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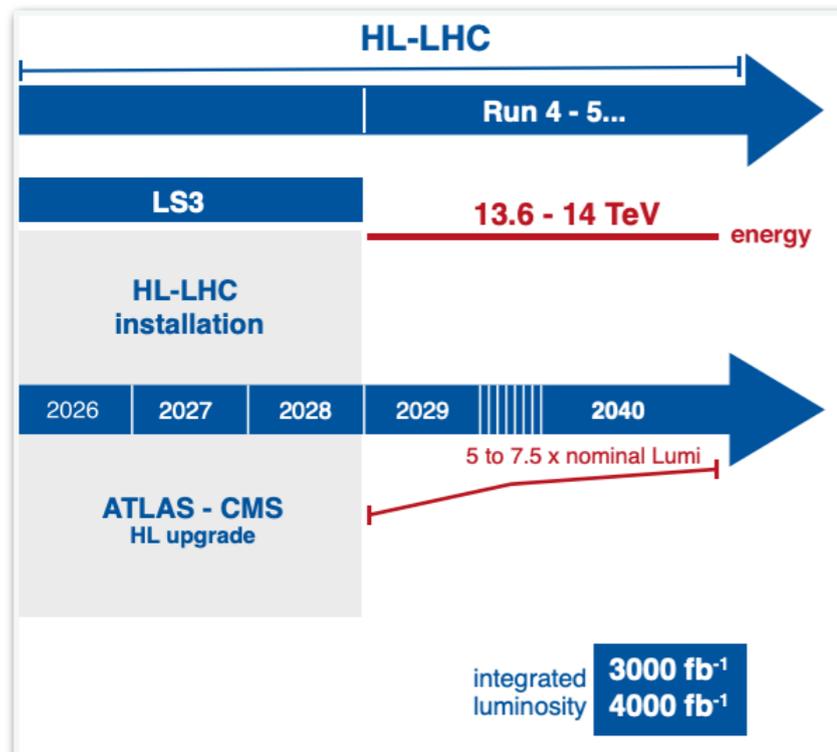
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# A Set of Physics Program at the ILC 250, 380, 1000

They can be a Killer Science in 2040?

( late 2030? )

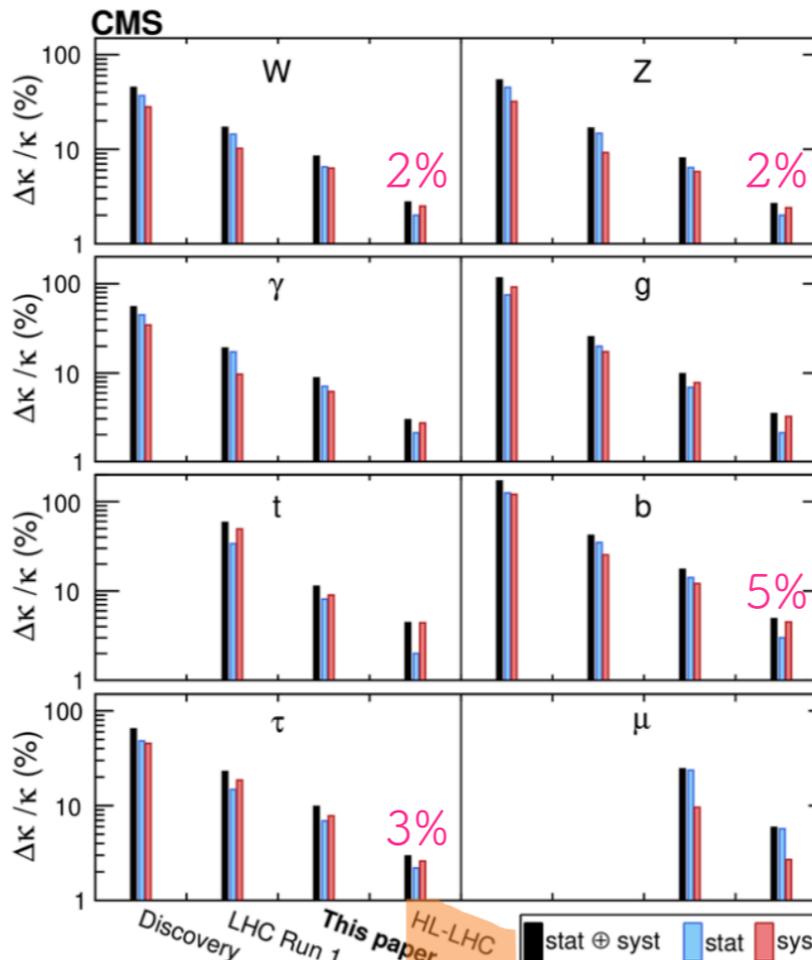
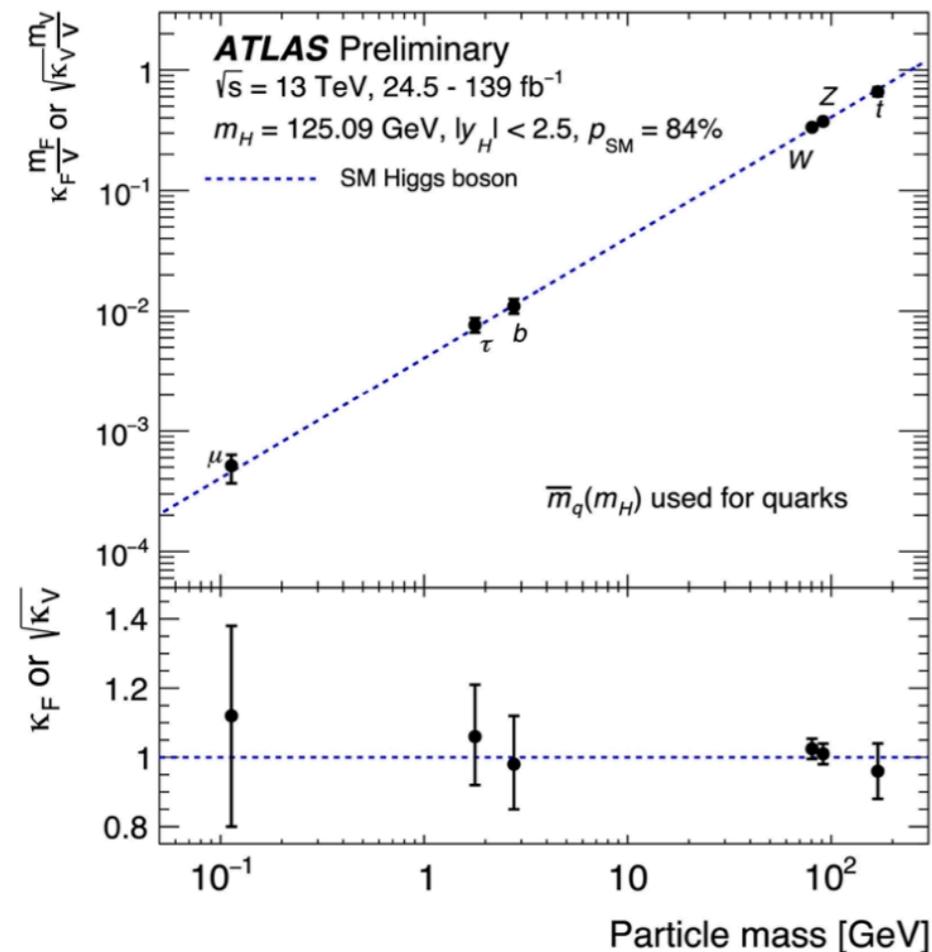


IDT view on the ILC project timeline  
-success oriented and assuming no major incident-



HL-LHC Run4 (Start 2029) for ~ 10+ years

## (HL-) LHC



$$\kappa = g_X / g_X^{\text{SM}} = 1 + \Delta\kappa$$

$$\Delta\kappa \sim \mathcal{O}(v^2/\Lambda^2)$$

e.g. A New Phys. at 1 TeV  
 Expected Deviation ~6%

## Precision in Run-2

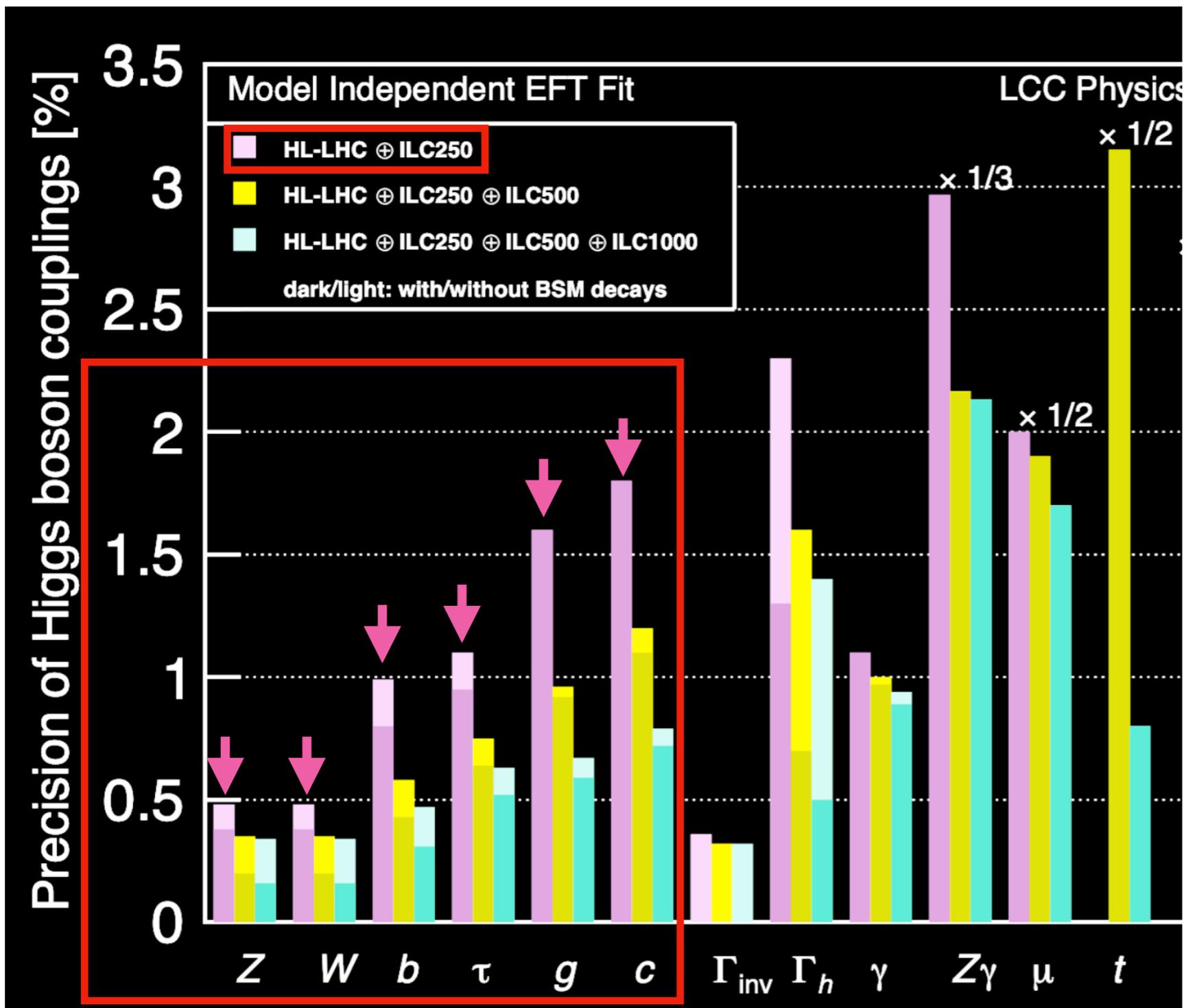
- W/Z: 5-10%
- 3rd Gen.: 10%
- 2nd Gen.: First Obs.



