

Determining Spin through Quantum Azimuthal-Angle Correlations

Matthew Buckley
Caltech
UC Berkeley/LBNL/IPMU

In Collaboration with H. Murayama, S. Choi, K. Mawatari

See previous works: 0711.0364, 0804.0476
with Heinemann, Klemm, Murayama, Rental.

The LHC Era

- Finally have access to TeV-scale physics
 - Solution to the Hierarchy Problem?
 - Dark Matter?
⇒ New Particles
 - SUSY, Extra-Dimensions, Little Higgs? Something totally different?



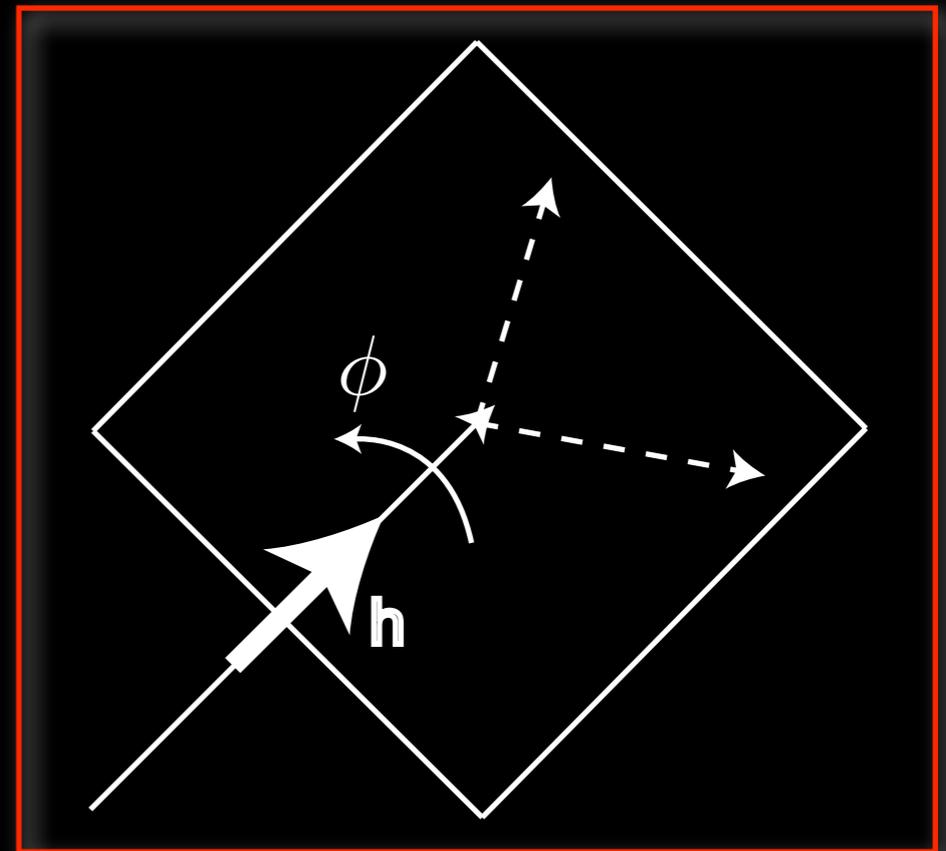
Spin Measurements

- Most techniques for next-generation colliders concentrate on distinguishing models:
 - Comparison of total cross section
 - Look for higher KK modes in UED
- At a linear collider can use threshold scans
 - Reconstruct production/polar decay angle
 - With long decay chains, can be used at hadron collider.

Spin and Quantum Interference

- Want a spin measurement with as few assumptions as possible.
- Back to Quantum Mechanics!
- Decay of particle with helicity h
 - Rotations about the z-axis (particle momentum) implies that

$$\mathcal{M}_{\text{decay}} \propto e^{iJ_z\phi} = e^{ih\phi}$$



Spin and Quantum Interference

- If particle is produced in multiple helicity states and then decays, then decay amplitudes interfere coherently:

$$\sigma \propto \left| \sum \mathcal{M}_{\text{prod.}} \mathcal{M}_{\text{decay}} \right|^2$$
$$\mathcal{M}_{\text{decay}}(h, \phi) = e^{ih\phi} \mathcal{M}_{\text{decay}}(h, \phi = 0)$$

- Sum runs over all helicities produced, generically $h = -s, \dots, s$ in which case

$$\sigma = A_0 + A_1 \cos \phi + \dots + A_n \cos n\phi, \quad n = 2s$$

New Physics

$$e^+e^- \rightarrow F^+F^- \rightarrow (\mu^+\chi)(\mu^-\chi) \rightarrow \mu^+\mu^- \cancel{E}$$

- i.e. $e^+e^- \rightarrow \tilde{\mu}^+\tilde{\mu}^- \rightarrow (\mu^+\tilde{\chi}_1^0)(\mu^-\tilde{\chi}_1^0)$ or
 $e^+e^- \rightarrow \mu_1^+\mu_1^- \rightarrow (\mu^+\gamma_1)(\mu^-\gamma_1)$

