Top electroweak couplings study using di-muonic state at \sqrt{s} = 500 GeV, ILC Full ILD detector simulation

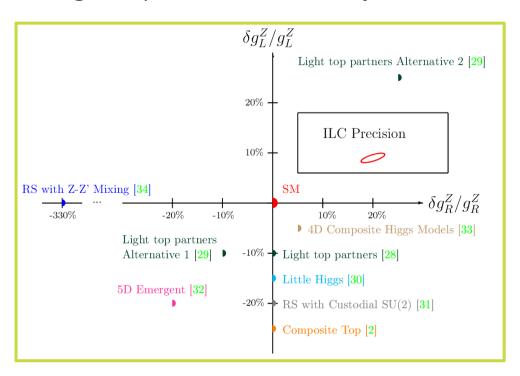
50th General meeting

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Top electroweak couplings

The top quark mass is comparable with the electroweak symmetry breaking scale. One can speculate that top quark plays a special role for the EWSB, for example such composite models. Therefore top quark electroweak couplings are good probes for New Physics.



Plot shows the predicted deviations from the Standard model of Z^0 couplings to t_L and t_R in composite models

Precision expected at the ILC will allow to distinguish between models. arXiv:1505.06020 [hep-ph]

Matrix element method

The most efficient method when all the kinematics can be reconstructed.

- Results of previous study show that 10 form factors can be fitted simultaneously at less than a percent precision.

Statistical uncertainties and correlation with the SM LO as normalization										
					Khaim	FK Kurik	ara Le Di	iherder: ar	Xiv:1503:0	4247
					Kheim, E.K. Kurihara, Le Diberder: arXiv:1503:04247					
$\mathcal{R} \in \delta \tilde{F}_{1V}^{\gamma}$	$\Re \delta \tilde{F}_{1V}^Z$	$\Re \delta \tilde{F}_{1A}^{\gamma}$	$\Re \delta \tilde{F}_{1A}^{Z}$	\mathcal{R} e $\delta \tilde{F}_{2V}^{\gamma}$	$\Re \delta \tilde{F}_{2V}^{Z}$	\mathcal{R} e $\delta \tilde{F}_{2A}^{\gamma}$	$\Re \delta \tilde{F}_{2A}^{Z}$	\mathcal{I} m $\delta \tilde{F}_{2A}^{\gamma}$	\mathcal{I} m $\delta \tilde{F}_{2A}^{Z}$	
0.0037	-0.18	-0.09	+0.14	+0.62	-0.15	0	0	0	0	
	0.0063	+.14	-0.06	-0.13	+0.61	0	0	0	0	
		0.0053	-0.15	-0.05	+0.09	0	0	0	0	
			0.0083	+0.06	-0.04	0	0	0	0	
				0.0105	-0.19	0	0	0	0	
					0.0169	0	0	0	0	
						0.0068	-0.15	0	0	
							0.0118	0	0	
								0.0069	-0.17	
									0.0100	
500 GeV&500 fb ⁻¹ Polarization 50/50 between $\pm 80\%$ and $\pm 30\%$										0 %

Emi Kou (LAL-Orsay) LFC 15, Trento, 7-11 Sep. 2015

This result is at parton level ignoring the detector effect, ISR and so on.

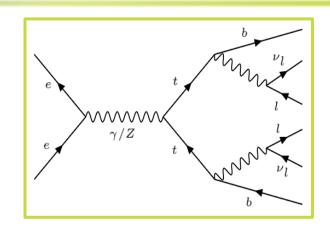
→ More realistic study is required!

Setting of this study

Sample

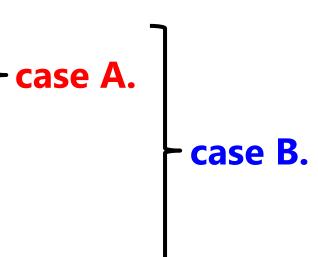
The top pair production di-muonic state;

$$t\bar{t} \rightarrow b\bar{b}\mu^+\mu^-\nu\bar{\nu}$$
 at $\sqrt{s}=500~{\rm GeV}$



Situation

- \checkmark The hadronization of b and \bar{b} quark
- ✓ The detector effects (= full simulation)
- ✓ ISR and beamsstrahlung
- ✓ Gluon emission from top quark
- √ γγ → hadrons background
- × Background events