## HL-LHC

- 2-5% for most of the channels

- Higher accuracy w.r.t. HL-LHC may be necessary
- Higher accuracy probes higher energy scale



$$\kappa = g_X / g_X^{\text{SM}} = 1 + \Delta\kappa$$

$$\Delta\kappa \sim O(v^2/\Lambda^2)$$

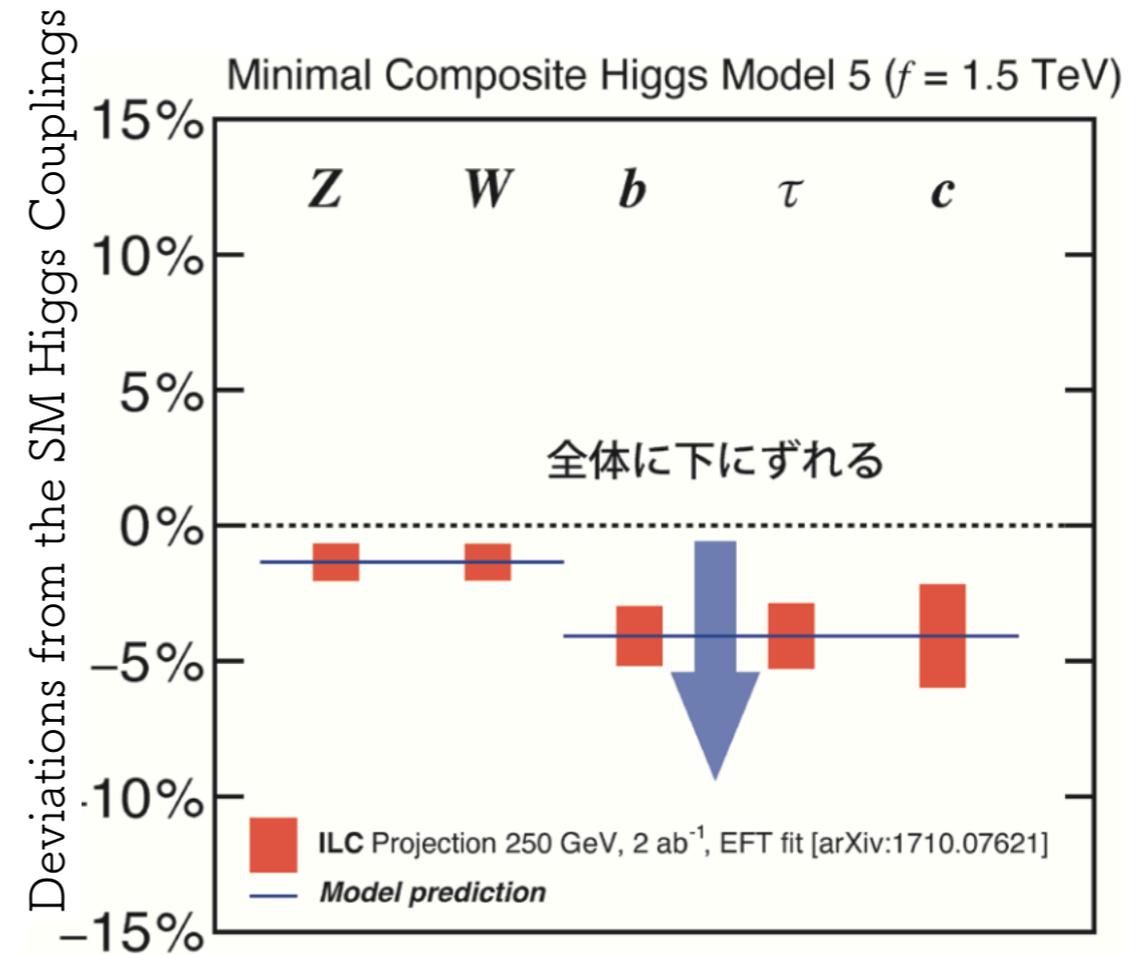
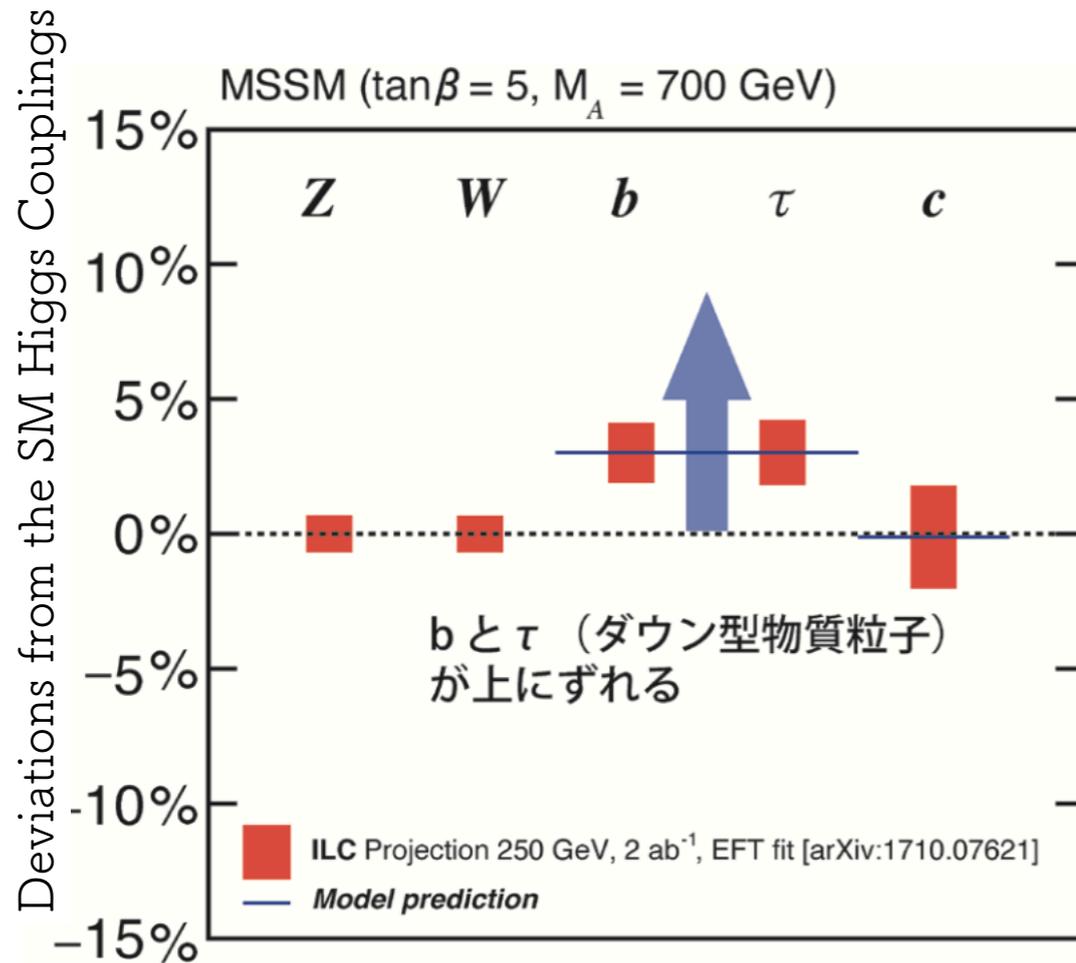
e.g. A New Phys. at 1TeV  
 Expected Deviation ~6%

### ILC250

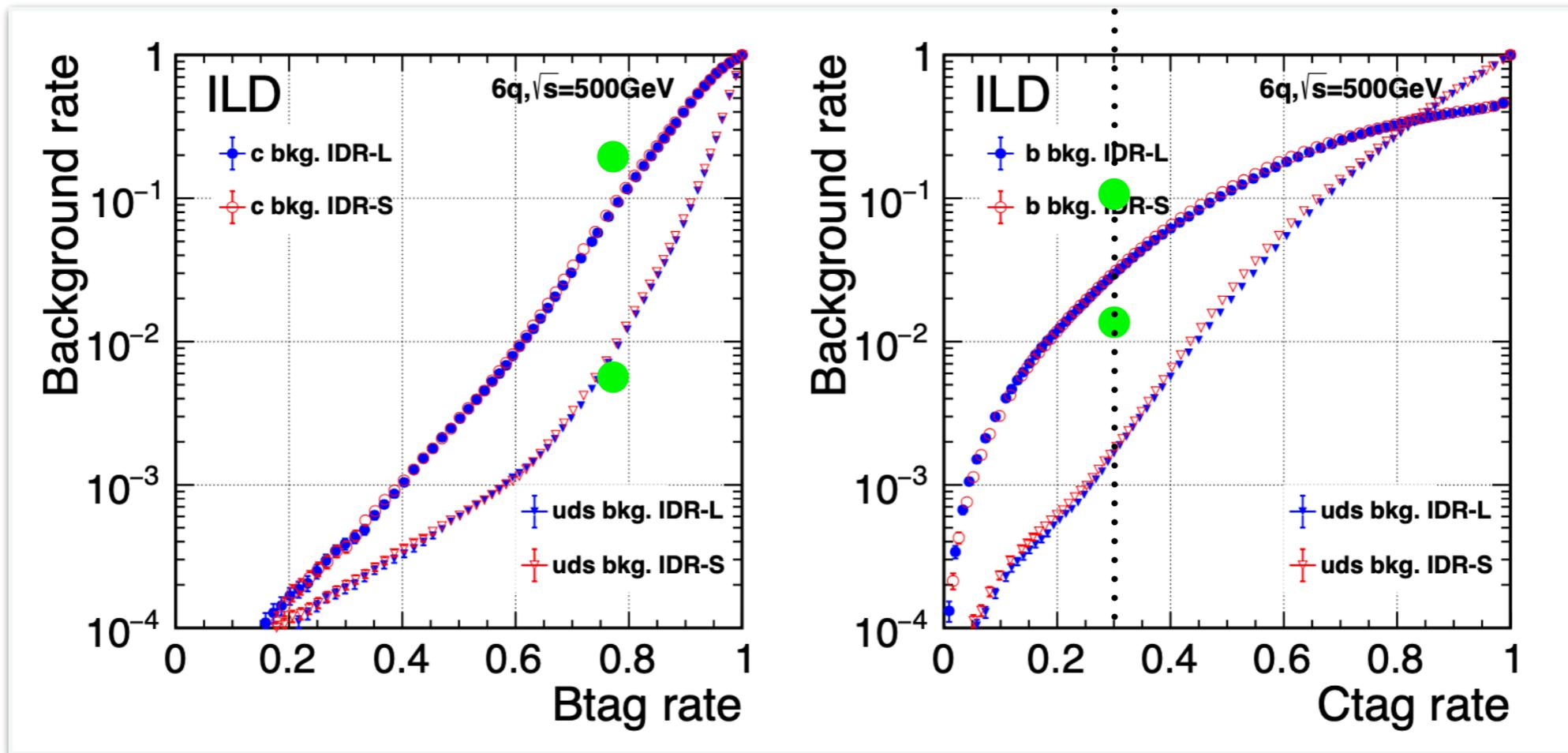
- 0.5-2% for Z, W, b, τ, g, c

The high precision measurement makes sense  
 to explore a New Phys. @ TeV Scale

# Two Example Scenarios of BSM



The precision of 5% (b) / 3% ( $\tau$ ) is not enough and better precision is highly demanded