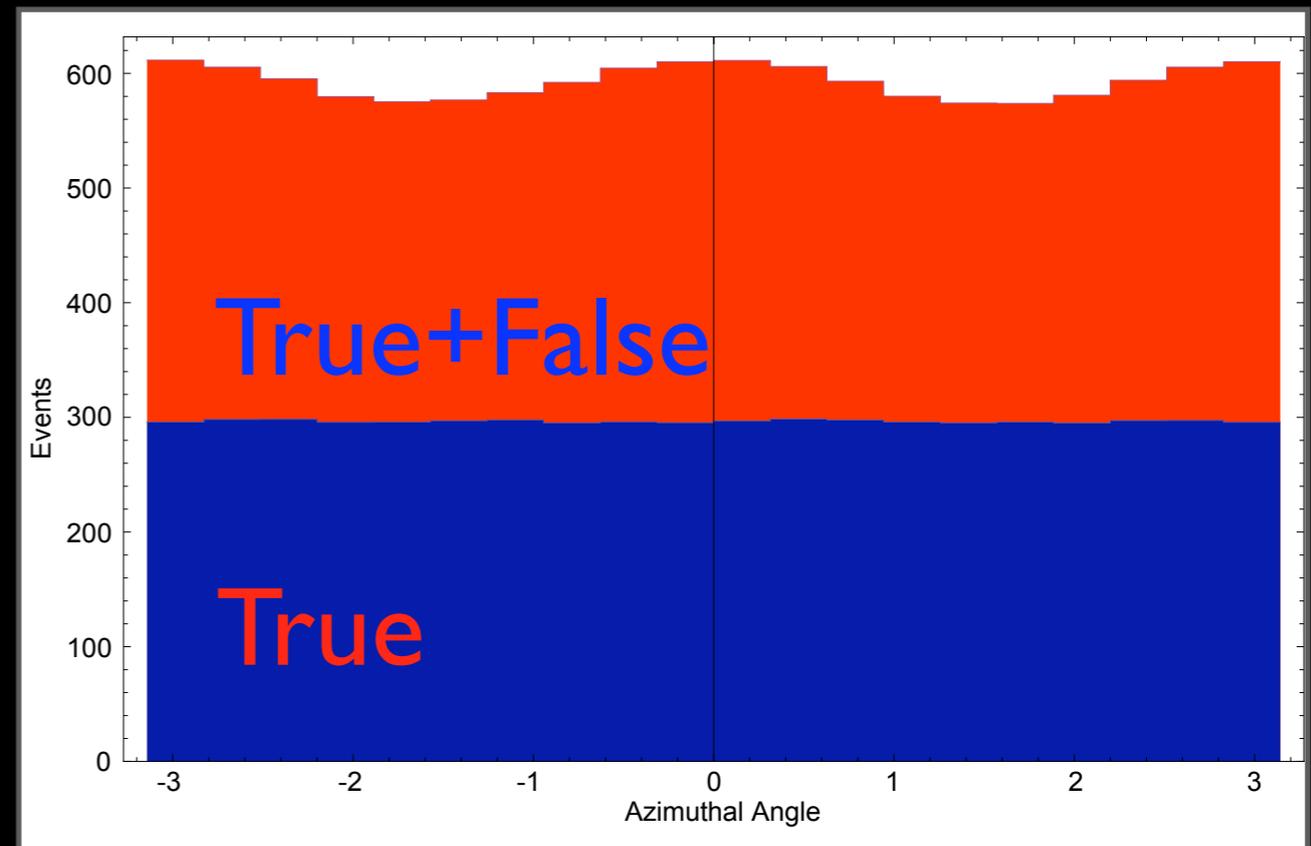
4+4 unknown momenta
-4 measured \cancel{p}
-4 mass relations

- 2-Fold ambiguity in reconstructing momenta & azimuthal angles ϕ_i (measured from production plane)

