

Search for dark photons at future e^+e^- colliders

Mikael Berggren¹, Sepideh Hosseini-Senvan^{1,2}

¹DESY, Hamburg, ²Universität Hamburg

European Physical Society Conference on High Energy Physics,
August 20-25, 2023
Universität Hamburg



CLUSTER OF EXCELLENCE
QUANTUM UNIVERSE



Universität Hamburg
DER FORSCHUNG | DER LEHRE | DER BILDUNG

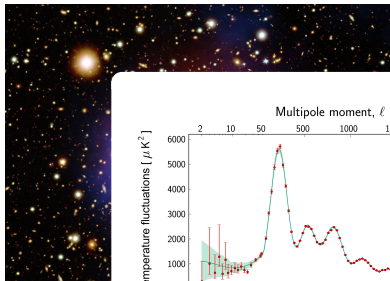


Introduction: Dark Matter

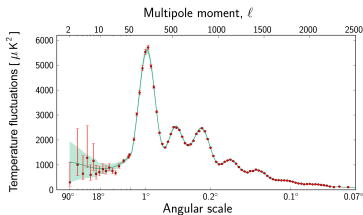


Bullet cluster

Introduction: Dark Matter

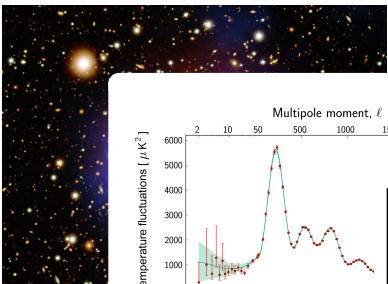


Bullet

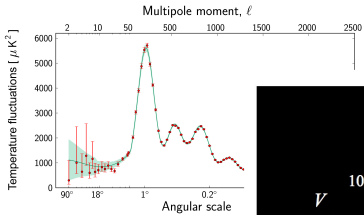


Planck CMB

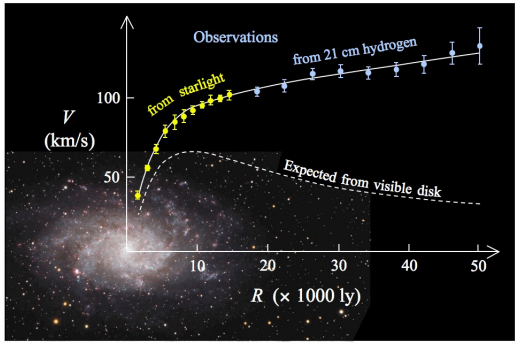
Introduction: Dark Matter



Bullet



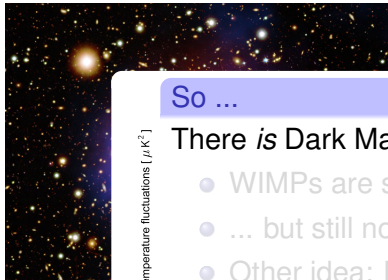
Planck CMB



M33 rotation curve



Introduction: Dark Matter



Bullet

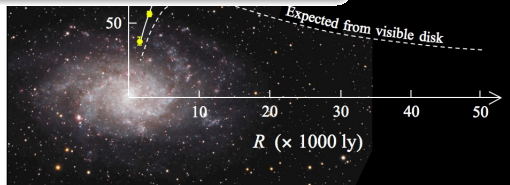
Temperature fluctuations [μK^2]

So ...

There *is* Dark Matter:

- WIMPs are still a good candidate....
- ... but still not seen.
- Other idea: Dark Matter lives in a dark, hidden sector \Rightarrow **FIPS**.

Plank CMB



M33 rotation curve

Introduction: Dark Matter



Bullet

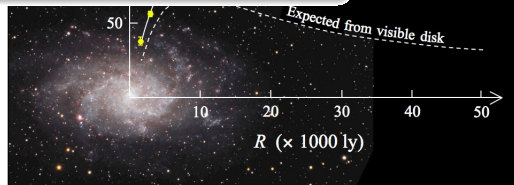
Temperature fluctuations [μK^2]

So ...

There *is* Dark Matter:

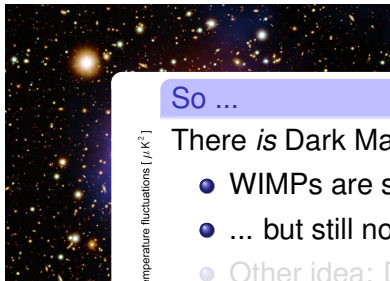
- WIMPs are still a good candidate....
- ... but still not seen.
- Other idea: Dark Matter lives in a dark, hidden sector \Rightarrow FIPS.

Plank CMB



M33 rotation curve

Introduction: Dark Matter



Bullet

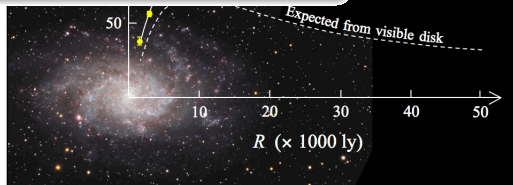
Temperature fluctuations [μK^2]

So ...

There *is* Dark Matter:

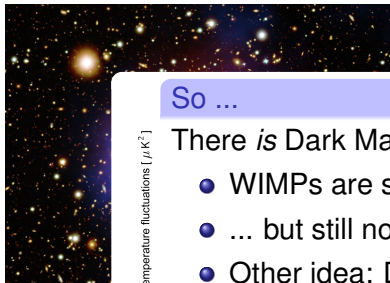
- WIMPs are still a good candidate....
- ... but still not seen.
- Other idea: Dark Matter lives in a dark, hidden sector \Rightarrow FIPS.

Plank CMB



M33 rotation curve

Introduction: Dark Matter



Bullet

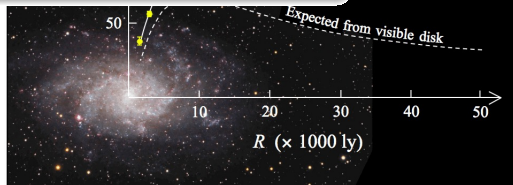
Temperature fluctuations [μK^2]

So ...

There *is* Dark Matter:

- WIMPs are still a good candidate....
- ... but still not seen.
- Other idea: Dark Matter lives in a dark, hidden sector \Rightarrow **FIPS**.

Planck CMB



M33 rotation curve

Introduction: FIPS

Feebly interacting particles is a class of models explaining dark matter and why it's not yet been seen in a **different way**.