b-tag 77%, light jet:  $1/170 (= 6E-3)$

c-jet:  $1/5$

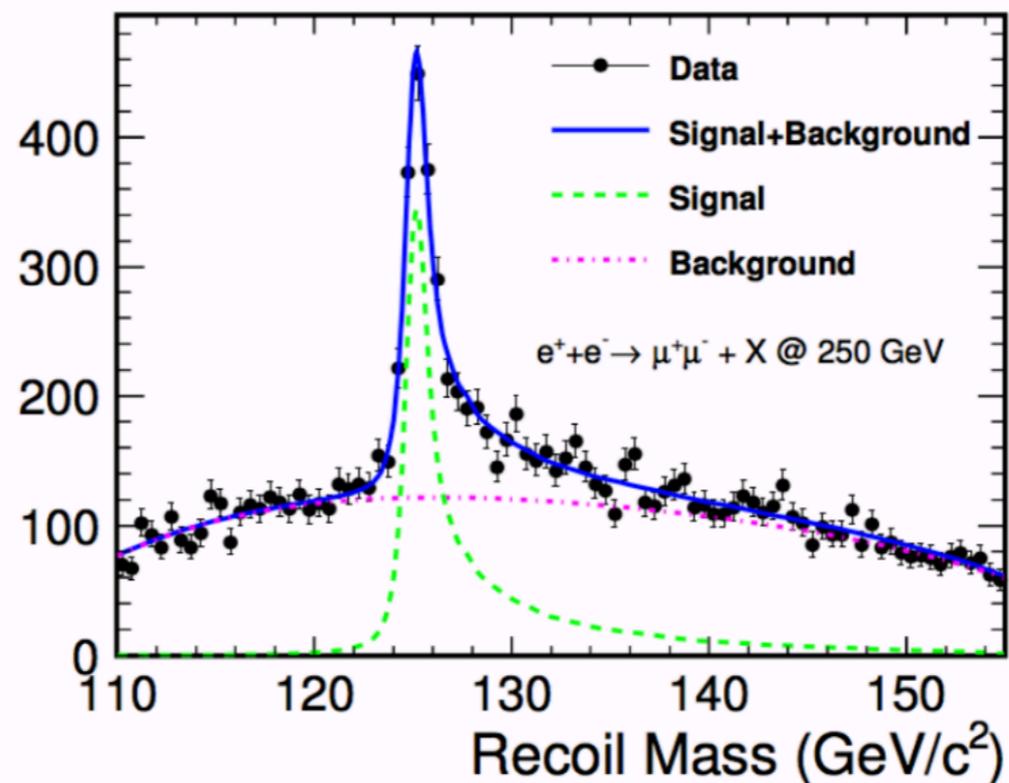
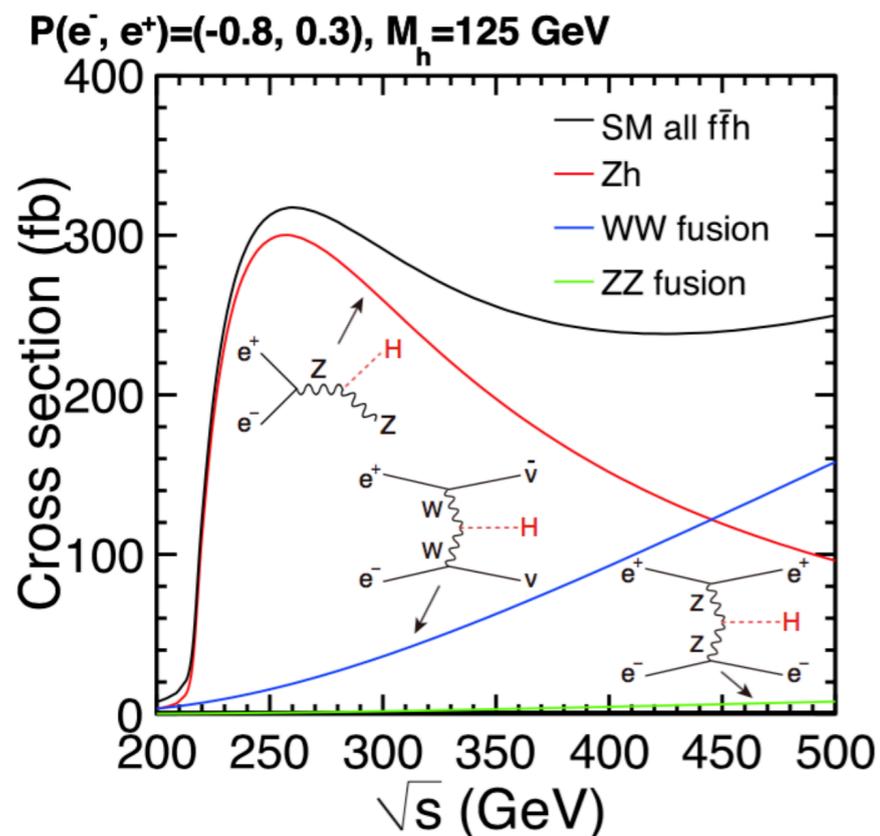
Large Room to improve b-tagging on ILC side?

c-tag 30%, light jet:  $1/70$

b-jet:  $1/9$

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**Abstract** The flavour-tagging algorithms developed by the ATLAS Collaboration and used to analyse its dataset of  $\sqrt{s} = 13\text{ TeV}$   $pp$  collisions from Run 2 of the Large Hadron Collider are presented. These new tagging algorithms are based on recurrent and deep neural networks, and their performance is evaluated in simulated collision events. These developments yield considerable improvements over previous jet-flavour identification strategies. At the 77%  $b$ -jet identification efficiency operating point, light-jet (charm-jet) rejection factors of 170 (5) are achieved in a sample of simulated Standard Model  $t\bar{t}$  events; similarly, at a  $c$ -jet identification efficiency of 30%, a light-jet ( $b$ -jet) rejection factor of 70 (9) is obtained.



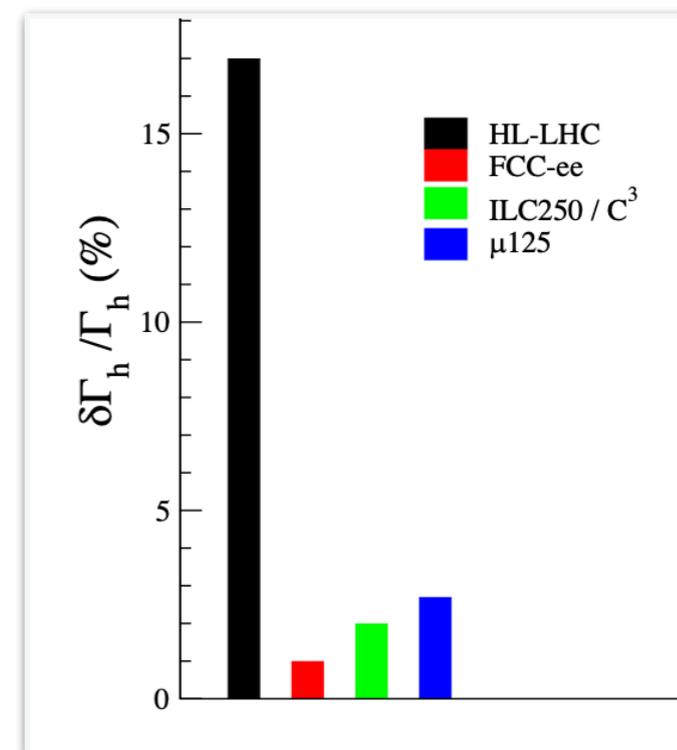
Higgs Production Zh:  $\sigma_{Zh} \propto (gh_{ZZ})^2$

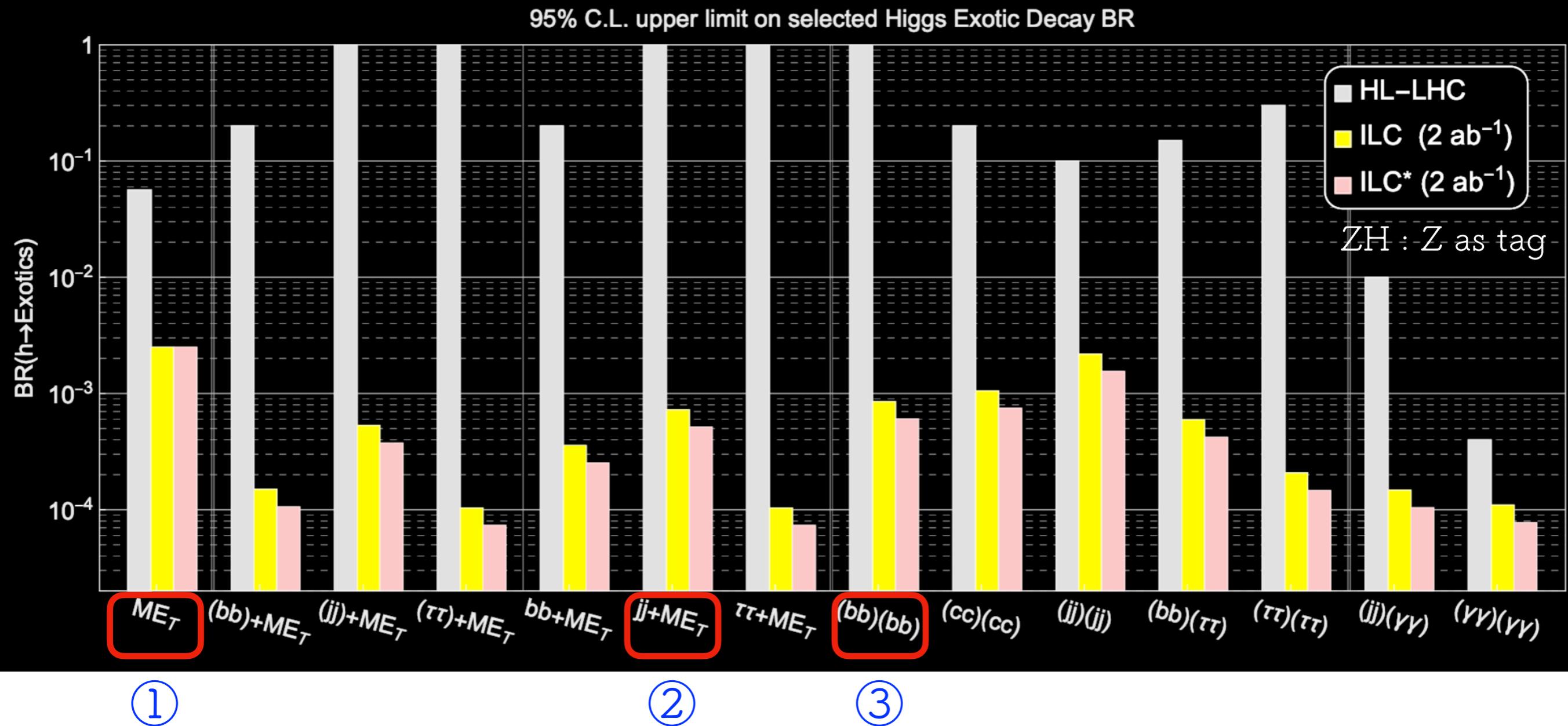
Partial Width  $\Gamma(h \rightarrow XX) \propto (gh_{XX})^2$

X-section measurement  
with Recoil Mass Method  
( @ Higgs Factory )

$$\begin{aligned} & \sigma_{ZH} \times \text{Br}(h \rightarrow XX) \\ &= \sigma_{ZH} \times \Gamma(h \rightarrow XX) / \Gamma_H \\ &\propto (gh_{ZZ})^2 (gh_{XX})^2 / \Gamma_H \end{aligned}$$

