

# W-boson pair production at lepton colliders in the Feynman-Diagram gauge

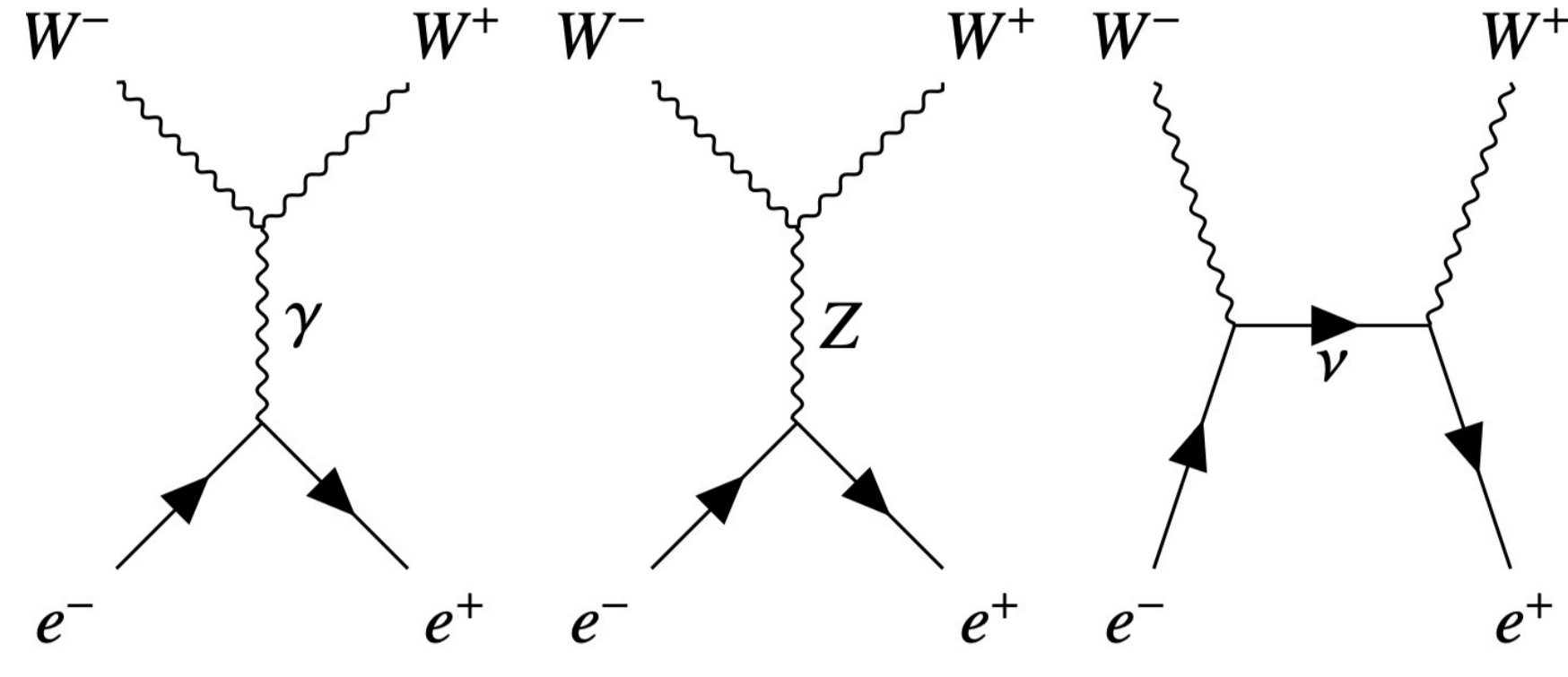


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## 1. Introduction



$$e^-(k, \sigma) + e^+(\bar{k}, \bar{\sigma}) \rightarrow W^-(p, \lambda) + W^+(\bar{p}, \bar{\lambda})$$

- We revisit  $e^-e^+ \rightarrow W^-W^+$  in the SM.
- This process has been thoroughly studied theoretically and experimentally in the LEP era.
- This process is still very important for precision test of the electroweak (EW) theory in future high-energy lepton collider, such as the ILC, CEPC and FCC-ee.
- The gauge cancellation among amplitudes is an obstacle to numerical evaluations, especially at high energy.
- causes cancellation of significant digits.
- Recently a new gauge fixing was proposed and it indicates gauge cancellation among the interfering amplitudes can be avoided!!