- Generically, FIPS are models where rather than having heavy new particles with sizeable couplings, the new physics might be **light**, but much **more weakly coupled**.
- So, the reason why the BSM has not yet been seen is **not the lack of energy**, but the **lack of precision** - be it **luminosity, background contamination or detector performance**.

Introduction: FIPS

Types of FIPS, and how to detect them

- The Higgs Portal: Dark Higgs
- The fermions Portal: Sterile Neutrinos.
- The Pseudoscalar Portal: Axions (and ALPS)

and

- The Vector Portal: Dark photons

which is what we will discuss here.

Introduction: FIPS

Types of FIPS, and how to detect them

- The Higgs Portal: Dark Higgs
- The fermions Portal: Sterile Neutrinos.
- The Pseudoscalar Portal: Axions (and ALPS)

and

- **The Vector Portal: Dark photons**

which is what we will discuss here.

The Vector Portal - Dark Photons, A_D

- Assume that there is a **dark sector** with a **dark U(1)** symmetry
- The relevant part of the Lagrangian is

$$\mathcal{L}_{gauge} = -\frac{1}{4} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} - \frac{1}{4} \hat{Z}_{D\mu\nu} \hat{Z}_D^{\mu\nu} + \frac{1}{2} \frac{\epsilon}{\cos\theta_W} \hat{Z}_{D\mu\nu} \hat{B}^{\mu\nu}.$$
 \hat{B} is the ordinary U(1) field-strength tensor, and \hat{Z}_D that of the dark U(1).
- The Dark Photon might mix with the photon by *kinetic mixing* - the $\hat{Z}_D \hat{B}$ term - , so that $e^+e^- \rightarrow A_D \rightarrow f\bar{f}$ is possible.
- The (arbitrary) **mixing parameter** ϵ must be small, so the coupling is weak. There will be few events, but the decay will form a very narrow peak, or even a displaced vertex.
- Note that the dark photon itself is **not** the dark matter, since it isn't stable ... Something else in the dark sector that is stable is needed in addition.

The Vector Portal - Dark Photons, A_D

- Assume that there is a **dark sector** with a **dark U(1)** symmetry
- The relevant part of the Lagrangian is

$$\mathcal{L}_{gauge} = -\frac{1}{4} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} - \frac{1}{4} \hat{Z}_{D\mu\nu} \hat{Z}_D^{\mu\nu} + \frac{1}{2} \frac{\epsilon}{\cos\theta_W} \hat{Z}_{D\mu\nu} \hat{B}^{\mu\nu}.$$
 \hat{B} is the ordinary U(1) field-strength tensor, and \hat{Z}_D that of the dark U(1).
- The Dark Photon might mix with the photon by **kinetic mixing** - the $\hat{Z}_D \hat{B}$ term - , so that $e^+e^- \rightarrow A_D \rightarrow f\bar{f}$ is possible.
- The (arbitrary) **mixing parameter ϵ** must be small, so the coupling is weak. There will be few events, but the decay will form a very narrow peak, or even a displaced vertex.
- Note that the dark photon itself is **not** the dark matter, since it isn't stable ... Something else in the dark sector that is stable is needed in addition.

The Vector Portal - Dark Photons, A_D

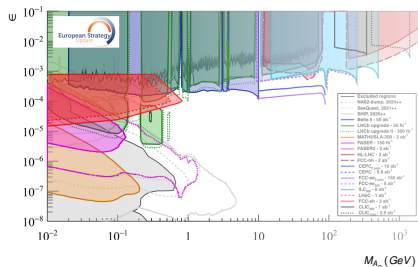
- Assume that there is a **dark sector** with a **dark U(1)** symmetry
- The relevant part of the Lagrangian is

$$\mathcal{L}_{gauge} = -\frac{1}{4} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} - \frac{1}{4} \hat{Z}_{D\mu\nu} \hat{Z}_D^{\mu\nu} + \frac{1}{2} \frac{\epsilon}{\cos\theta_W} \hat{Z}_{D\mu\nu} \hat{B}^{\mu\nu}.$$
 \hat{B} is the ordinary U(1) field-strength tensor, and \hat{Z}_D that of the dark U(1).
- The Dark Photon might mix with the photon by **kinetic mixing** - the $\hat{Z}_D \hat{B}$ term - , so that $e^+e^- \rightarrow A_D \rightarrow f\bar{f}$ is possible.
- The (arbitrary) **mixing parameter** ϵ must be small, so the coupling is weak. There will be few events, but the decay will form a very narrow peak, or even a displaced vertex.
- Note that the dark photon itself is **not** the dark matter, since it isn't stable ... Something else in the dark sector that **is** stable is needed in addition.

Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

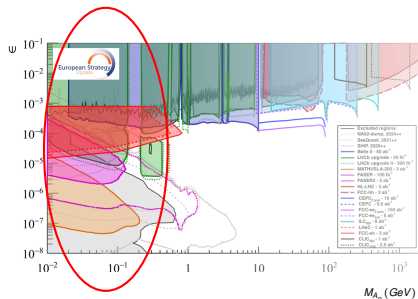
- All experiment, log mass-scale (from the EPPSU briefing book.)
- Masses up to ~ 1 GeV: LLPs detected in **Beam-dump** experiments. Sensitive to very small couplings
- Beyond that: colliders
 - Up to 10 GeV: B factories - extremely high luminosity.
 - Then: e^+e^- up to their maximum energy
 - ... and beyond that pp colliders



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

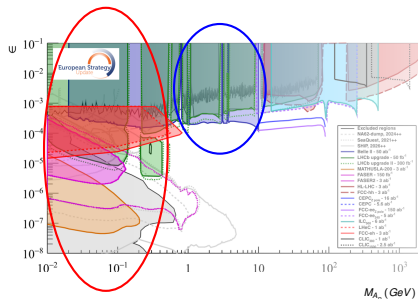
- All experiment, log mass-scale (from the EPPSU briefing book.)
- Masses up to ~ 1 GeV: LLPs detected in **Beam-dump** experiments. Sensitive to very small couplings
- Beyond that: colliders
 - Up to 10 GeV: B factories - extremely high luminosity.
 - Then: e^+e^- up to their maximum energy
 - ... and beyond that pp colliders



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

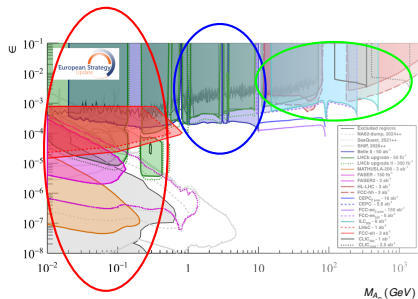
- All experiment, log mass-scale (from the EPPSU briefing book.)
- Masses up to ~ 1 GeV: LLPs detected in **Beam-dump** experiments. Sensitive to very small couplings
- Beyond that: colliders
 - Up to 10 GeV: **B factories** - extremely high luminosity.
 - Then: e^+e^- up to their maximum energy
 - ... and beyond that pp colliders.