$X = Z, W \rightarrow$  extract  $\Gamma_h$

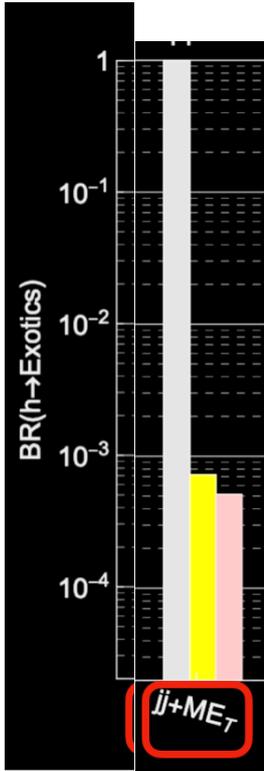
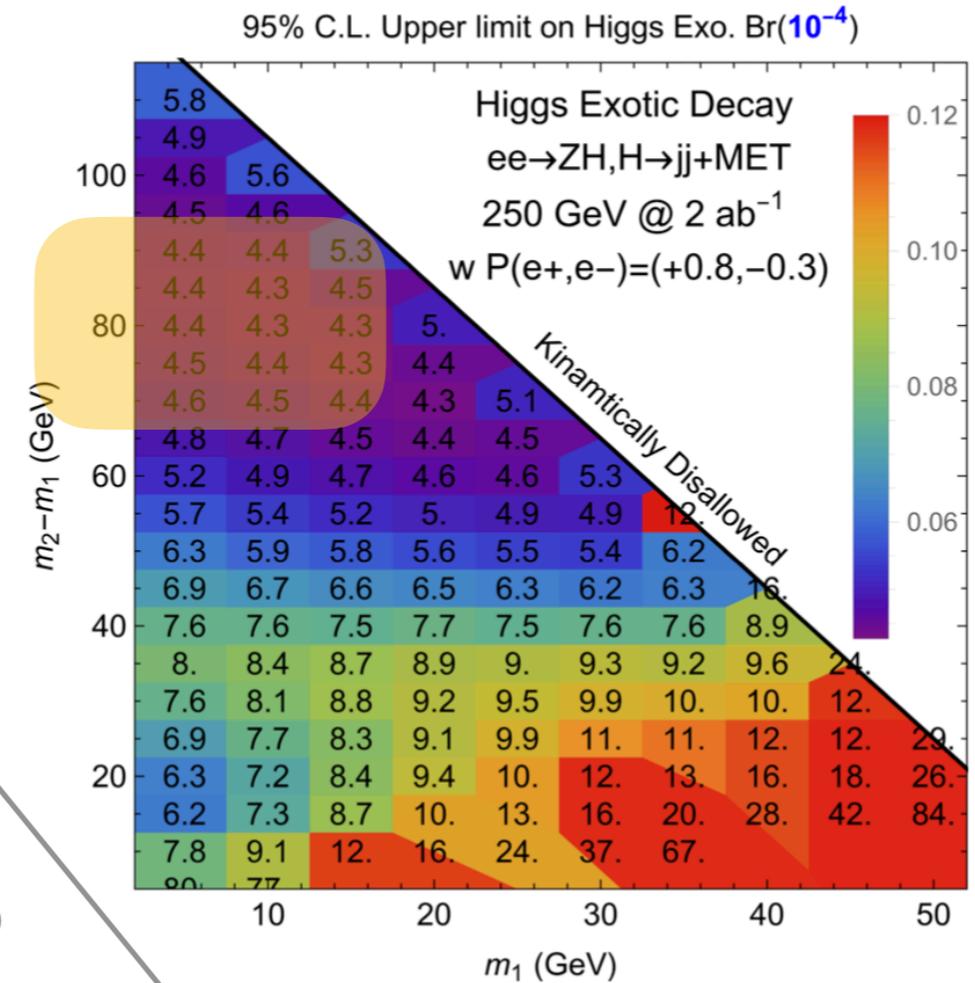




- ① Fully Invisible Higgs Decay:  
 $H \rightarrow 4\nu$  (0.1%), 95% Confidence Upper Limit on Br = 0.16%

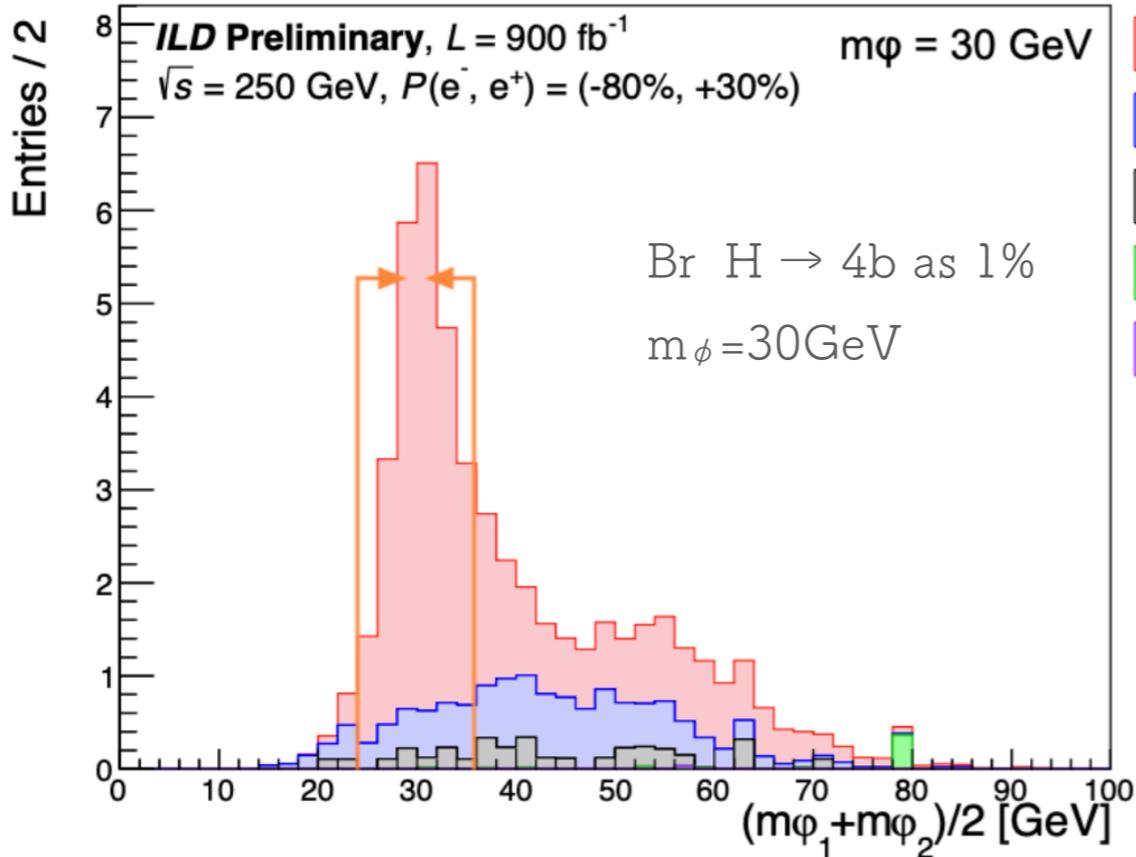
# 9 Exotic Higgs Decay - Partially Visible Case

②  $H \rightarrow X1 + (X2) \rightarrow X1 + (X1 + 2\text{-jets})$   
 BG.:  $H \rightarrow ZZ^* \rightarrow qq \nu \nu$   
 95% C.L. Upper Limit Br. =  $4 E-4 \sim 8E-4$

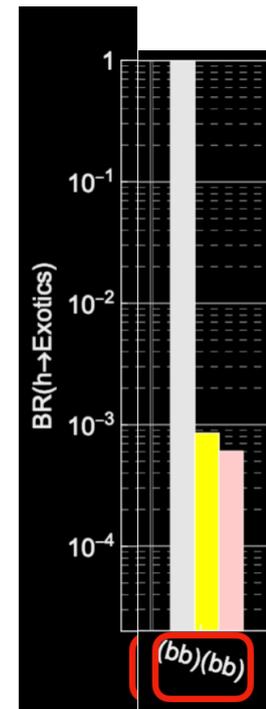


③  $H \rightarrow \phi\phi \rightarrow (b\bar{b})(b\bar{b})$

ZH



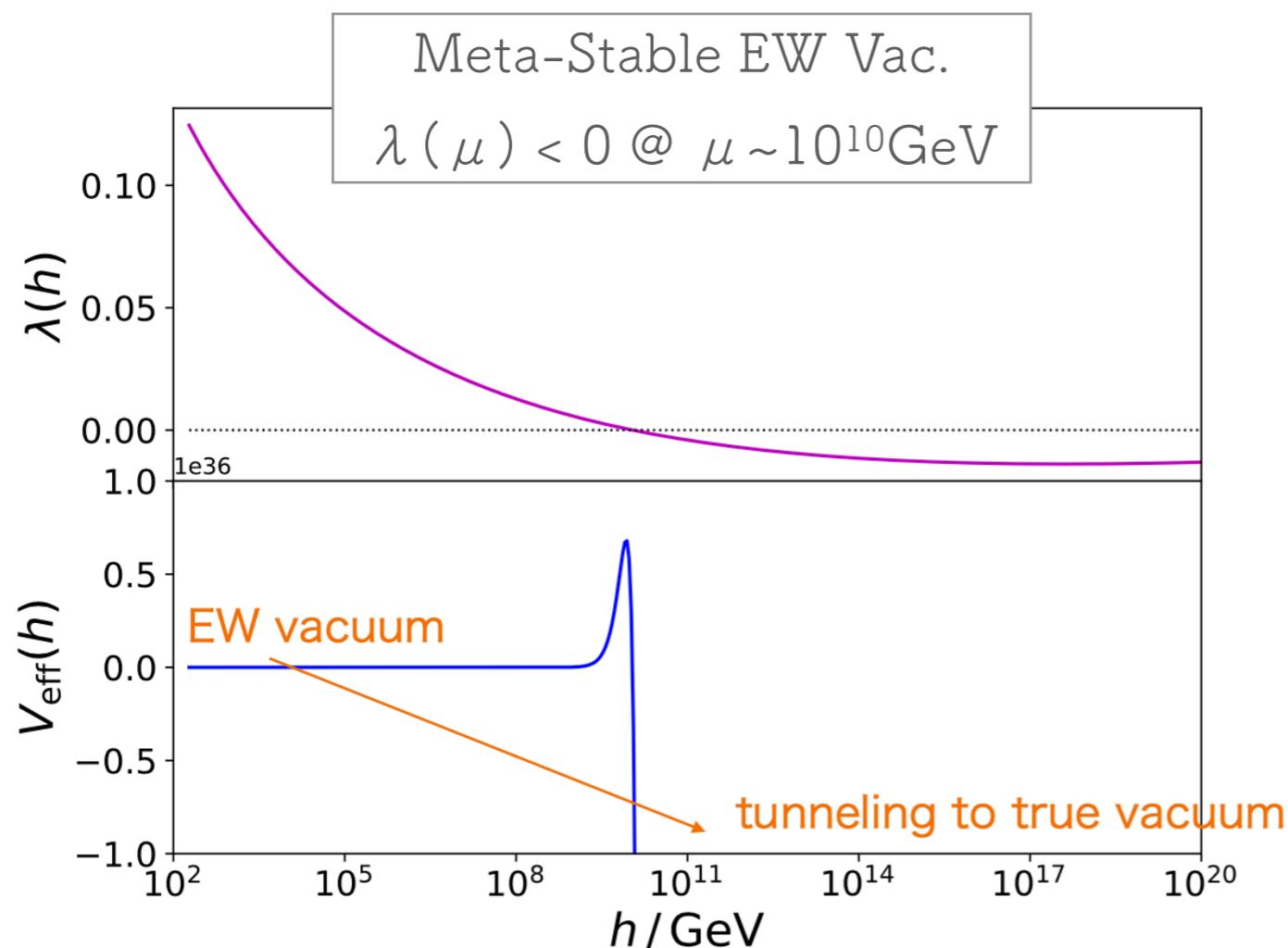
95% C.L Upper Limit is expected as 0.1%  
 ( $1E-3$ )



one-loop RG equation of  $\lambda$  and  $y_t$  useful to understand  $\lambda(\mu)$

$$16\pi^2 \frac{d\lambda}{d \ln \mu} \Big|_{\text{one-loop}} = 12\lambda \left( 2\lambda + y_t^2 - \frac{g_Y^2 + g_2^2}{4} - \frac{g_2^2}{2} \right) - 6y_t^4 + 6 \left( \frac{g_Y^2 + g_2^2}{4} \right)^2 + 12 \left( \frac{g_2^2}{4} \right)^2$$

