Heavy Neutral Leptons (HNL) at e^+e^- colliders — theory perspective













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CLUSTER OF EXCELLENCE QUANTUM UNIVERSE



Jürgen R. Reuter

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The quest for Heavy Neutral Leptons (HNL)





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The neutrino mystery

- Neutrinos masses is already physics beyond the standard model G
- Simple extension of SM: just add ν_R and Yukawa couplin
- $-M_{\nu} \overline{\nu^{C}} \nu$ Singlet allows for a Majorana mass term:



ngs
$$\nu_R = (\mathbf{1}, \mathbf{1}, 1) - m_{\nu}(\overline{\nu}_L \nu_R + h \cdot c.) \left(1 + \frac{h}{v}\right)$$

[Minkowski, 1977; Mohapatra/Senjanovic, 1980; Yanagida, 1981] Dedicated "seesaw" models for neutrino physics: type I (singlet fermion), type II (triplet scalar), type III (triplet fermion)







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Explanations for tiny masses

- Suppression by large scale O
- Accidentally small numbers O
- Symmetry protection:

("classic seesaw") (neutrino Yukawas, "tuned")

$$F = \begin{pmatrix} F_e(1+\epsilon_e) & iF_e(1-\epsilon_e) & F_e\epsilon'_e \\ F_\mu(1+\epsilon_\mu) & iF_\mu(1-\epsilon_\mu) & F_\mu\epsilon'_\mu \\ F_\tau(1+\epsilon_\tau) & iF_\tau(1-\epsilon_\tau) & F_\tau\epsilon'_\tau \end{pmatrix}, \quad M_M = \begin{pmatrix} \overline{M}(1-\mu) & 0 & 0 \\ 0 & \overline{M}(1+\mu) & 0 \\ 0 & 0 & M' \end{pmatrix}$$
Neutrino-Yukawa couplings

Ineutino-rukawa coupings

Model realisations:

Ş $\varepsilon, \varepsilon' \ll \mu \ll 1$ Inverse seesaw type Ş $\mu \ll \varepsilon, \varepsilon' \ll 1$ Linear seesaw type c Ģ uMSM type $\varepsilon, \varepsilon', \mu \ll 1$ Ş "Mass communist" type $\mu \ll 1, M' \longrightarrow M$

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e.g. B—L, flavor/discrete symmetries, Froggatt-Nielsen type, e.g. Hagedorn et al., 1408.7118

Mohapatra 1986; Mohapatra/Valle 1986 Akhmedov/Lindner/Schnapka/Valle 1995 Asaka/Shaposhnikov 2005/06; Kersten/Smirnov, 2007







Mixing of light ("flavored") and heavy ("sterile") neutrinos

 $3 \quad au$

$$\nu_{L\ell} = \sum_{k=1}^{3} U_{\ell k} \nu_k + \sum_{k'=4}^{n_R+3} V_{\ell k'} N_{k'}.$$
 Direct se

$$\begin{split} \Delta \mathcal{L} &= -\frac{g_W}{\sqrt{2}} W_{\mu}^+ \sum_{k=1}^3 \sum_{\ell}^{\prime} \left[\overline{\nu}_k U_{\ell k}^* \gamma^{\mu} P_L \ell \right] \\ &- \frac{g_W}{\sqrt{2}} W_{\mu}^+ \sum_{k'=4}^{n_R+3} \sum_{\ell}^{\tau} \left[\overline{N}_{k'} V_{\ell k'}^* \gamma^{\mu} P_L \ell \right] + \text{H.c.} \end{split}$$

Direct constraints: absolute mass scale (KATRIN, β -decay kinematics)

Constraints from cosmology (large scale structure, CMB)

Neutrino oscillation measurements of
$$U_{\ell'k}$$
:

Searches for charged lepton number violation:



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 $\mu \to e\gamma, \tau \to e\gamma, \tau \to \mu\gamma, \mu \to eee, \tau \to \{e, \mu\}^3$ 1605.05081, 0908.2381, 1001.3221





Simplified neutrino model (HNL model)

Simplified model with right-handed (ν SM) and sterile neutrinos After EWSB heavy (sterile) neutrinos do mix with ν SM neutrinos G Lagrangian: $\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_N + \mathcal{L}_{WN\ell} + \mathcal{L}_{ZN\nu} + \mathcal{L}_{HN\nu}$ $\mathcal{L}_N = \xi_{\nu} \cdot \left(\bar{N}_k i \partial N_k - m_{N_k} \bar{N}_k N_k \right) \quad \text{for } k = 1, 2, 3$ $\mathcal{L}_{WN\ell} = -\frac{g}{\sqrt{2}} W^+_{\mu} \sum_{k=1}^3 \sum_{l=e}^\tau \bar{N}_k V^*_{lk} \gamma^{\mu} P_L \ell^- + \text{ h.c.}, \qquad \qquad \bigvee_{\mathbf{N}} W$ k=1 l=eN