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

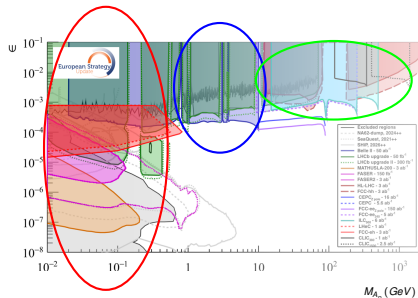
- All experiment, log mass-scale (from the EPPSU briefing book.)
- Masses up to ~ 1 GeV: LLPs detected in **Beam-dump** experiments. Sensitive to very small couplings
- Beyond that: colliders
 - Up to 10 GeV: **B factories** - extremely high luminosity.
 - Then: e^+e^- up to their maximum energy
 - ... and beyond that pp colliders.



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

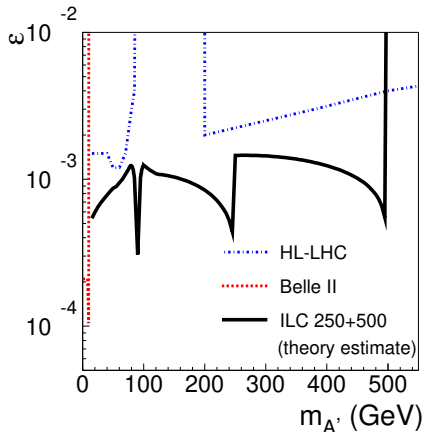
- All experiment, log mass-scale (from the EPPSU briefing book.)
- Masses up to ~ 1 GeV: LLPs detected in **Beam-dump** experiments. Sensitive to very small couplings
- Beyond that: colliders
 - Up to 10 GeV: **B factories** - extremely high luminosity.
 - Then: e^+e^- up to their maximum energy
 - ... and beyond that pp colliders.



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

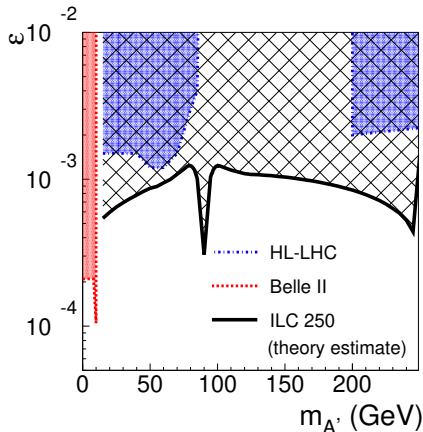
- BELLE II, ILC 250+500 and HL-LHC on **linear mass scale**.
- ... and zoomed to Higgs factory reach.



Dark Photon limits from the EPPSU

Current projections from the European Particle Physics Strategy Update of 2019

- BELLE II, ILC 250+500 and HL-LHC on **linear mass scale**.
- ... and zoomed to Higgs factory reach.



Dark Photon limits from the EPPSU

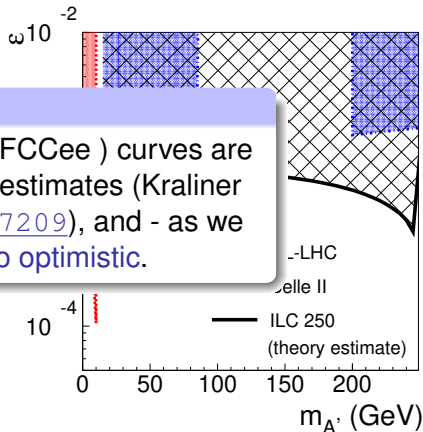
Current projections from the European Particle Physics Strategy Update of 2019

- BELLE II, ILC 250+500 and HL-LHC on **linear mass scale**.

- ... and zero mass dark photon factory re

Attention !

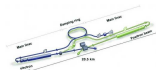
The "ILC" (and CepC, FCCee) curves are very simplistic theory estimates (Kraliner & al. [arxiv:1503.07209](https://arxiv.org/abs/1503.07209)), and - as we will see - are much **too optimistic**.



Higgs factories and beyond

The Bestiary of proposed future e^+e^- colliders, and their detectors

ILC



CEPC



FCC-ee



CLIC



This talk
uses ILD

ILD



CEPC Baseline



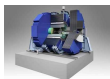
IDEA



CLICdp



SiD



FST



CEPC 4th concept



CLD



slide stolen from B. Dudar

The circular machines are Higgs (and Z) factories, the linear ones can extend far beyond in energy.

Higgs factories and beyond

The Bestiary of proposed future e^+e^- colliders, and their detectors

ILC

CEPC

FCC-ee

CLIC

Why dark photons at Higgs factories ?

... when LEP II reached almost as high energies ?

- At least 1000 times the luminosity !
- and polarisation, triggerless running, 40 years for detector development ...
- \Rightarrow Enormous **increase in sensitivity** !

This talk uses ILD

CLICdp

SiD

slide stolen from B. Dudar

The circular machines are Higgs (and Z) factories, the linear ones can extend far beyond in energy.

Higgs factories and beyond

The Bestiary of proposed future e^+e^- colliders, and their detectors

ILC

CEPC

FCC-ee

CLIC

Why dark photons at Higgs factories ?

... when LEP II reached almost as high energies ?

- At least 1000 times the luminosity !
- and polarisation, triggerless running, 40 years for detector development ...
- ⇒ Enormous increase in sensitivity !

This talk uses ILD

CLICdp

SiD

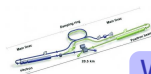
slide stolen from B. Dudar

The circular machines are Higgs (and Z) factories, the linear ones can extend far beyond in energy.


Higgs factories and beyond

The Bestiary of proposed future e^+e^- colliders, and their detectors


ILC




CEPC



FCC-ee



CLIC




Why dark photons at Higgs factories ?


... when LEP II reached almost as high energies ?

- At least 1000 times the luminosity !
- and polarisation, triggerless running, 40 years for detector development ...
- \Rightarrow Enormous **increase in sensitivity** !


This talk uses ILD



SiD



CLICdp



The circular machines are Higgs (and Z) factories, the linear ones can extend far beyond in energy.