Reconstruction

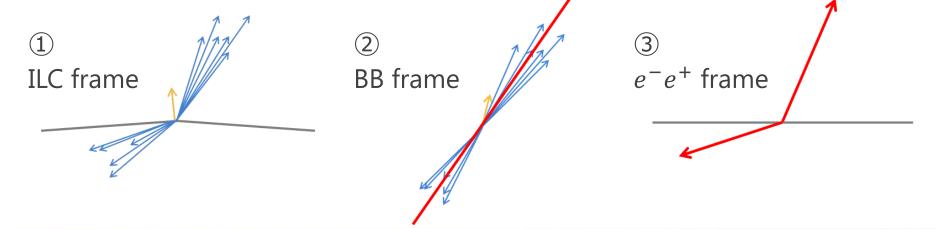
Process of the reconstruction

- Isolated leptons tagging
- B-jets reconstruction using the Thrust axis method
- Kinematical reconstruction

Thrust axis method

- Collect all hadronized particles in the ILC frame
- 2 Boost them into the BB frame and calculate thrust axis
- 3 Boost the vectors along the thrust axis into the rest frame of e^-e^+
- case A. using same parameters with ②

case B. considering ISR/BS effects using \rightarrow introduce the k_e -, k_e +



Considering ISR/BS effects

Collinear approximation:

Photons are emitted on the beam directions by ISR/BS

$$\vec{e} - \hat{\eta}_{e} - E_{e} - (1)$$

$$\vec{e}^{\,+} = \hat{\eta}_{e^+} E_{e^+} \tag{2}$$

with,

$$\hat{\eta}_{e^{-}} = (\sin \theta_c, 0, \cos \theta_c) \tag{3}$$

$$\hat{\eta}_{e^+} = (\sin \theta_c, 0, -\cos \theta_c) \tag{4}$$

$$E_{e^{\pm}} = E = 250 \text{ GeV} \tag{5}$$

where θ_c is the beam crossing angle, $\theta_c = 7$ mrad.

In this approximation, the directions are not changed but only the energies are changed. Then the electron and positron thre-momenta become:

$$(\vec{e}^{\,-})^* = \hat{\eta}_{e^-} E_{e^-}^* = \hat{\eta}_{e^-} E(1 - k_{e^-}) \text{ with } k_{e^-} = \frac{E - E_{e^-}^*}{E}$$

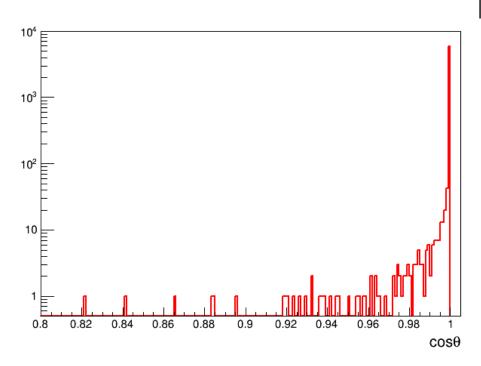
$$(\vec{e}^{+})^{*} = \hat{\eta}_{e^{+}} E_{e^{+}}^{*} = \hat{\eta}_{e^{+}} E(1 - k_{e^{+}}) \text{ with } k_{e^{+}} = \frac{E - E_{e^{+}}^{*}}{E}$$

$$k_{e^{-}} = \frac{E - E_{e^{-}}^{*}}{E} \tag{6}$$

$$k_{e^{+}} = \frac{E - E_{e^{+}}^{*}}{E} \tag{7}$$

where $E_{e^{\pm}}^{*}$ is the energy of electron or positron just before collision.

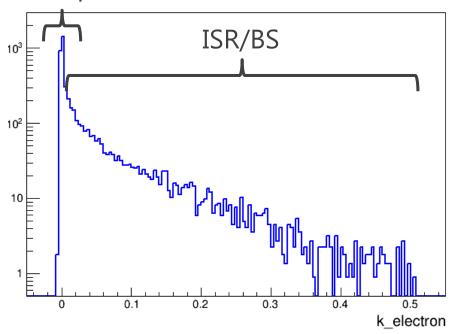
Distribution of $\cos\theta$ and k_{e^-} (MC truth)



 $\cos \theta$

 $(\theta : angle between e^{-*} and beam direction)$

Beam spread



$$\frac{250 - E_{e^-}^*}{250} (= k_{e^-})$$

(Indeed, it is not same as k_e because $\cos\theta$ has width)

Kinematical reconstruction

Strategy of the kinematical reconstruction

- There are 8 unknown kinematics in this state
 - (= the momenta of two neutrinos and energy of two b-jets)

and (k_{e^-}, k_{e^+}) in case B.

