notin [80.100] \text{GeV}$ and > 30 GeV/c^2
- Masses from edges. Beam-energy spread dominates error.

$$\Delta(M_{ ilde{\chi}_1^0}) = 920 \mathrm{MeV}/c^2$$

 $\Delta(M_{ ilde{\mu}_\mathrm{L}}) = 100 \mathrm{MeV}/c^2$

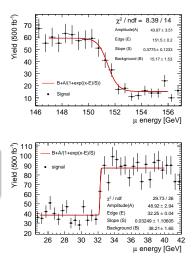


$\tilde{\mu}_{\rm L}\tilde{\mu}_{\rm L}$

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$$\Delta(M_{ ilde{\chi}_1^0}) = 920 {
m MeV}/c^2 \ \Delta(M_{ ilde{\mu}_{
m L}}) = 100 {
m MeV}/c^2$$



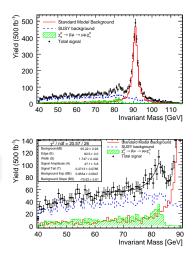


Selections

- $\theta_{missingp} \in [0.2\pi; 0.8\pi]$
- $p_{Tmiss} > 40 {
 m GeV}/c$
- β of μ system > 0.6.
- $E_{miss} \in [355, 395]$ GeV

Masses from edges. Beam-energy spread dominates error.

 $\Delta(M_{\tilde{\chi}^0_2}) = 1.38 {
m GeV}/c^2$

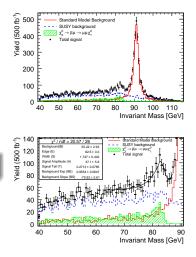




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$\tilde{\mu}_{\mathrm{R}}$ threshold scan

From these spectra, we can estimate $M_{\tilde{e}_R}$, $M_{\tilde{\mu}_R}$ and $M_{\tilde{\chi}_1^0}$ to < 1 GeV.



$\tilde{\mu}_{\rm R}$ threshold scan

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- So: Next step is $M_{\tilde{\mu}_{\rm R}}$ from threshold:
 - 10 points, 10 fb $^{-1}$ /point.

• Luminousity $\propto E_{CMS}$, so this is $\Leftrightarrow 170 \text{ fb}^{-1} @ E_{CMS} = 500 \text{ GeV}.$

Error on $M_{\tilde{\mu}_{\mathrm{R}}}$ = 197 MeV

Channels with μ :s

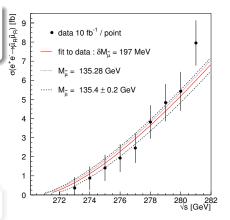
$\tilde{\mu}_{\mathrm{R}}$ threshold scan

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Channels with μ :s

$\tilde{\mu}_{\mathrm{R}}$ threshold scan

From these spectra, we can 5(e⁺e`→й_Rй_R) [fb] 2 8 6 estimate $M_{\tilde{e}_{P}}$, $M_{\tilde{\mu}_{P}}$ and $M_{\tilde{\chi}_{1}^{0}}$ to < data 10 fb⁻¹ / point 1 GeV. fit to data : $\delta M_{\tilde{u}} = 197 \text{ MeV}$ M. = 135.28 GeV So: Next step is $M_{\tilde{\mu}_{\rm R}}$ from 6 $M_{\widetilde{\mu}}=~135.4\pm0.2~GeV$ threshold: 5 • 10 points, 10 fb $^{-1}$ /point. 4 з • Luminousity $\propto E_{CMS}$, so this is \Leftrightarrow 170 fb⁻¹ @ *E*_{CMS}=500 GeV. 2 0 280 28 √s [GeV] Error on $M_{\tilde{\mu}_{\rm R}}$ = 197 MeV 272 274 276 278 282