

Top FCNC decays: review of theoretical expectations

Aleksander Filip Żarnecki



Faculty of Physics, University of Warsaw

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Standard Model

On the tree level only charged current top decays are allowed in the Standard Model

$$t \rightarrow W^+ b \quad \text{dominant, BR} = 91\%$$

$$t \rightarrow W^+ s/d \quad \text{CKM suppressed}$$

FCNC top decays are only possible on loop level.

Four two-particle final states can be considered:

$$t \rightarrow q\gamma, qZ, qg, qH \quad q = u, c$$

Current experimental limits are (RPP2014):

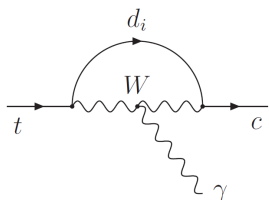
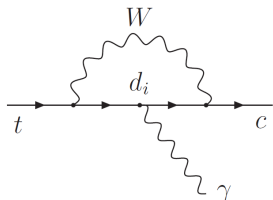
$$BR(t \rightarrow \gamma q) < 5.9 \cdot 10^{-3} \quad 95\% \text{ CL}$$

$$BR(t \rightarrow Z q) < 2.1 \cdot 10^{-3}$$

Leading order diagrams for FCNC decay $t \rightarrow c \gamma$

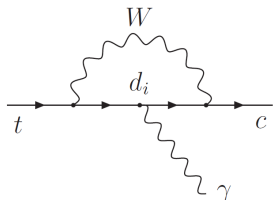
When neglecting down quark masses, the decay amplitude is suppressed (GIM):

$$\mathcal{M} \sim \sum_{d_i} V_{td_i}^* V_{cd_i} = 0$$



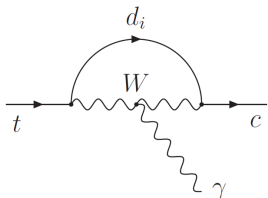
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However, taking into account quark masses, GIM cancelation is not perfect

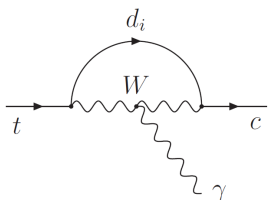
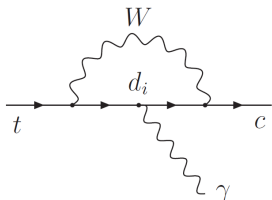


$$\mathcal{M} \sim \sum_{d_i} V_{td_i}^* V_{cd_i} \mathcal{F}(x_{d_i})$$

$$x_{d_i} = \frac{m_{d_i}^2}{M_W^2}$$



Leading order diagrams for FCNC decay $t \rightarrow c \gamma$



Assuming $m_d = m_s \ll m_b$ the leading contribution is:

$$\mathcal{M} \sim V_{tb}^* V_{cb} [\mathcal{F}(x_b) - \mathcal{F}(0)]$$

Resulting decay width:

$$\Gamma(t \rightarrow c \gamma) \sim |V_{bc}|^2 \alpha_{em}^3 m_t \left(\frac{m_b}{M_W} \right)^4$$

Double suppression due to

- CKM: $|V_{bc}| \sim 0.04$
- GIM: $\frac{m_b}{M_W} \sim 0.04$

Standard Model expectations for the FCNC top decays (Snowmass 2013):

$$BR(t \rightarrow c \gamma) \sim 5 \cdot 10^{-14}$$

$$BR(t \rightarrow c Z) \sim 1 \cdot 10^{-14}$$

$$BR(t \rightarrow c g) \sim 5 \cdot 10^{-12}$$

$$BR(t \rightarrow c H) \sim 3 \cdot 10^{-15}$$

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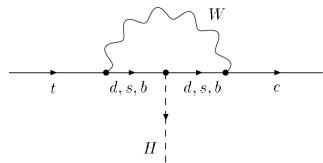
$$BR(t \rightarrow c g) \sim 5 \cdot 10^{-12}$$

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Only for $t \rightarrow c H$ channel, GIM mechanism is not applicable (in one of the diagrams) due to Higgs coupling proportional to mass.

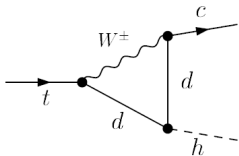
But the contribution of this diagram is still suppressed by $\frac{m_b}{M_W}$ (Higgs coupling)



Two Higgs Doublet Model

Probably the simplest possible extension of the SM.