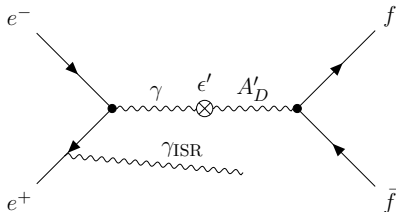
Dark Photons at e^+e^- beyond SuperKEKB

- Signal process:

$e^+e^- \rightarrow \gamma_{ISR} A_D \rightarrow \mu^+ \mu^- \gamma_{ISR}$,
 where E_{ISR} is such that the
 recoil-mass against the ISR is
 M_{A_D}

- Both σ and Γ scales with ϵ^2 .

- One could hope to exclude $\sigma > \mathcal{O}(1 \text{ fb})$
- For the corresponding ϵ^2 , Γ is $\mathcal{O}(10 \text{ keV})$ to $\mathcal{O}(10 \text{ MeV})$.
- \Rightarrow detector resolution will determine the peak-width
- \Rightarrow decay is prompt ($cr < 1 \text{ nm}$).



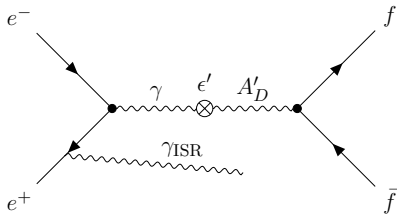
Dark Photons at e^+e^- beyond SuperKEKB

- Signal process:

$e^+e^- \rightarrow \gamma_{ISR} A_D \rightarrow \mu^+ \mu^- \gamma_{ISR}$,
 where E_{ISR} is such that the
 recoil-mass against the ISR is
 M_{A_D}

- Both σ and Γ scales with ϵ^2 .

- One could hope to exclude $\sigma > \mathcal{O}(1 \text{ fb})$
- For the corresponding ϵ^2 , Γ is $\mathcal{O}(10 \text{ keV})$ to $\mathcal{O}(10 \text{ MeV})$.
- \Rightarrow detector resolution will determine the peak-width
- \Rightarrow decay is prompt ($c\tau < 1 \text{ nm}$).



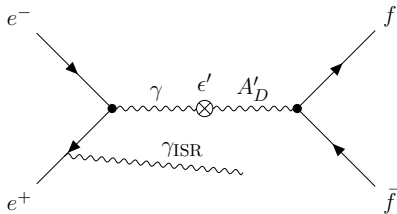
Dark Photons at e^+e^- beyond SuperKEKB

- Signal process:

$e^+e^- \rightarrow \gamma_{ISR} A_D \rightarrow \mu^+ \mu^- \gamma_{ISR}$,
 where E_{ISR} is such that the recoil-mass against the ISR is M_{A_D}

- Both σ and Γ scales with ϵ^2 .

- One could hope to exclude $\sigma > \mathcal{O}(1 \text{ fb})$
- For the corresponding ϵ^2 , Γ is $\mathcal{O}(10 \text{ keV})$ to $\mathcal{O}(10 \text{ MeV})$.
- \Rightarrow detector resolution will determine the peak-width
- \Rightarrow decay is prompt ($c\tau < 1 \text{ nm}$).



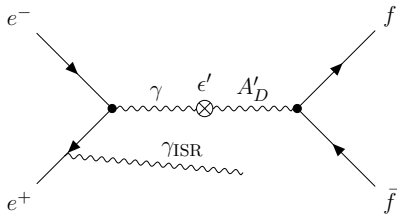
Dark Photons at e^+e^- beyond SuperKEKB

- Signal process:

$e^+e^- \rightarrow \gamma_{ISR} A_D \rightarrow \mu^+ \mu^- \gamma_{ISR}$,
 where E_{ISR} is such that the recoil-mass against the ISR is M_{A_D}

- Both σ and Γ scales with ϵ^2 .

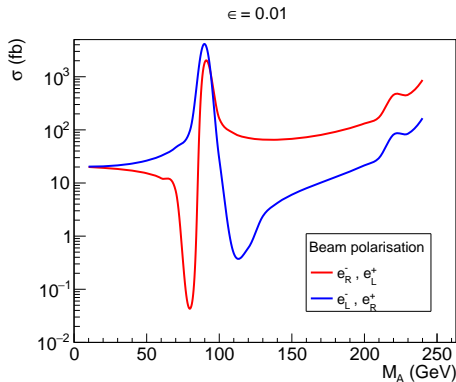
- One could hope to exclude $\sigma > \mathcal{O}(1 \text{ fb})$
- For the corresponding ϵ^2 , Γ is $\mathcal{O}(10 \text{ keV})$ to $\mathcal{O}(10 \text{ MeV})$.
- \Rightarrow detector resolution will determine the peak-width
- \Rightarrow decay is prompt ($c\tau < 1 \text{ nm}$).



Properties of Dark Photon production and decay

Generate events using the UFO files describing the model of Curtin & al. ([arxiv:1412.0018](https://arxiv.org/abs/1412.0018)) interfaced to Whizard 3.0.

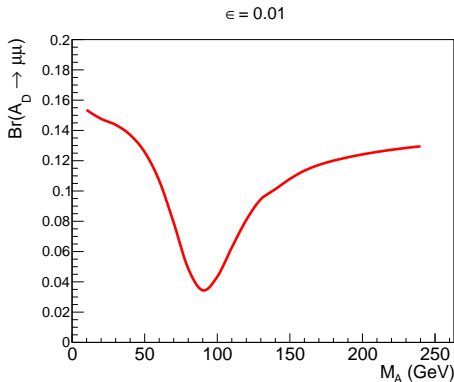
- Production cross-section σ for fully polarised beams.
- $\text{BR}(A_D \rightarrow \mu\mu)$
- “Effective” $\sigma \times \text{BR}$ (meaning: $N_{\text{events}} = \sigma_{\text{eff}} \times \int \mathcal{L}$) for
 - ILC expected polarisations, and
 - unpolarised beams.
- Γ_{tot}



Properties of Dark Photon production and decay

Generate events using the UFO files describing the model of Curtin & al. ([arxiv:1412.0018](https://arxiv.org/abs/1412.0018)) interfaced to Whizard 3.0.

