

ILD tracker performance: the prize for a smaller detector

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¹DESY, Hamburg

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Outline

- 1 Introduction
- 2 Effects of Tracking geometry
- 3 Geometry used
- 4 Results: helix parameters
- 5 Results: Higgs recoil-mass @ 350 GeV
- 6 Conclusions

Present optimisation studies

Detector-component optimisation in ILD (post DBD):

- Presently
- Mainly has been about ECal
- Aimed at cost-reduction.
- Only considers JER as metric - mainly for highest energy jets.
- Studies on:
 - Sensitive detector technology
 - Number of layers
 - Radius and/or length

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Effects of Tracking geometry

Reminder:

- $\Delta(1/p_T) \propto L^{-2.5}$ (2 purely geometric + (\geq) 0.5 because of less points in TPC).
 - But only linear in σ_{point} and B-field
- Please note: Stored energy in B-field $\propto B^2 V$, so at equal stored energy, a smaller detector can have a higher field.
- Also: $\sigma_{point,TPC}^2 = \sigma_0^2(\sin \phi) + \frac{C_d^2(B)}{N_{eff}(\sin \theta)} Z$, $C_d^2(B) \propto 1/(1 + (\mu B)^2) \Rightarrow$ complicated relation, but gets better with shorter drift-length and higher B .
- Also: Higher B-field \Rightarrow possible to have smaller beam-pipe/vertex-detector \Rightarrow better IP-resolution.

Detailed estimation: SGV

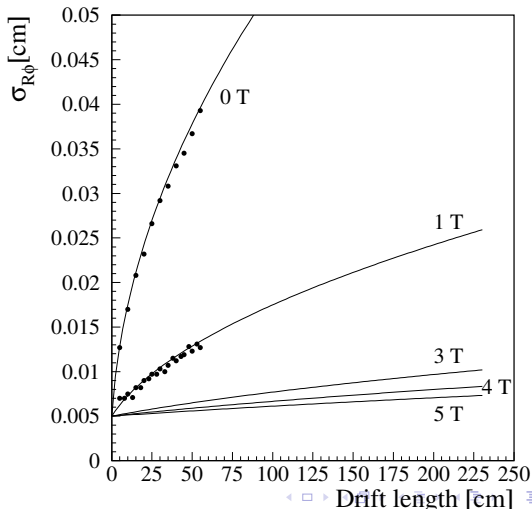
- The description of the point-errors in the TPC have been extended to include all the terms in the description. Inputs:
 - $\sigma_{R\phi}$ and at σ_Z zero drift length.
 - Zero B-field diffusion in $R\phi$.
 - Mobility.
 - Track-radial direction angle dependence of $\sigma_{R\phi}$.
 - Ratio of diffusion in Z and $R\phi$, default 2.

Numbers (mostly) from Ron Settles for T2K gas.

- Replace the default simplified TPC layer structure (pad-rows grouped by 9) to the full 225 layers ILD to simplify scaling.
- Script to scale the default ILD.

TPC point-resolution vs. Z and B in SGV and DBD

- Points: Prototype measurements (from DBD/DBD SVN)
- Lines: Formula used in SGV.

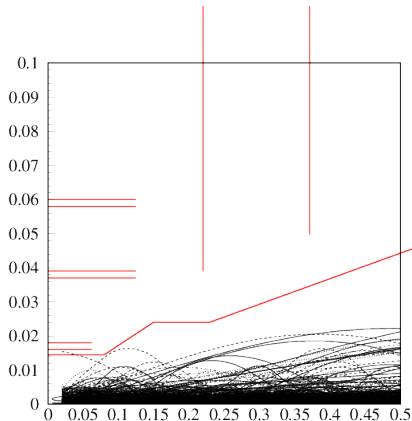


Vertex-detector size and B

- If the B field goes from
 - ... 3.5 T ... to 5 T,
 - the cone of bs pairs get squeezed.
 - The "cone" is a parabola \Rightarrow
 - The radius of the edge of the "cone" $\propto 1/\sqrt{B}$ at a given Z.
 - However:
 - The particles directly hitting the VTX are **not** from the "cone", and have no $P_T - \theta$ correlation.
 - The back-scatters from the BCal that hits the VTX are **not** produced at a given Z if the detector shrinks.

