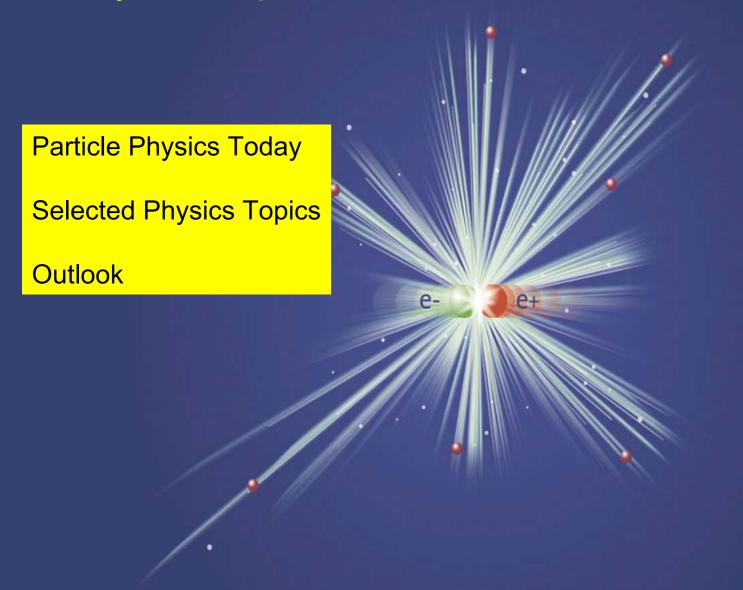
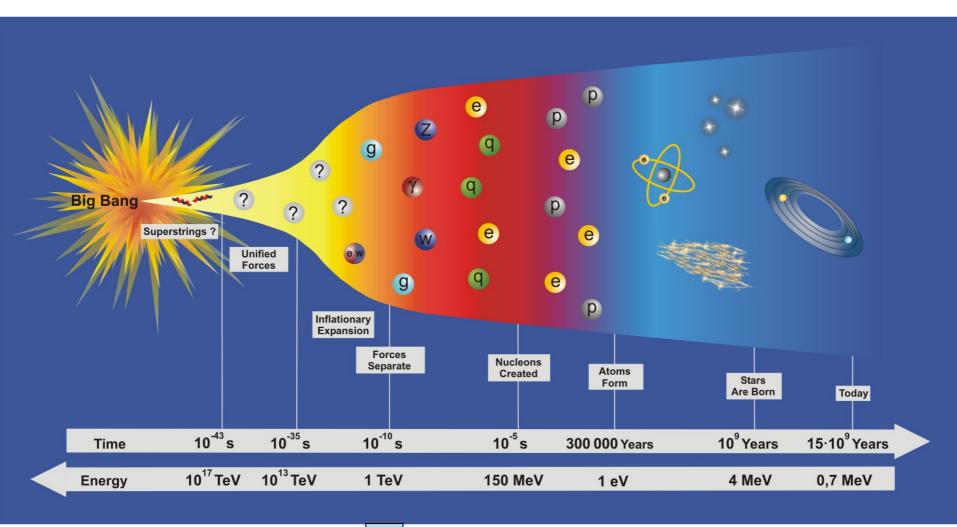
Physics topics at the Linear Collider



History of the Universe





Particle Physics Today or Status of the Standard Model

Past few decades

"Discovery" of Standard Model

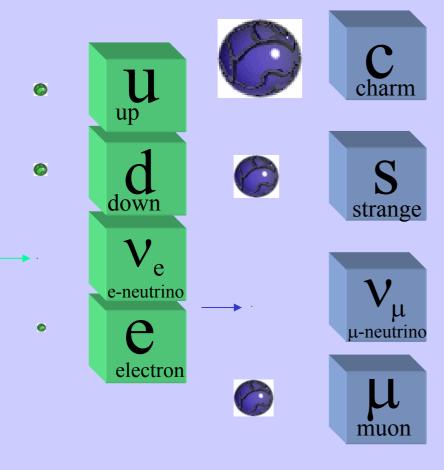
through synergy of

hadron - hadron colliders (e.g. Tevatron)

lepton - hadron colliders (HERA)

lepton - lepton colliders (e.g. LEP)

Matter Particles

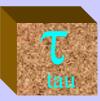


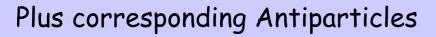












Matter Particles

Matter (Stars <=> living organisms) consists of 3 families of Quarks and Leptons

Matter around us: only 1 of the 3 families

Matter at high energies:

,democratic', all 3 families present

- Situation fraction of seconds after the creation of the Universe
- —> Study of Matter at High Energies knowledge about Early Universe

Standard Model

 $e^+e^- \rightarrow Z^0 \rightarrow f \overline{f}$ where f=q,l,v

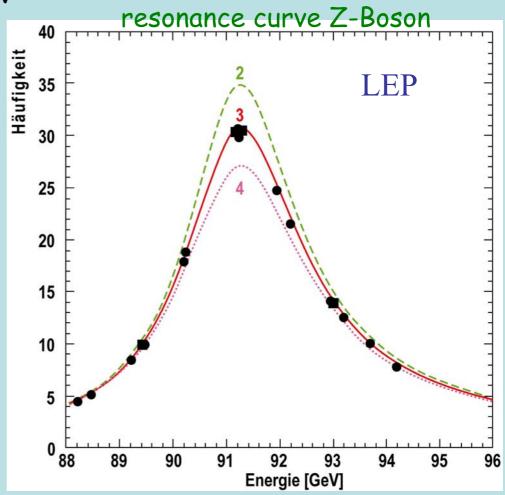
 σ_Z and Γ_Z depend on number of (light) neutrinos



number of families:

N = 2.984 + 0.008

Nobel Prize 2008: Kobayashi-Maskawa)



Forces

acts on/in

relationath force corriers

tyne

Type	rei.strength	Torce curriers	acis on/in
Strong Force	1	Gluons g m = 0	Quarks Atomic Nucleus
Electro-magnet. Force	~ 1/1000	Photon γ m = 0	Electric Charge Atoms, Chemistry
Weak Force	~ 10 -5	W, Z Bosons m = 80 , 91 GeV	Leptons, Quarks Radioactive Decays (β-decay)
Gravitation	~ 10 -38	Graviton? m = 0	Mass, Energy

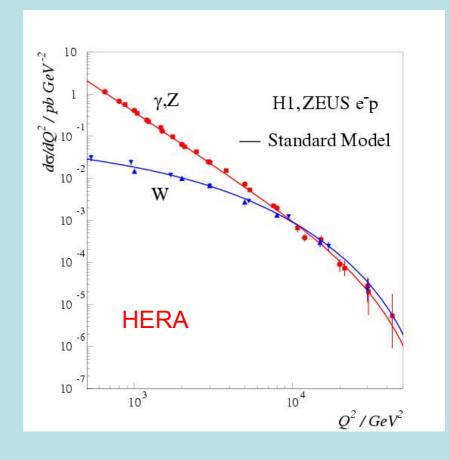
Force Carriers (Bosons) exchange interactions

Forces

Four fundamental *Forces* act between *Matter Particles* through *Force Carriers* (Gluons, W[±] und Z⁰, γ, Graviton)

forces in our energy regime:
 different strengths
 forces at high energies:
 democratic......unification

—>Situation immediately after creation of the Universe



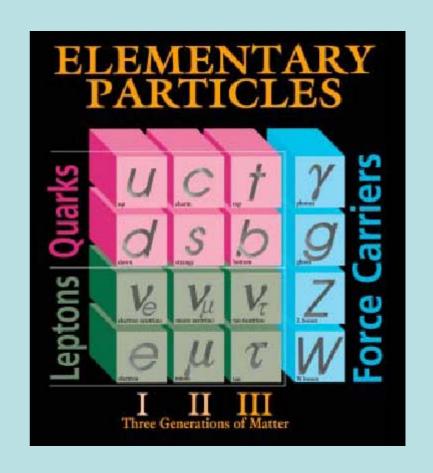
What have we learned the last 50 years or Status of the **Standard Model**

The physical world is

composed of Quarks and Leptons

interacting via force carriers (Gauge Bosons)

Last entries: top-quark 1995 tau-neutrino 2000



Standard Model

- mathematical description of all interactions, involving weak, electromagnetic, strong forces, through closely related symmetry principles (gauge symmetries)
- Symmetries are of fundamental importance for describing the dynamics in particle physics

```
Noether-Theorem: Symmetry <---> Conservation Law
e.g. Rotation angular momentum
Mirror image Parity
gauge transformation charge, Baryon #...
```

Local gauge symmetry —> Invariance under local phase transformation

(QED)

1967/68 Glashow - Salam - Weinberg: gauge theory to unify el.magn. and weak forces



Standard Model of electroweak interaction

Problem: gauge invariance only possible for <u>massless</u> gauge bosons (m=0, R->00 => Phase trafo can be compensated through gauge trafo <u>everywhere</u> in space)

Massive gauge bosons => Violation of gauge invariance

Solution:

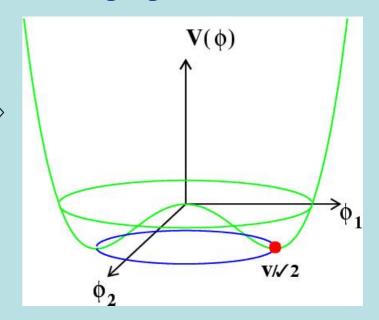
Introduction of a scalar background field

(Higgs-Field)

Vacuum expectation value v

(Analogy: super conductivity)

Nobel Prize 2008: Nambu



The Higgs mechanism

Paradigm: All (elementary) particles are massless

- \Rightarrow gauge principle works
- ⇒ renormalizable theory (finite cross sections)

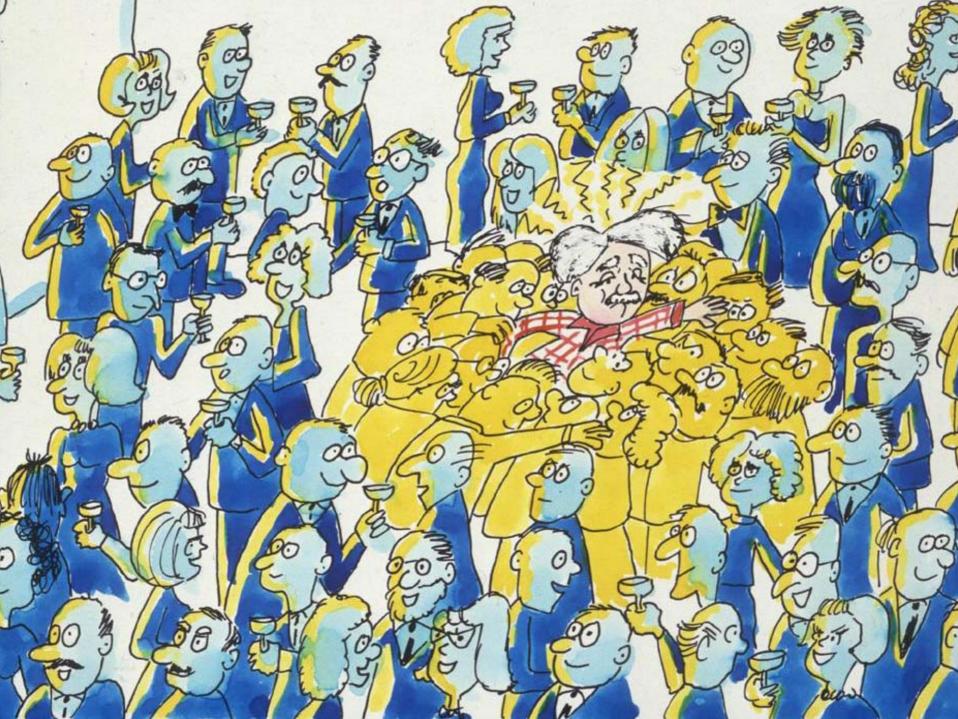
Permanent interaction of particles with a scalar Higgs field acts as if the particles had a mass (effective mass):

$$= \frac{1}{1/\cancel{q}} + \frac{1}{\cancel{q}} \underbrace{\begin{pmatrix} g_f v / \sqrt{2} \\ H \times \end{pmatrix}}_{1/\cancel{q}} + \frac{1}{\cancel{q}} \underbrace{\begin{pmatrix} g_f v \\ \sqrt{2} \end{pmatrix}}_{H \times } + \cdots + \cdots$$

$$\frac{1}{\cancel{q}} + \frac{1}{\cancel{q}} \left(\frac{g_f v}{\sqrt{2}} \right) \frac{1}{\cancel{q}} + \cdots = \frac{1}{\cancel{q}} \sum_{n=0}^{\infty} \left[\left(\frac{g_f v}{\sqrt{2}} \right) \frac{1}{\cancel{q}} \right]^n = \frac{1}{\cancel{q}} - \left(\frac{g_f v}{\sqrt{2}} \right)$$

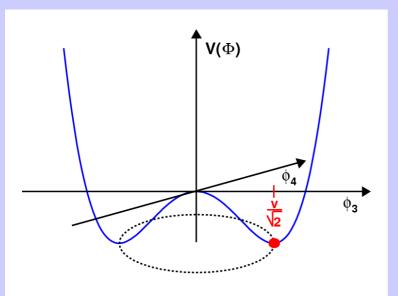






The Higgs mechanism

How to add such a field in a gauge invariant way?