Decay channel $t \rightarrow c h$ is affected by modified Higgs couplings:



hdd :

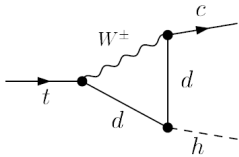
$$g = g_{SM} \times (\sin(\beta - \alpha) - \tan \beta \cdot \cos(\beta - \alpha))$$

possible enhancement at large $\tan \beta$

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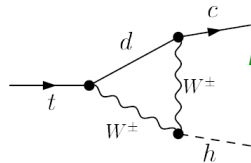
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hWW :

$$g = g_{SM} \times \sin(\beta - \alpha)$$

no enhancement possible

Two Higgs Doublet Model

New contributions to $t \rightarrow c h$ (as well as to $t \rightarrow c\gamma$, cZ , cg) from diagrams with H^\pm in the loop (instead of W^\pm).

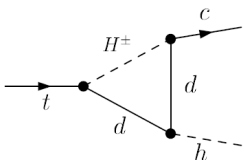
In case of 2HDM(II) (as an example):

$H^+ b \bar{t} :$
 $\frac{ig}{2\sqrt{2}M_W} V_{tb} [m_b(1 + \gamma_5)\tan\beta + m_t(1 - \gamma_5)\cot\beta]$
 $\tan\beta$ in all 3 vertexes !

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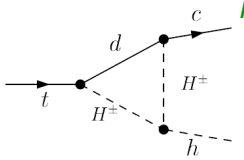


The diagram shows a top quark (t) entering from the left, splitting into a charm quark (c) and a down quark (d). A dashed line representing a charged Higgs boson (H^\pm) forms a loop with the down quark line. The loop then splits into a charm quark (c) and a Higgs boson (h).

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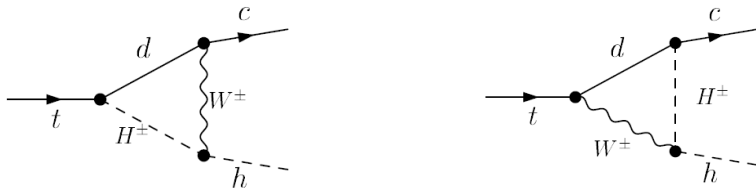
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$H^\pm H^\pm h$:

$$\frac{-ig}{M_W} \left[\frac{M_h^2}{\sin 2\beta} (\cos^3\beta \cos\alpha - \sin^3\beta \sin\alpha) - M_{H^\pm}^2 \sin(\alpha - \beta) \right]$$

enhancement possible for both large and small $\tan\beta$

One also has to consider diagrams with both H^\pm and W^\pm :



In the “standard” 2HDM scenarios, loop contributions can be enhanced significantly. However, FCNC remain suppressed at the tree level due to assumed flavour diagonal Higgs couplings.

However, one can also consider “non standard” scenarios, as 2HDM(III) or “Top 2HDM”, where one of Higgs doublets couple to top quark only, where tree level FCNC couplings are possible!...

Expected maximal branching ratios for different scenarios

Significant differences between different papers - overall limit ranges given

Model	$BR(t \rightarrow c h)$	$BR(t \rightarrow c \gamma)$	$BR(t \rightarrow c g)$	$BR(t \rightarrow c Z)$
SM	$3 \cdot 10^{-15}$	$5 \cdot 10^{-14}$	$5 \cdot 10^{-12}$	10^{-14}
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\mathcal{R} SUSY	$10^{-9} - 10^{-6}$	$10^{-9} - 10^{-5}$	$10^{-5} - 10^{-3}$	$10^{-6} - 10^{-4}$

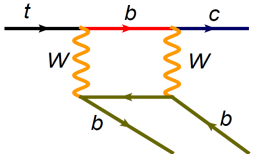
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\mathcal{R} SUSY	$10^{-9} - 10^{-6}$	$10^{-9} - 10^{-5}$	$10^{-5} - 10^{-3}$	$10^{-6} - 10^{-4}$
Little Higgs	10^{-5}	$1.3 \cdot 10^{-7}$	$1.4 \cdot 10^{-2}$	$2.6 \cdot 10^{-5}$
Quark Singlet	$4.1 \cdot 10^{-5}$	$7.5 \cdot 10^{-9}$	$1.5 \cdot 10^{-7}$	$1.1 \cdot 10^{-4}$
Randal-Sundrum	10^{-4}	10^{-9}	10^{-10}	10^{-3}

Remark

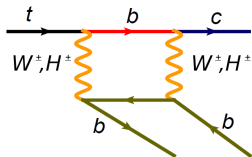
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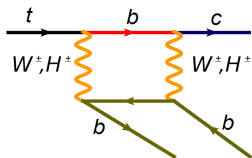


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However, for other models it can be important and has to be taken into account.