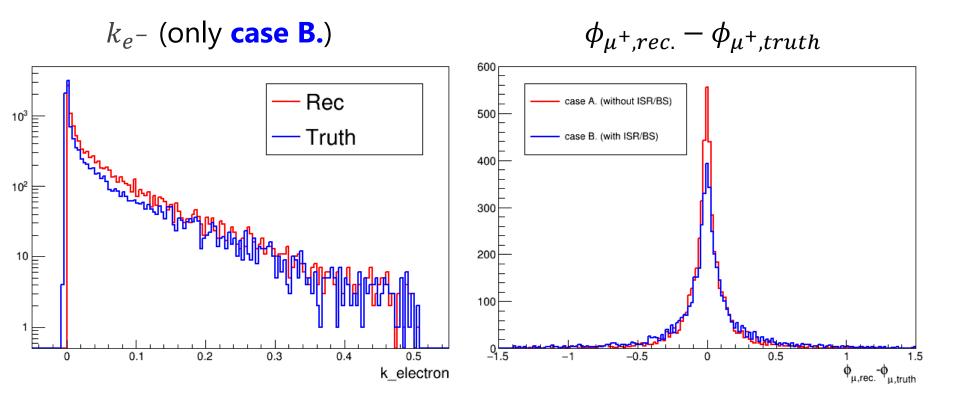
- Impose 8 constraints (= initial state constraints and $m_t, m_{\bar{t}}, m_{W^+}, m_{W^-}$)
 - \rightarrow Solutions are obtained in terms of $(\theta_t, \phi_t) \rightarrow (\theta_t, \phi_t, m_t, m_{\bar{t}}, m_{W^+}, m_{W^-})$

But the equation is nonlinear. Furthermore an ambiguity of b-charge remains.

- → Typically 4 solutions per event.
- Select the optimal solution

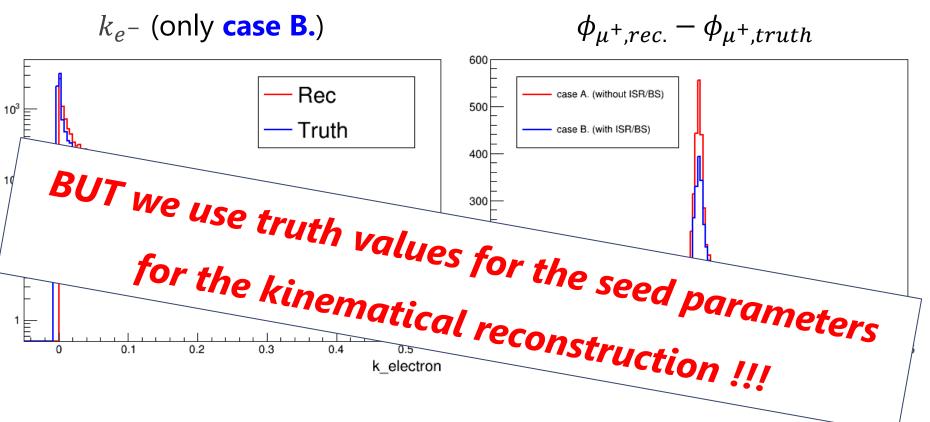
Compare $E_b^{\text{meas.}}$ (by thrust axis method) and $E_b^{\text{rec.}}$ (by kinematical reconstruction).

Kinematical reconstruction: Results



- **case B.** is not that bad comparing with **case A.**
 - → Need to check if it is enough for the analysis using MEM

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Status of the 10 form factors fit (case A.)