## “Feynman-Diagram Gauge”

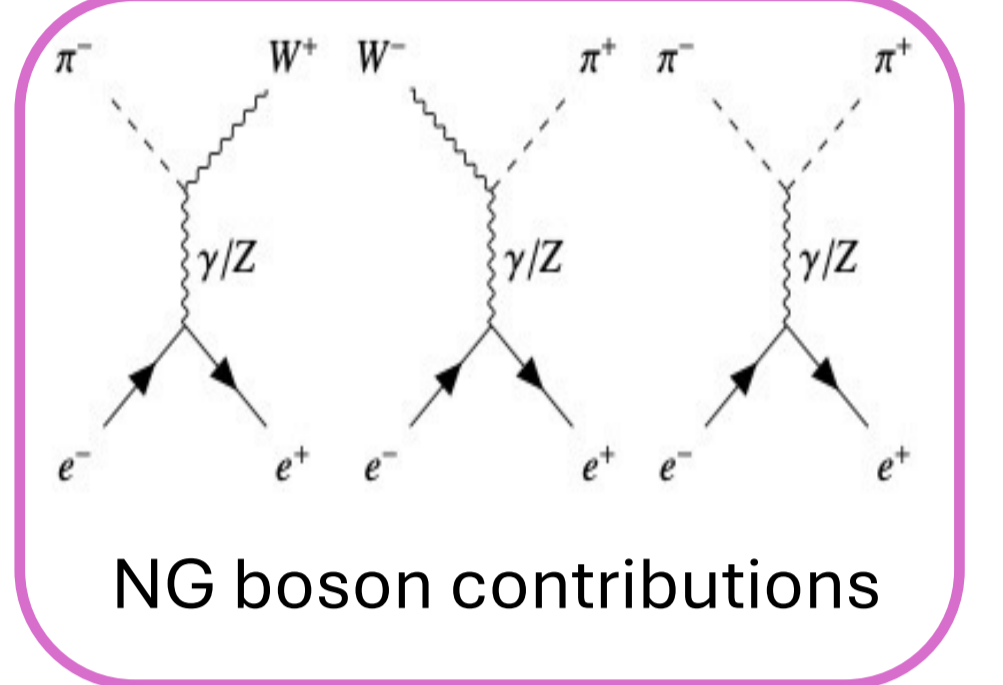
[1],[2],[3]

## 3. Helicity amplitudes

$$\mathcal{M}_{i\sigma}^{\lambda\bar{\lambda}} = \sqrt{2} e^2 c_i \tilde{\mathcal{M}}_i^{\lambda\bar{\lambda}}(\theta) \varepsilon d_{\Delta\sigma, \Delta\lambda}^{J_0}(\theta) P_i(\theta)$$

### U gauge

$\Delta\lambda$	$(\lambda, \bar{\lambda})$	$\tilde{\mathcal{M}}_\gamma^{\lambda\bar{\lambda}}(\theta)$	$\tilde{\mathcal{M}}_Z^{\lambda\bar{\lambda}}(\theta)$	$\tilde{\mathcal{M}}_\nu^{\lambda\bar{\lambda}}(\theta)$
0	(0, 0)	$-2\gamma^2\beta + \beta$	$2\gamma^2\beta - \beta$	$2\gamma^2(\beta - \cos\theta) - \beta$
+1	(+, 0), (0, -)	$-2\gamma\beta$	$2\gamma\beta$	$2\gamma(\beta - \cos\theta) - \frac{1}{\gamma}$
-1	(0, +), (-, 0)	$-2\gamma\beta$	$2\gamma\beta$	$2\gamma(\beta - \cos\theta) + \frac{1}{\gamma}$
0	(±, ±)	$-\beta$	$\beta$	$\beta - \cos\theta$
±2	(±, ∓)	0	0	$-\sqrt{2}$