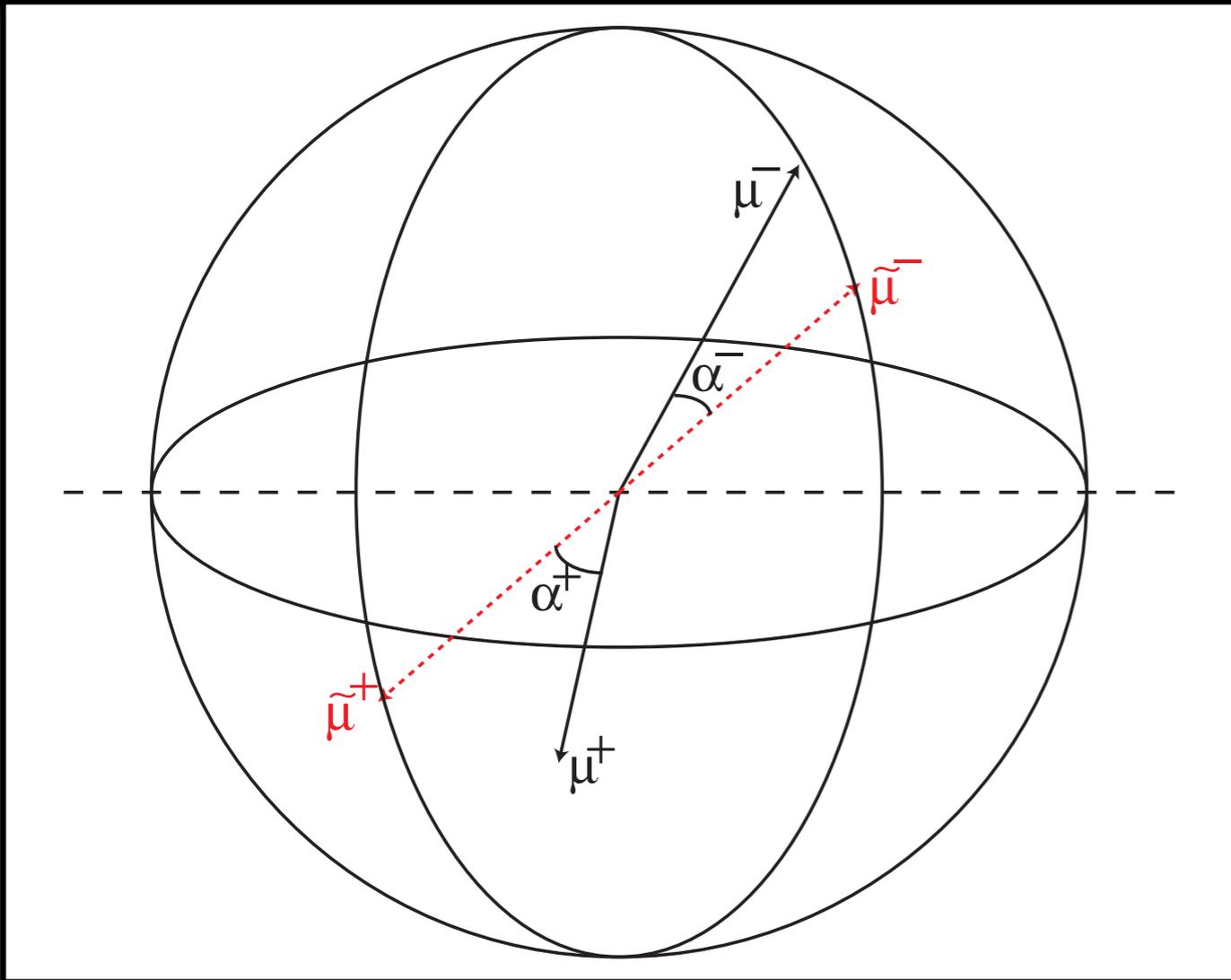
False Solutions

- Plotting both true and false distribution gives spurious high-frequency noise in distributions ϕ_i
- ϕ_1, ϕ_2 are not observable, but $\phi \equiv \Delta\phi$ is.

Scalar decay:



ϕ Reconstruction



Opening angles α^\pm
defined by

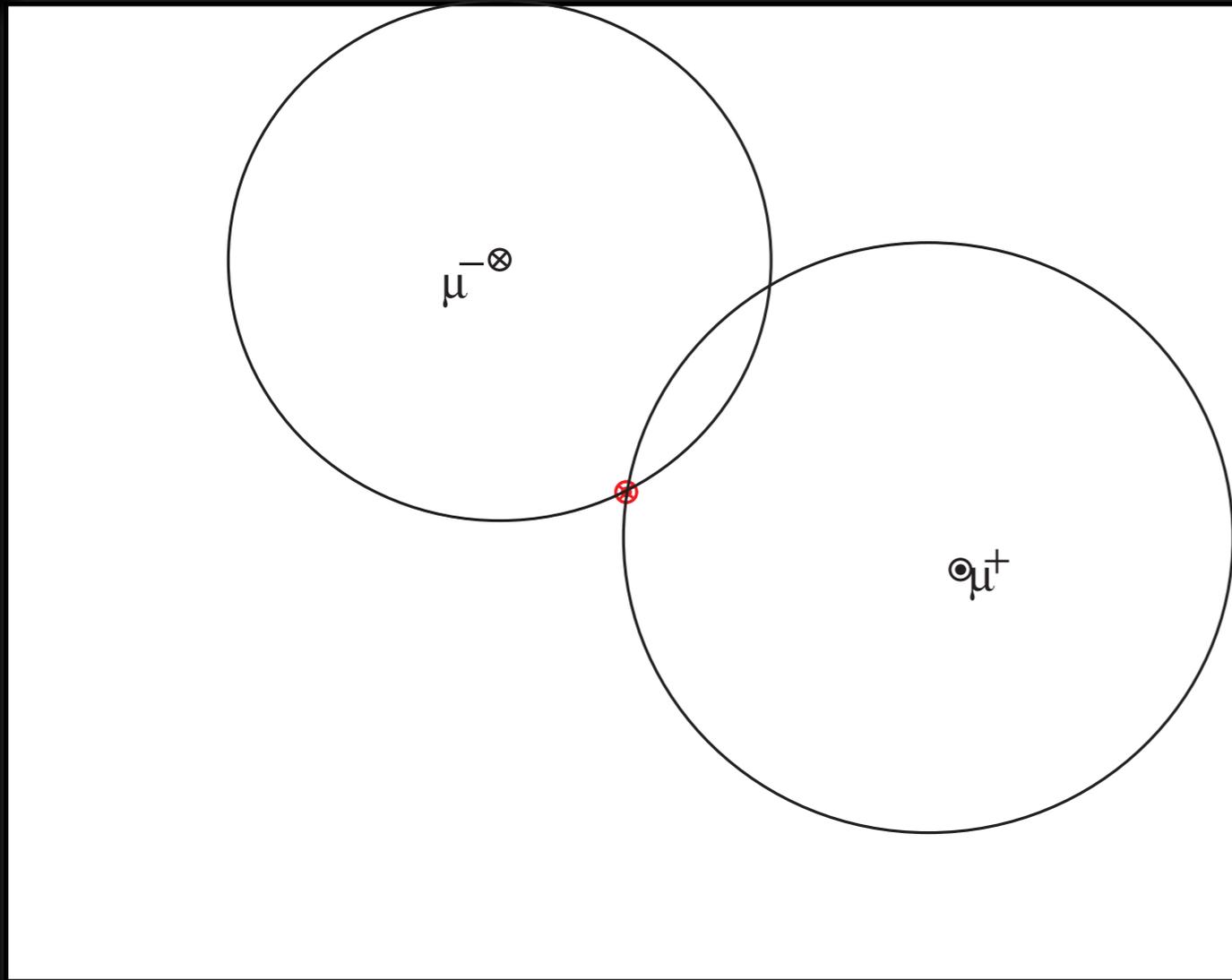
$$m_{\tilde{\mu}^\pm}^2 - m_{\tilde{\chi}}^2 = \sqrt{s} E_{\tilde{\mu}^\pm} (1 - \beta_{\tilde{\mu}^\pm} \cos \alpha^\pm)$$

Straightforwardly,

$$\phi \equiv \phi_T = \phi_F$$

Since interference argument
only needs *some* reference plane,
we expect same expansion
in $\cos n\phi_i$ and $\cos n\phi$

ϕ Reconstruction



Opening angles
defined by

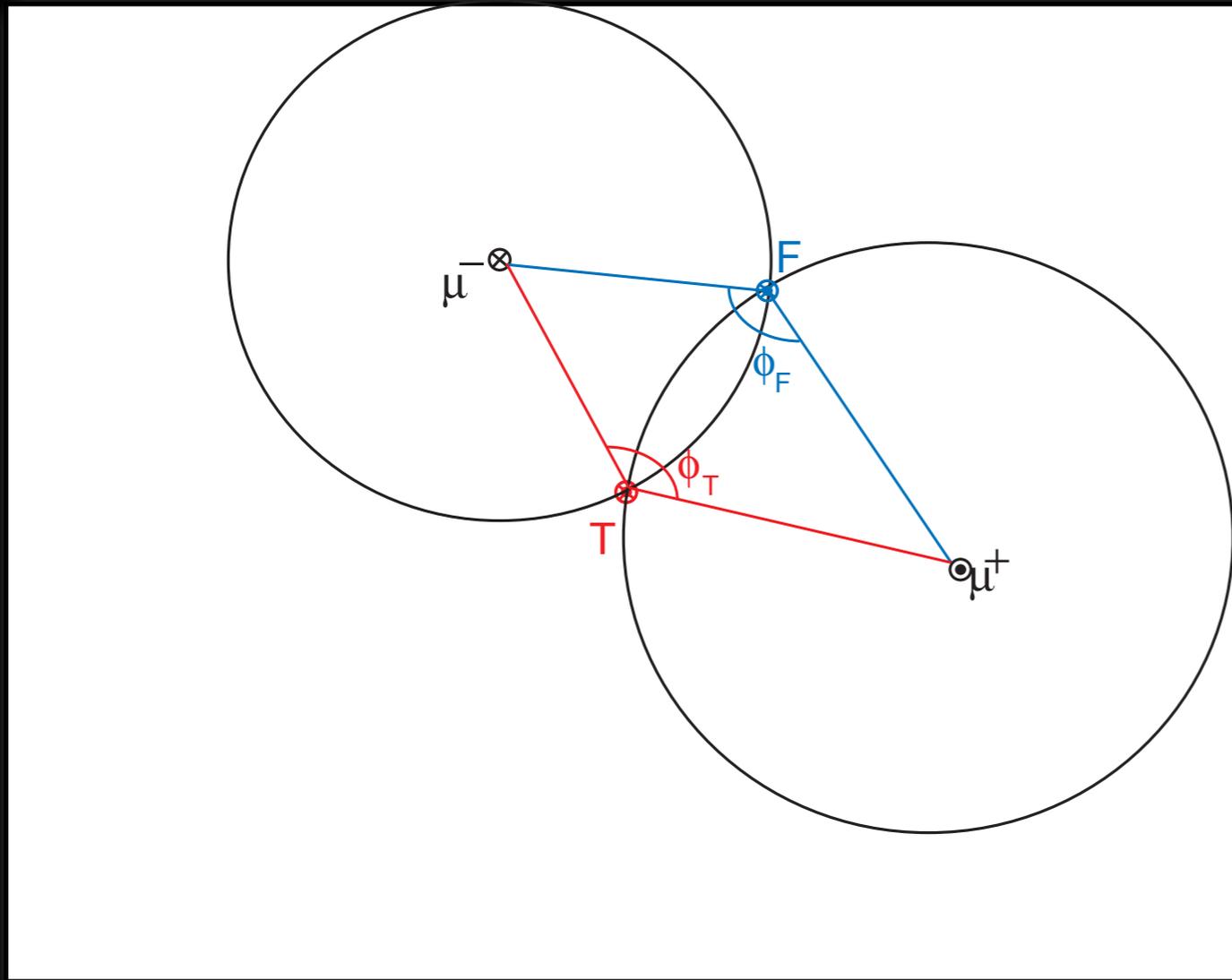
$$m_{\tilde{\mu}^{\pm}}^2 - m_{\tilde{\chi}}^2 = \sqrt{s} E_{\tilde{\mu}^{\pm}} (1 - \beta_{\tilde{\mu}^{\pm}} \cos \alpha^{\pm})$$