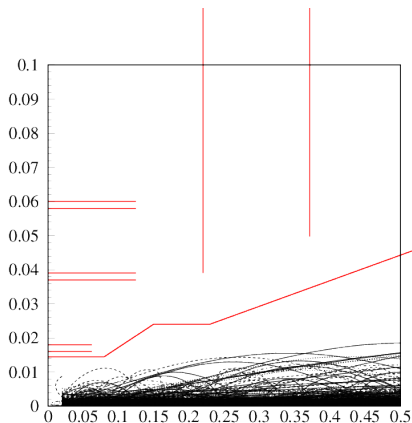
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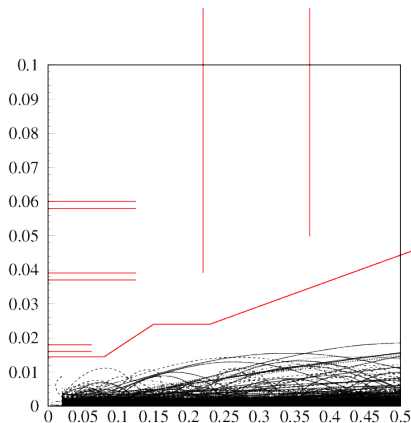
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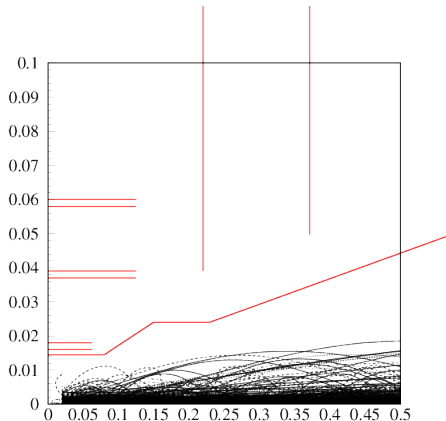
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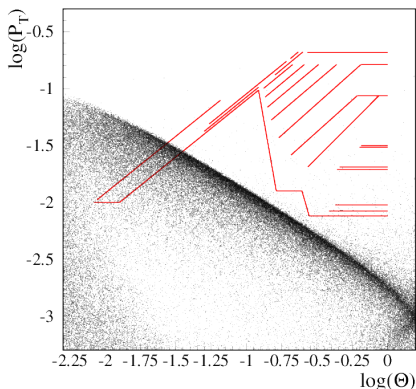
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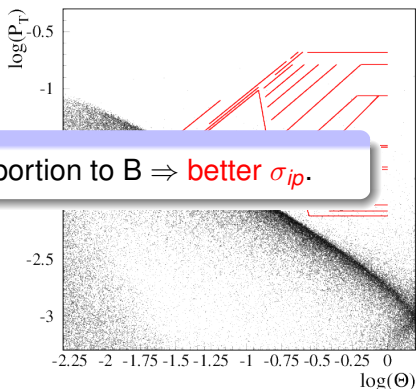
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Vertex-detector size and B

- If the B field goes from
- ... 3.5 T ... to 5 T,
- the cone of bs pairs get **squeezed**.
- The "cone" is a parabola
- This \Rightarrow Reduce $R_{VTX-inner}$ in proportion to B \Rightarrow **better σ_{ip}** .
- "cone" $\propto 1/\sqrt{B}$ at a given Z.
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Geometries used

I studied five different ways to change the ILD baseline geometry. For each of these I did modifications in 5 steps:

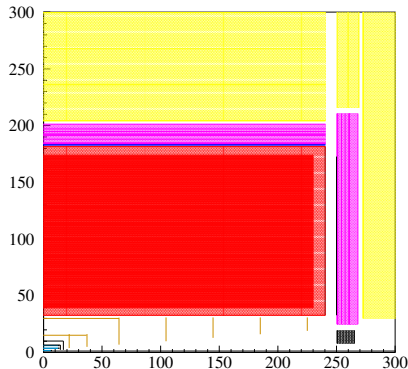
- 1 Keep baseline aspect ratio.
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(All showing the largest modification using SGV:s detector description visualiser)