Δ_{F1}	-0.0067 ± 0.0082				
Δ_{F2}	0.035 ± 0.017				
Δ_{F3}	-0.056 ± 0.012				
Δ_{F4}	0.035 ± 0.018				
Δ_{F5}	-0.022 ± 0.026				
Δ_{F6}	0.042 ± 0.045				
Δ_{F7}	-0.0081 ± 0.015				
Δ_{F8}	0.010 ± 0.032				
Δ_{F9}	0.013 ± 0.024				
Δ_{F10}	-0.010 ± 0.022				

(Preliminary)

5000 events, after cut on the $\chi^2_{\rm tot.}$ to keep ~83% of the events without ISR/BS

Some small biases are observed (eg. Δ_{F3}) at few percent level

→ No show stopper yet !!!

→ Use DBD samples for the matrix element method

Summary & Plan

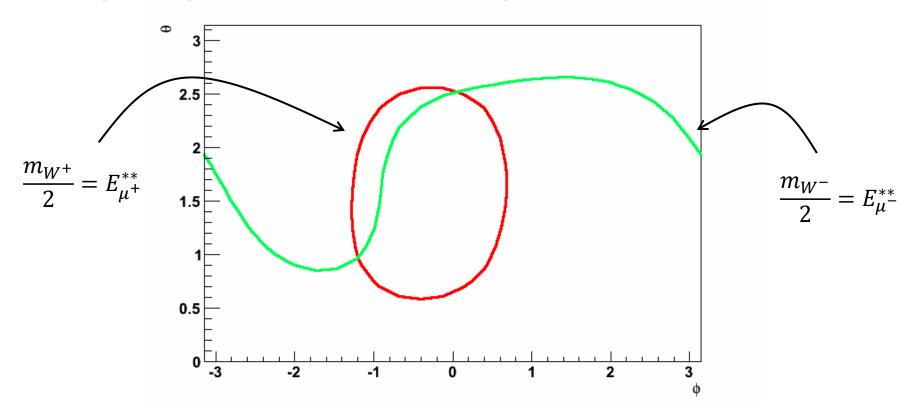
- Started realistic study using DBD samples
 - Accuracy of kinematical reconstruction is not that bad comparing the samples without ISR/BS
- Analysis with Matrix element method
 - 10 form factors can be fitted at percent precision in case A.
 - → Apply Matrix element method on DBD samples (case B.)

- Visit François at LAL to learn MEM and statistics in March
- Borys, master student from Ukraine will join us to work on seeds issue at LAL.

Back up

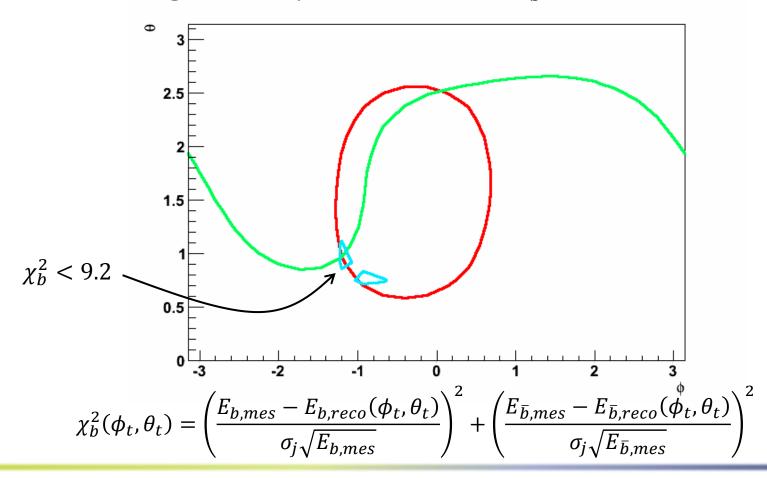
Kinematical constraints

In the W rest frame, the energy of isolated lepton is equal to $m_W/2$ (with ignoring ISR and bremsstrahlung)



Measurements of b-quark energies

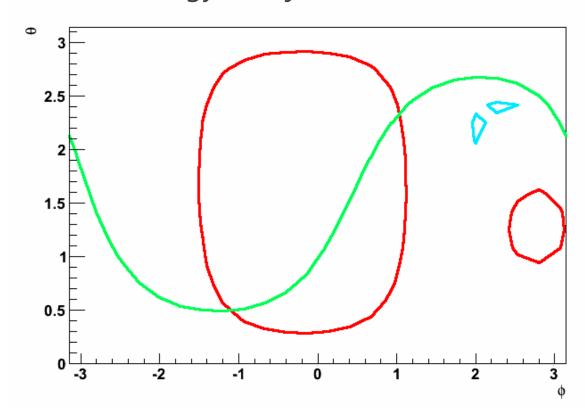
To select a right solution, we can use the measurement of b-jets energy. (Because this figure is at parton level, the χ_b^2 doesn't make sense.)