Straightforwardly,

$$\phi \equiv \phi_T = \phi_F$$

Since interference argument
only needs *some* reference plane,
we expect same expansion
in $\cos n\phi_i$ and $\cos n\phi$

ϕ Reconstruction



Opening angles
defined by

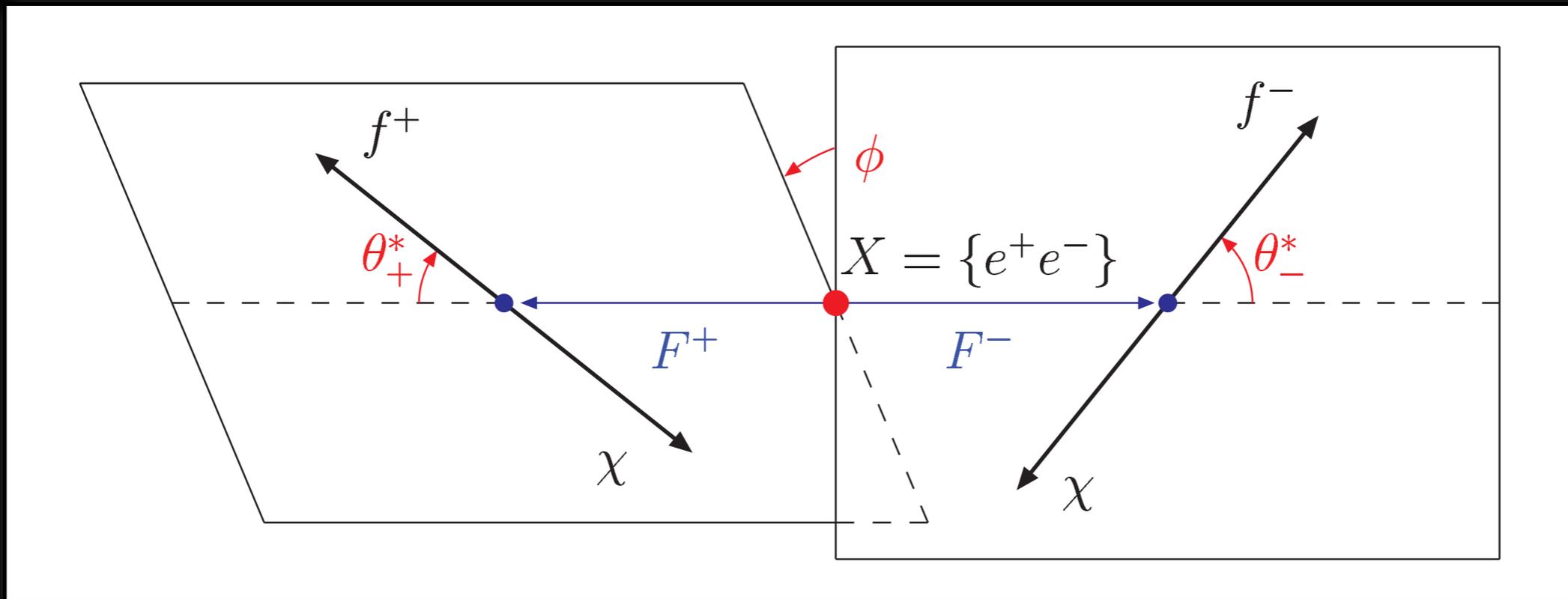
$$m_{\tilde{\mu}^\pm}^2 - m_{\tilde{\chi}}^2 = \sqrt{s} E_{\tilde{\mu}^\pm} (1 - \beta_{\tilde{\mu}^\pm} \cos \alpha^\pm)$$

