# [157] Possible effects of the composite dark matter at the linear collider

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# Abstract

We consider the model of the composite dark matter based on the QCD-like SU(N) hidden color confining gauge theory.

Our meson-like dark matters (call it dark mesons) are a bounded state of dark quark ( $\psi$ ) and anti-dark quark ( $\overline{\psi}$ ) pairs. This dark matter sector connects to the Standard model via a real singlet scalar particle. The SU(N) hidden color gauge sector in the dynamical chiral symmetry breaking generates Nambu-Goldstone bosons ( $\tilde{\pi}$  dark meson) and massive composite scalar bosons ( $\tilde{\sigma}$  dark meson) simultaneously. A real singlet scalar particle gives the current mass for the dark quark. The current mass for dark quark breaks explicitly chiral symmetry. Nambu-Goldstone bosons are massive, meaning they are dark matter candidates. In the dark mesons of SU(N) hidden color, the chiral partner of the  $\tilde{\pi}$  meson is the  $\tilde{\sigma}$  meson (iso scalar-scalar).  $\tilde{\sigma}$  is also a candidate as dark matter. We will investigate the invisible process for the Higgs boson's decay at the future linear collider.

## The Model

The hidden sector Lagrangian:

$$\mathcal{L}_{H} = -\frac{1}{2} \mathrm{Tr} F^{2} + \mathrm{Tr} \overline{\psi_{i}} (i \gamma^{\mu} \partial_{\mu} + g_{H} \gamma^{\mu} G_{\mu} + g' Q \gamma^{\mu} B_{\mu} - y_{i} \phi) \psi_{i}$$
  
$$B_{\mu} = \cos \theta_{W} A_{\mu} - \sin \theta_{W} Z_{\mu}, g' = e / \cos \theta_{W}, i = 1,2,3$$

If the diagonal component of the Yukawa coupling,

$$y_1 = y_2 \neq y_3.$$

⇒Flavour symmetry breaking:  $SU(3) \rightarrow SU(2)_V \times U(1)_{\tilde{B}}$ In our present analysis, we omitted  $\psi_3$ . We assume that the mesons of which  $\psi_3$  is a constituent are heavy mass. The Lagrangian of the Higgs portal is as follows.

$$\mathcal{L}_{portal} = -\frac{1}{2}\mu_{H}^{2}(H^{+}H) - \frac{1}{2}\mu_{s}^{2}\phi^{2} - \lambda_{H}(H^{+}H)^{2} - \frac{1}{4}\lambda_{s}\phi^{4} - \frac{1}{2}\lambda\phi^{2}(H^{+}H)$$

Dark quarks are connected to the SM through the Yukawa couplings of Higgs and a singlet scalar particle  $\phi$ .

#### Park hadrons can be observed by the collider experiments.

We have assumed that the dark matter section is simply a scaled-up version of QCD.

A part of the effective Lagrangian depending on  $\phi$ :

 $\mathcal{L}_{int} = g_{\pi} \phi \left( \widetilde{\pi^{+}} \widetilde{\pi^{-}} + \widetilde{\pi^{-}} \widetilde{\pi^{+}} + \widetilde{\pi^{0}} \widetilde{\pi^{0}} \right) + g_{\sigma} \phi \widetilde{\sigma} \widetilde{\sigma}$ 

In the dark mesons of SU(N) hidden color, the chiral partner of the  $\tilde{\pi}$  meson is the  $\tilde{\sigma}$  meson (iso scalar-scalar).  $\tilde{\sigma}$  is also a candidate as dark matter.

Decay modes & decay width:

$$-\frac{H}{-} \times -\frac{\phi}{-} - \underbrace{\overline{\psi}\psi \to \widetilde{\pi^{+}}, \widetilde{\pi^{-}}, \widetilde{\pi^{0}}, \widetilde{\sigma}}_{(\overline{\psi}\overline{\psi}) \to \widetilde{\pi^{-}}, \widetilde{\pi^{+}}, \widetilde{\pi^{0}}, \widetilde{\sigma}}$$
ume:  $g_{\pi} = g_{\sigma} \equiv g$ 

$$\Gamma_{\phi \to Dark \ mesons} \propto \frac{g^2}{4\pi} \frac{1}{M_{\phi}^2} \sqrt{\left(\frac{M_{\phi}}{2}\right)^2 - (\widetilde{m}_{Dark \ mesons})^2}$$

Total decay width of  $\phi$ :

 $\Gamma^{\phi}_{\text{total}} = \Gamma_{\phi \to \widetilde{\pi^{+}} \widetilde{\pi^{-}}} + \Gamma_{\phi \to \widetilde{\pi^{0}} \widetilde{\pi^{0}}} + \Gamma_{\phi \to \widetilde{\sigma} \widetilde{\sigma}}$ 

### **Method** (The case of SU(3) hidden color gauge sector)

We use an effective theory for dark matter interactions in the framework of the linear sigma model. We scale up the dark meson mass  $m_{\widetilde{DM}}$ , by giving the hadron vacuum expectation value v, and the dark matter energy scale  $v_{DM}$ .

$$m_{\widetilde{DM}} = m \frac{v_{DM}}{v}$$

In hadron physics, the experimental value of  $\sigma$  is not fixed, so the results of the LATTICE calculations are used as a reference. This result does not fully incorporate the effect of chiral symmetry since the simulation is based on the Wilson action. For the other mesons, the median experimental value was used [1]. A recent analysis shows that the dark pion has a mass of 1.8 TeV [2].

S. Navas et al. (Particle Data Group), to be published in Phys. Rev. D 110, 030001 (2024)
 T. Abe, R. Sato, T. Yamanaka, arXiv:2404.03963v1 [hep-ph] (2024)



### **Brunching ratio**

v <sub>DM</sub> [TeV]	$M_{\phi}$ [TeV]	$\widetilde{m}_{\pi}$ [TeV]	$\widetilde{m}_{\sigma}$ [TeV]	$\frac{\Gamma_{\phi \to \widetilde{\pi^+} \widetilde{\pi^-}}}{\Gamma_{\text{total}}}$	$\frac{\Gamma_{\phi \to \widetilde{\pi^0} \widetilde{\pi^0}}}{\Gamma_{total}}$	$\frac{\Gamma_{\phi \to \widetilde{\sigma} \widetilde{\sigma}}}{\Gamma_{\text{total}}}$
1	10	1.08	4.615	0.5893	0.2946	0.1161
1	10	1.08	2.35	0.5123	0.2561	0.2316
1.67	16	1.8	7.715	0.6113	0.3057	0.0830
1.67	16	1.8	3.94	0.5137	0.2569	0.2294

### Summary

Ass

The decay mode  $\phi \to \widetilde{\pi^0} \, \widetilde{\pi^0}$  is a visible process because it is  $\widetilde{\pi^0} \to 2\gamma$ . The other decay modes are invisible processes. We assume that  $\tilde{\sigma}$  decay into  $\widetilde{\pi^+} \, \widetilde{\pi^-}$  and  $\widetilde{\pi^0} \, \widetilde{\pi^0}$ , do not decay into  $2\gamma$  (In hadron physics, it is not yet known whether  $\sigma$  meson decay into photons or not).  $\widetilde{\pi^+}$ , and  $\widetilde{\pi^-}$  could be stable dark matter. We are supposed to identify dark matter through an invisible process at the future linear collider.

This work is supported by "Joint Usage/Research Center for Interdisciplinary Large-scale Information Infrastructures (JHPCN)" and "High Performance Computing Infrastructure (HPCI)" in Japan (Project ID: jh230030 and jh240034).