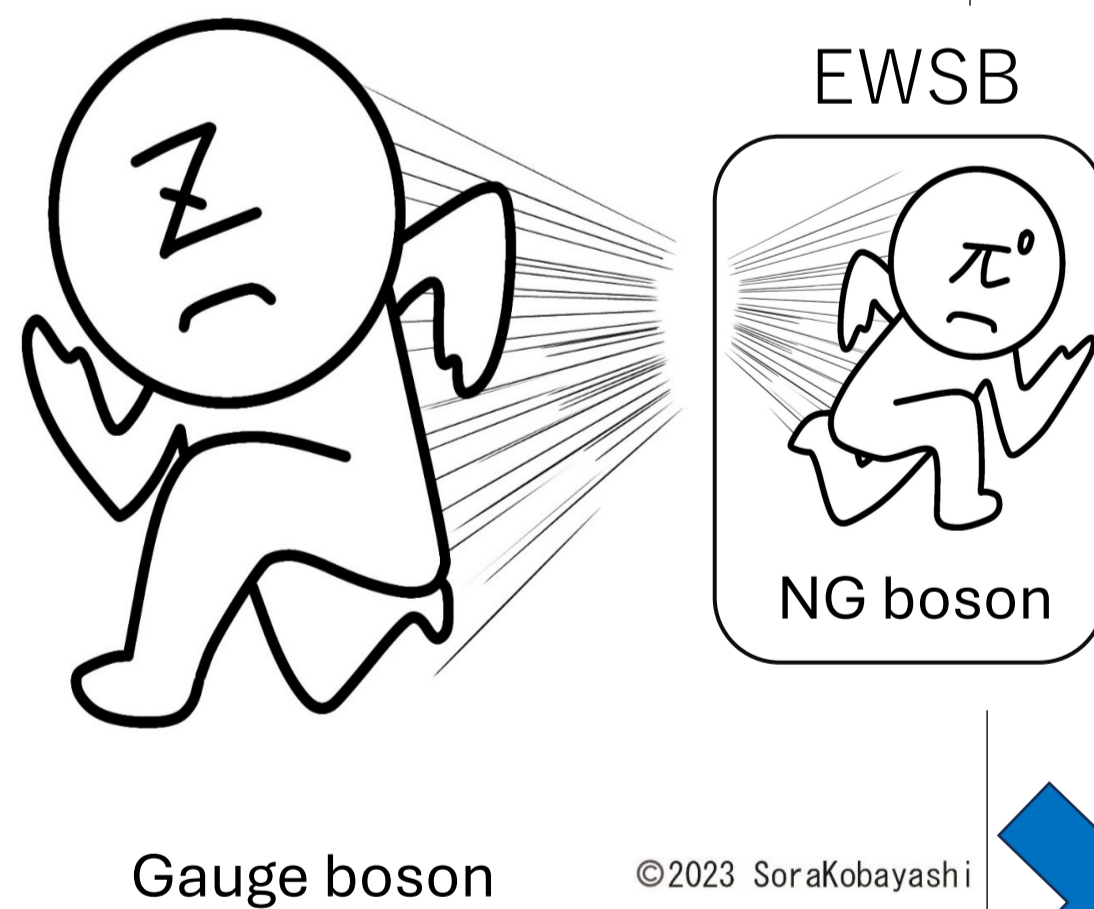


### FD gauge

$\Delta\lambda$	$(\lambda, \bar{\lambda})$	$\tilde{\mathcal{M}}_\gamma^{\lambda\bar{\lambda}}(\theta)$	$\tilde{\mathcal{M}}_Z^{\lambda\bar{\lambda}}(\theta)$	$\tilde{\mathcal{M}}_\nu^{\lambda\bar{\lambda}}(\theta)$
0	(0, 0)	$\frac{1}{\gamma^2} \frac{3+\beta}{(1+\beta)^2} + 1$	$-\frac{1}{\gamma^2} \frac{3+\beta}{(1+\beta)^2} - \frac{s_W^2}{c_W^2} \left( \frac{\beta}{2s_W^2} - 1 \right)$	$-\frac{1}{\gamma^2} \frac{2}{(1+\beta)^2} (1 + \cos\theta)$
+1	(+, 0), (0, -)	$\frac{1}{2\gamma} \left( \frac{3-\beta}{1+\beta} + 1 \right)$	$-\frac{1}{2\gamma} \left( \frac{3-\beta}{1+\beta} - \frac{s_W^2}{c_W^2} \right)$	$-\frac{1}{\gamma} \frac{2}{1+\beta} (1 + \cos\theta)$
-1	(0, +), (-, 0)	$\frac{1}{2\gamma} \left( \frac{3-\beta}{1+\beta} + 1 \right)$	$-\frac{1}{2\gamma} \left( \frac{3-\beta}{1+\beta} - \frac{s_W^2}{c_W^2} \right)$	$\frac{1}{\gamma} \frac{2}{1+\beta} (\beta - \cos\theta)$