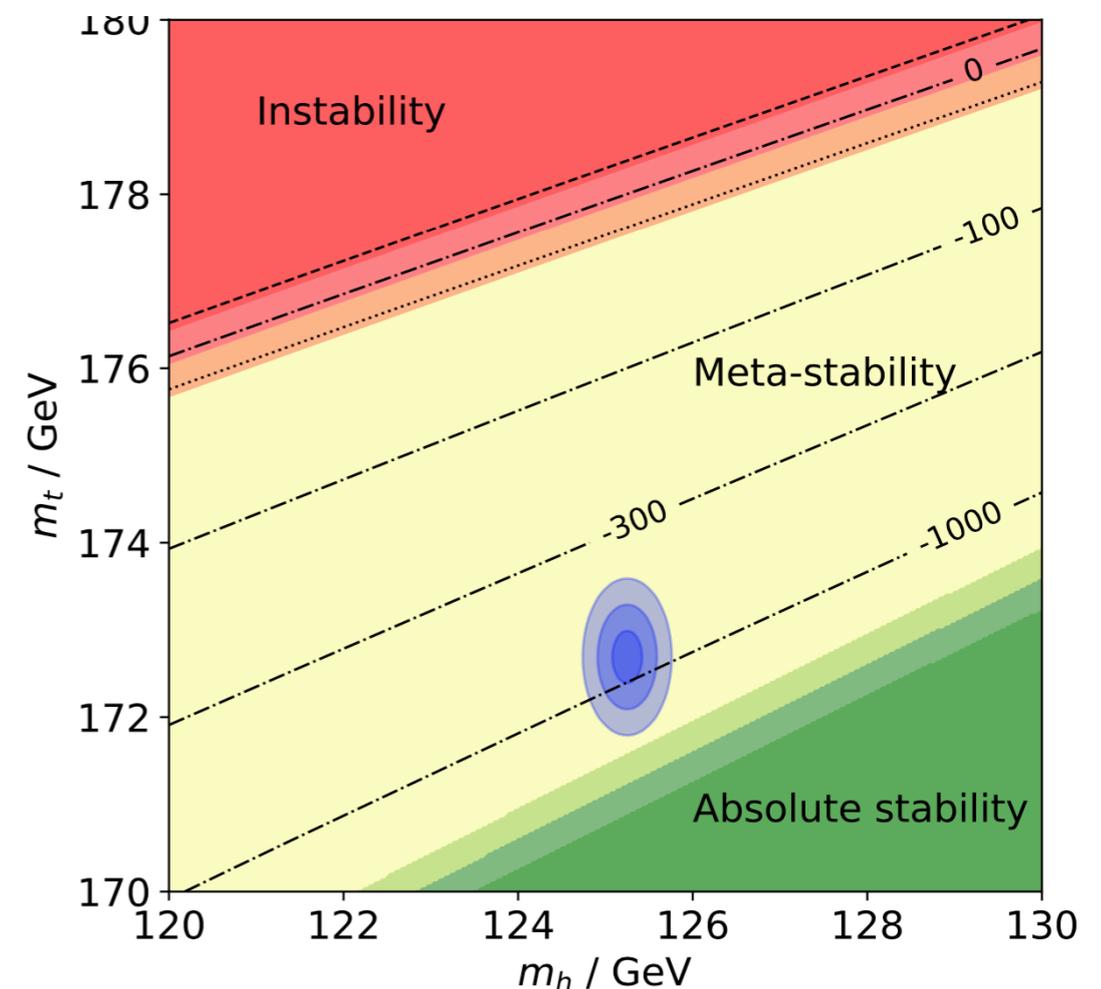
Larger Higgs mass  $\rightarrow$  being stable  
 Larger Top mass  $\rightarrow$  being unstable

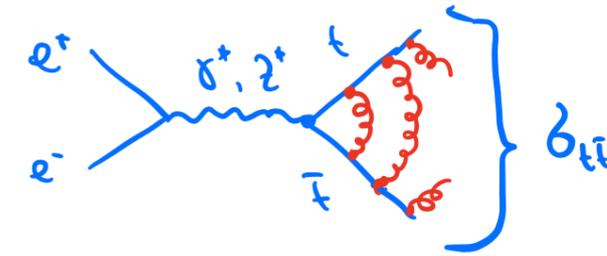
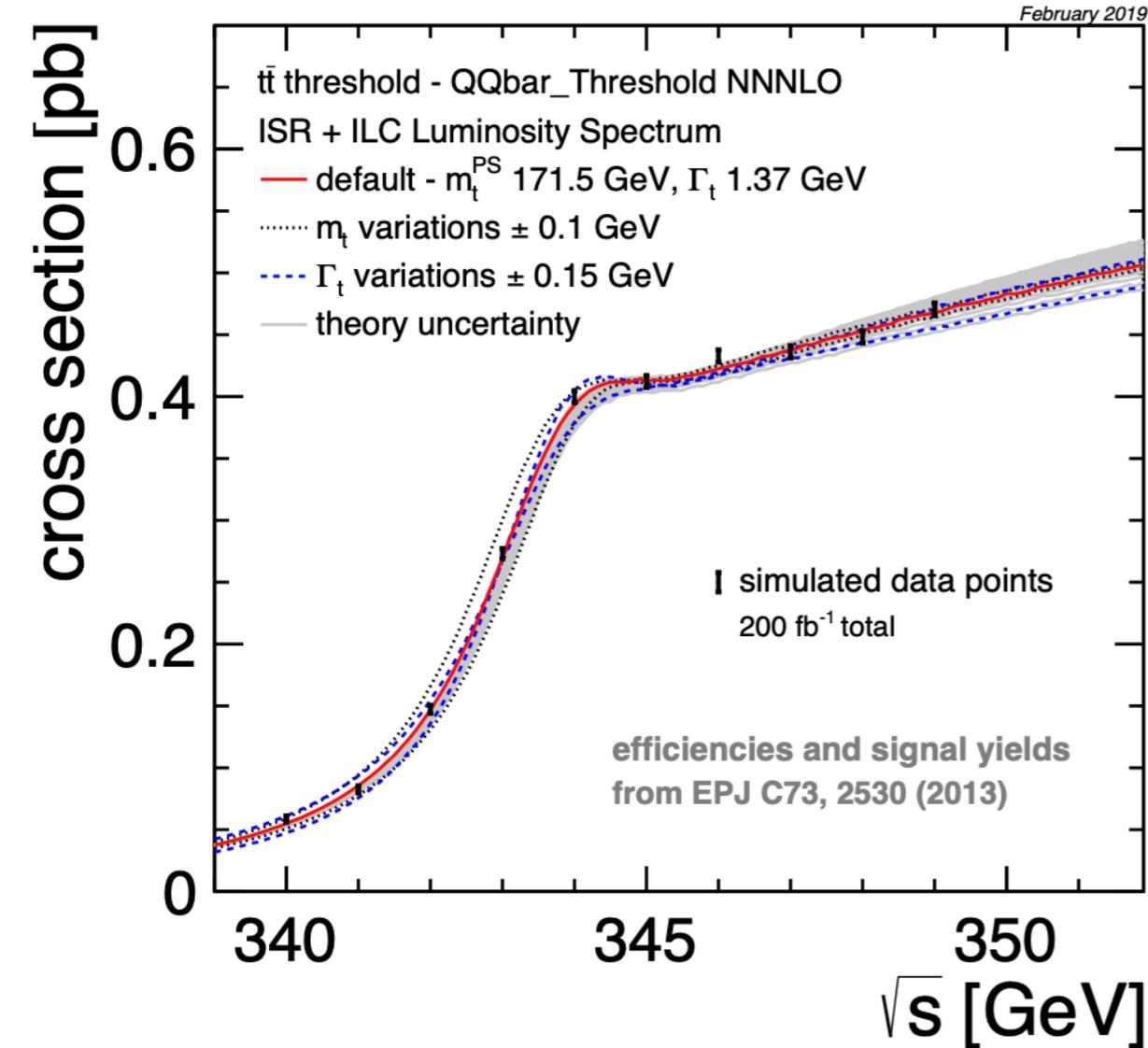


EW vacuum decay rate

$$\log_{10} (\gamma \times \text{Gyr Gpc}^3) = -785^{+45}_{-49} + 155^{+181}_{-222} + 181^{+181}_{-276}$$

$\Delta m_h$      $\Delta m_t$      $\Delta \alpha_s$





## Threshold Scan

→ short-distance mass (theoretically well defined)

$\Delta m \sim 20 \text{ MeV}$  (statistical)

$\Delta m \sim 50 \text{ MeV}$  (sys.: higher-order loop,  $\alpha_s$ )

EPJ C73, 2530 (2013)

enough precision to resolve EW Vac. Stability

(c.f. LHC  $\sigma(t\bar{t})$  n-differential, NLO  $173.2 \pm 1.6 \text{ GeV}$ )

a few other methods ( $t\bar{t} \gamma$ , boosted top, angular scale, ...) are proposed at ILC

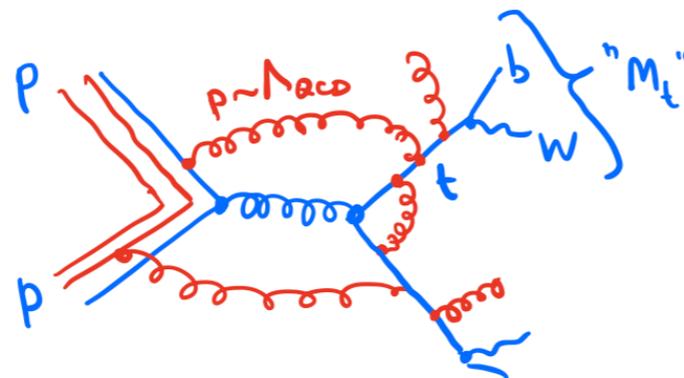
## Reconstructed Top invariant mass

Need to model soft-gluon radiation/exchanges  
 (theoretically not well defined)

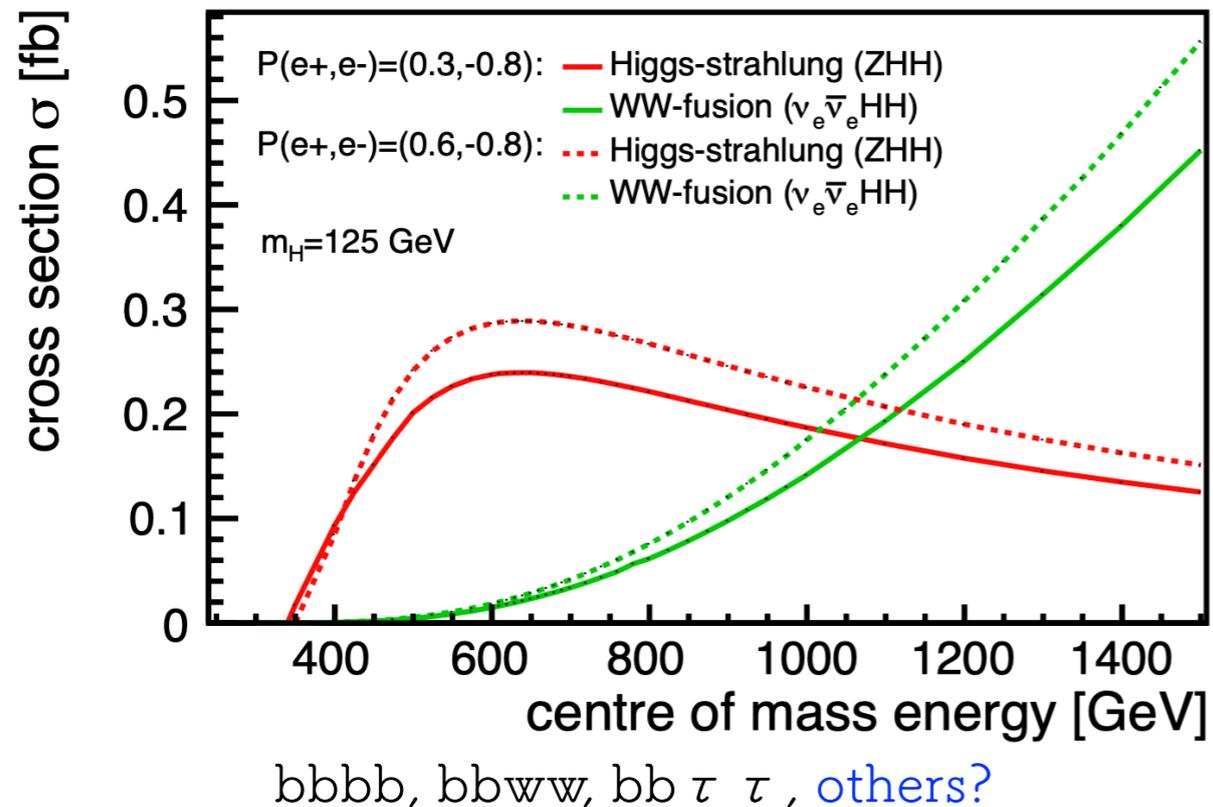
LHC MC template method:  $\Delta m \sim 600 \text{ MeV}$

(HL-LHC:  $\Delta m \sim$  a few hundred MeV??)