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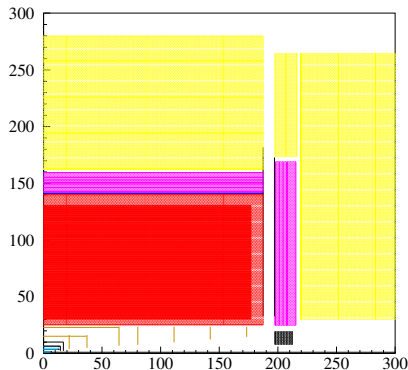


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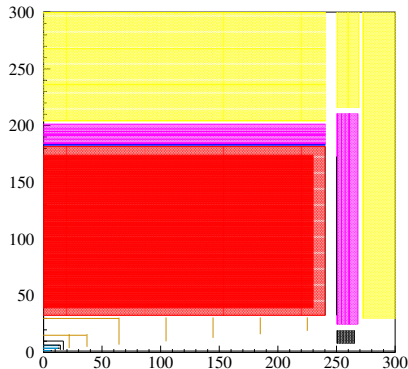


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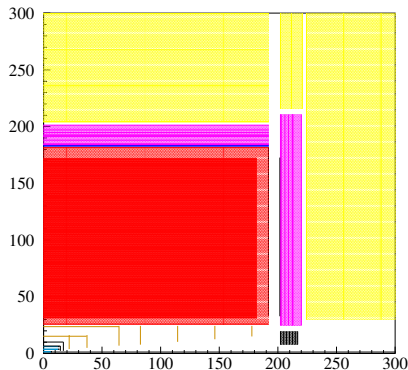


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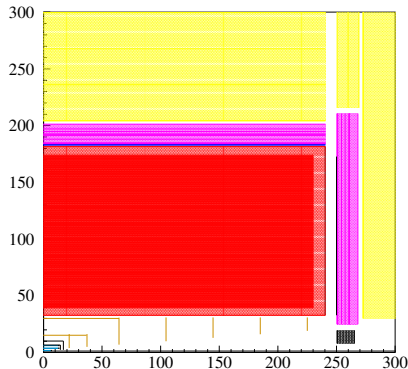


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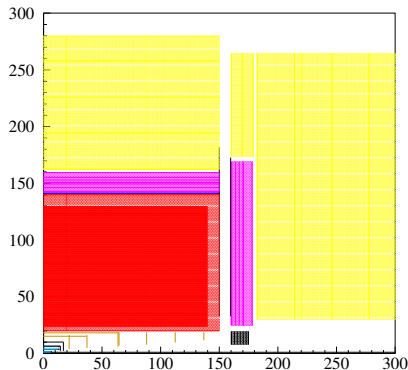


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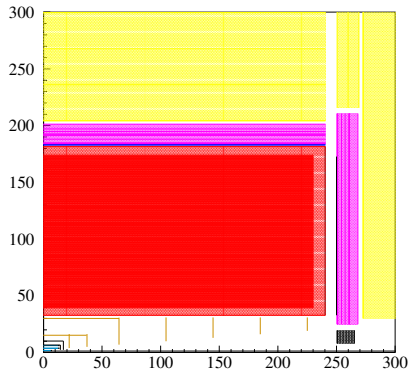


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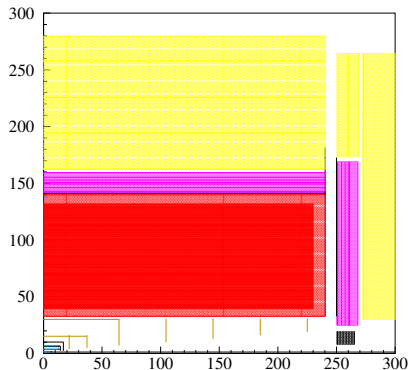


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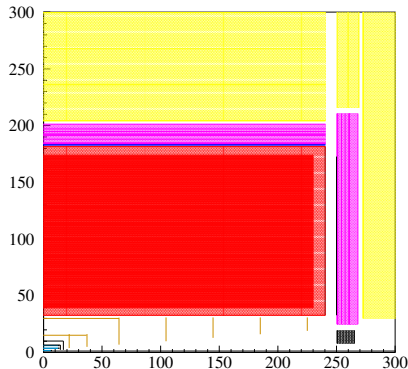