Introduction of SU(2)xU(1) invariant Mexican hat potential

$$V(\Phi) = -\mu^2 |\Phi|^2 + \lambda |\Phi|^4$$

Simplest case (SM):

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi_4 \end{pmatrix}$$

complex doublet of weak iso-spin

This is only the most economic way. Many more possibilities exist, e. g. two doublets (minimal SUSY), triplets, ...

Higgs mechanism requires the existence of at least one scalar, massive Higgs boson.





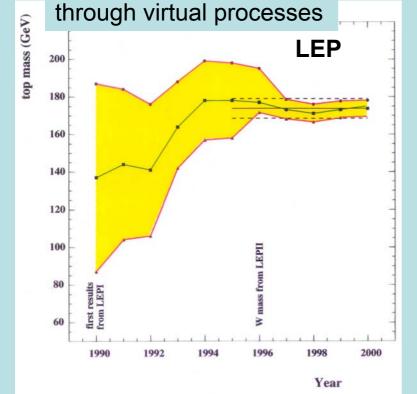
Status of the Standard Model

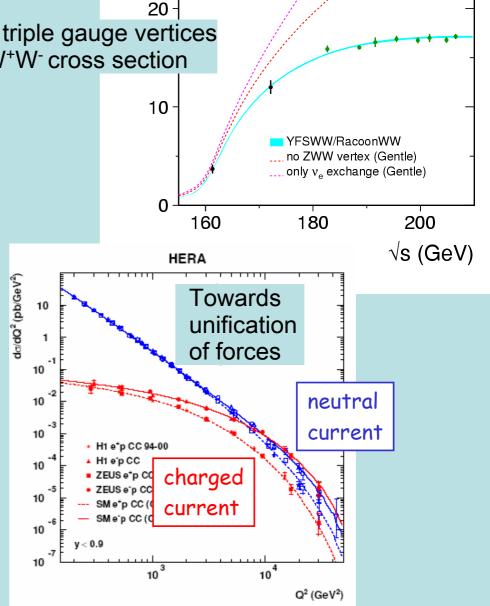
Verification of triple gauge vertices from e⁺e⁻ → W⁺W⁻ cross section

30

o_{ww} (bb)

Indirect determination of the top quark mass: Proves high energy reach

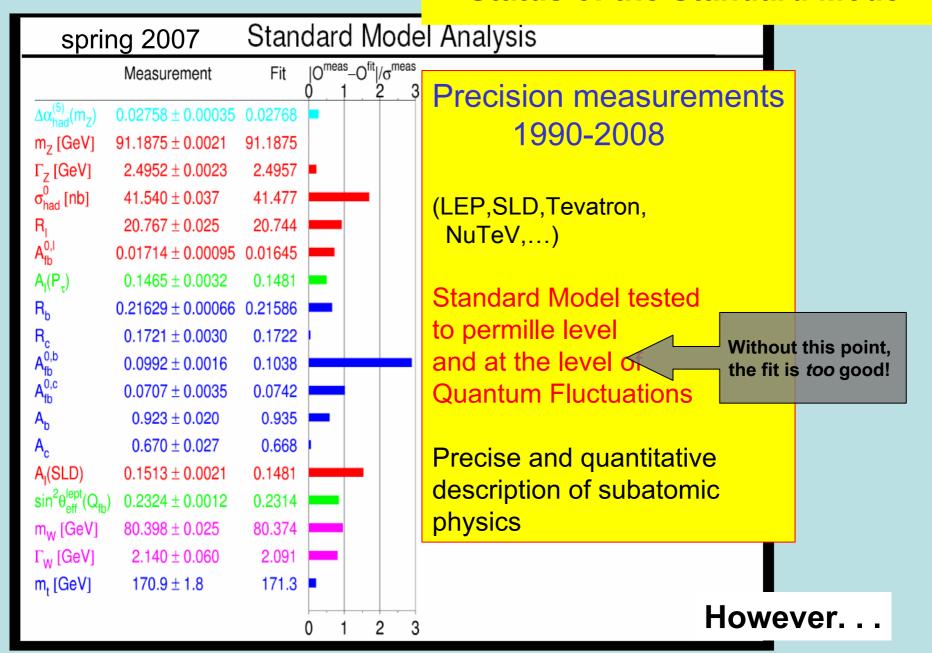


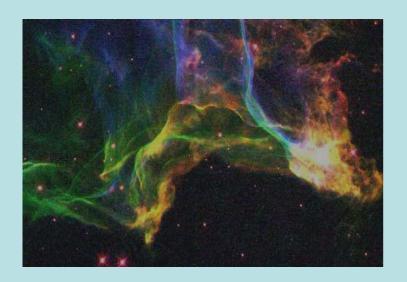


_EP

PRELIMINARY

Status of the Standard Model

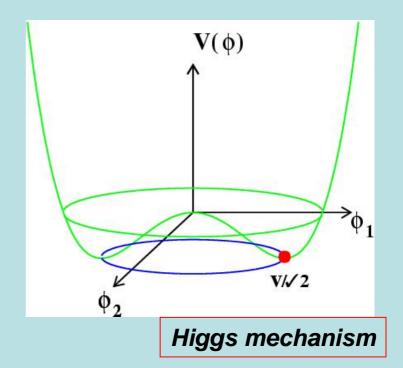


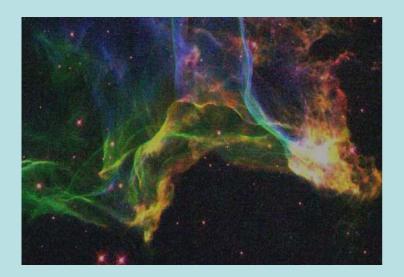


Standard Model

 What is the origin of mass of elementary particles or

why are carriers of weak force so heavy while the photon is massless



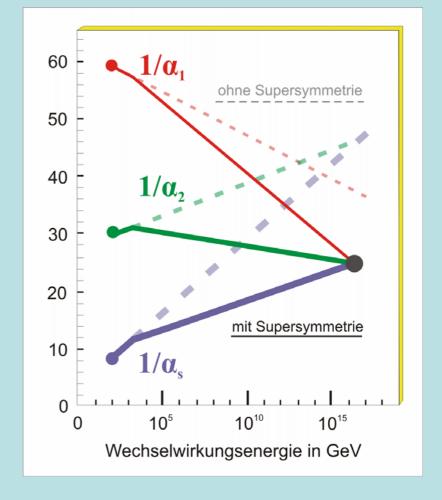


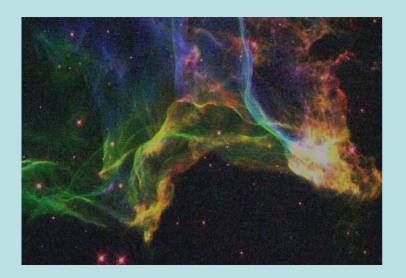
Ultimate Unification

- Do the forces unify, at what scale
- Why is gravity so different
- Are there new forces

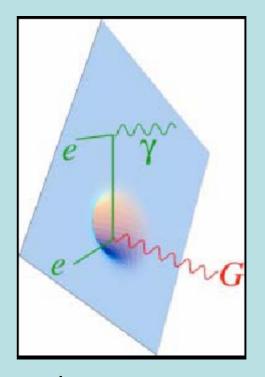
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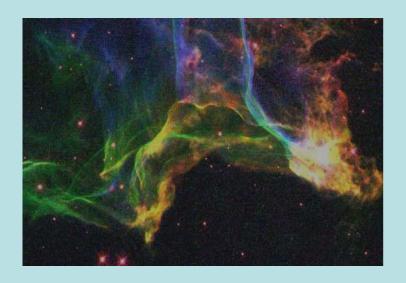




Hidden Dimensions or Structure of Space -Time

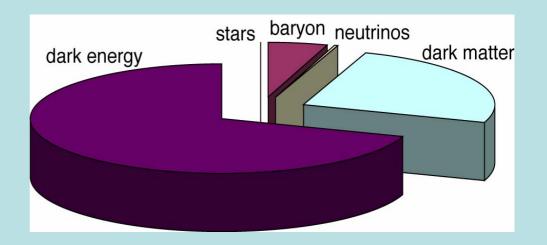


- Are there more than four space-time dimensions
- What is the quantum theory of gravity
- •
- •
- •



Cosmic Connections

- What is dark matter
- What is dark energy
- •
- What happened to antimatter
- •
- •



The next steps at the energy frontier

Interplay of Hadron and Lepton Colliders

There are two distinct and complementary strategies for gaining new understanding of matter, space and time at future particle accelerators

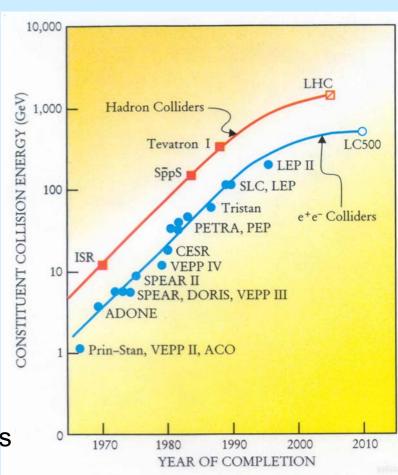
HIGH ENERGY

or

direct discovery of new phenomena i.e. accelerators operating at the energy scale of the new particle

HIGH PRECISION

Access to new physics at high energies through the precision measurement of phenomena at lower scales

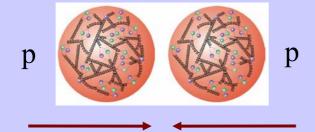


Both strategies have worked well together

→ much more complete understanding than from either one alone

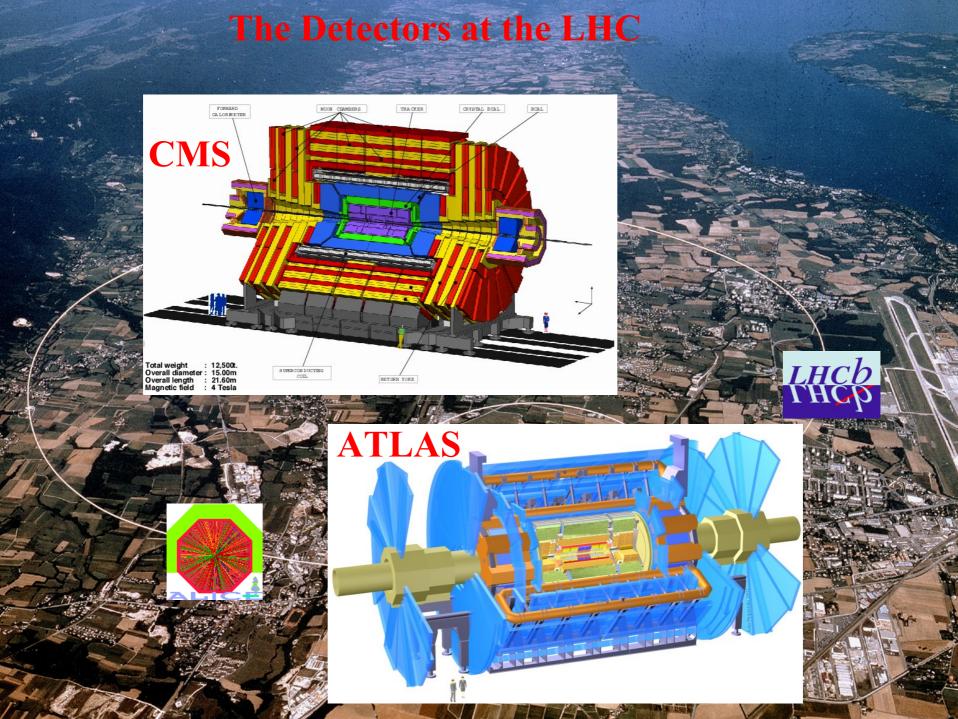
Hadron Collider

Lepton Collider

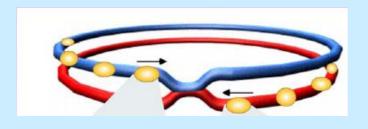


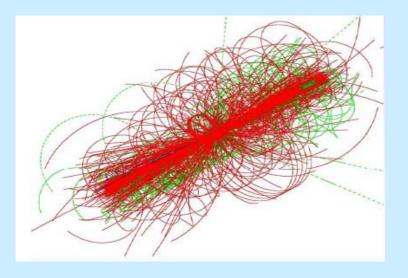


- p = composite particle:
 unknown √s of IS partons,
 no polarization of IS partons,
 parasitic collisions
- p = strongly interacting:
 huge SM backgrounds,
 highly selective trigger needed,
 radiation hard detectors needed
- e = pointlike particle:
 known and tunable √s of IS particles,
 polarization of IS particles possible,
 kinematic contraints can be used
- e = electroweakly interacting low SM backgrounds, no trigger needed, detector design driven by precision