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Incomplete literature:

Aguilar-Saavedra ea., hep-ph/0502189; hep-ph/0503026; Shaposhnikov, 0804.4542; Das/Okada, 1207.3734; Banerjee ea., 1503.05491; Antusch, Cazzato, Fischer, 1612.0272; Cai, Han, Li, Ruiz, 1711.02180; Pascoli, Ruiz, Weiland, 1812.08750





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- "Light"/"heavy" HNL decay width scales as $\Gamma_N \sim V_{\ell N}^2 M_N^5 G_F^2$, $V_{\ell N}^2 M_N^3 G_F$ Decay length in lab frame: $\lambda_N = \frac{\beta \gamma}{\Gamma_N} \sim \frac{|\dot{p}|}{V_{\ell_N}^2 M_N^6 G_F^2}$ 0 3 regimes: prompt decays, displaced vertices, long-lived 0
- Neutrino widths: $\Gamma_N \gtrsim \mathcal{O}(1 \text{ keV})$ prompt decays only, 0 no LLP signature, displaced vertices possible for $M_N \lesssim 10 \, {\rm GeV}$



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 - no LLP signature, displaced vertices possible for M_N





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$$M_N^5 G_F^2, V_{\ell N}^2 M_N^3 G_F$$

 $\lesssim 10 \, \text{GeV}$
 $\int_{-\infty}^{\infty} 10^{-2} \int_{-\infty}^{\infty} 10^{-2} \int_{-\infty}^{\infty} 10^{-2} \int_{-\infty}^{\infty} 10^{-2} \int_{-\infty}^{\infty} 1000 - 2000 - 3000$

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Same-sign LNV searches at LHC

from arXiv:2011.02547









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Single production quadratic in mixing angle:

$$\sigma(pp \to N\ell^{\pm} + X) \equiv |V_{\ell N}|^2 \times \sigma_0(pp \to N\ell^{\pm} + X).$$

Pair production/*t*-channel exchange quartic in mixing angle:

$$\sigma(pp \to \ell_i^{\pm} \ell_j^{\pm} + X) \equiv |V_{\ell_i N} V_{\ell_j N}|^2 \times \sigma_0(pp \to \ell_i^{\pm} \ell_j^{\pm} + X).$$

Gigantically smaller background for LNV pair production process

Two cases of hadron collider superiority: rate at *W,Z* decays, same-sign processes

Might be motivation for e^-e^- running (or μ TRISTAN, μ^+e^+ , $\mu^+\mu^+$





Charged current decay



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Charged current decay



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"Light HNL:" Displaced vertex searches at the Z pole





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At lepton colliders, single production work much better than at LHC:

Associated production: $\ell^+\ell^- \rightarrow \nu N$





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Equivalent to LHC (NC/CC) Drell-Yan production: for heavy HNL propagator-suppressed

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- At lepton colliders, single production work much better than at LHC:
- Associated production: $\ell^+\ell^- \rightarrow \nu N$

- Enhancement due to *W*-electron fusion
- Vector boson fusion: $\ell^+\ell^- \to \bar{\nu}\nu N + \ell^+\ell^- N$ (less important)







Equivalent to LHC (NC/CC) **Drell-Yan production: for** heavy HNL propagator-suppressed

absent at LHC! By far dominant process for heavy HNL!



At lepton colliders, single production work much better than at LHC: Associated production: $\ell^+\ell^- \to \nu N$

Enhancement due to *W*-electron fusion Vector boson fusion: $\ell^+\ell^- \to \bar{\nu}\nu N + \ell^+\ell^- N$ (less important)

- At lepton colliders: optimal single channel is $\ell^+\ell^- \to N\nu \to \ell^\pm j j \nu$
- HNL mass reconstructable as resonance peak
- Major backgrounds: $\ell^+\ell^- \rightarrow jj\ell^\pm\nu$, $\ell\ell\ell'\ell'$, $\{jj, jjjj\}\ell\ell$, $jj\ell^+\nu\ell^-\bar{\nu}$
- QED-ISR/beamstrahlung: CLIC-3 vs. MuC-3
- Off-shell processes extend sensitivity beyond collider energy!



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Equivalent to LHC (NC/CC) **Drell-Yan production: for** heavy HNL propagator-suppressed

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absent at LHC! By far dominant process for heavy HNL!



ILC 500 GeV, (-80%, +30%), $m_N = 300$ GeV



Reach for HNLs





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K. Mękała/JRR/A.F. Żarnecki, 2301.02602

cf. next talk by K. Mękała

LHC analysis [1812.08750], diff. assumption: $V_{eN} = V_{\mu N} eq V_{\tau N} = 0$

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High-energy lepton colliders outperform high-energy hadron colliders over the whole mass range!