0	(±, ±)	$-\beta$	$\beta$	$\beta - \cos\theta$
±2	(±, ∓)	0	0	$-\sqrt{2}$

## 2. Gauge choices

- The treatment of Nambu-Goldstone (NG) bosons depend on gauge choices.
- In the early universe (before EW symmetry breaking), gauge bosons run at the speed of light (mass 0).



When EW symmetry breaks, NG bosons appear.

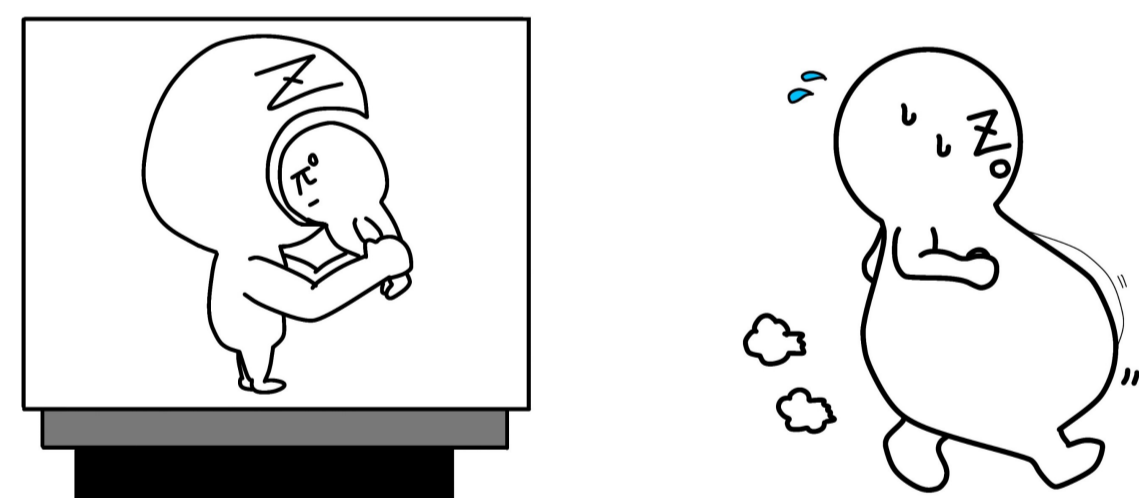
$$q^\mu = (E, 0, 0, q),$$

$$\beta = \frac{q}{E} = \sqrt{1 - \frac{m^2}{E^2}},$$

$$\gamma = \frac{E}{m} = \frac{1}{\sqrt{1 - \beta^2}}.$$

### Unitary (U) gauge

Gauge bosons eat NG bosons and become fat (run slower = get the mass).



$\mu, \nu = 0$  to 3,

$$G_{\mu\nu}(q) = \frac{i}{q^2 - m^2 + i\epsilon} \left( -g_{\mu\nu} + \frac{q_\mu q_\nu}{m^2} \right)$$

$$\epsilon^\mu(q, 0) = \frac{1}{m} (q, 0, 0, E) = \gamma(\beta, 0, 0, 1)$$

### Feynman-Diagram (FD) gauge

Gauge bosons and NG bosons run together (run slower = get the mass).