- Production cross-section σ for fully polarised beams.
- $\text{BR}(A_D \rightarrow \mu\mu)$
- “Effective” $\sigma \times \text{BR}$ (meaning: $N_{events} = \sigma_{eff} \times \int \mathcal{L}$) for
 - ILC expected polarisations, and
 - unpolarised beams.
- Γ_{tot}

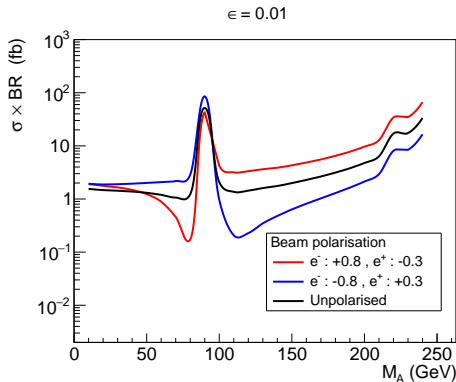


Properties of Dark Photon production and decay

Generate events using the UFO files describing the model of Curtin & al. ([arxiv:1412.0018](https://arxiv.org/abs/1412.0018)) interfaced to Whizard 3.0.

- Production cross-section σ for fully polarised beams.
- $\text{BR}(A_D \rightarrow \mu\mu)$
- “Effective” $\sigma \times \text{BR}$ (meaning: $N_{\text{events}} = \sigma_{\text{eff}} \times \int \mathcal{L}$) for
 - ILC expected polarisations, and
 - unpolarised beams.

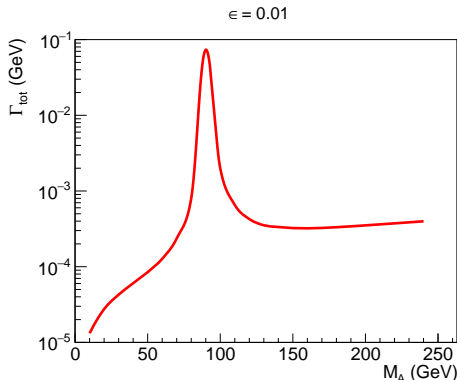
• Γ_{tot}



Properties of Dark Photon production and decay

Generate events using the UFO files describing the model of Curtin & al. ([arxiv:1412.0018](https://arxiv.org/abs/1412.0018)) interfaced to Whizard 3.0.

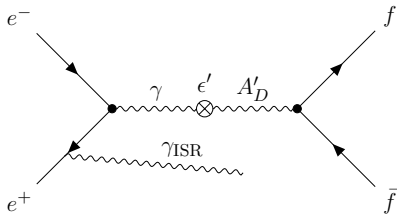
- Production cross-section σ for fully polarised beams.
- $\text{BR}(A_D \rightarrow \mu\mu)$
- “Effective” $\sigma \times \text{BR}$ (meaning: $N_{\text{events}} = \sigma_{\text{eff}} \times \int \mathcal{L}$) for
 - ILC expected polarisations, and
 - unpolarised beams.
- Γ_{tot}



Dark Photons in a real detector

Pass such generated events through the full `Geant4`-based simulation (`ddsim`) and reconstruction (`Marlin`) of ILD.

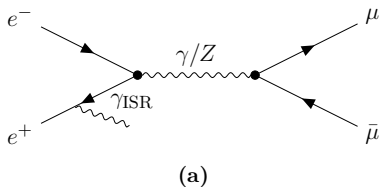
- Select events with **two muons**, and possibly an **isolated photon** - nothing else.
- Include **all** (fully simulated) SM background.
- Look for an **arbitrarily small peak** in the $M(\mu\mu)$ distribution, with natural width $\ll \delta_{det}(M)$, over the SM background
- ... which varies with M_A , and is not only $e^+e^- \rightarrow \mu^+\mu^- + ISR$
- The target.



Dark Photons in a real detector

Pass such generated events through the full Geant4-based simulation (ddsim) and reconstruction (Marlin) of ILD.

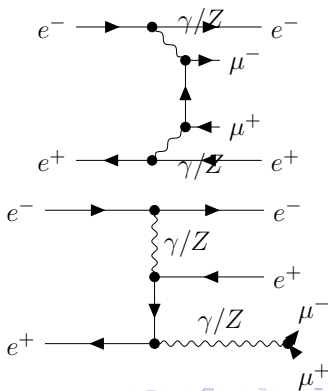
- Select events with **two muons**, and possibly an **isolated photon** - nothing else.
- Include **all** (fully simulated) SM background.
- Look for an **arbitrarily small peak** in the $M(\mu\mu)$ distribution, with natural width $\ll \delta_{det}(M)$, over the SM background
- ... which varies with M_A , and is not only $e^+e^- \rightarrow \mu^+\mu^- + ISR$
- The target.



Dark Photons in a real detector

Pass such generated events through the full `Geant4`-based simulation (`ddsim`) and reconstruction (`Marlin`) of ILD.

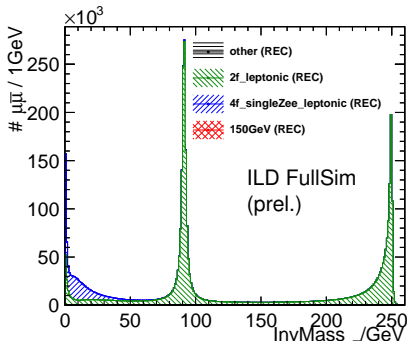
- Select events with **two muons**, and possibly an **isolated photon** - nothing else.
- Include **all** (fully simulated) SM background.
- Look for an **arbitrarily small peak** in the $M(\mu\mu)$ distribution, with natural width $\ll \delta_{det}(M)$, over the SM background
- ... which varies with M_A , and is not only $e^+e^- \rightarrow \mu^+\mu^- + ISR$
- The target.



Dark Photons in a real detector

Pass such generated events through the full `Geant4`-based simulation (`ddsim`) and reconstruction (`Marlin`) of ILD.

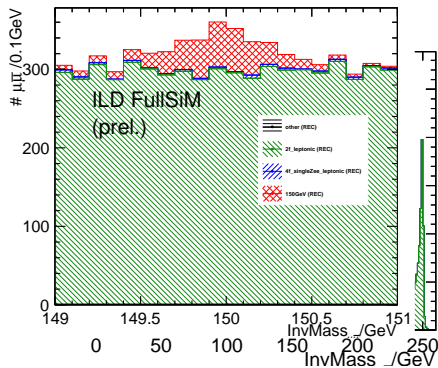
- Select events with **two muons**, and possibly an **isolated photon** - nothing else.
- Include **all** (fully simulated) SM background.
- Look for an **arbitrarily small peak** in the $M(\mu\mu)$ distribution, with natural width $\ll \delta_{det}(M)$, over the SM background
- ... which varies with M_A , and is not only $e^+e^- \rightarrow \mu^+\mu^- + ISR$
- The target.