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- Heavy neutrinos can be Dirac, Majorana or a mixture of both ("pseudo-Dirac")
- HHNL: Lepton collider discriminant: combine CP information (charge + lepton decay angle)
- Solution We were substant of LNC and LNV dilepton events by ratio of SS / OS: $R_{\ell\ell} = \#(\ell^{\pm}\ell^{\pm})/\#(\ell^{+}\ell^{-})$







- Heavy neutrinos can be Dirac, Majorana or a mixture of both ("pseudo-Dirac")
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$$N_{\ell} = \frac{1}{\sqrt{2}}(N_{+} - iN_{-})$$
$$N_{\bar{\ell}} = \frac{1}{\sqrt{2}}(N_{+} + iN_{-})$$

Interaction eigenstates

Mass eigenstates



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- Heavy neutrinos can be Dirac, Majorana or a mixture of both ("pseudo-Dirac")
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LHC separation of LNC and LNV dilepton events by ratio of SS / OS: $R_{\ell\ell} = \#(\ell^{\pm}\ell^{\pm})/\#(\ell^{+}\ell^{-})$ LHNL:

$$g_{+}(t) = e^{-iMt} e^{-\frac{\Gamma}{2}t} \cos\left(\frac{\Delta M}{2}t\right), \qquad N_{\ell}(t) = g_{+}(t)N_{\ell}$$
$$g_{-}(t) = i e^{-iMt} e^{-\frac{\Gamma}{2}t} \sin\left(\frac{\Delta M}{2}t\right) \qquad N_{\bar{\ell}}(t) = g_{-}(t)N_{\ell}$$



 $V_{\ell} + g_{-}(t)N_{\bar{\ell}}$ $N_{\ell} + g_+(t) N_{\bar{\ell}}$



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- $g_{+}(t) = e^{-iMt} e^{-\frac{\Gamma}{2}t} \cos\left(\frac{\Delta M}{2}t\right), \qquad R_{ll} = \frac{\int_{0}^{\infty} |g_{-}|^{2} dt}{\int_{0}^{\infty} |g_{+}|^{2} dt} = g_{-}(t) = i e^{-iMt} e^{-\frac{\Gamma}{2}t} \sin\left(\frac{\Delta M}{2}t\right) \qquad R_{ll} = \frac{\int_{0}^{\infty} |g_{+}|^{2} dt}{\int_{0}^{\infty} |g_{+}|^{2} dt} = \frac{\int_{0}^{\infty} |g_{+}|^{2} dt}{\int_{0}^{\infty} |g_{+}|^{2} dt}$



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LHC separation of LNC and LNV dilepton events by ratio of SS / OS: $R_{\ell\ell} = \#(\ell^{\pm}\ell^{\pm})/\#(\ell^{+}\ell^{-})$

$$= \frac{\Delta M^2}{2\Gamma^2 + \Delta M^2}$$



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LHC separation of LNC and LNV dilepton events by ratio of SS / OS: $R_{\ell\ell} = \#(\ell^{\pm}\ell^{\pm})/\#(\ell^{+}\ell^{-})$

$$= rac{\Delta M^2}{2\Gamma^2 + \Delta M^2}$$
 $R_{ll}
ightarrow 0$ as $(\Gamma/\Delta M)^{-1}
ightarrow 0$ Dirac line $R_{ll}
ightarrow 1$ as $\Gamma/\Delta M
ightarrow 0$ Majoran

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mit na limit

Heavy neutrinos can be Dirac, Majorana or a mixture of both ("pseudo-Dirac") 6 HHNL: Lepton collider discriminant: combine CP information (charge + lepton decay angle) LHC separation of LNC and LNV dilepton events by ratio of SS / OS: $R_{\ell\ell} = \#(\ell^{\pm}\ell^{\pm})/\#(\ell^{+}\ell^{-})$ LHNL:



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mit





- Dominant *t*-channel production (*W* exchange):
- **On-shell production**
- Off-shell more difficult: need to scan each parameter point



$$\sigma \propto \frac{|V_{\ell_{in}}N|^2 \cdot |V_{\ell_{out}N}|^2}{|V_{eN}|^2 + |V_{\mu N}|^2}$$

cf. also talk by Krzysztof Mękała

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Conclusions & Outlook



- Image Meavy Neutrinos Leptons: connections to BAU (?), dark matter (?), flavor symmetries (?)
- Tiny neutrino masses: (1) scale suppression, (2) small parameters ("tuned"), (3) symmetry
- Without more experimental guidance (e.g. DUNE): beware of model prejudices
- Hadron colliders do cover same-sign LNV-/ u-less double beta decay signatures
- Lepton collider LHNL: long-lived particles / displaced vertices (luminosity rules)
- Lepton collider HHNL: superior due to t-channel enhancement (W-lepton fusion)
- Mass peak of HHNL reconstructable in hadronic final states at lepton colliders
- Lepton collider clean environments allow Majorana/Dirac discrimination
- \mathbf{M} Interesting flavor complementarities between hadron, ee and $\mu\mu$ colliders



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Heavy Duty Neutrino

Neutrino

Dr





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SUPER STRONG