$M, N = 0$  to 4,

$$n^\mu = (\text{sgn}(q^0), -\vec{q}/|q|)$$

$$G_{MN}(q) = \frac{i}{q^2 - m^2 + i\epsilon} \begin{pmatrix} -g_{MN} + \frac{q_M n_N + n_M q_N}{n \cdot q} & i \frac{m n_M}{n \cdot q} \\ -i \frac{m n_N}{n \cdot q} & 1 \end{pmatrix}$$

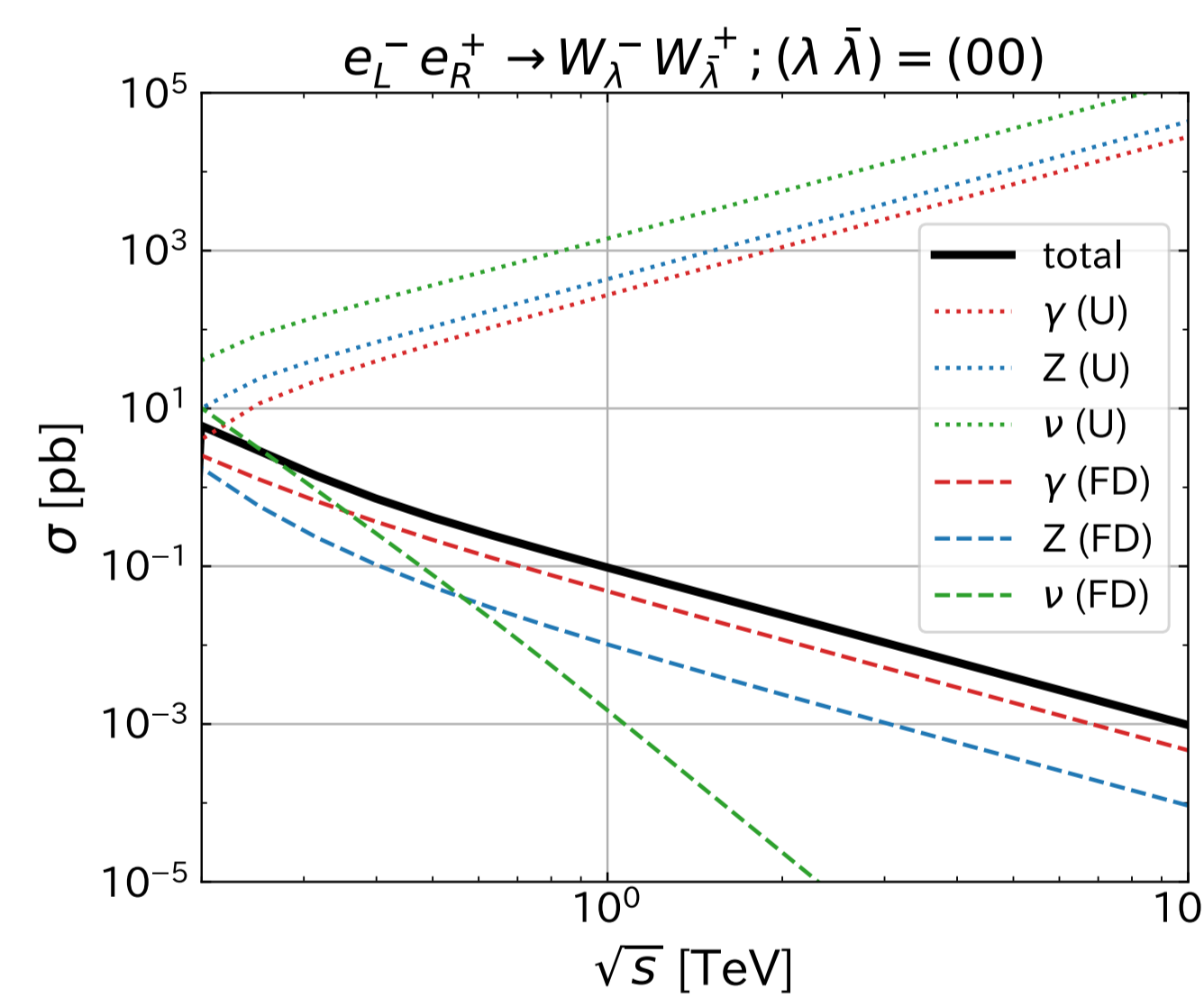
$$\epsilon^M(q, 0) = (\tilde{\epsilon}^M(q, 0), i)$$