For hadronic channels, one also has to consider dominant $t \rightarrow b W^+$ decay.

For the $cb\bar{b}$ final state all contributions should be considered together:

- $t \rightarrow c h \rightarrow c b\bar{b}$
- $t \rightarrow c Z \rightarrow c b\bar{b}$
- $t \rightarrow b W^+ \rightarrow b c\bar{b}$
- $t \rightarrow c b\bar{b}$ (non-resonant)

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Decay $t \rightarrow c h$ in 2HDM is an interesting scenario, allowing for very large enhancement of signal with minimal extension of the model and no new final states - well suited for the case study.

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Test configuration of the model:

- $m_{h_1} = 125$ GeV
- $\text{BR}(t \rightarrow ch_1) = 10^{-3}$
- $\text{BR}(h \rightarrow b\bar{b}) = 100\%$

Generated samples at $\sqrt{s}=500$ GeV

- $e^+e^- \rightarrow t\bar{t}$ (2HDM/SM)
- $e^+e^- \rightarrow ch_1\bar{t}, t\bar{c}h_1$ (2HDM)
- $e^+e^- \rightarrow cb\bar{b}l^+\nu$ (SM)

2HDM events generated with CIRCE1 spectra, but no ISR (!)

All results on parton level only! No hadronization or detector effects.

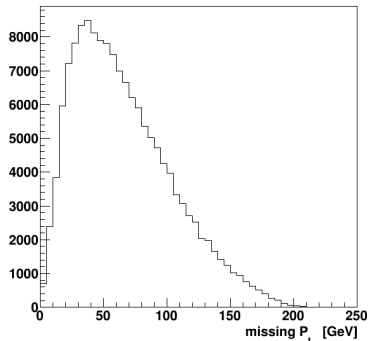
All results VERY PRELIMINARY

Final state selection

Considered final state: 4 jets + lepton + missing p_t

$t \rightarrow ch_1 + \text{higgs decay to } b\bar{b} + \text{leptonic } W^\pm \text{ decay for second top}$

- Missing $p_t > 15 \text{ GeV}$

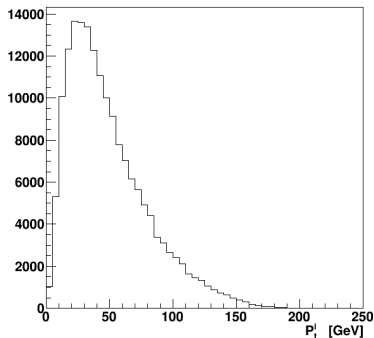


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- Lepton with $p_t > 15$ GeV
 $|\cos\theta_l| < 0.995$

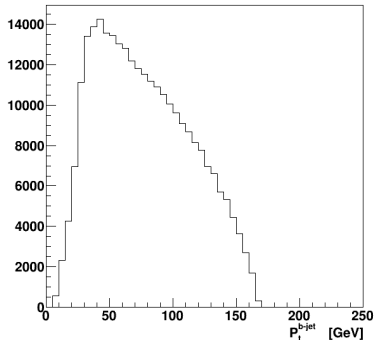


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- **3 b-jets** with $p_t > 10$ GeV
 $|\cos\theta_b| < 0.975$