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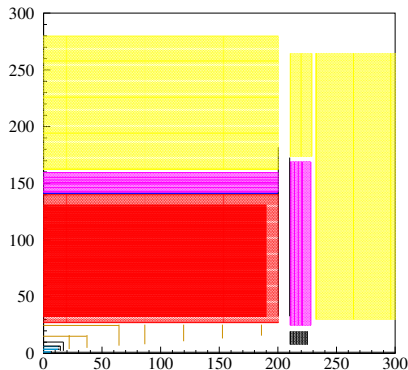


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Other options

Apart from the pure modifications of the geometry, I considered

- 1 Only Outer extent of the TPC modified. Everything outside was also moved, as was the FTD strip-discs. VTX, SIT and FTD pixels unchanged.
- 2 Also modify B, keeping $B^2 V$ constant (V =volume of solenoid).
- 3 Keep B fixed, but modify TPC inner radius (and hence the outer layer of the SIT and the outer radius of the FTD discs.)
- 4 Both 2 and 3.
- 5 In addition to 4, also scale beam-pipe and VTX-inner with B.
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 $2 \times 5 \times 5 \times 6 = 300 \text{ cases !}$
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Effects of Tracking geometry on helix parameters

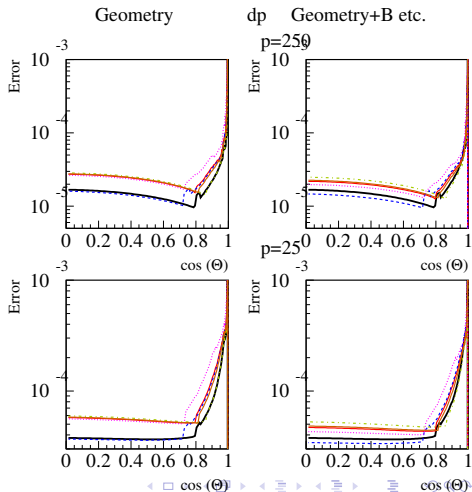
Check $\Delta(1/p)$ and $\Delta(ip_{R\phi})$ at different p and $\cos\theta$ ($\Delta(ip_Z)$ similar to $\Delta(ip_{R\phi})$, angles not relevant - the other uncertainties dominate)

- Red/magenta: fixed aspect-ratio, decrease size.
- Blue: fixed R, decrease Z.
- Green/orange: fixed Z, decrease R.
- Black: TDR detector.
- $\Delta(1/p)$ vs. $\cos\theta$
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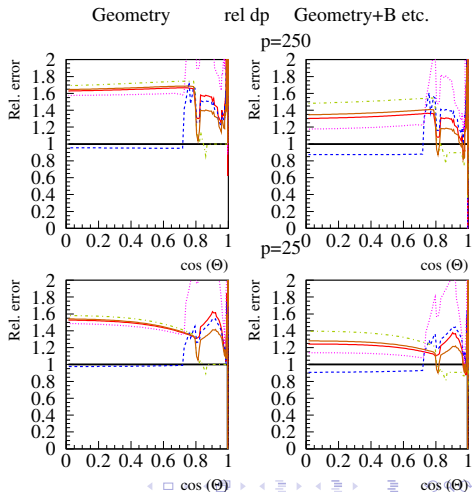
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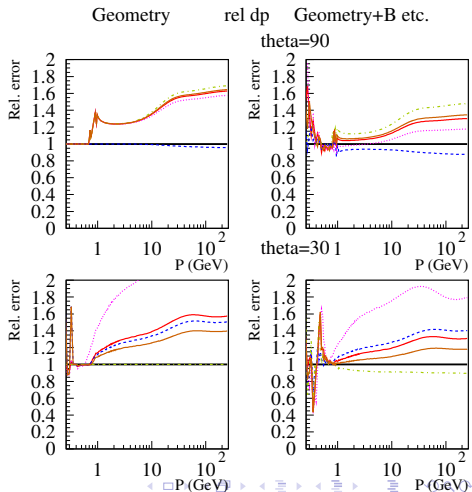
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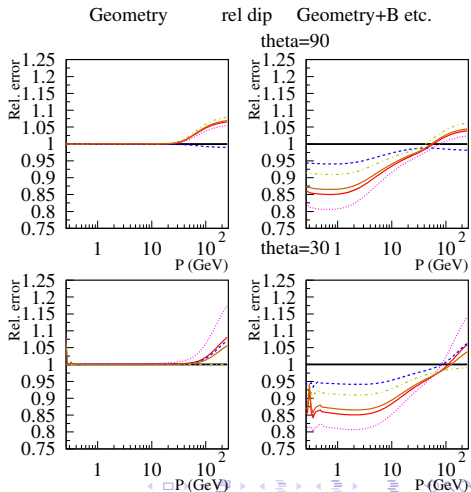
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Effects of Tracking geometry on helix parameters