Dark Photons in a real detector

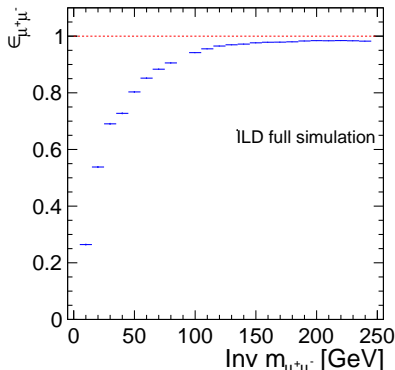
Pass such generated events through the full `Geant4`-based simulation (`ddsim`) and reconstruction (`Marlin`) of `ILD`.

- Select events with **two muons**, and possibly an **isolated photon** - nothing else.
- Include **all** (fully simulated) SM background.
- Look for an **arbitrarily small peak** in the $M(\mu\mu)$ distribution, with natural width $\ll \delta_{det}(M)$, over the SM background
- ... which varies with M_A , and is not only $e^+e^- \rightarrow \mu^+\mu^- + ISR$
- The target.



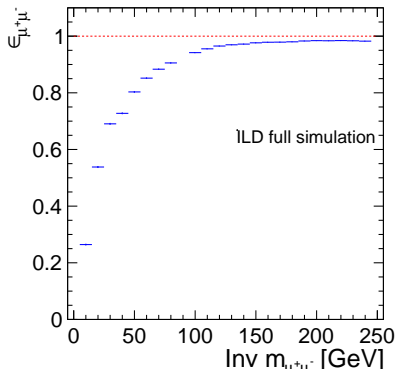
Dark Photons in a real detector

- Efficiency to find two muons.
- Why so low ? ILD track-finding is 100 % efficient down to $p_T \sim 300$ MeV and angles to the beam above $\sim 10^\circ$!?
- Here's why: **Angular distribution** of the muons - we need to see both to get a pair, obviously!



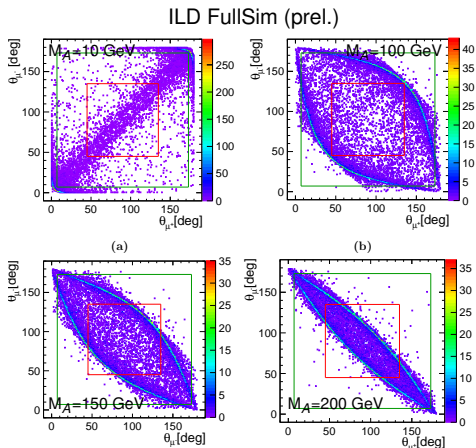
Dark Photons in a real detector

- Efficiency to find two muons.
- **Why so low ?** ILD track-finding is 100 % efficient down to $p_T \sim 300$ MeV and angles to the beam above $\sim 10^\circ$!?
- Here's why: **Angular distribution** of the muons - we need to see both to get a pair, obviously!



Dark Photons in a real detector

- Efficiency to find two muons.
- **Why so low** ? ILD track-finding is 100 % efficient down to $p_T \sim 300$ MeV and angles to the beam above $\sim 10^\circ$!?
- Here's why: **Angular distribution** of the muons - we need to see both to get a pair, obviously!



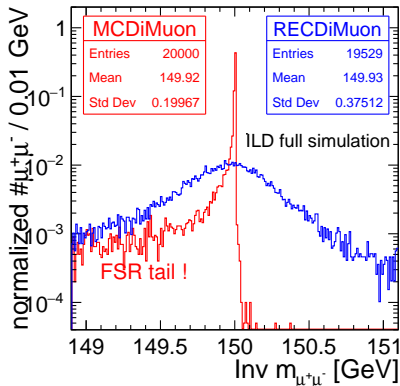
Dark Photons in a real detector

- Mass resolution:

$M = p_1 p_2 (1 - \cos \theta_{12})$, and the ISR is along the beam and $\sigma(1/p_T)$ vs. p is constant, so error-propagation gives $\sigma_M \propto M^2$, right ?

- Wrong.

- Due to M.S., for $p \gtrsim 100$ GeV, $\sigma(1/p_T)$ is not constant, rather $\propto p^{-1}$.
- Strong dependence on θ in the forward region.
- and most muons are below 100 GeV
- and are not in the barrel
- and are not on the curve for ISR at zero angle



Dark Photons in a real detector

- Mass resolution:

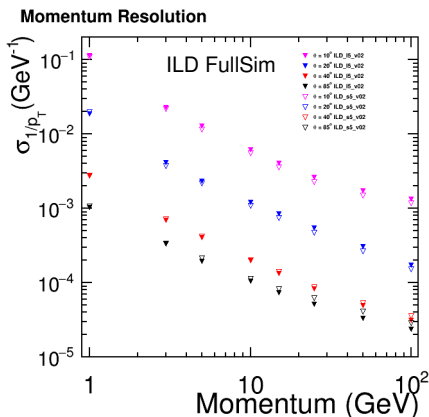
$M = p_1 p_2 (1 - \cos \theta_{12})$, and the ISR is along the beam and $\sigma(1/p_T)$ vs. p is constant, so error-propagation gives $\sigma_M \propto M^2$, right ?

- **Wrong.**

- Due to M.S., for $p \gtrsim 100$ GeV, $\sigma(1/p_T)$ is **not** constant, rather $\propto p^{-1}$.

Strong dependence on θ in the forward region.

- and most muons are below 100 GeV
- and are not in the barrel
- and are not on the curve for ISR at zero angle



Dark Photons in a real detector

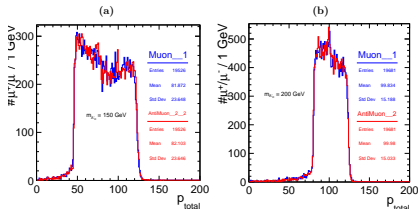
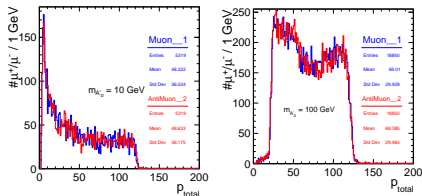
- Mass resolution:

$M = p_1 p_2 (1 - \cos \theta_{12})$, and the ISR is along the beam and $\sigma(1/p_T)$ vs. p is constant, so error-propagation gives $\sigma_M \propto M^2$, right?