Proton-Proton Collisions at the LHC

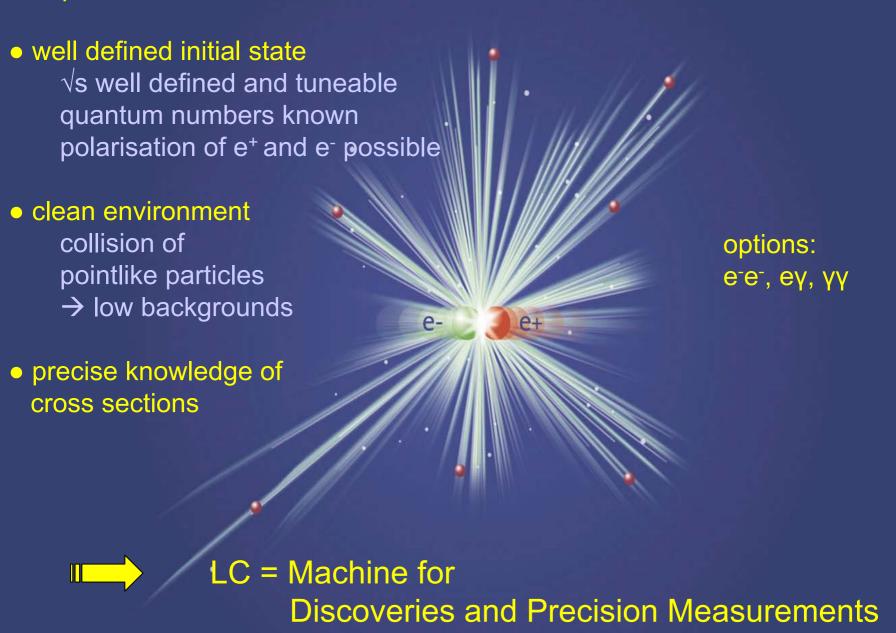




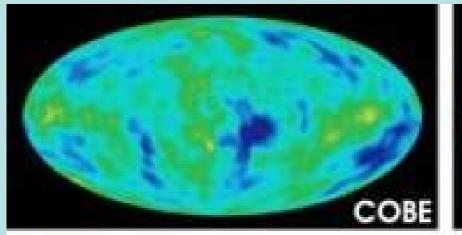
- 2835 + 2835 proton bunches separated by 7.5 m
- → collisions every 25 ns= 40 MHz crossing rate
- 10¹¹ protons per bunch
- at 10^{34/}cm^{2/}s
 ≈ 35 pp interactions per crossing pile-up
- $\rightarrow \approx 10^9$ pp interactions per second !!!
- in each collision≈ 1600 charged particles produced

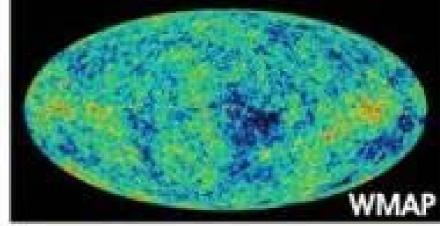
enormous challenge for the detectors

The power of an Electron-Positron Linear Collider



An Analogy: What precision does for you ...

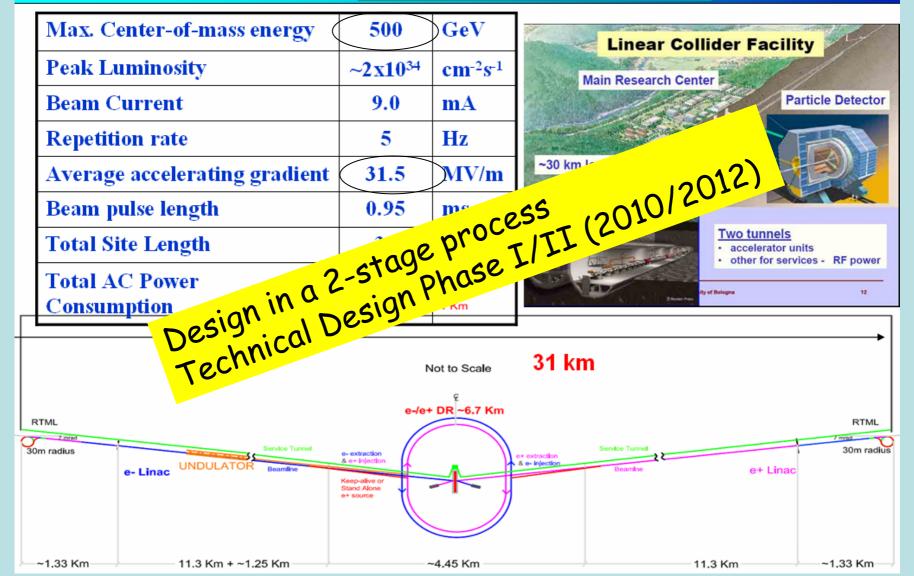




High Energy Colliders: ILC (E_{cm} up to ~ 1TeV)

ILC @ 500 GeV

ILC web site: http://www.linearcollider.org/cms/



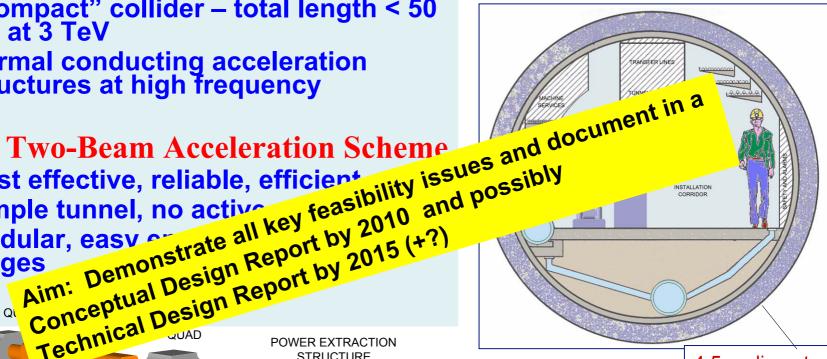
High Energy Colliders: CLIC $(E_{cm} up to \sim 3TeV)$

BPM

- High acceleration gradient: ~ 100 MV/m
 - "Compact" collider total length < 50 km at 3 TeV
 - Normal conducting acceleration structures at high frequency
- **Novel Two-Beam Acceleration Scheme**
 - Cost effective, reliable, efficient
 - Simple tunnel, no active
 - Modular, easy er stages

Technical Design Report by 2015 (+?) 12 GHz - 140 MW **ACCELERATING** Main beam - 1 A, 200 ns from 9 GeV to 1.5 Te)

CLIC TUNNEL CROSS-SECTION

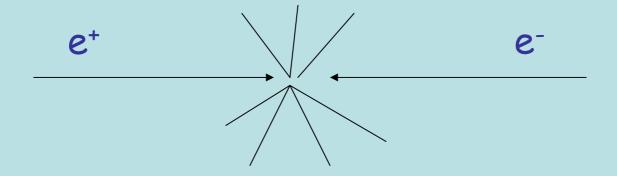


4.5 m diameter

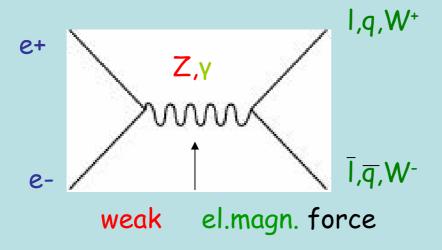
Drive beam - 95 A, 300 ns from 2.4 GeV to 240 MeV



Electron - Positron - Reactions

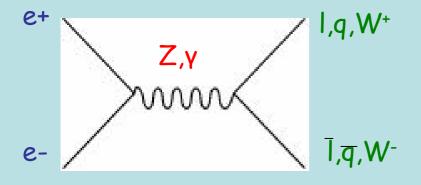


Description in particle physics:

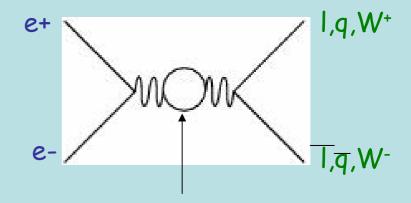


Electron - Positron - Reactions

direct measurement:

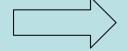


indirect measurement:



Heisenberg: ΔΕΔt > h

=> (extremely) short fluctuations
to high energies (masses) possible
("Quantum fluctuations")



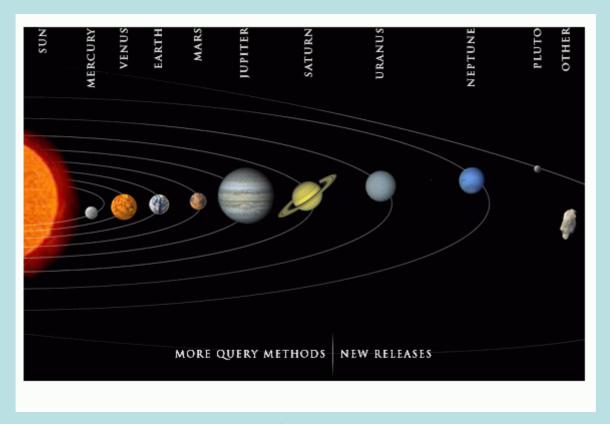
Modifications of rate and properties of reactions

Effect of high masses indirectly measureable

example:

Precision + Model

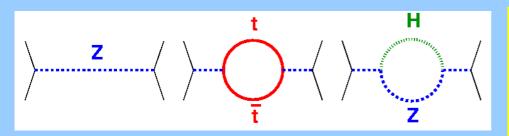
planets



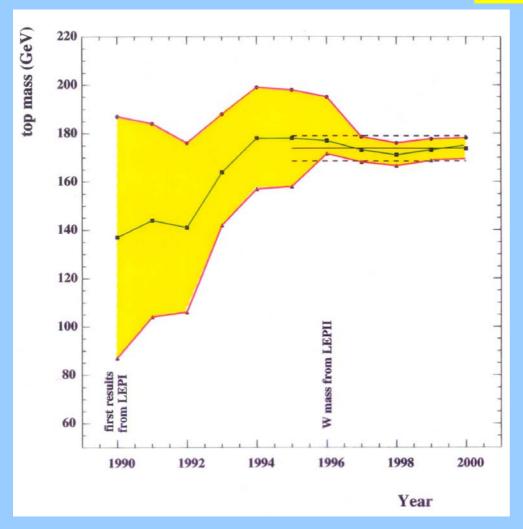
Precision measurements of Uranus orbit

—--> deviation from model calculations

—--> prediction of Neptune



Standard Model: Testing Quantum Fluctuations



LEP:

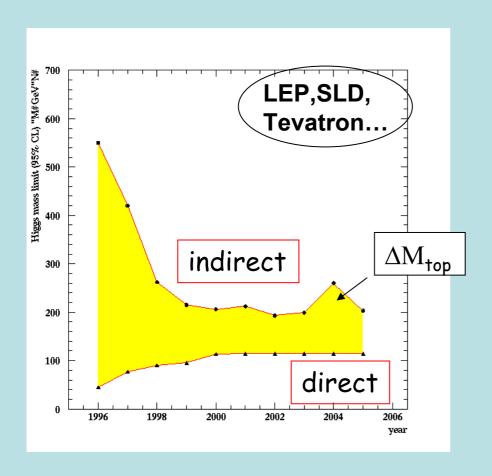
Indirect determination of the top mass

possible due to

- precision measurements
- known higher order electroweak corrections

$$\propto (\frac{M_t}{M_W})^2, \ln(\frac{M_h}{M_W})$$

Synergy of colliders:



Time evolution of experimental limits on the Higgs boson mass

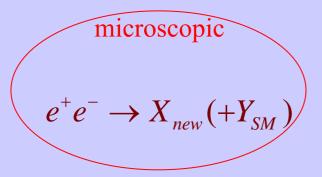
knowledge obtained only through combination of results from different accelerator types

in particular: Lepton and Hadron Collider

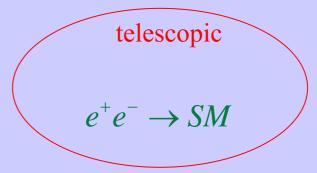
*M*_H between 114 and ~200 GeV

The Role of the LC

Explore new Physics through high precision at high energy



Study the properties of new particles (cross sections, BR's, quantum numbers)



Study known SM processes to look for tiny deviations through virtual effects (needs ultimate precision of measurements and theoretical predictions)

- →precision measurements will allow
- -- distinction of different physics scenarios
- -- extrapolation to higher energies

The (I)LC Physics Case

or

Relation of Hadron Collider and Linear Collider

- 1. Since the (I)LC will start after the start of LHC, it must <u>add</u> significant amount of information. This is the case! (see e.g. TESLA TDR, Snowmass report, ACFA study etc.)
- 2. Neither (I)LC nor LHC can draw the whole picture alone. An (I)LC will
- add new discoveries and
- precision of (I)LC will be essential for a better understanding of the underlying physics
- 3. There are probably pieces which can only be explored by the LHC due to the higher mass reach. <u>Joint interpretation</u> of the results will improve the overall picture
- 4. Overlapping running of both machines will further increase the potential of both machines and might be mandatory, depending on the physics scenario realized

International Linear Collider Parameters global consensus (Sept. 2003)

(1) baseline machine

200 GeV < \sqrt{s} < 500 GeV integrated luminosity ~ 500 fb⁻¹ in 4 years electron polarisation ~ 80%

(2) energy upgrade

to $\sqrt{s} \sim 1 \text{ TeV}$ integrated luminosity $\sim 1 \text{ ab}^{-1}$ in 3 years

(3) options

positron polarisation of $\sim 50\%$ high luminosity running at M_Z and W-pair threshold e^-e^- , $e\gamma$, $\gamma\gamma$ collisions

(4) concurrent running with LHC desired

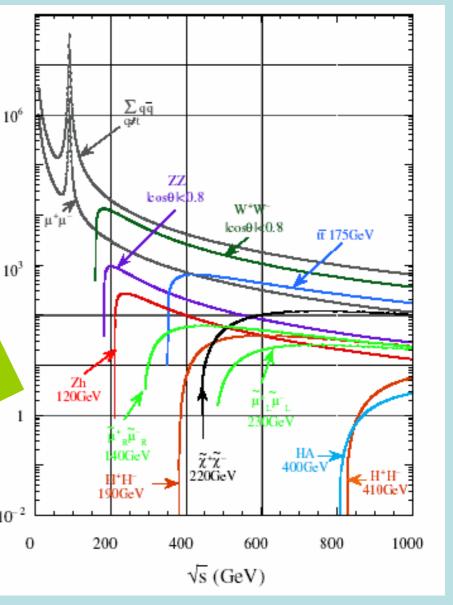
! Times quoted for data taking cover only part of program!