The point of the exercise is to reduce the size (=area) of the calorimeters (in particular ECal). Here I show the performance as a function of $A_{ECal}/A_{ECal, TDR}$

- Circles/triangles: fixed aspect-ratio.
- Squares: fixed R.
- Stars/inv. triangles: fixed Z.
- $\Delta(1/p)$, barrel.
- $\Delta(1/p)$, endcap.
- $\Delta(ip_{R\phi})$, barrel.
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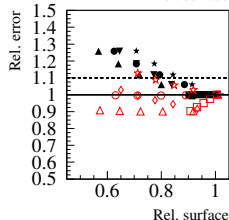
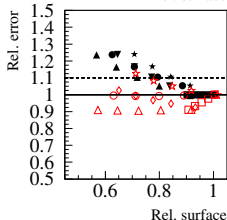
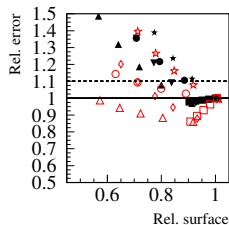
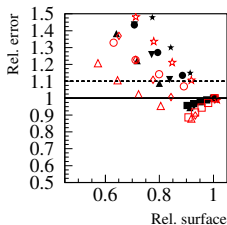
Filled: Only modify geometry; Open: also do other changes. 

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$dp, \theta=90$

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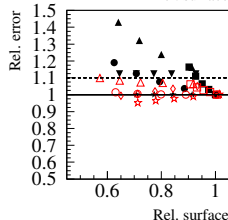
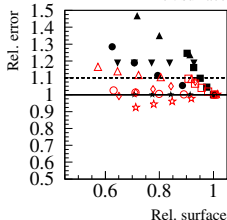
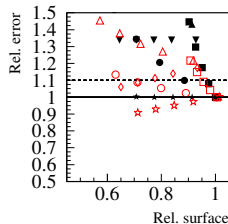
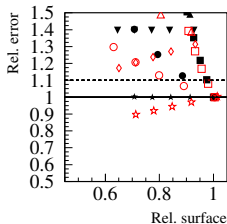
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Effects of Tracking geometry on helix parameters

The point of the exercise is to reduce the size (=area) of the calorimeters (in particular ECal). Here I show the performance as a function of $A_{ECal}/A_{ECal, TDR}$

dp, theta=30

- Circles/triangles: fixed aspect-ratio.
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- $\Delta(ip_{R\phi})$, barrel.
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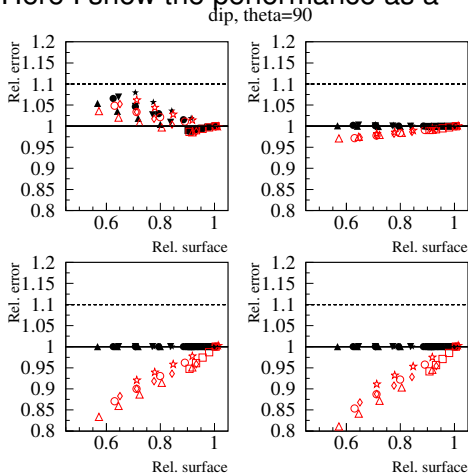