- Wrong.**

- Due to M.S., for $p \gtrsim 100$ GeV, $\sigma(1/p_T)$ is **not** constant, rather $\propto p^{-1}$. **Strong** dependence on θ in the forward region.
- and most muons **are** below 100 GeV
- and **are not** in the barrel
- and **are not** on the curve for ISR at zero angle

ILD FullSim (prel.)



Dark Photons in a real detector

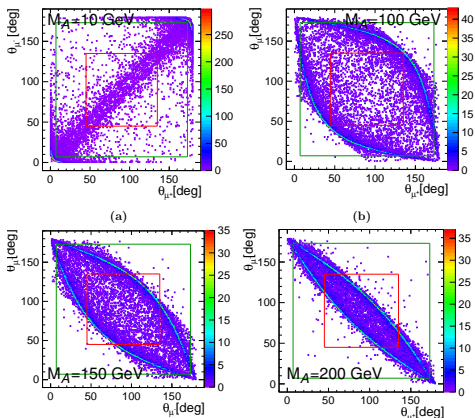
- Mass resolution:

$M = p_1 p_2 (1 - \cos \theta_{12})$, and the ISR is along the beam and $\sigma(1/p_T)$ vs. p is constant, so error-propagation gives $\sigma_M \propto M^2$, right ?

- **Wrong.**

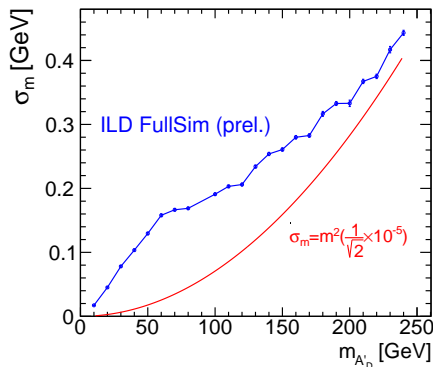
- Due to M.S., for $p \gtrsim 100$ GeV, $\sigma(1/p_T)$ is **not** constant, rather $\propto p^{-1}$. **Strong** dependence on θ in the forward region.
- and most muons **are** below 100 GeV
- and **are not** in the barrel
- and **are not** on the curve for ISR at zero angle

ILD FullSim(prel.)



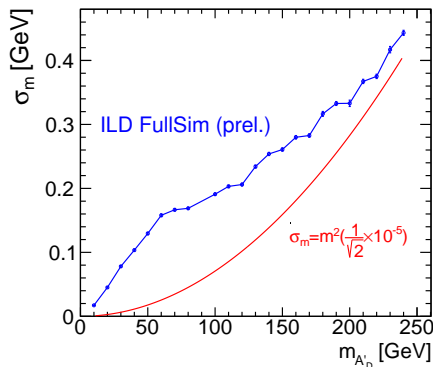
Dark Photons in a real detector

- Bottom line: **None** of the assumptions on the mass-resolution - the **red** curve - used for the EPPSU curve are valid. The **correct** full simulation values are the **blue** curve.
- The resolution will **vary a lot event-by-event** - with angle and momentum of the muons, and the angle of the ISR.
- \Rightarrow **Event-by-event simulation is essential.**



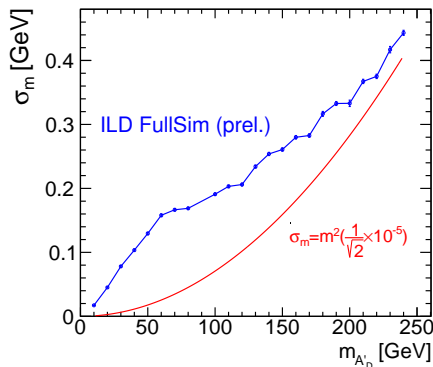
Dark Photons in a real detector

- Bottom line: **None** of the assumptions on the mass-resolution - the **red** curve - used for the EPPSU curve are valid. The **correct** full simulation values are the **blue** curve.
- The resolution will **vary a lot event-by-event** - with angle and momentum of the muons, and the angle of the ISR.
- \Rightarrow **Event-by-event simulation is essential.**



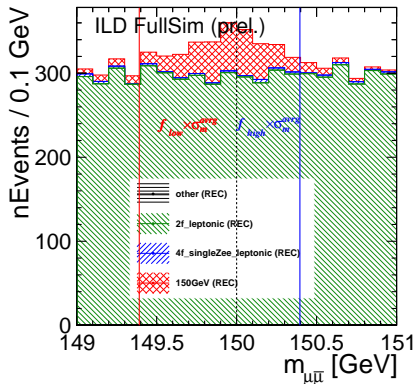
Dark Photons in a real detector

- Bottom line: **None** of the assumptions on the mass-resolution - the **red** curve - used for the EPPSU curve are valid. The **correct** full simulation values are the **blue** curve.
- The resolution will **vary a lot event-by-event** - with angle and momentum of the muons, and the angle of the ISR.
- \Rightarrow **Event-by-event simulation is essential.**



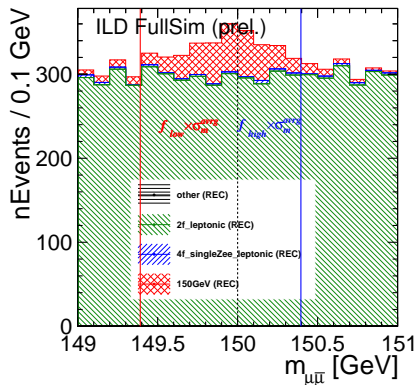
Dark Photons in a real detector

- However, the uncertainty is known, event-by-event, since the track-fit covariance matrix is output from the fit !
- Use this to optimise the search:
 - Define the signal-window as a factor times the event-specific σ_m .
 - factors are different above and below the tested mass, because of FSR.
 - Optimise these factors for sensitivity at each tested mass.



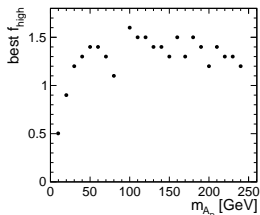
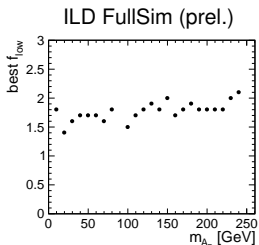
Dark Photons in a real detector

- However, the uncertainty is known, event-by-event, since the track-fit covariance matrix is output from the fit !
- Use this to optimise the search:
 - Define the signal-window as a factor times the event-specific σ_m .
 - factors are different above and below the tested mass, because of FSR.
 - Optimise these factors for sensitivity at each tested mass.