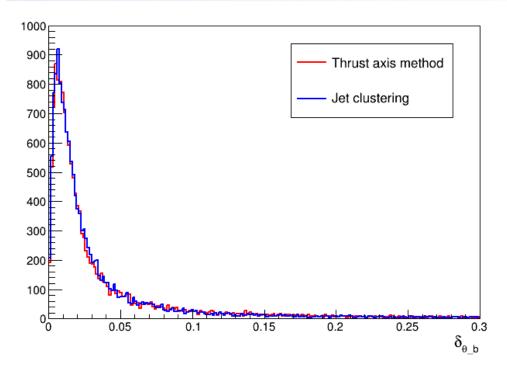
16

Miss combination of b-quarks

When we use the anti-b direction for the top reconstruction, the measurements of energy of b-jets excludes this combination.



Comparison Thrust axis method and Jet clustering



The angle between truth direction and reconstructed direction of b quark

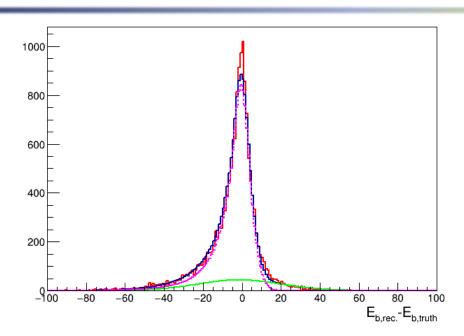
Red: Thrust axis method

Blue: Jet clustering (LCFIPlus)

Two methods produce almost same precision for direction of b-quark.

→We select the thrust axis method for this study so far.

Assessment of energy of b-jet



Deviation of energy of b-jet

(using thrust axis method)

Red: Original

Blue: Fitted

(Blue is a sum of Pink and Green)

To estimate the b-jet energy resolution we use multiple Crystal Ball functions for fitting.

$$CB(x|\alpha, n, \mu, \sigma) = \begin{cases} N \exp\left(-\frac{(x-\mu)^2}{2\sigma^2}\right), & \frac{x-\mu}{\sigma} > -\alpha \\ N \cdot A \cdot \left(B - \frac{x-\mu}{\sigma}\right)^{-n}, & \frac{x-\mu}{\sigma} < -\alpha \end{cases}$$

So far we use $\sigma_{iet}E_b^K$ with two parameters σ_{iet} and K for σ .

χ^2 algorithm : General

1. Define the χ_{μ}^2 ;

$$\chi_{\mu}^{2} = \chi_{\mu^{+}}^{2} + \chi_{\mu^{-}}^{2}, \quad \chi_{\mu^{\pm}}^{2} = \left(\frac{E_{\mu^{\pm}}^{**}(\theta_{t},\phi_{t},m_{t},m_{ar{t}},m_{W^{+}},m_{W^{-}}) - m_{W^{\pm}}/2}{\sigma[E_{\mu^{\pm}}^{**}]}\right)^{2}$$

The energy of μ^{\pm} in the W^{\pm} rest frame, $E_{\mu^{\pm}}^{**}$, must be equal to $m_{W^{\pm}}/2$.

2. Define the δ_h^2 ;

$$\delta_b^2 = -2\log L_b - 2\log L_{\overline{b}}, \ L_b = \operatorname{CB}\left(E_b^{\text{meas.}} - E_b^{\text{rec.}}(\theta_t, \phi_t, m_t, m_{\overline{t}}, m_{W^+}, m_{W^-})\right)$$

The likelihood function is obtained from the assessment of b-jets energy.

3. Compound $\chi^2_{\mathrm{tot.}}$; $\chi^2_{\mathrm{tot.}} = \chi^2_{\mu} + \delta^2_b$

One minimizes the $\chi^2_{\text{tot.}}$ to obtain the optimal solution; $(\theta_t, \phi_t, m_t, m_{\bar{t}}, m_{W^+}, m_{W^-})$.