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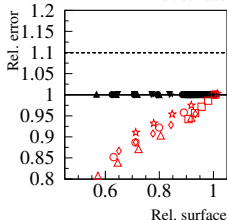
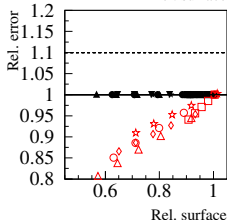
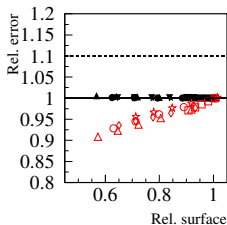
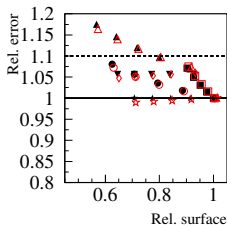
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Effects of Tracking geometry on helix parameters, 4T

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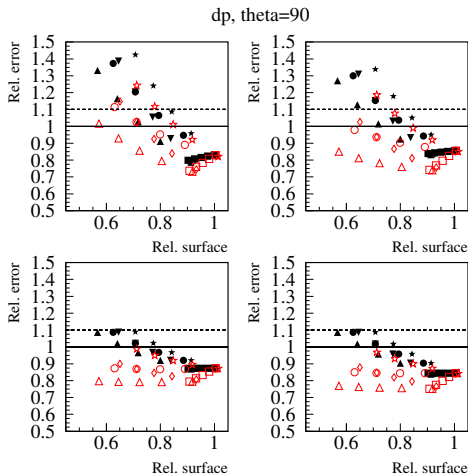
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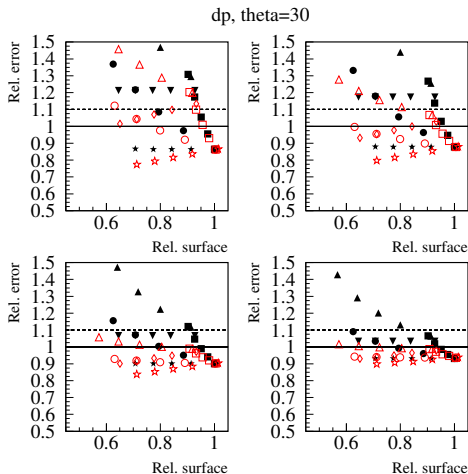


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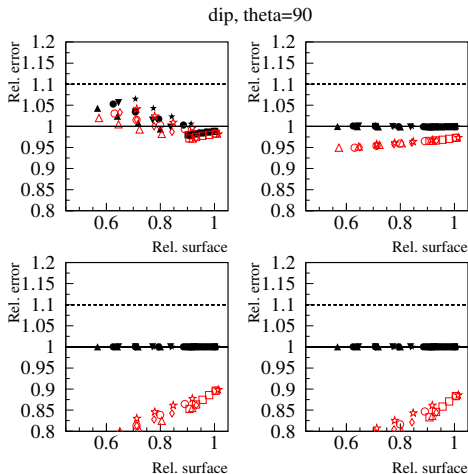


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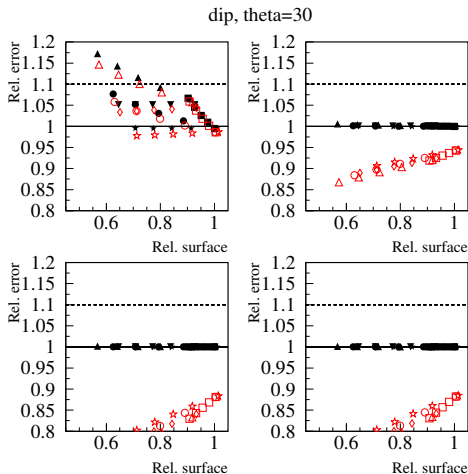


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Higgs recoil-mass @ 350 GeV: measurables