Straightforwardly,

$$\phi \equiv \phi_T = \phi_F$$

Since interference argument
only needs *some* reference plane,
we expect same expansion
in $\cos n\phi_i$ and $\cos n\phi$

ϕ Reconstruction



$$\cos \phi = \frac{\hat{n}_+ \cdot \hat{n}_- + \cos \alpha_+ \cos \alpha_-}{\sin \alpha_+ \sin \alpha_-}$$

$$(m_{\pm}^2 - m_0^2) = \sqrt{s} E_{f_{\pm}} \left(1 - \sqrt{1 - \frac{4m_{\pm}^2}{s}} \cos \alpha_{\pm} \right)$$

Spinor-Scalar Measurement

$$e^+e^- \rightarrow \tilde{\mu}_R^+ \tilde{\mu}_R^- \rightarrow (\mu^+ \tilde{\chi}_1^0)(\mu^- \tilde{\chi}_1^0)$$
$$e^+e^- \rightarrow \mu_{R1}^+ \mu_{R1}^- \rightarrow (\mu^+ \gamma_1)(\mu^- \gamma_1)$$

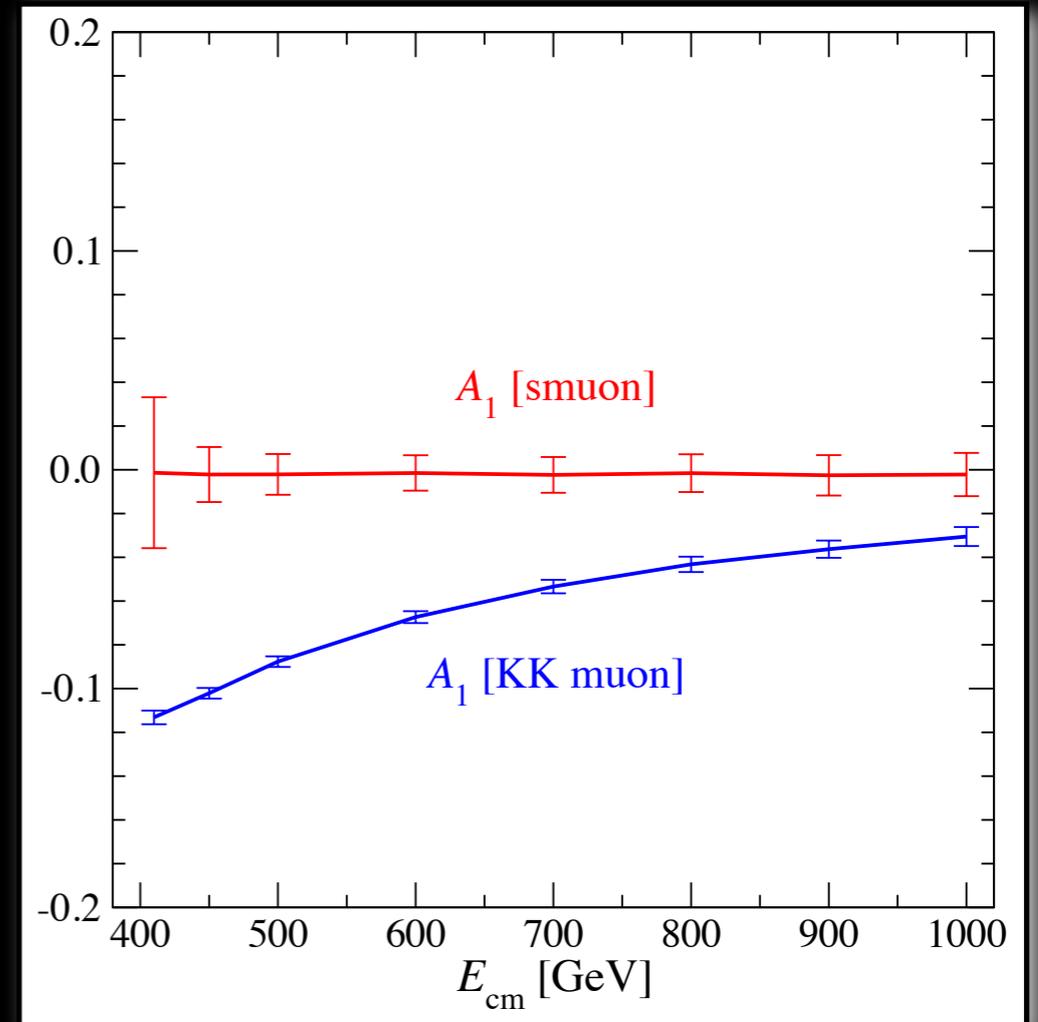
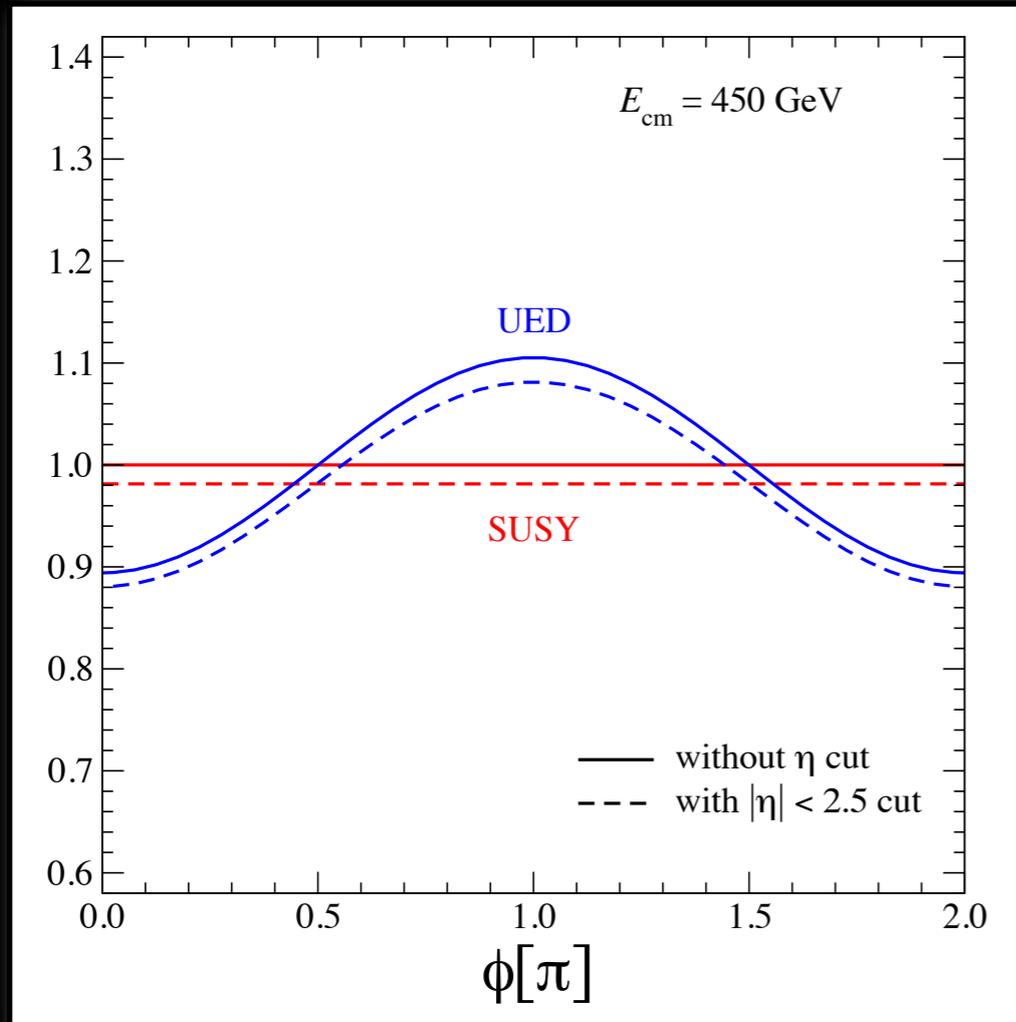
- Choose mass spectrum

$$m_{\pm} = m_{\tilde{\mu}_R^{\pm}} = m_{\mu_{R1}^{\pm}} = 200 \text{ GeV}$$

$$m_0 = m_{\tilde{\chi}_1^0} = m_{\gamma_1} = 50 \text{ GeV}$$

- Assume 500 fb^{-1} of luminosity and $\sqrt{s} \leq 1 \text{ TeV}$
- Model detector acceptance cuts with $|\eta_{\mu}|, |\eta_{\cancel{E}}| \leq 2.5$
- Simulated using HELAS/BASES

Spinor-Scalar Measurement



$$A_1 = \frac{\pi^2 m_{\mu_{R1}^\pm}^2}{8(s + 2m_{\mu_{R1}^\pm}^2)} \left(\frac{1 - 2m_{\gamma_1}^2/m_{\mu_{R1}^\pm}^2}{1 + 2m_{\gamma_1}^2/m_{\mu_{R1}^\pm}^2} \right)^2 \leq \pi^2/48 \approx 0.206$$