Physics at the LC

Comprehensive and high precision coverage of energy range from M_7 to ~ 1 TeV

Selected Physics Topics

- Higgs Mechanimostly from
 Supersymples mostly from
 Streexamples mostly from
 In the street of the at lower energies



cross sections few fb to few pb \rightarrow e.g. O(10,000) HZ/yr

Physics Examples

Electroweak Symmetry Breaking

- Higgs mechanism
- no Higgs scenarios

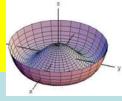
Supersymmetry

- unification of forces
- dark matter

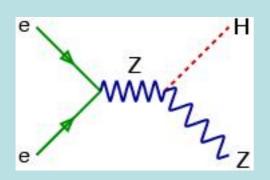
Precision tests of the Standard Model

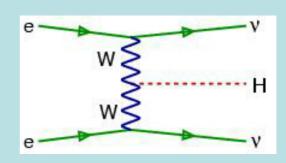
- top quark properties
- high luminosity running at the Z-pole

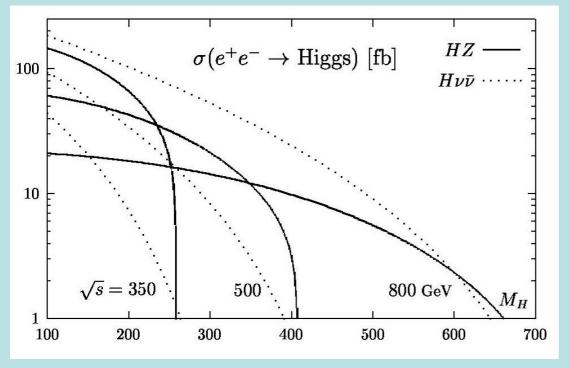
The Higgs: Key to Understanding Mass



Dominant production processes at LC:







Task at the LC:

- determine properties of the Higgs-boson
- establish Higgs mechanism responsible for the origin of mass
- . . . together with LHC

Tasks at the LC

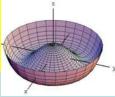
Establishing the Higgs mechanism as being responsible for EW symmetry breaking requires more than discovering one or more Higgs bosons and measuring its/their mass(es).

Precision measurements must comprise:

- Mass
- Total width
- •Quantum numbers J^{CP} (Spin, CP even?)
- •Higgs-fermion couplings (∝ mass?)
- •Higgs-gauge-boson couplings (W/Z masses)
- •Higgs self-coupling (spontaneous symmetry breaking)

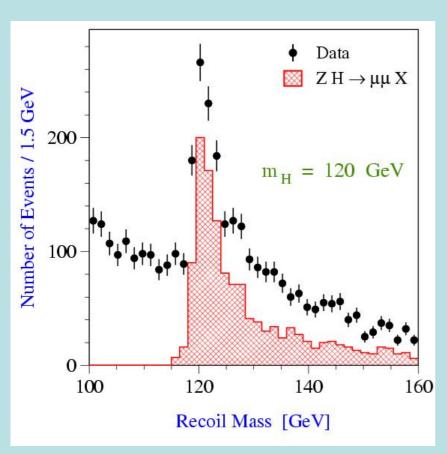
Precision should be sufficient to distinguish between different models (e. g. SM/MSSM, effects from XD, ...)

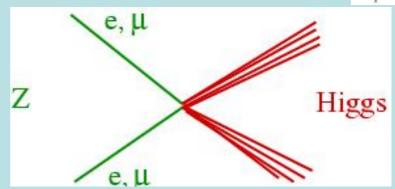
The Higgs: Key to Understanding Mass



Recoil mass spectrum

ee -> HZ with Z -> 1+1-



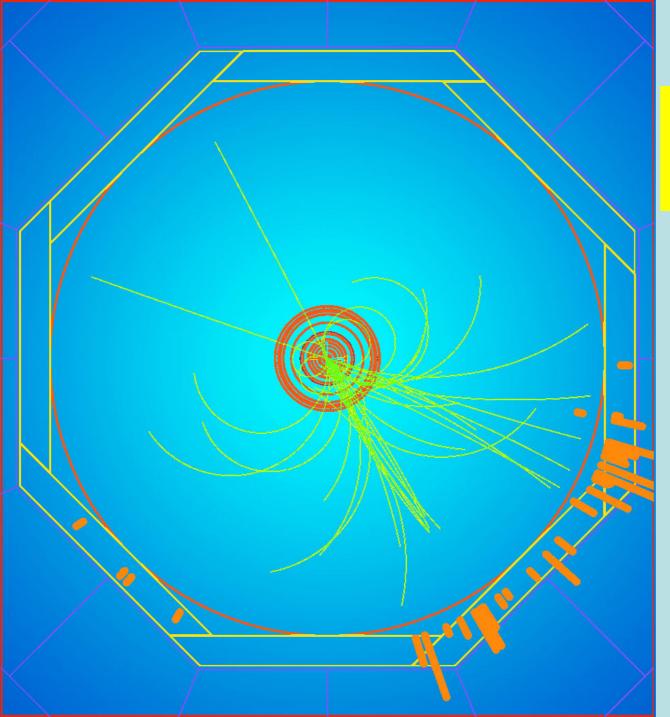


 $\Delta \sigma \sim 3\%$

model independent measurement

∆m ~ 50 MeV

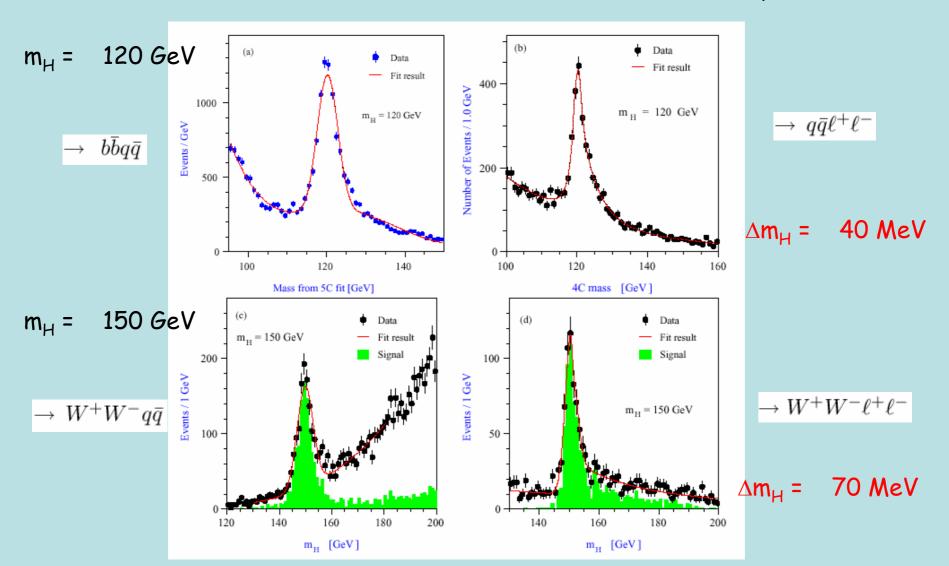
sub-permille precision



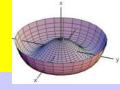
ee -> HZ Z -> | | H -> qq

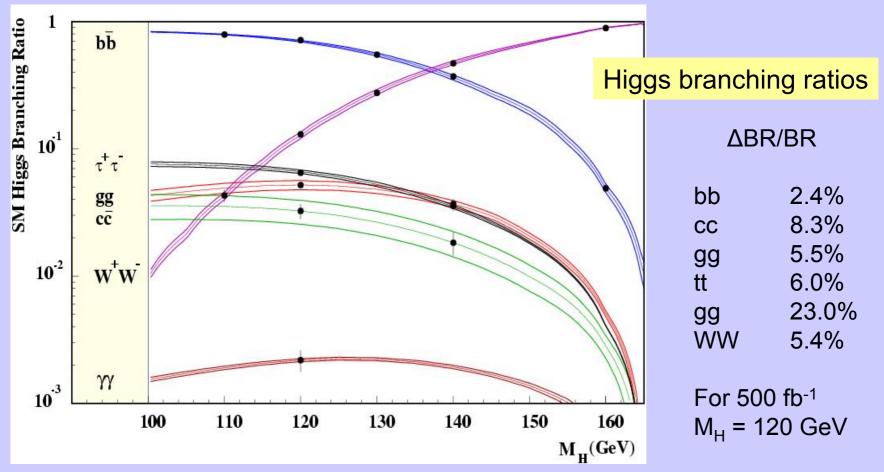
Precision physics of Higgs bosons

ee -> HZ diff. decay channels



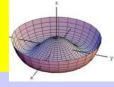
The Higgs: Key to Understanding Mass





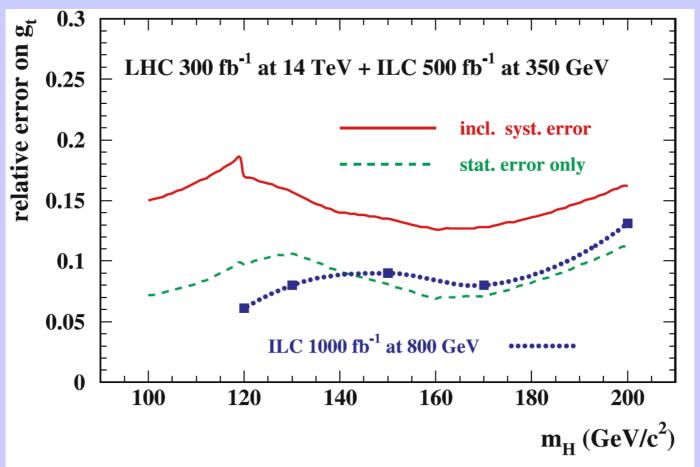
model-independent measurements at %-level possible

Example: Top Yukawa Coupling

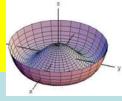


LHC sensitive to top Yukawa coupling of light Higgs through tth production. LC BR measurement (h→bb and h→WW) turns rate measurement into an absolute coupling measurement

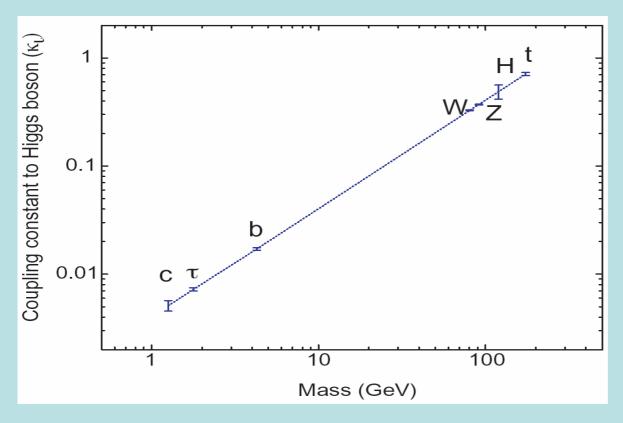
LC direct measurement only at high energy (> 800 GeV)



The Higgs: Key to Understanding Mass



Testing the Yukawa couplings...