MC mass  $\Leftrightarrow$  Theoretical mass mapping studies on-going



## X-section: HH production



Direct measurement of Higgs cubic self-coupling

[ 500 GeV ] ZHH ( $\Delta \lambda \sim 25\%$ )

[ 1 TeV ] WW-fusion ( $\Delta \lambda \sim 10\%$ )

SM predicts  $\lambda_3$  (w/  $m_H, v_{ev}$ )

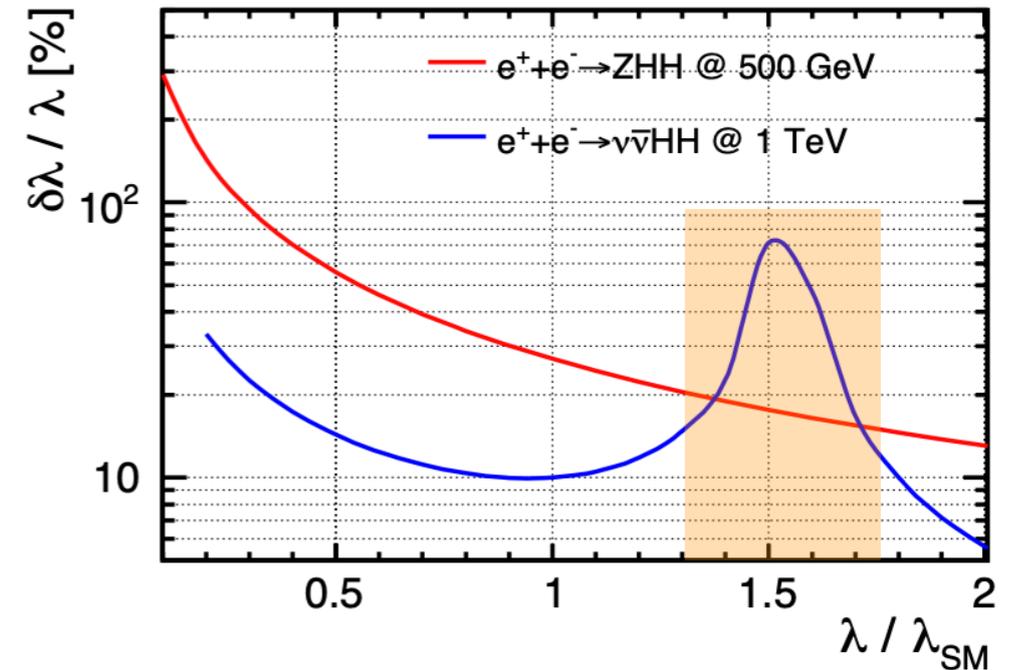
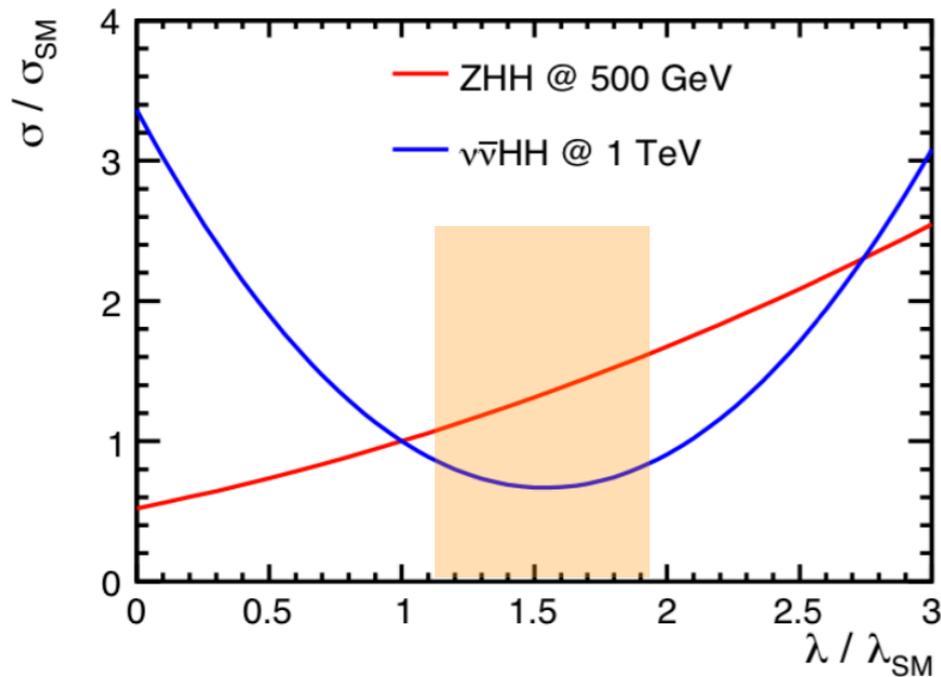
→ Good Test of **New Phys. in Higgs Sector**

[ Example: Deviation of  $\lambda$  ]

• 2HDM (Yukawa Type-I)  $\lambda: -0.5 \sim 1.5$

• EW Baryogenesis Model  $\lambda: 1.5 \sim 2.5$

By the way, there are two di-Higgs processes ZH & WW-fusion and they interfere ⇒



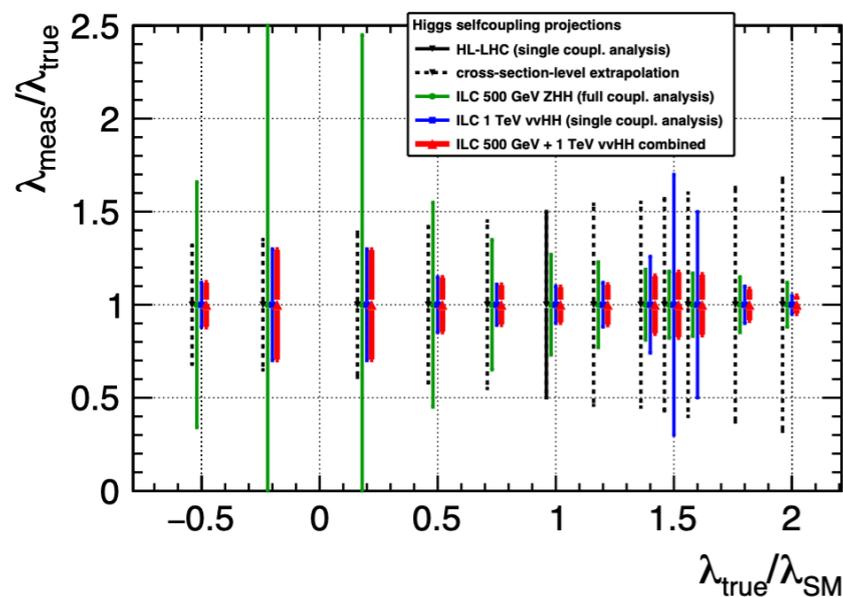
The 2 X-sections behave oppositely

- **ZHH**      **Constructive** interference
- **WW-fusion**      **Destructive** interference

Combination of two measurements

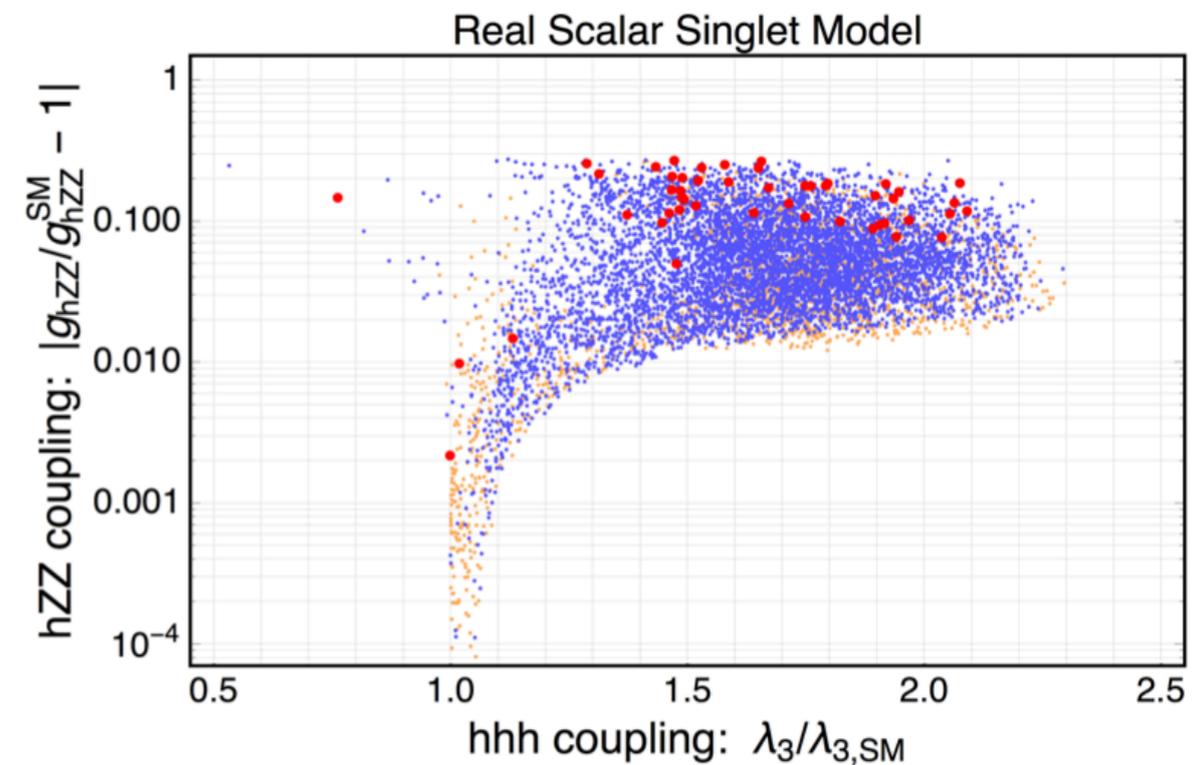
⇒ **a 20% precision in  $\lambda_3$** , even if it deviates from the SM by 50-100%

important to run at the both 500 and 1000 GeV



LHC ⇒ gluon fusion dominant  
 ⇒ Destructive Interference  
 ⇒ suffered for enhanced  $\lambda_3$