χ^2 algorithm : Optional

The direction of b-jets is obtained by thrust axis method.

- \rightarrow Add 4 angles $(\theta_b, \phi_b, \theta_{\bar{b}}, \phi_{\bar{b}})$ to the minimization parameters
- \rightarrow Add constraints of 4 angles $(\theta_b, \phi_b, \theta_{\bar{b}}, \phi_{\bar{b}})$ to $\chi^2_{\text{tot.}}$ as follows;

$$\chi_{\text{direction}}^2 = \chi_{\theta_b}^2 + \chi_{\phi_b}^2 + \chi_{\theta_{\overline{b}}}^2 + \chi_{\phi_{\overline{b}}}^2, \qquad \chi_{\theta_b}^2 = \left(\frac{\theta_b^{\text{meas.}} - \theta_b}{\sigma[\theta_b^{\text{meas.}}]}\right)^2$$

 $(\chi^2_{\phi_b}, \chi^2_{\theta_{\overline{b}}}, \chi^2_{\phi_{\overline{b}}}$ are same as $\chi^2_{\theta_b})$

$$(\chi^2_{\text{tot.}})' = \chi^2_{\text{tot.}} + \chi^2_{\text{direction}}$$

We can use $(\chi^2_{\text{tot.}})'$ instead of $\chi^2_{\text{tot.}}$ to get the optimal solution, which is written in $(\theta_t, \phi_t, m_t, m_{\bar{t}}, m_{W^+}, m_{W^-}, \theta_b, \phi_b, \theta_{\bar{b}}, \phi_{\bar{b}})$

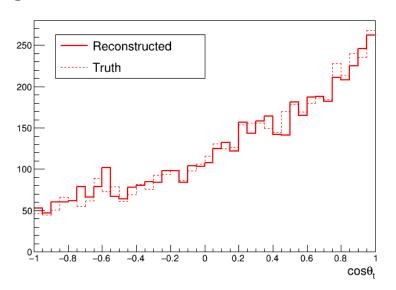
Kinematical reconstruction: Results

Reconstructed particles → 9 helicity angles :

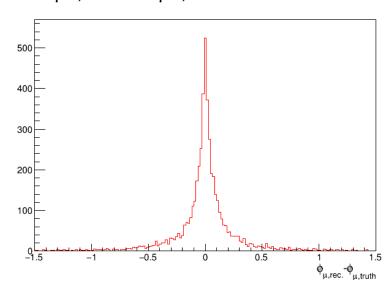
$$\cos heta_t$$
 , $\cos heta_{W^+}$, ϕ_{W^+} , $\cos heta_{\mu^+}$, ϕ_{μ^+} , $\cos heta_{W^-}$, ϕ_{W^-} , $\cos heta_{\mu^-}$, ϕ_{μ^-}

(→ Matrix element squared → Fit Form Factors)

eg.)
$$\cos \theta_t$$



eg.)
$$\phi_{\mu^+,rec.} - \phi_{\mu^+,truth}$$



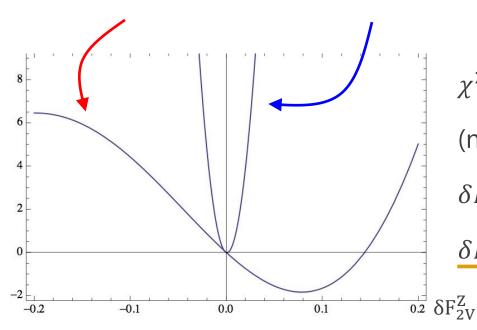
The ratio of wrong assignment of b-quark is only 2.1 %!

(cf. ~16 % for the semi-leptonic analysis, Yo Sato Top@LC 2016)

Analysis with Matrix element method

Illustration of analysis: 5000 unpolarized events

• using only $\cos \theta_t$ • using the complete set of 9 helicity angles



$$\chi^2(\delta F_{2V}^Z)$$
 function

(normalized such that $\chi^2(0) = 0$)

$$\delta F_{2V}^Z = 0.080 \pm 0.05$$
 by $\cos \theta_t$

$$\delta F_{2V}^Z = 0.001 \pm 0.01$$
 by the complete set

Number of events not used on purpose to compare the intrinsic power of these two to determine a single form factor, F_{2V}^Z .