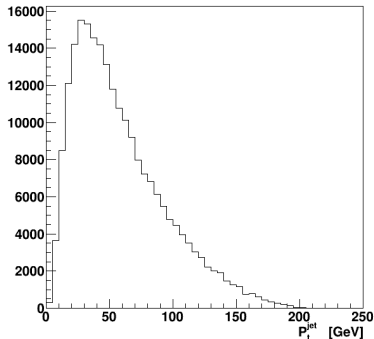


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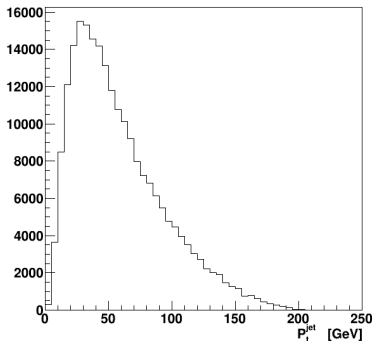


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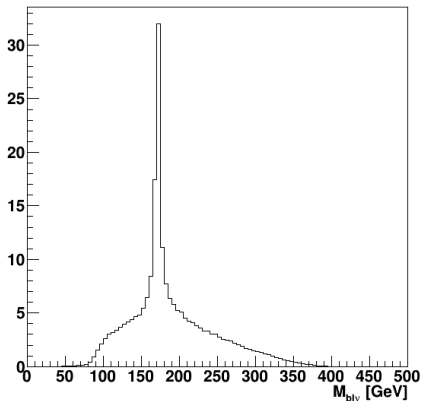
Cut efficiency for considered final state $\sim 75\%$ (without detector effects!)

Event selection

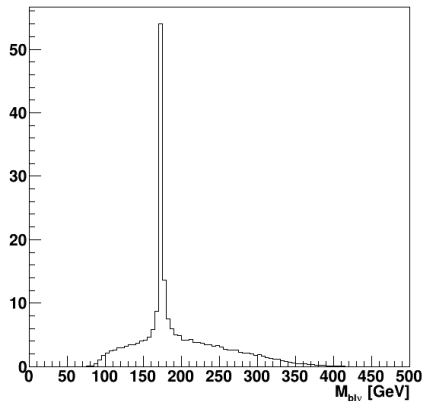
Top reconstruction

With 3 jets b-tagged, there are 3 possible quark combination.

Reconstruction of $t \rightarrow b\nu$



$$e^+e^- \longrightarrow cb\bar{b}bl^+\nu \text{ (SM)}$$



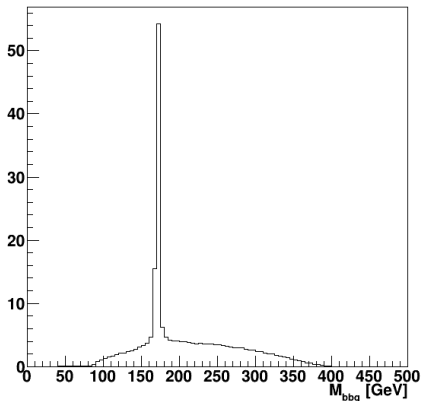
$$e^+e^- \longrightarrow ch_1\bar{t}, t\bar{c}h_1 \text{ (2HDM)}$$

Event selection

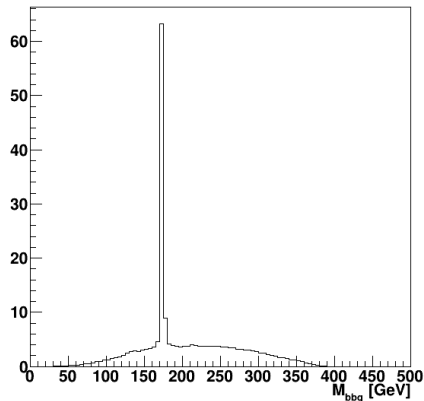
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Reconstruction of $t \rightarrow b\bar{b}c$



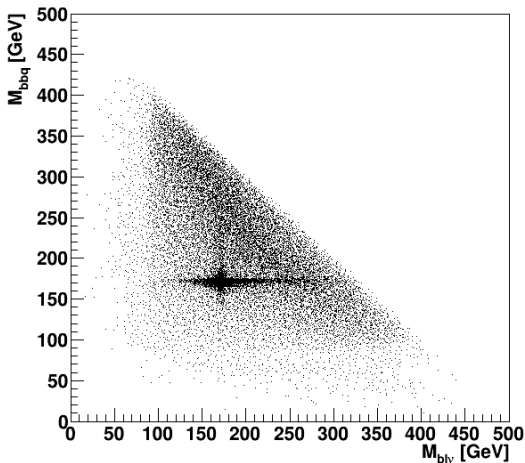
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Top reconstruction

Best matching of b-jets to top and anti-top can be selected by minimizing the χ^2 of mass constraint.