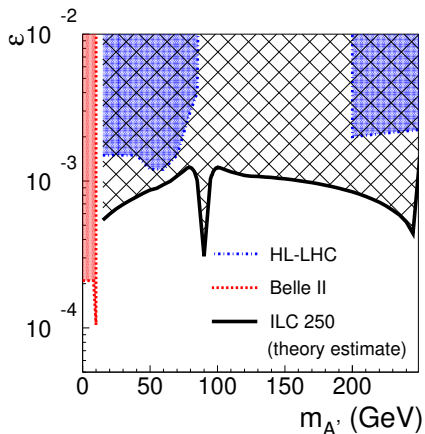
Dark Photons in a real detector

- However, the uncertainty is known, event-by-event, since the track-fit covariance matrix is output from the fit !
- Use this to optimise the search:
 - Define the signal-window as a factor times the event-specific σ_m .
 - factors are different above and below the tested mass, because of FSR.
 - Optimise these factors for sensitivity at each tested mass.



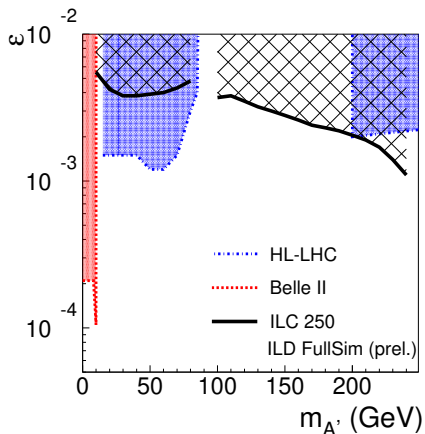
Dark Photons in a real detector

- Compared to the theory curve ...
- ... this is **the (current) result with full simulation**.
- At the highest mass, the correct limit is a factor two higher, a factor four at 100 GeV.
- This is **due to the correct estimate of the error**.
- Below M_Z , the difference is larger, and HL-LHC limits are expected to be stronger.
- Here, the reason is both **the correct error-estimate**, but also **the much larger background from non- $Z \rightarrow \mu\mu$ processes**.



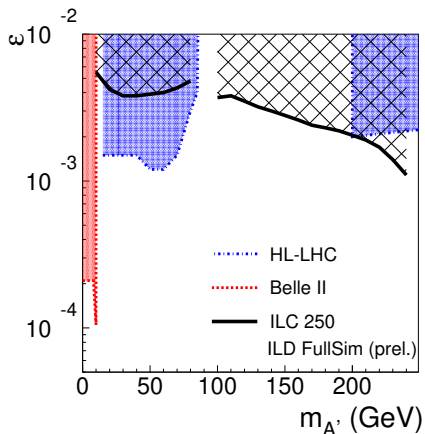
Dark Photons in a real detector

- Compared to the theory curve ...
- ... this is **the (current) result with full simulation**.
- At the highest mass, the correct limit is a factor two higher, a factor four at 100 GeV.
- This is **due to the correct estimate of the error**.
- Below M_Z , the difference is larger, and HL-LHC limits are expected to be stronger.
- Here, the reason is both **the correct error-estimate**, but also **the much larger background from non- $Z \rightarrow \mu\mu$ processes**.



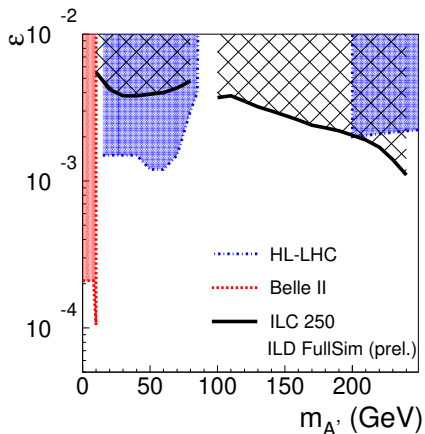
Dark Photons in a real detector

- Compared to the theory curve ...
- ... this is **the (current) result with full simulation**.
- At the **highest mass**, the **correct limit** is a factor two higher, a factor four at 100 GeV.
- This is **due to the correct estimate of the error**.
- Below M_Z , the difference is larger, and HL-LHC limits are expected to be stronger.
- Here, the reason is both **the correct error-estimate**, but also **the much larger background** from $\text{non-}Z \rightarrow \mu\mu$ processes.



Dark Photons in a real detector

- Compared to the theory curve ...
- ... this is **the (current) result with full simulation**.
- At the **highest mass**, the **correct limit** is a factor two higher, a factor four at 100 GeV.
- This is **due to the correct estimate of the error**.
- **Below M_Z** , the difference is larger, and **HL-LHC limits** are expected to be stronger.
- Here, the reason is both **the correct error-estimate**, but also the **much larger background** from $\text{non-}Z \rightarrow \mu\mu$ processes.



Conclusion and outlook

Uptake:

- Even for - or maybe in particular for - the most simple topology **full simulation is needed**.
- Because in these cases, **precision** is the most important aspect.
- Even though the correctly evaluated reach is significantly less than the theory estimate, e^+e^- colliders **will probe lower dark photon couplings** than HL-LHC, at least for masses above M_Z .

Outlook:

- Several non-trivial ameliorations are possible

Conclusion and outlook

Uptake:

- Even for - or maybe in particular for - the most simple topology **full simulation is needed**.
- Because in these cases, **precision** is the most important aspect.
- Even though the correctly evaluated reach is significantly less than the theory estimate, e^+e^- colliders **will probe lower dark photon couplings** than HL-LHC, at least for masses above M_Z .

Outlook:

- Several non-trivial ameliorations are possible

Conclusion and outlook

Uptake:

- Even for - or maybe in particular for - the most simple topology **full simulation is needed**.
- Because in these cases, **precision** is the most important aspect.
- Even though the correctly evaluated reach is significantly less than the theory estimate, e^+e^- colliders **will probe lower dark photon couplings** than HL-LHC, at least for masses above M_Z .

Outlook:

- Several non-trivial **ameliorations** are possible
 - **LR weighting** of the samples with different polarisations.
 - Include $A_D \rightarrow e^+e^-$: Need methods to compensate for brems-strahlung to get good enough mass-resolution.
 - No use of **the ISR photon** made. Can it be used? Background reduction at low M_A , or even better resolution?
 - Use event-by-event error better: **un-binned Maximum Likelihood**.
 - Spend some running-time **scanning E_{CMS}** .
 - ...

Thank You !

Backup

BACKUP SLIDES