Vector-Spinor Measurement

$$e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow (\ell^+ \tilde{\nu}_\ell)(\ell^- \tilde{\nu}_\ell^*)$$

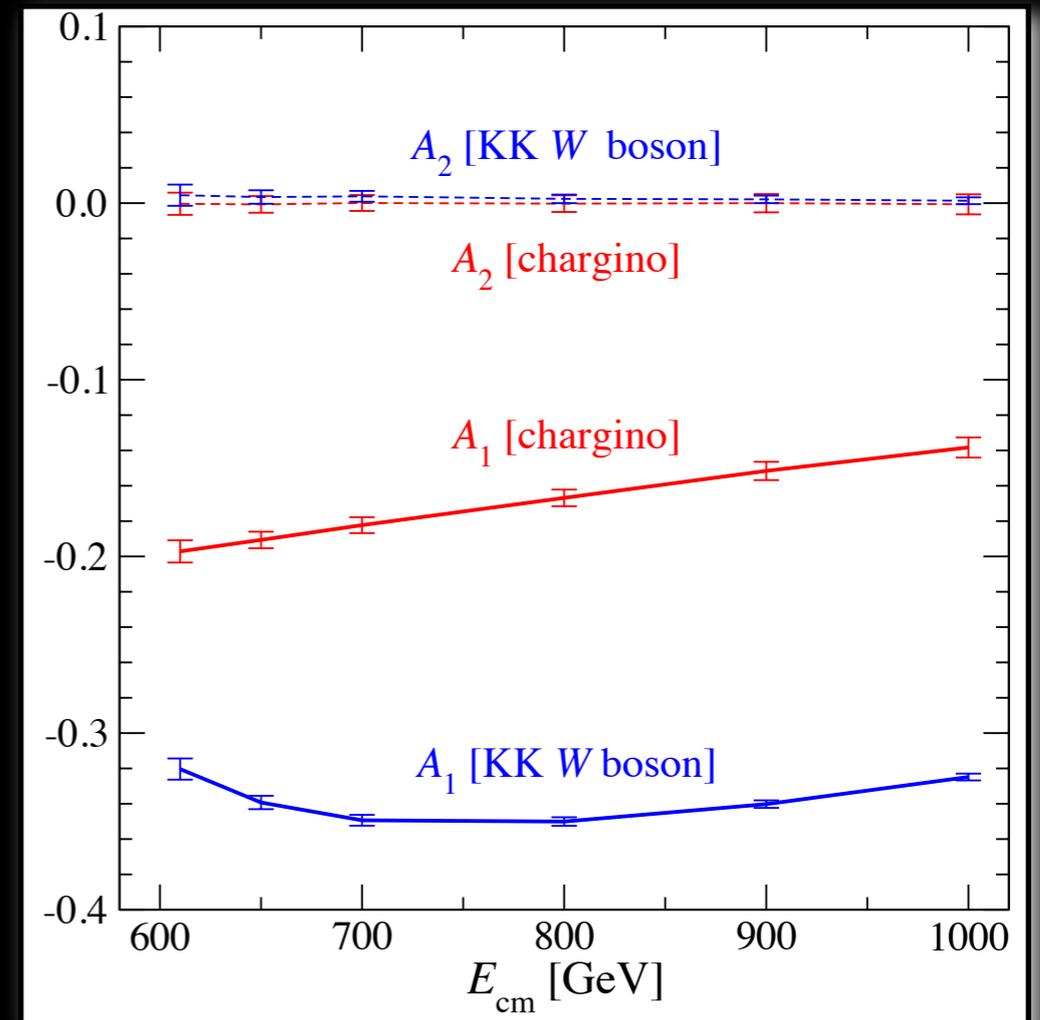
$$e^+e^- \rightarrow W_1^+ W_1^- \rightarrow (\ell^+ \nu_{\ell 1})(\ell^- \bar{\nu}_{\ell 1})$$

- Choose

$$m_{\pm} = m_{\tilde{\chi}_1^{\pm}} = m_{W_1^{\pm}} = 300 \text{ GeV}$$

$$m_0 = m_{\tilde{\nu}_\ell} = m_{\nu_{\ell 1}} = 200 \text{ GeV}$$

Production amplitudes suppressed
by $\sim m_{W_1^{\pm}}^2/s$ leads to $A_2 \lesssim 0.5\%$



Conclusions

- Measurement of azimuthal angular dependence offers model-independent measurement of spin.
- Reconstruction of azimuthal angles ϕ_i impossible for most interesting new physics.
- We have demonstrated that $\Delta\phi \equiv \phi$ is both measurable and contains spin information
 - This technique successful in discerning spin 0 and spin-1/2 particles
 - Spin-1/2 vs. spin-1 more difficult.