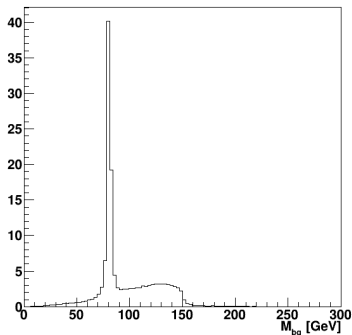


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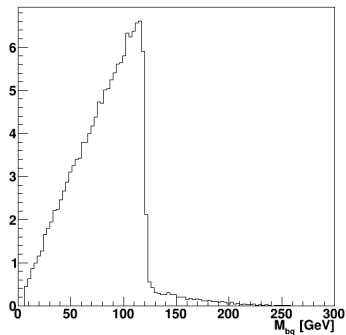
W^\pm veto

Irreducible SM background can be suppressed by reconstructing second W

Invariant mass of two jets from hadronic top decay (one b-tagged)



$$e^+e^- \longrightarrow cb\bar{b}b l^+\nu \text{ (SM)}$$



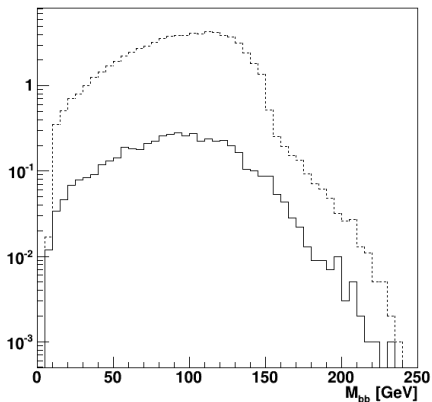
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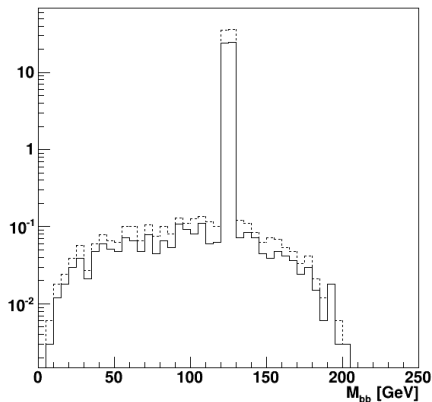
Candidate events

W^\pm veto used: events with $70 < M_{bq} < 90$ GeV rejected

Invariant mass of two b-jets before (dashed) and after (solid) W^\pm veto



$$e^+e^- \longrightarrow cb\bar{b}l^+\nu \text{ (SM)}$$



$$e^+e^- \longrightarrow ch_1\bar{t}, t\bar{c}h_1 \text{ (2HDM)}$$

Expected events

For 500 fb^{-1} , after taking into account $BR(h \rightarrow b\bar{b}) \approx 57\%$ for signal

	Background	Signal
	SM	2HDM(III)
All	272'000	300
Semi-leptonic	108'000	85
4 jets, 3 b-tags	73	75
W veto	5.0	50.6
H mass ± 10 GeV	0.8	48.9

Total of about 540 FCNC top decays expected in the considered model, 180 with leptonic decay of second top, 105 with $h \rightarrow b\bar{b}$.

\Rightarrow cut efficiency $\sim 45\%$ overall efficiency $\sim 10\%$

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Fully hadronic channel is more challenging, but with 3 b-tagged jets background suppression should still be possible.

Selected references for theoretical BR estimates in different models:

- *Snowmass 2013: Top quark working group report*, arXiv:1311.2028
- B.Yang, N.Liu¹, J.Han, arXiv:1308.4852
- I.Bauma, G.Eilamb, S.Bar-Shalomb, arXiv:0802.2622
- Hou Hong-Sheng, arXiv:hep-ph/0703067
- J.J.Cao, G.Eilam, M.Frank, et al., arXiv:hep-ph/0702264
- J.L.Diaz-Cruz, C.Pagliarone, arXiv:hep-ph/0612120
- J.A.Aguilar-Saavedra, arXiv:hep-ph/0409342
- S.Bejar, J.Guasch, J.Sola, arXiv:hep-ph/0101294
- Barbara Mele, arXiv:hep-ph/0003064