- First-order EWPT requires BSM fields, with significant coupling to Higgs and mass  $\sim$  weak scale
- A singlet scalar is the simplest benchmark:
 
$$V = m_S^2 S^2 + \frac{\kappa}{2} S^2 |H|^2$$
- Invisible to the LHC as long as  $m_S > m_h/2$
- Models with first-order EWPT typically predict  **$\sim 20-120\%$  enhancement** in  $h^3$

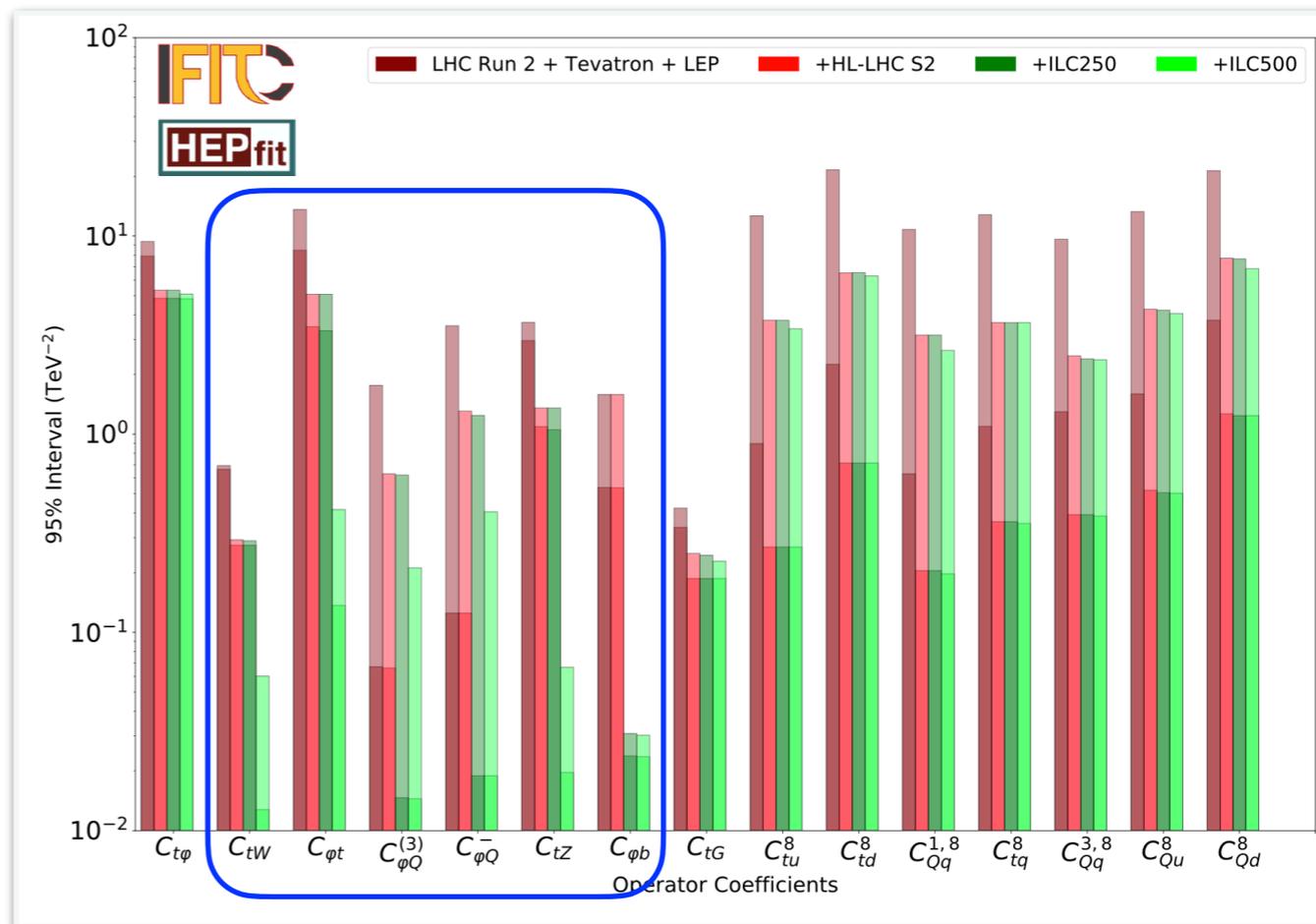


[Huang, Long, L-T. Wang]

## Top EFT Operator

- 500 GeV (or above) Top EW Couplings
- Deviations: Parametrized by EFT
- **1/10 - 1/100 better than HL-LHC**

$pp \rightarrow ttZ/tt\gamma$  or  $pp \rightarrow tZq/t\gamma q$   
 Precision @ HL-LHC: 5-10%