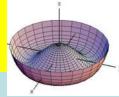


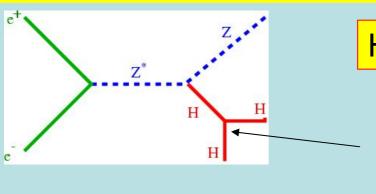
Precision ~ % level

...through the measurement of absolute BRs:

$$BR(H \to X) = \frac{\left[\sigma(HZ) \cdot BR(H \to X)\right]^{meas}}{\sigma(HZ)^{meas}}$$

The Higgs: Key to Understanding Mass



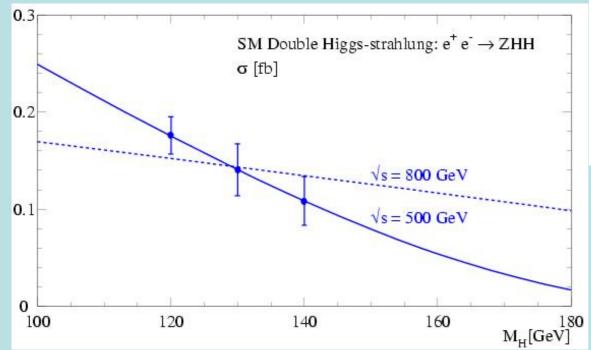


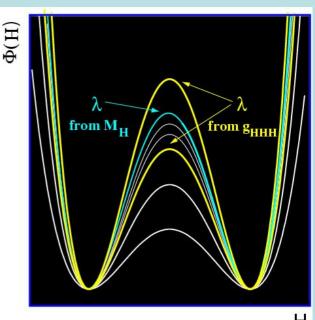
Higgs self coupling

9ннн

$$\Phi(H)=\Lambda v^2H^2 + \Lambda vH^3 + 1/4\Lambda H^4$$

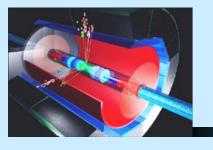
SM: $g_{HHH}=6\Lambda v$, fixed by M_H



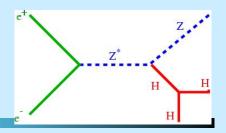


 $\Delta \lambda/\lambda \sim 10-20 \%$ for 1 ab⁻¹

requires excellent calorimeter resolution

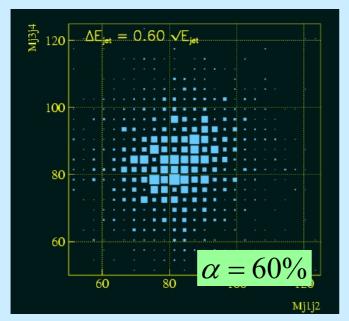


Jet energy resolution

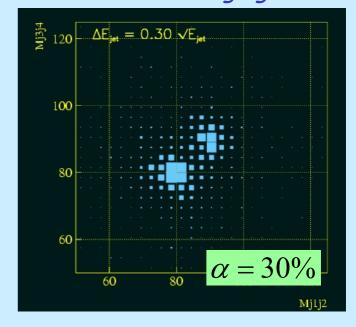


- Dijet masses in WWvv, ZZvv events (no kinematic fit possible):
- Challenge: separate W and Z in their hadronic decay mode

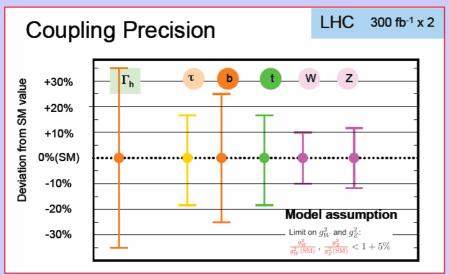
LEP-like detector

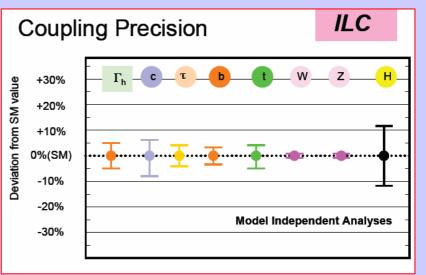


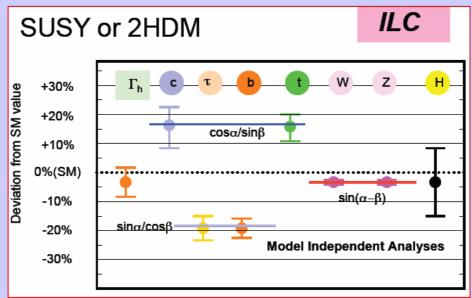
LC design goal



e.g. Coupling Precision and New Physics





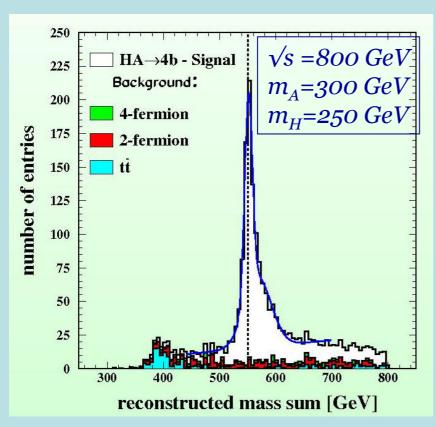


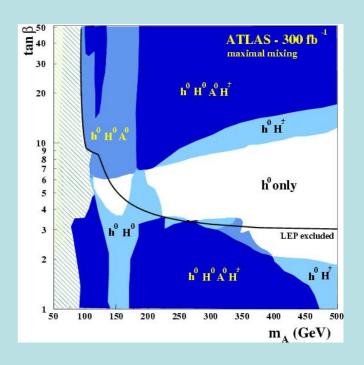
Yamashita

Heavy SUSY-Higgs

Heavy SUSY Higgs bosons:

observation and mass/BR/width(?) measurements
deep into the LHC wedge region at 800-1000 GeV LC





 $HA \rightarrow bbbb$ and $HA \rightarrow bb\tau\tau/\tau\tau bb$ observable

HA: 5σ discovery possible up to $\Sigma m = \sqrt{s} - 30 \text{ GeV}$

Precision physics of Higgs bosons

Conclusion

precision measurements at the LC together with the results from LHC are crucial to establish the Higgs mechanism responsible for the origin of mass and for revealing the character of the Higgs boson

if the electroweak symmetry is broken in a more complicated way then foreseen in the Standard Model the LC measurements strongly constrain the alternative model

The Higgs is Different!

All the matter particles are spin-1/2 fermions. All the force carriers are spin-1 bosons.

Higgs particles are spin-0 bosons.

The Higgs is neither matter nor force;

The Higgs is just different.

This would be the first fundamental scalar ever discovered.

The Higgs field is thought to fill the entire universe. Could give some handle of dark energy (scalar field)?

Many modern theories predict other scalar particles like the Higgs. Why, after all, should the Higgs be the only one of its kind?

LHC and LC can search for new scalars with precision.

LHC and LC results will allow to study the Higgs mechanism in detail and to reveal the character of the Higgs boson

This would be the first investigation of a scalar field

This could be the very first step to understanding dark energy

Beyond the Higgs

Why are electroweak scale (10² GeV) and the Planck scale (10¹⁹ GeV) so disparate?

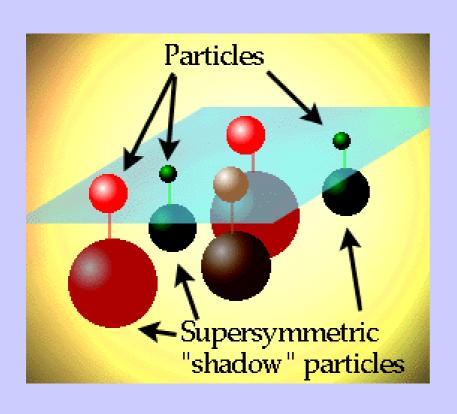
Are there

new particles ? → supersymmetry

new forces ? → strong interactions

hidden dimensions?

Introduction of an additional symmetry to the SM:



boson ↔ fermion symmetry

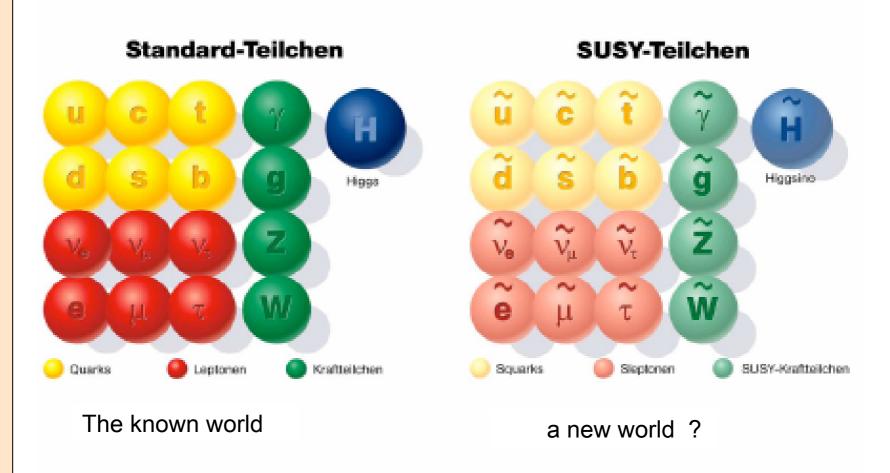
Each SM particle gets a SUSY partner whose spin differs by 1/2. All other quantum numbers are equal.

But so far no SUSY particle seen (SUSY symmetry broken)

but

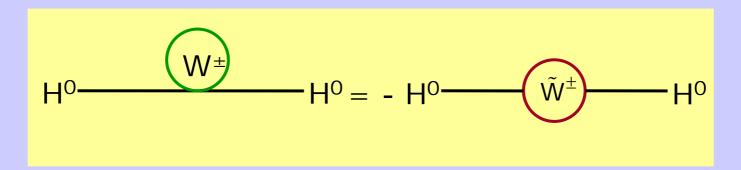
SUSY well motivated theory

Supersymmetrie



Solution to hierarchy problem

Motivation 1: It solves the hierarchy problem

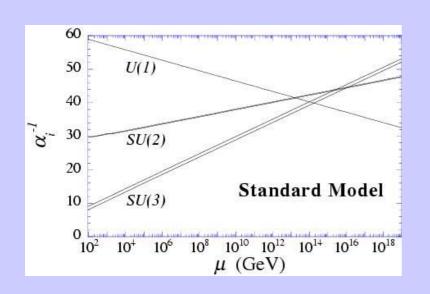


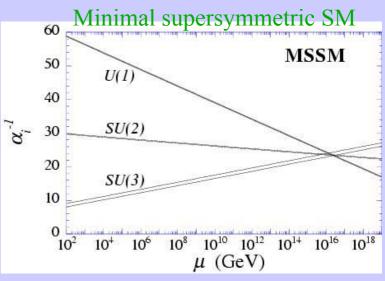
The divergence in the Higgs mass corrections is cancelled exactly for unbroken SUSY.

If it is not broken too strongly (i. e. if the SUSY partners are at $< \sim 1$ TeV), there is no fine tuning necessary.

Unification of gauge couplings

Motivation 2: Gauge coupling constants unify





(Requires light (< TeV) partners of EW gauge bosons)

This is achieved for $\sin^2\theta_W^{SUSY} = 0.2335(17)$ Experiment: $\sin^2\theta_W^{exp} = 0.2315(2)$

More good reasons ...

Motivation 3: Provides cold dark matter candidate

If lightest SUSY particle is stable, it is an excellent dark matter candidate

Motivation 4: Link to gravity

SUSY offers the theoretical link to incorporate gravity. Most string models are supersymmetric.

Motivation 5: Predicts light Higgs boson

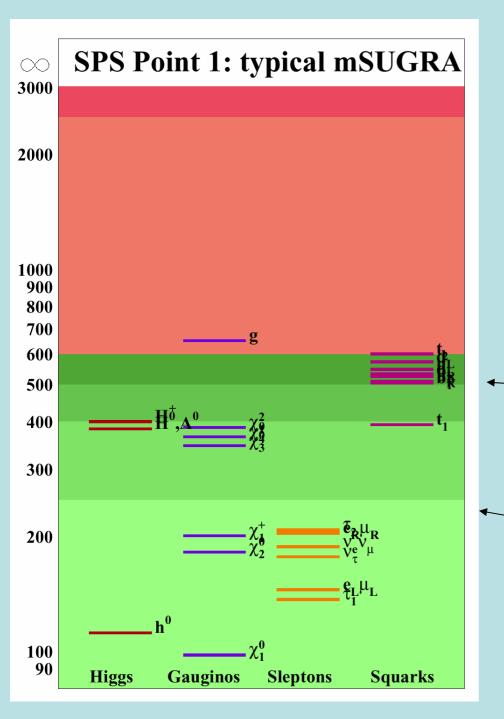
SUSY predicts a light (< 135 GeV) Higgs boson as favored by EW precision data.

- best motivated extension of SM
 grand unification connection to gravity light Higgs sin²Θ_W
 dark matter candidate
- mass spectrum depends on the unknown breaking scheme
- LC task for SUSY
 reconstruction of kinematically accessible sparticle spectrum
 i.e. measure sparticle properties (masses, Xsections, spin-parity)

extract fundamental parameters (mass parameters, mixings, couplings) at the weak scale

extrapolate to GUT scale using RGEs

→ determine underlying supersymmetric model

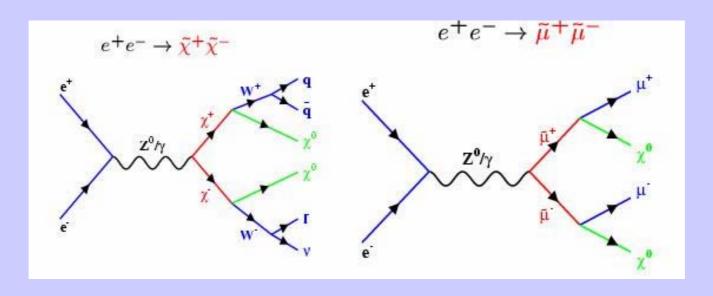


Mass spectra depend on choice of models and parameters...