- Look at $e^+e^- \rightarrow ZH$,
 $Z \rightarrow \mu^+\mu^-$, $H \rightarrow X$.
- Signal only, perfect μ finding, SGV.
- Recoil-mass =
$$\sqrt{(E_Z - E_{CMS})^2 - \bar{p}_Z^2}$$
 where

$$E_Z = E_{\mu^+} + E_{\mu^-}, \bar{p}_Z = \bar{p}_{\mu^+} + \bar{p}_{\mu^-},$$

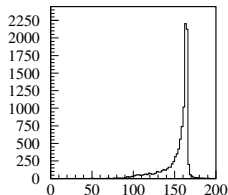
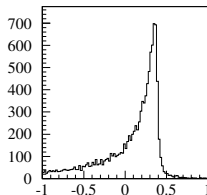
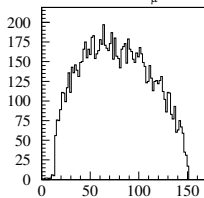
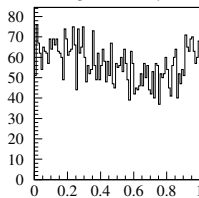
$$E_{CMS} = \text{nominal} = 350.$$
- So, it's all about measuring the μ :s !
- Note: E range 20 to 150, θ in barrel.

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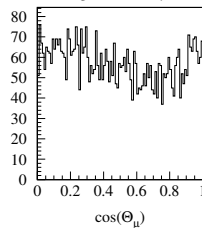
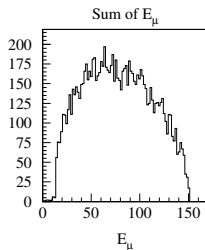
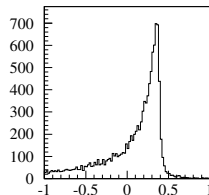
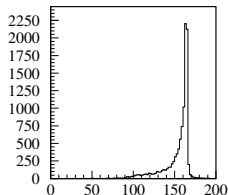
Sum of E_μ Angle between μ 's E_μ  $\cos(\Theta_\mu)$

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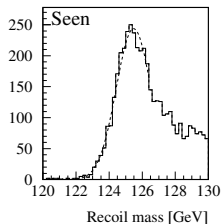
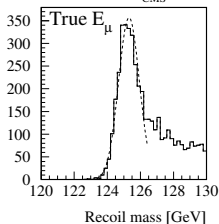
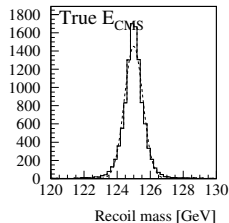
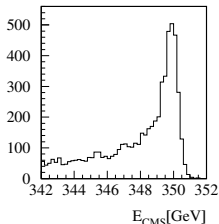
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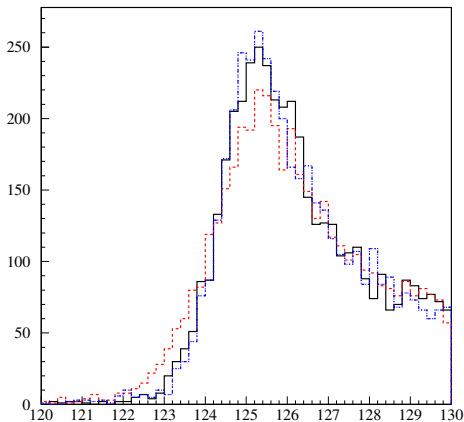
Higgs recoil-mass @ 350 GeV: The recoil mass

- $E_{CMS} \neq \text{nominal}$, due to beam spectrum.
- Assume E_{CMS} known \Rightarrow see effect of **detector alone**.
- Or: Assume μ :s perfectly measured \Rightarrow see effect of **beam-spectrum alone**.
- Fold the two: the **observable distribution**.



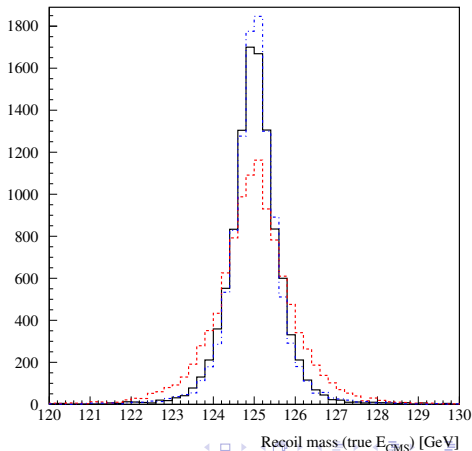
Higgs recoil-mass @ 350 GeV: The good, the bad, the ugly

- This shows observable recoil-mass for the nominal ILD (black), the worst case (red) and the best case (blue)
- ... and this shows the case if E_{CMS} would be known, ie. the pure detector effect.