## Top FCNC Operator

Forbidden at SM tree level

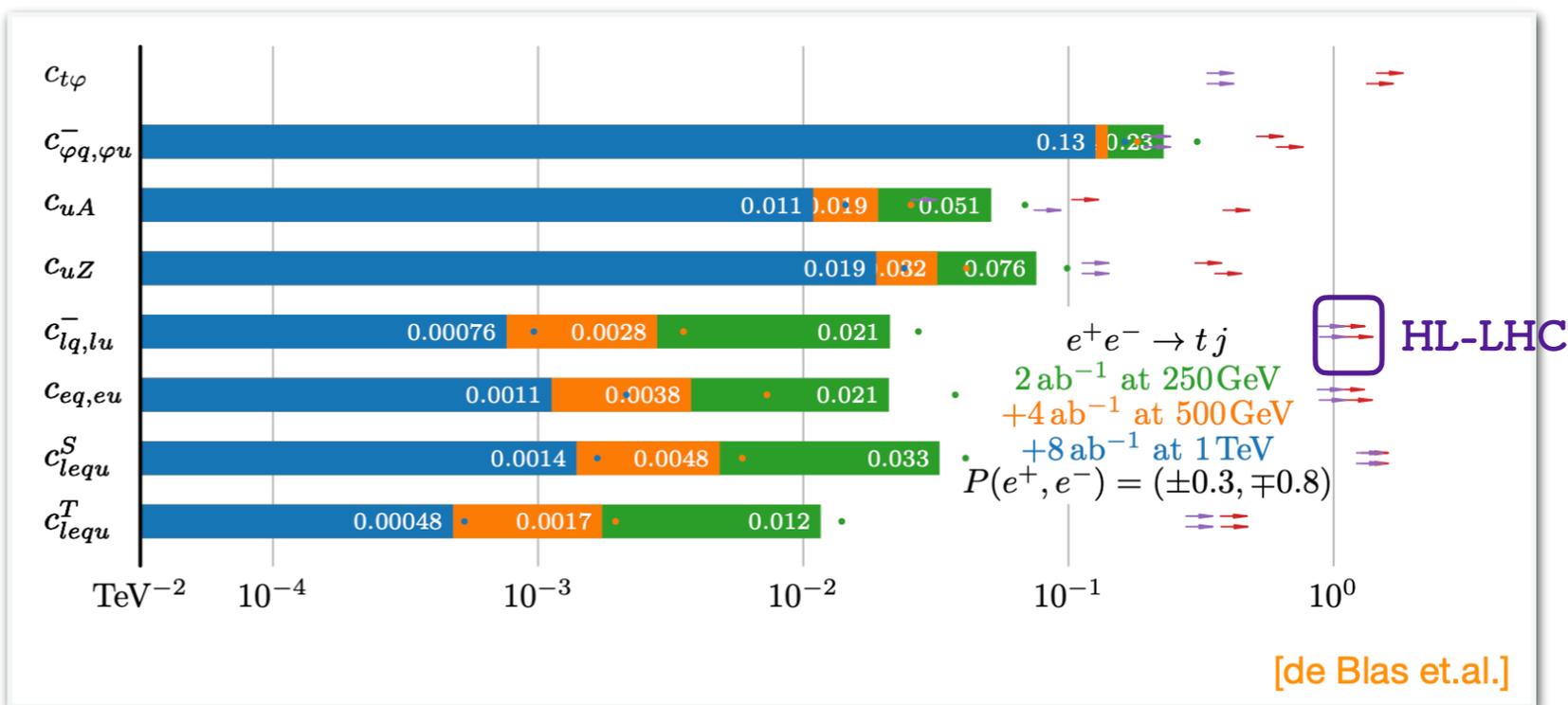
Suppressed by GIM mechanism

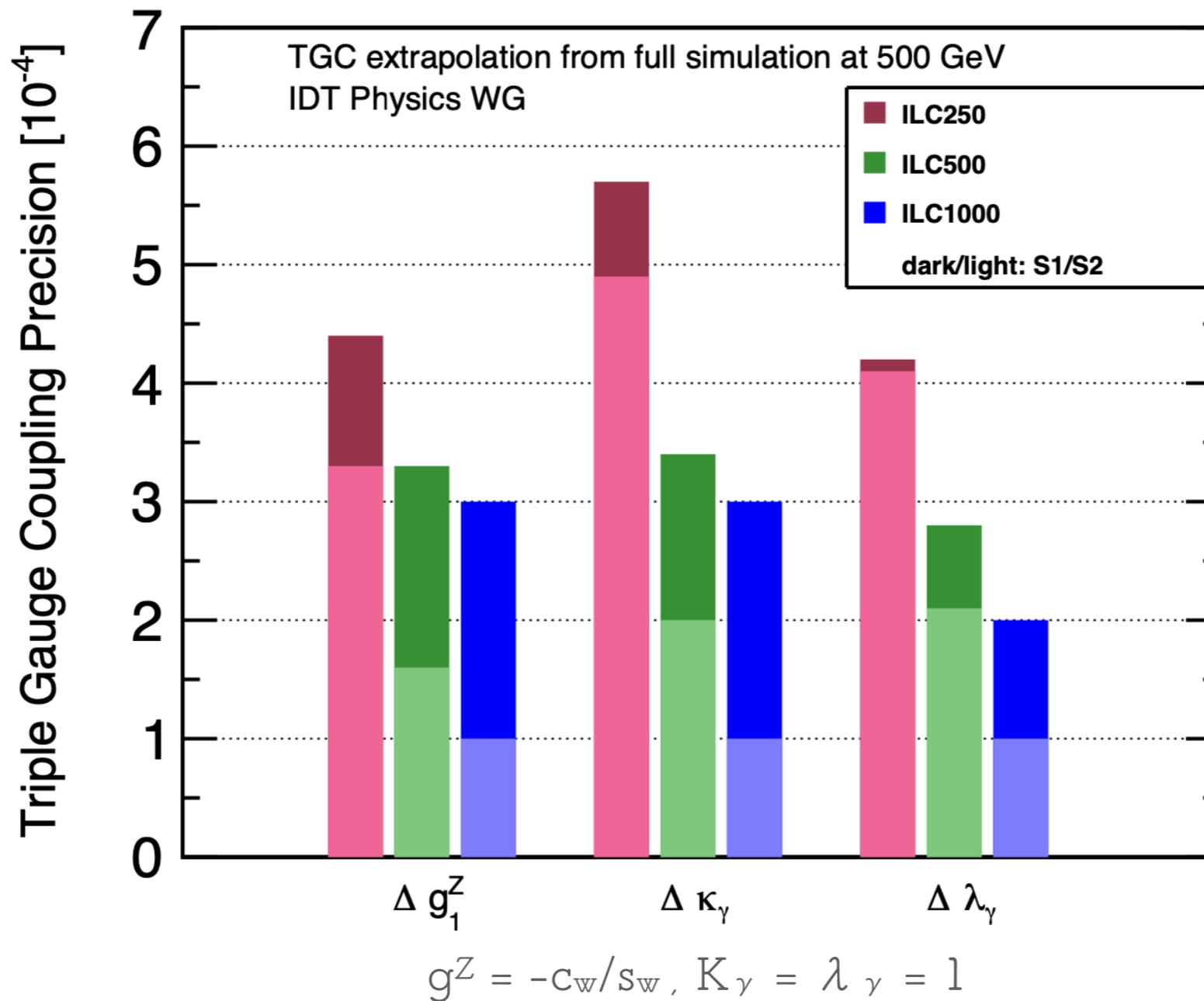
HL-LHC:  $1E-3 \sim 1E-5$

Higher Energy  $e^+e^-$  is powerful

e.g. 4-f Operator  $ee \rightarrow tq$

improvement of 3 orders

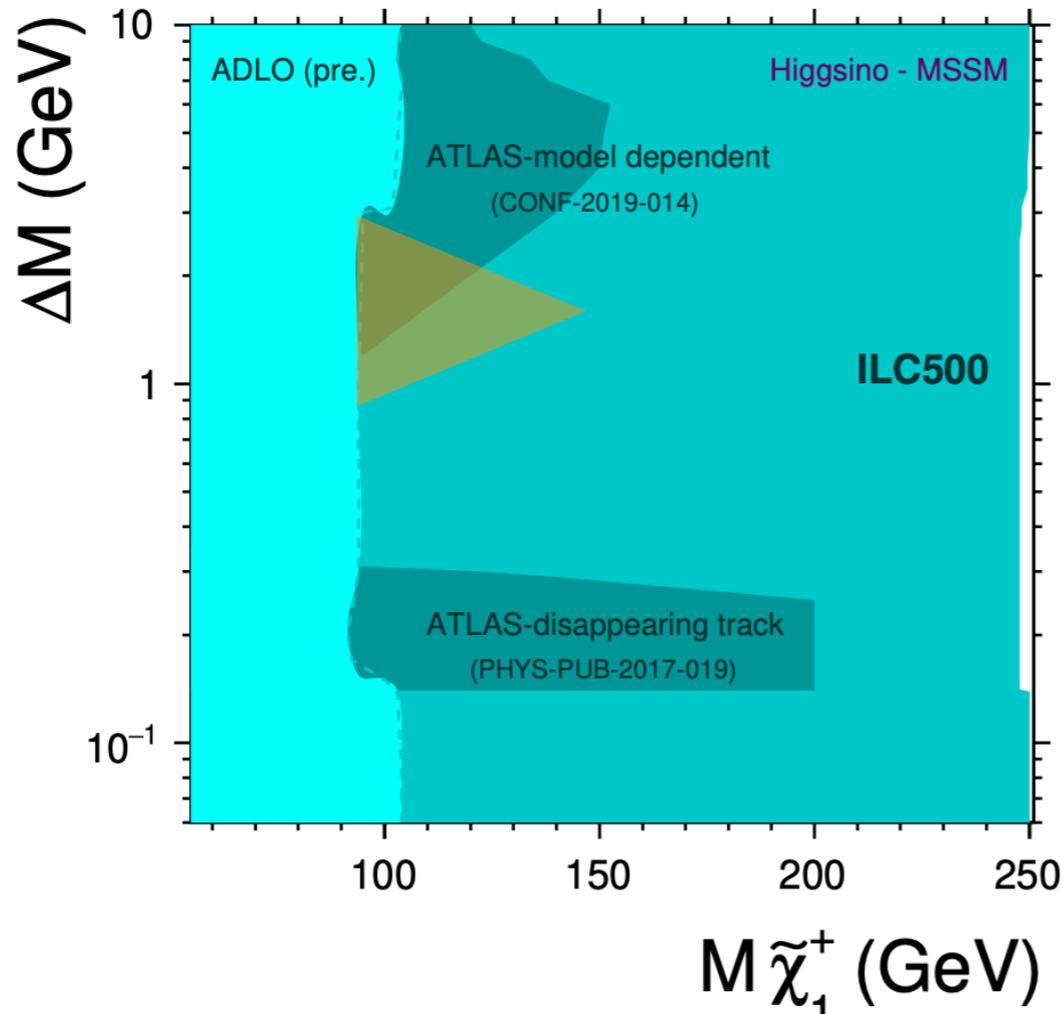




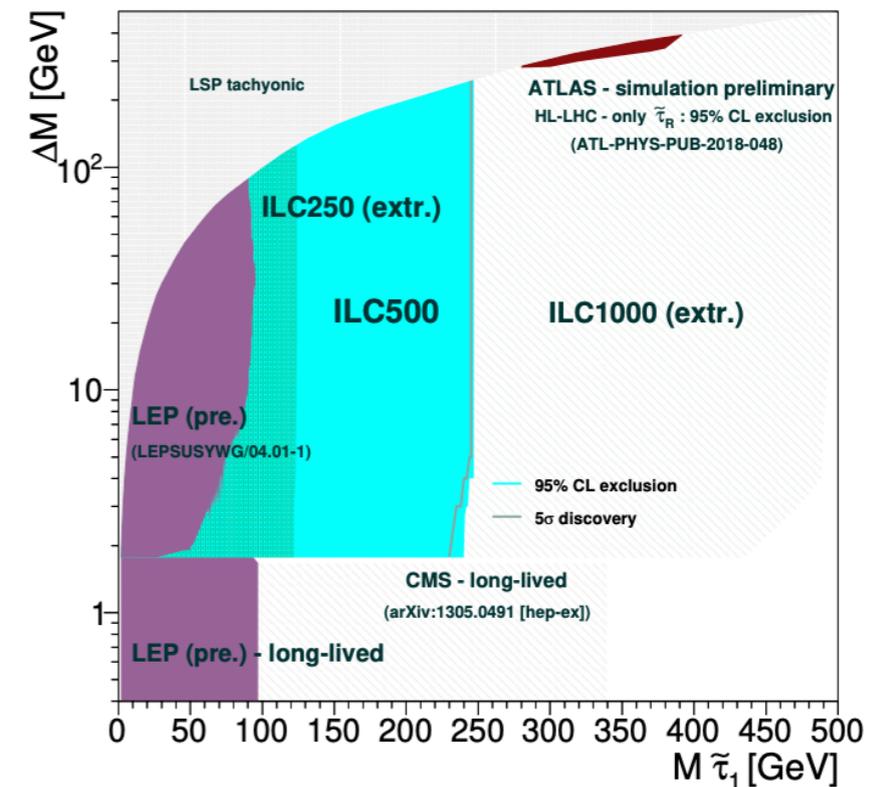
The precision of  $O(1E-3)$  for HL-LHC (p.175)

The improvement of "one order" is expected at higher  $\sqrt{s}$

$\Delta m$  (Chargino - LSP) : small mass gap



$\Delta m$  (stau, LSP) v.s.  $m_{\text{stau}}$



Low  $p_T$  lepton tracking: difficult at LHC  
 Very big efforts, but relatively small gain...

Great Potential to cover whole the interesting region by Higher Energy  $e^+e^-$  Machine

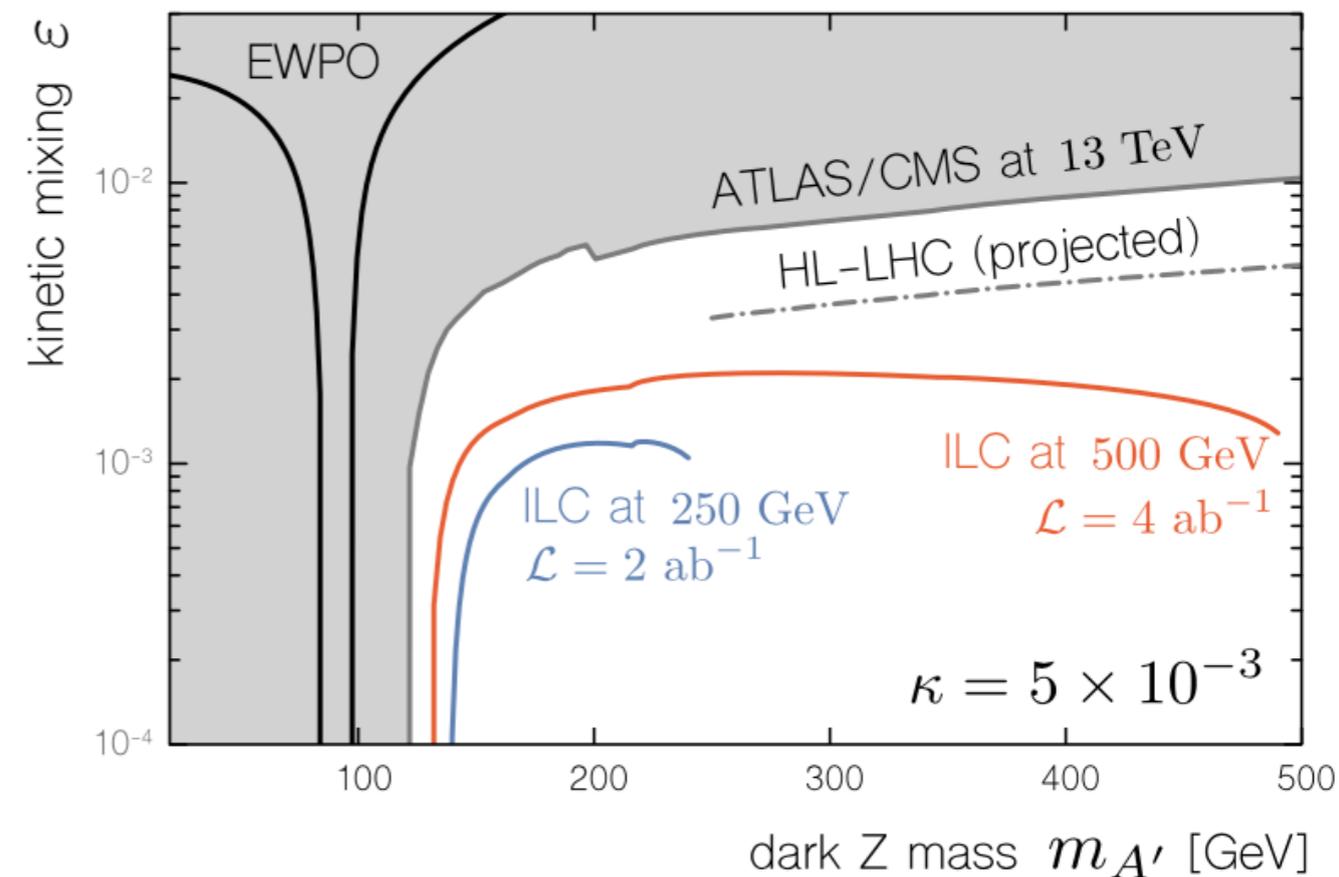
Some BSM predicts an alternative U(1) symmetry  
 If no SM field is charged, it could be "Dark Sector"

Maxim's Slide in LCWS 2022

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} + \frac{\varepsilon}{2c_W}B_{\mu\nu}A'^{\mu\nu} - \frac{1}{4}A'_{\mu\nu}A'^{\mu\nu}$$

Dominant coupling to SM is a loop effect

As long as Dark Z is within the kinematic reach of an e+e- collider, the region inaccessible by HL-LHC can be probed



# Summary

A Set of Physics Program at the ILC 250, 380, 1000

They can be a Killer Science in 2040?

- **Higgs**

- To explore a New Physics at 1 TeV, the measurement of **Higgs Couplings (0.5-2%)** is a strong approach and having advantage w.r.t. HL-LHC in 2040
- Exotic Higgs Decays: LC has great advantage (3 orders)
- Higgs-Self Coupling:  $\lambda_3$  (  $\Delta \lambda \sim 10\%$  ) LC has great advantage

- **m<sub>Top</sub>**  $\rightarrow y_{\text{top}}$ : **Energy Scan** at LC is a unique approach to determine the **pole-mass (20 MeV stat., 50 MeV sys. )**  $\Rightarrow$  conclude EW Vac. Stability

- **EW Precision Test** (1) Top (2) TGC : **1 - 3 orders** of advantage

- **New Particle Search**, e.g. EW SUSY @ mass degenerated region: Great advantage w.r.t. HL-LHC

**YES**