——— well measureable at LHC

precise spectroscopy at the Linear Collider

Production and decay of supersymmetric particles at e⁺e⁻ colliders (ILC)



charginos

s-muons

Lightest supersymmetric particle stable in most models



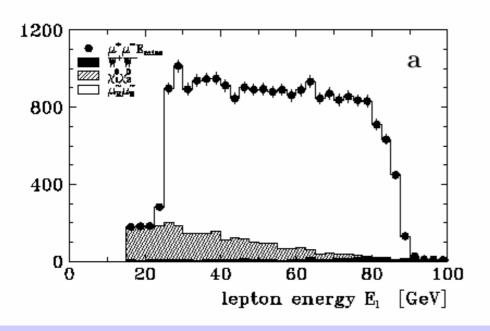
candidate for dark matter

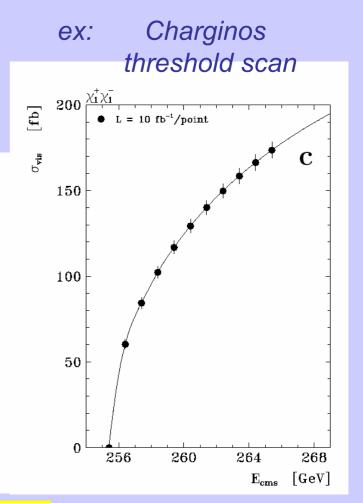
Experimental signature: missing energy

Measurement of sparticle masses

ex: Sleptons

lepton energy spectrum in continuum

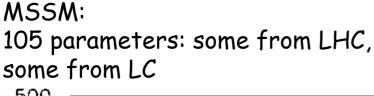


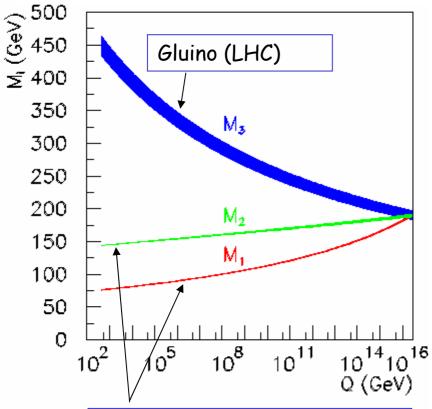


achievable accuracy:

 $\delta m/m \sim 10^{-3}$

Test of Unification





SUSY partners of electroweak bosons and Higgs

Extrapolation of SUSY parameters from weak to GUT scale (e.g. within mSUGRA)

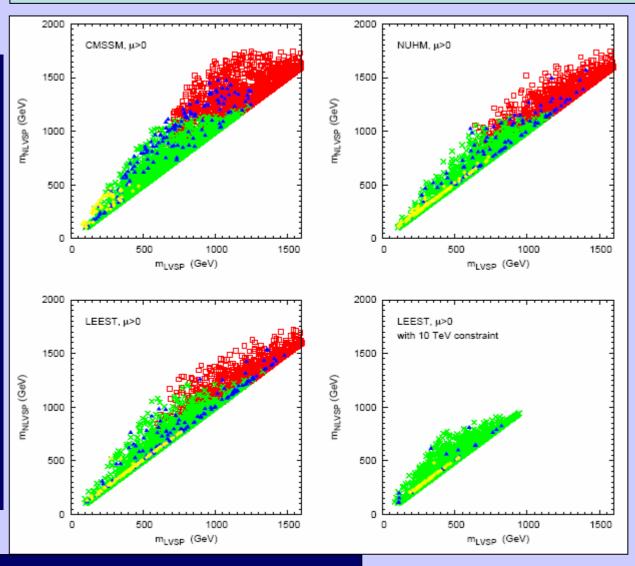
Gauge couplings unify at high energies,

Gaugino masses unify at same scale

Precision provided by LC for sleptons, charginos and neutralinos will allow to test if masses unify at same scale as forces



Sparticles may not be very light

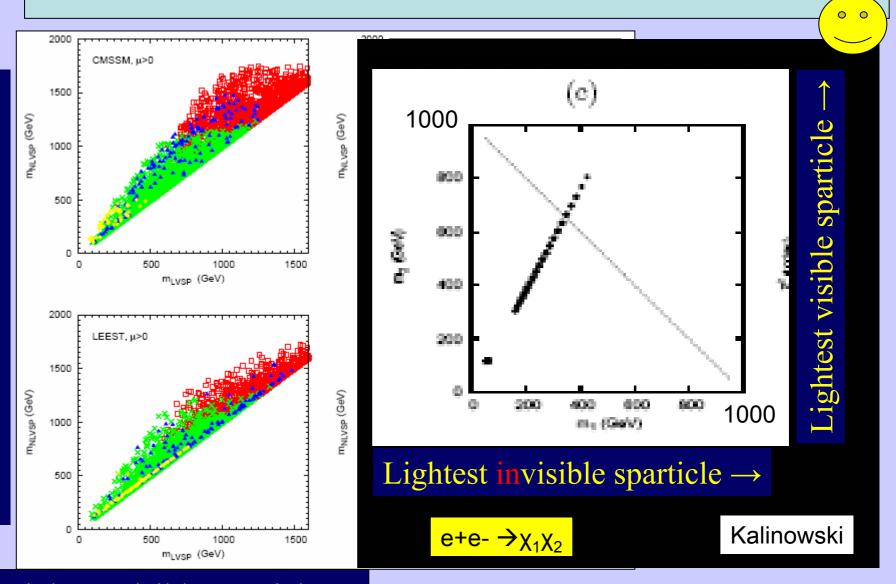


Lightest visible sparticle →

JE + Olive + Santoso + Spanos



LSP light in most cases

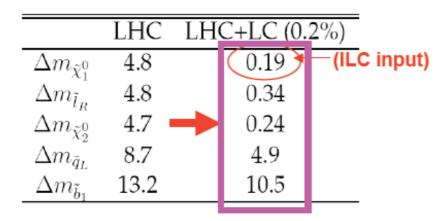


Lightest visible sparticle →

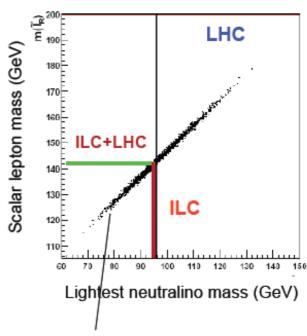


Using the $M(\chi^0_1)$ from ILC

300 fb-¹@LHC ∆M values in GeV

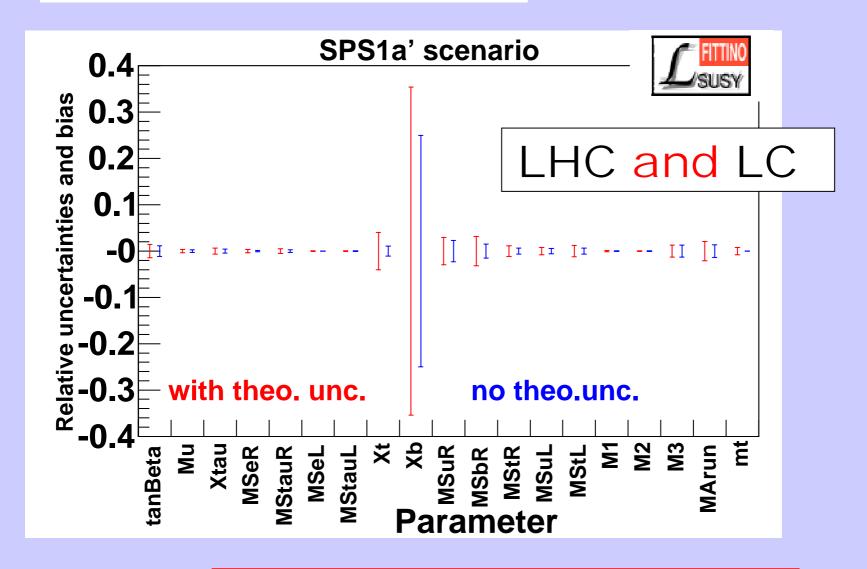


Significant improvements even if only $m(x^0)$ is measured at ILC



Strong correlation at LHC
An input from ILC resolve this correlation

MSSM parameters from global fit

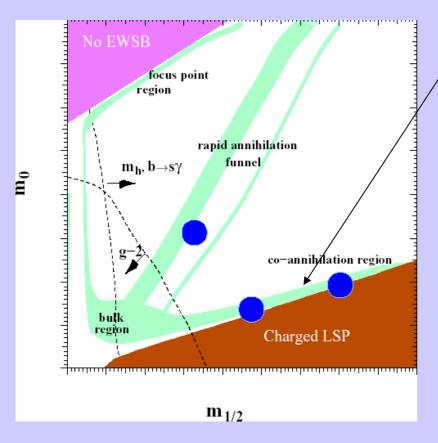


→ only possible with information from BOTH colliders

Dark Matter

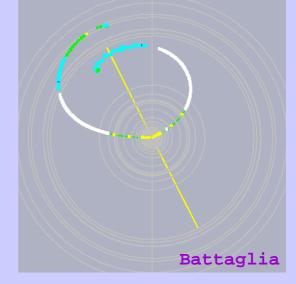
If SUSY LSP responsible for Cold Dark Matter, need accelerators to show that its properties are consistent with CMB data

- Future precision on $\Omega h^2 \sim 2\%$ (Planck) -> match this precision!
- WMAP points to certain difficult regions in parameter space:

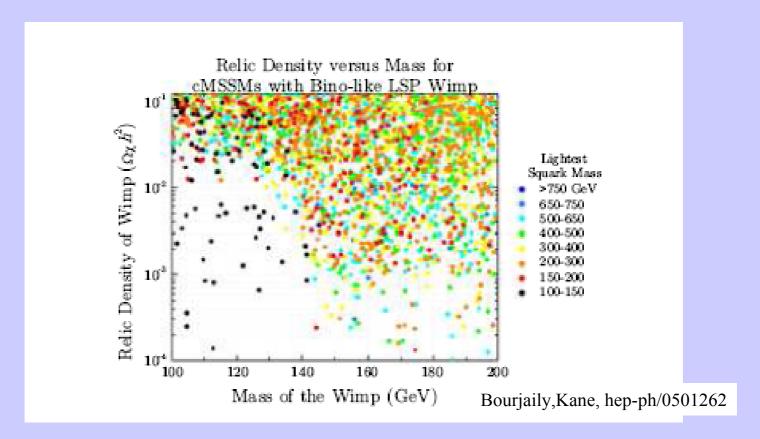


small
$$\Delta M = M_{\tilde{\ell}} - M_{\chi_1^0}$$

e.g. smuon pair production at 1TeV only two very soft muons! need to fight backgrounds



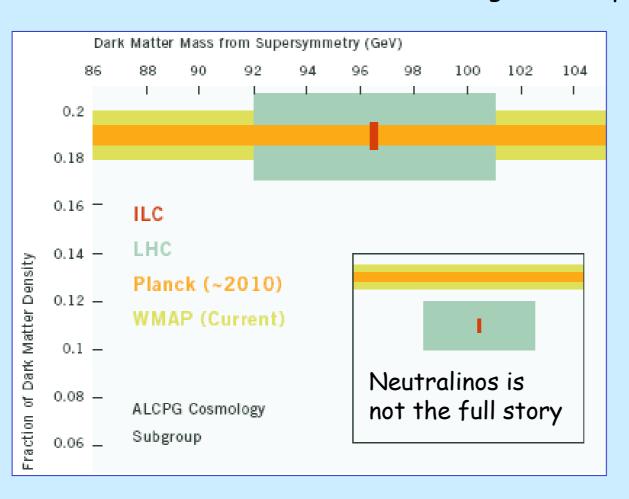
LSP responsible for relic density Ω_{CDM} ?



→ need to measure many parameters, in particular coupling to matter

Dark Matter and SUSY

Is dark matter linked to the Lightest Supersymmetric Particle?



Accel. and sat. data (WMAP and Planck):

complementary views of dark matter.

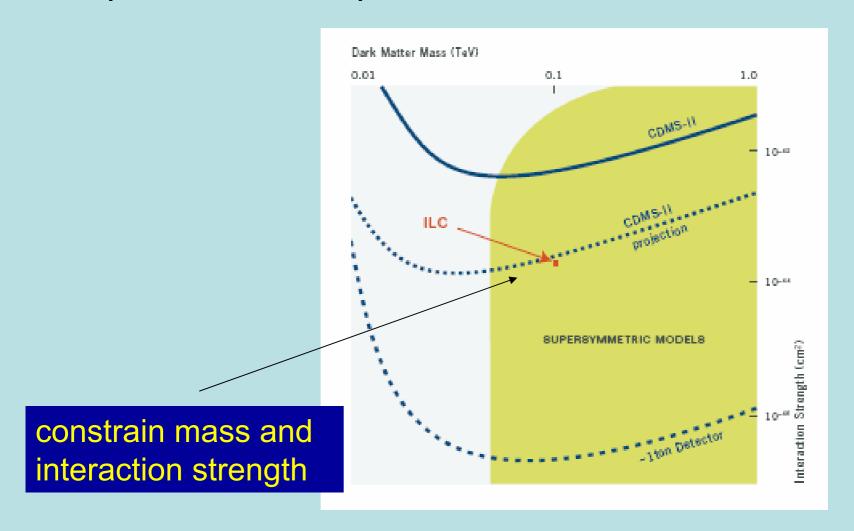
LHC/ILC: identify DM particle, measures its mass;

WMAP/Planck:

sensitive to total density of dark matter.