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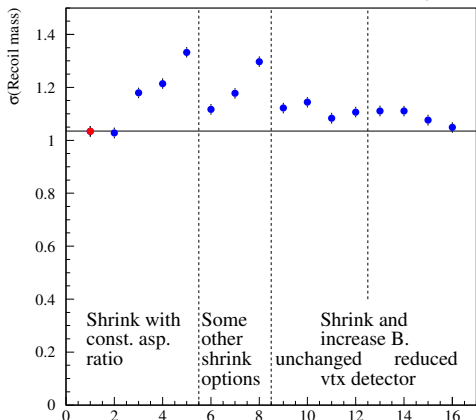
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Higgs recoil-mass @ 350 GeV: fits with different options

To substantiate: Fit the recoil-mass (Gaussian from 120 to 126.5 GeV in the observable case, free Gaussian in the E_{CMS} -known case)

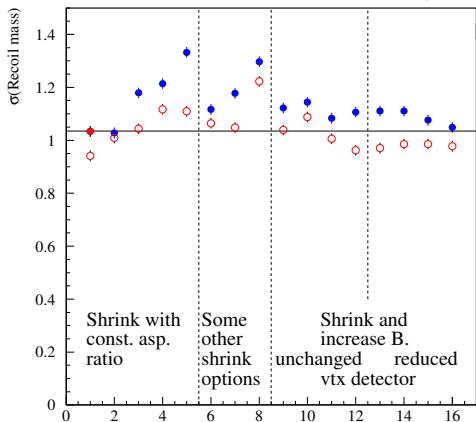
- $\sigma_{M-recoil}$ for a representative set of options.
- Same, but in the “4T” case.
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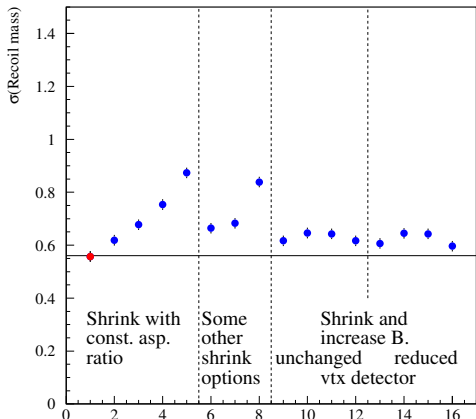
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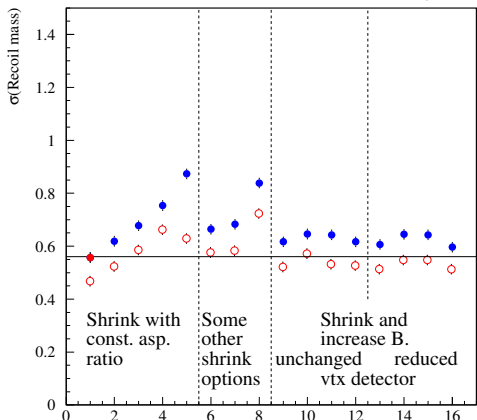
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Conclusions

- A large number of possible ways to **reduce the size** of the ILD tracking system were studied.
- A number of **auxiliary changes** that a reduced size would allow for were also studied: Increased **B-field**, changes of the **inner radius** of the TPC and/or the vertex detector.
- The errors of the **basic helix parameters** were evaluated for all of these scanning in momentum at a few fixed θ angles or in θ at a few fixed momenta.
- In addition, the precision on M_H from the **recoil-mass** method was evaluated with with SGV for a sub-set of the options.
- All taken together, the option with $R_{TPC}=160$ cm and/or $Z_{max,TPC}$ between 230 and 190 cm would be a viable option, **provided the B-field is increased**. Corresponds to $N_{paddrows} = 225 \rightarrow \sim 190$.
- These values correspond to a reduction of the **ECal area of between 15 and 25 %**, and a **B-field between 3.7 and 3.9 T** (or 4.2 to 4.4, if the baseline field would be increased to 4 T)

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