Together they establish the nature of dark matter.

Comparison with expectations from direct searches



Supersymmetry

Conclusions

The Linear Collider will be a unique tool for high precision measurements

- model independent determination of SUSY parameters
- determination of SUSY breaking mechanism
- extrapolation to GUT scale possible

LHC and LC together could shed first light into the Dark Universe

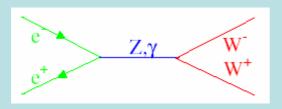
but what if

No Higgs boson(s) found....

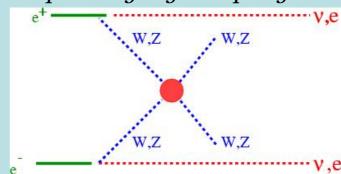
- ightarrow divergent W_L W_L \rightarrow W_L W_L amplitude in SM at $\Lambda^2 = o \left(\frac{4\pi\sqrt{2}}{G_F} \right) \approx (1.2 TeV)^2$
- → SM becomes inconsistent unless a new strong QCD-like interaction sets on
- → Goldstone bosons ("Pions") = W states ("technicolor")
- → no calculable theory until today in agreement with precision data

Experimental consequences: deviations in

triple gauge couplings



quartic gauge couplings:



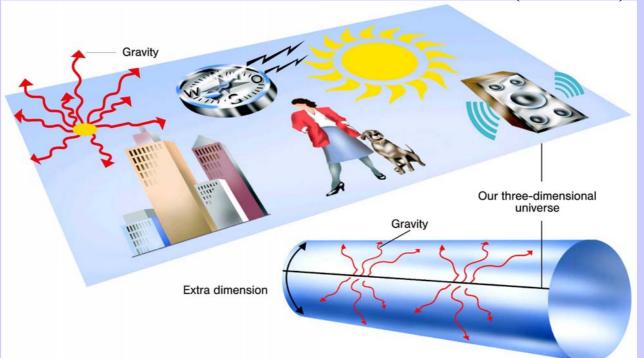
LC (800 GeV): sensitivity to energy scale Λ:

triple gauge couplings: ~ 8 TeV

quartic gauge couplings: ~ 3 TeV

⇒ complete threshold region covered

- Completely alternative approach to solve hierarchy problem: "There is no hierarchy problem"
- Suppose the SM fields live in "normal" 3+1 dim. space
- Gravity lives in $4 + \delta$ dimensions
- δ extra dimensions are curled to a small volume (radius R)



classical

$$V(r) = \frac{m_1 m_2}{M_{Pl}^2} \frac{1}{r}$$

$$G_N = 1/M_{Pl}^2$$

ADD-model:

 δ = new space dimension with radius R, which only communicates through gravity

$$V(r) \propto \frac{m_1 m_2}{M_D^{2+\delta} R^{\delta}} \frac{1}{r} \qquad r >> R$$

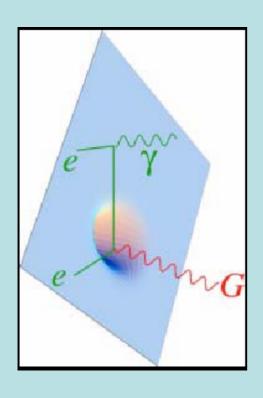
compare 4-dim and $4+\delta$ V(r):

$$M_{Pl}^2 = 8\pi R^{\delta} M_D^{\delta+2}$$

example $M_D = 1 \text{ TeV}$: for $\delta = 2(3) \rightarrow R = 1 \text{ mm(nm)}$

Extra dimensions provide an explanation for the hierarchy problem

String theory motivates brane models in which our world is confined to a membrane embedded in a higher dimensional space



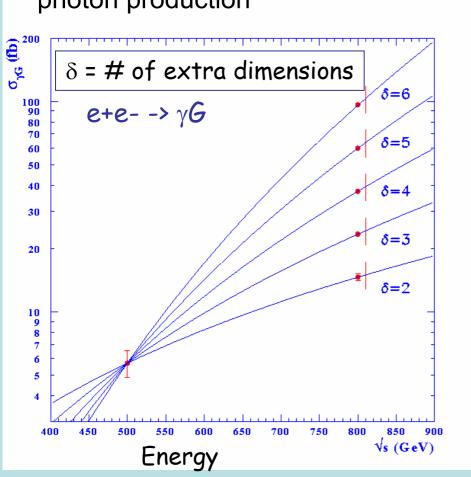
e.g. large extra dimensions:

Emission of gravitons into extra dimensions

Experimental signature

single photons

cross section for anomalous single photon production



measurement of cross sections at different energies allows to determine number and scale of extra dimensions

(500 fb-1 at 500 GeV, 1000 fb-1 at 800 GeV)

How does Polarisation help?

Polarisation reduces main background source: e+e-->vvy

=> heavy reduction of background

with $P(e^-) > 0$

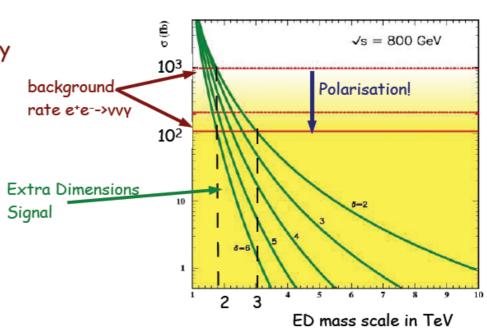
& $P(e^{+}) < 0$

Example: Extra Dimensions ete-->yG [TESLA-TDR]

precision of polarimetry important:

sensitivity to signal

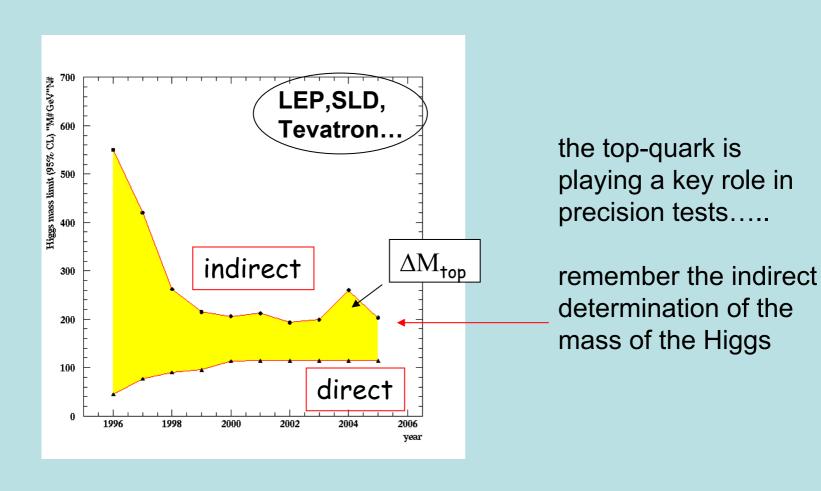
- uncertaintyof backgroundexpectation
- ~ uncertainty of polarisation!



W

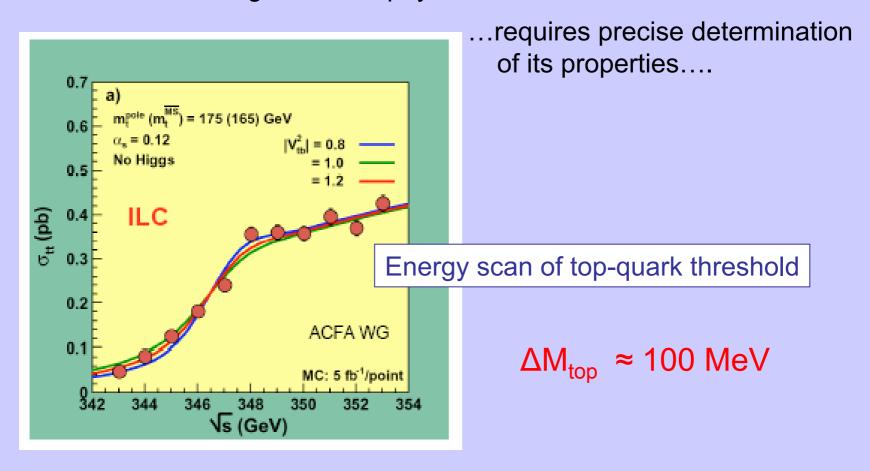
Precision measurements of SM processes

Precision electroweak tests

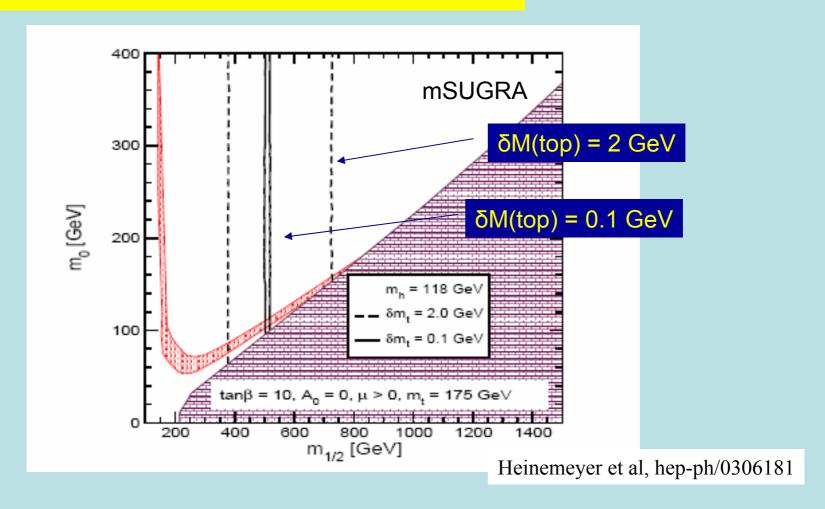


Precision electroweak tests

As the heaviest quark, the top-quark could play a key role in the understanding of flavour physics.....

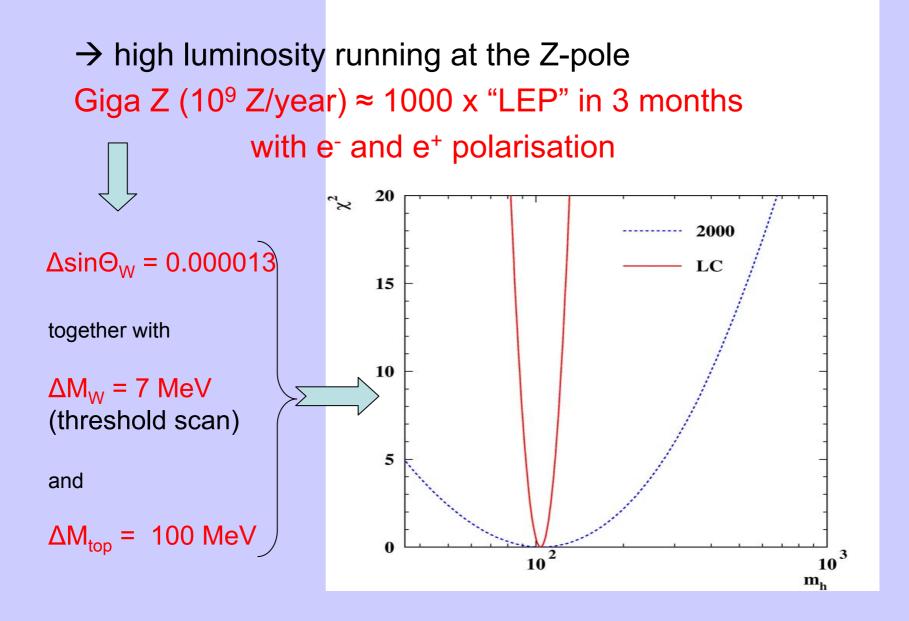


Precision electroweak tests



→ constrain allowed parameter space

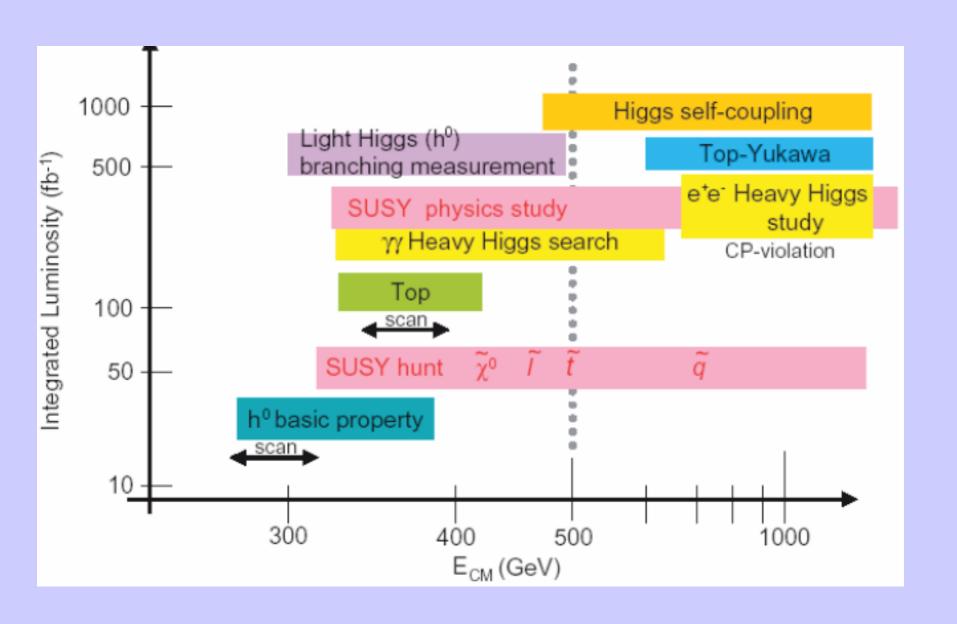
Precision Electroweak Tests



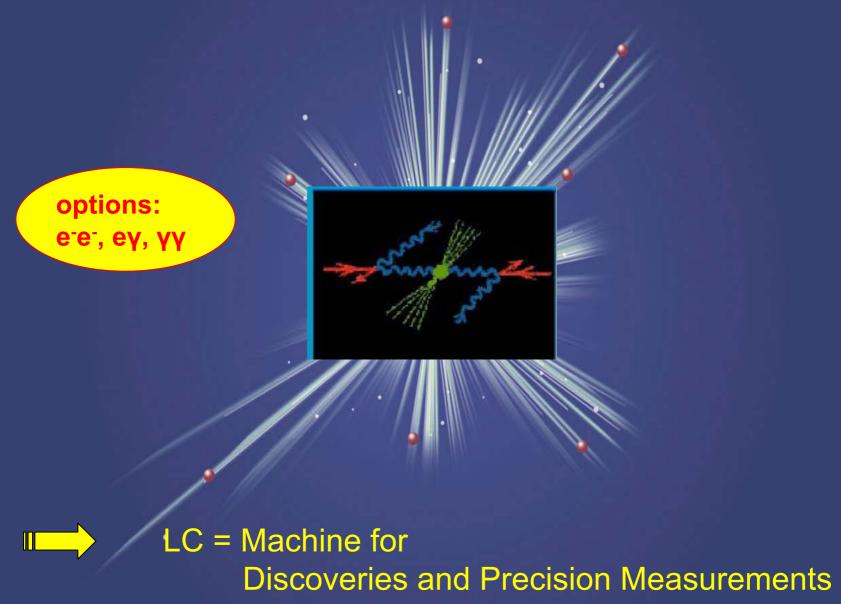
The (I)LC physics case

- 0. Top quark at threshold
- 'Light' Higgs (consistent with precision EW)
 ⇒ verify the Higgs mechanism is at work in all elements
- 2. 'Heavy' Higgs (inconsistent with precision EW)
 ⇒ verify the Higgs mechanism is at work in all elements
 ⇒ find out why prec. EW data are inconsistent
- 3. 1./2. + new states (SUSY, XD, little H, Z', ...)
 ⇒ precise spectroscopy of the new states
 ⇒ precision measurements of couplings of SM&new states
 properties of new particles above kinematic limit
- 4. No Higgs, no new states (inconsistent with precision EW)
 ⇒ find out why precision EW data are inconsistent
 ⇒ look for threshold effects of strong/delayed EWSB

Intermezzo: (I)LC Physics Reach



The power of an Electron-Positron Linear Collider



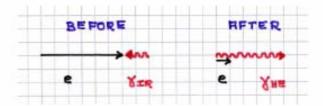
P.Zerwas, PLC05

The Collider Principle

4

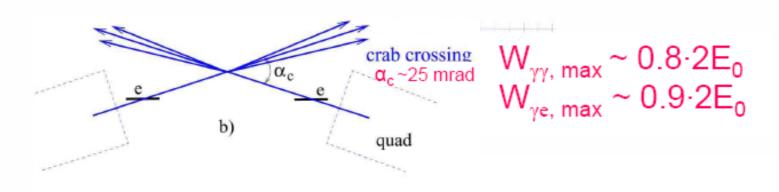
LC: conversion of e^{\pm} energy to γ energy by

Compton back-scattering of laser light



PIONEERS: I.F.Ginzburg, G.L.Kotkin, V.G.Serbo, V.I.Telnov

Pizma ZhETF 34 (1981) 514; JETP Lett. 34 (1982) 491; Nucl. Instr. Meth. 47 (1983) 205; (+ Panfil) ibid. A219 (1984) 5



P.Zerwas, PLC05

Luminosities

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$$\bullet$$
 $e\gamma$ luminosity

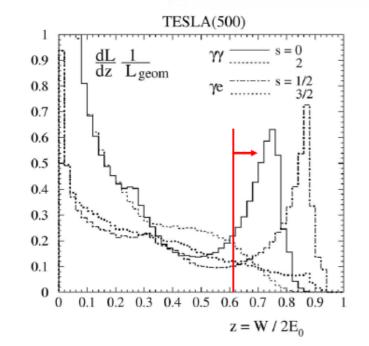
$$d\mathcal{L}/dm_{e\gamma}^2 = F(y = m_{e\gamma}^2)$$

$$m_{e\gamma} = M_{e\gamma}/\sqrt{s}$$

$$\gamma \gamma$$
 luminosity

$$d\mathcal{L}/dm_{\gamma\gamma}^2 = \int_{m_{\gamma\gamma}^2}^1 \frac{dy}{y} F(y) F(m_{\gamma\gamma}^2/y)$$

$$m_{\gamma\gamma} = M_{\gamma\gamma}/\sqrt{s}$$



$$\mathcal{LC}$$
 : $P_{-}=0.85$ / Laser : $P_{c}=-1$

\sqrt{s}_{ee} [GeV]	500	800
$m_{e\gamma} \ge 0.8 m_{e\gamma}^{max}$	0.9	1.3
$m_{\gamma\gamma} \ge 0.8 m_{\gamma\gamma}^{max}$	1.1	1.7
e^+e^-	3.4	5.8

$$\mathcal{L}_{\gamma\gamma}(\geq 0.8) \sim \frac{1}{3}\mathcal{L}_{ee}$$
 [1]

$$[10^{34} \mathrm{cm}^{-2} \mathrm{s}^{-1}]$$

P.Zerwas, PLC05

 $\gamma\gamma$ Cross sections

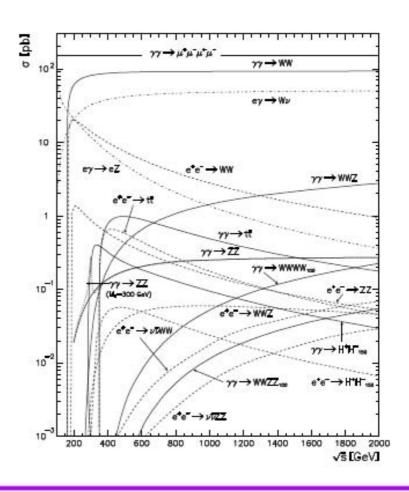
pointlike : $\gamma\gamma \sim 3$ to $10 \times e^+e^-$

examples : $t, W^{\pm}, H^{\pm}, \tilde{e}, \tilde{\chi}^{\pm}, \dots$

large size : 10 to 10^5 fb \Rightarrow

 10^3 to 10^7 evts

for $\mathcal{L} = 100 \text{ fb}^{-1}$



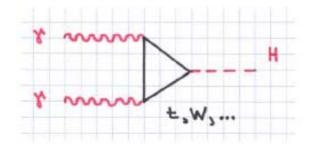
3

P.Zerwas, PLC05

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2. ELECTROWEAK SYMMETRY BREAKING

a) LIGHT HIGGS IN $\gamma\gamma$ COLLISIONS



$$\sigma_{\gamma\gamma} = \Gamma_{\gamma\gamma} \,\hat{\sigma} \, d\mathcal{L} / dm_{\gamma\gamma}^2(M_H^2)$$

- sharp onset for polarized beams
- helicities $\lambda_1 = \lambda_2$ enh. signal / sup. bkgd

$$\Gamma_{\gamma\gamma}(H)$$

$$\leq 140 \text{ GeV}: H \rightarrow bb$$

F

error :
$$\delta\Gamma_{\gamma\gamma}/\Gamma_{\gamma\gamma} \simeq 1.7\%$$

lifetime:
$$\tau_H = BR_{\gamma\gamma}/\Gamma_{\gamma\gamma} \sim 15\%$$

$$\geq 140 \; \mathrm{GeV} : H \to ZZ$$

F

V.Telnov, PLC05

Supersymmetry in $\gamma\gamma$

In supersymmetric model there are 5 Higgs bosons:

 h^0 light, with $m_h < 130$ GeV

 H^0, A^0 heavy Higgs bosons;

 H^+, H^- charged bosons.

 $M_H \approx M_A$, in e⁺e⁻ collisions H and A are produced in pairs (for certain param. region), while in $\gamma\gamma$ as the single resonances, therefore:

in e⁺e⁻ collisions $M_{H,A}^{max} \sim E_0$ (e⁺e⁻ $\rightarrow H + A$) in $\gamma\gamma$ collisions $M_{H,A}^{max} \sim 1.6E_0$ ($\gamma\gamma \rightarrow H(A)$)

For some SUSY parameters H,A can be seen only in γγ (but not in e+e- and LHC)

P.Zerwas, PLC05

3. SUPERSYMMETRY

• selectron production : $e^+e^- \to \tilde{e}\tilde{e} \quad \tilde{m} \leq \frac{1}{2}\sqrt{s}$

extended channel : $e\gamma \to \tilde{e}\tilde{\chi}^0$ $\tilde{m} \leq \sqrt{s} - \tilde{m}_{\chi}$

[i.g. larger than $\frac{1}{2}\sqrt{s}$

■ sneutrino production : $e^+e^- \to \tilde{\nu}\tilde{\nu}$ invisible for light ν [cf. SPS1a']

exploit $\tilde{\chi}^{\pm}$ decays in pairs : $\tilde{\chi}_1^{\pm} \to \ell^{\pm} \tilde{\nu}_{\ell}$ / difficult bkgds

<u>alternative channel</u>: $e\gamma \to \tilde{\nu}_e \tilde{\chi}_1^{\pm}$

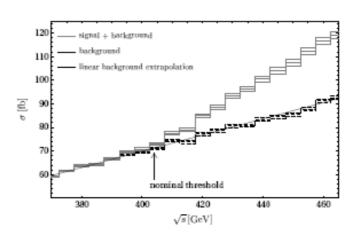
$$\rightarrow \tilde{\nu}_e \tilde{\nu}_\mu \, \mu$$

 $\sigma \sim \beta$: sharp onset \Rightarrow

F: Freitas, Porod, pmz

<u>threshold scan</u>: $\sqrt{s_{\gamma\gamma}} \ge m_{\tilde{\nu}} + m_{\tilde{\chi}}$

 $m_{\tilde{\nu}_e} = 169.8 \pm 3.2 \text{ GeV}$



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P.Zerwas, PLC05

6. SUMMARY

Central problems in high energy $\gamma\gamma$ physics:

- $\gamma\gamma \to H$: virtual high masses $H^{\pm}, KK, ...$ extended H, A search range* CP violation
- $e\gamma \rightarrow \tilde{e}, \, \tilde{\nu}:$ extended selectron / sneutrino search range*
- Majorana ν : signal in $e^-\gamma \to e^+WW$
- \bigcirc QCD: $\sigma_{\gamma\gamma}(had)$ and quark/gluon structure of photon

LC $\gamma\gamma$: valuable and unique complements to LHC and LC e^+e^- mode

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Outlook

Turn on of LHC entering an exciting phase of particle physics at the highest collision energies ever

Expect

- revolutionary advances in understanding the microcosm
- changes to our view of the early Universe

Results from LHC will guide the way

Expect

- period for decision taking on next steps in 2010 to 2012

Need

- R&D and technical design work now to enable these decisions
- ongoing for several projects
- global collaboration and stability on long timescales (reminder: first workshop on LHC was 1984)
- intensified efforts

Collaboration in network of HEP laboratories/institutes in Europe, Americas, Asia

Planning and execution of HEP projects today need global partnership

Use the exciting times ahead to establish such a partnership

Past decades saw precision studies of 5 % of our Universe → Discovery of the Standard Model

The LHC will soon deliver data

Preparations for the LC as a global project are under way

We are just at the beginning of exploring 95 % of the Universe

Past decades saw precision studies of 5 % of our Universe → Discovery of the Standard Model

The LHC will soon deliver data

Preparations for the LC as a global project are under way

We are just at the beginning of exploring 95 % of the Universe

the future is bright in the Dark Universe