$$\tilde{\epsilon}^\mu(q, 0) = \epsilon^\mu(q, 0) - \frac{q^\mu}{Q}$$

$$= \frac{1}{\gamma(1+\beta)} (-1, 0, 0, 1)$$

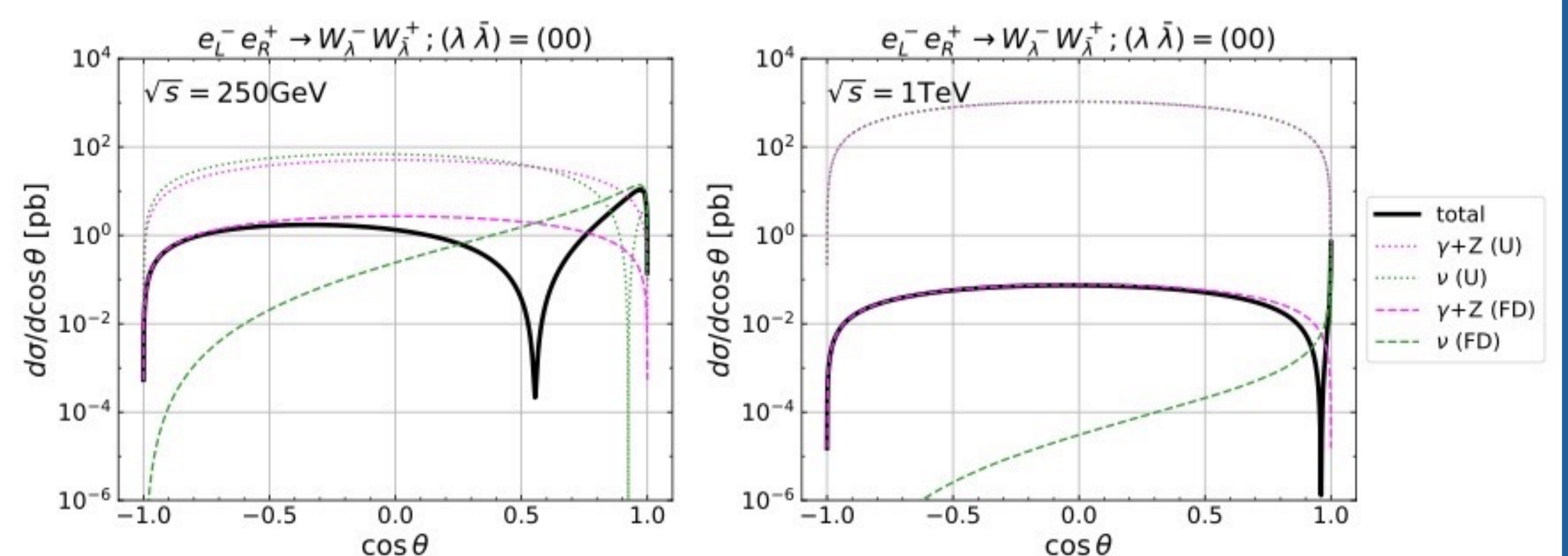
## 4. Cross sections

### Total cross section



- The solid black line is the physical observable.  $total \propto |M_\gamma + M_Z + M_\nu|^2$
- Dotted lines (U gauge) and dashed lines (FD gauge) are unphysical quantities.  $\gamma \propto |M_\gamma|^2, Z \propto |M_Z|^2, \nu \propto |M_\nu|^2$

### Differential cross section



### U gauge:

- Each amplitude has energy growth. → Large cancellation among the three amplitudes ( $\gamma, Z, \nu$ ) at high energies.
- No clue for the physical distribution from each contribution.

### FD gauge:

- No energy growth of individual amplitudes at all. → No unphysical cancellation.
- The contribution from the associated Goldstone bosons is manifest.
- Clear indication to the physical distribution from each contribution.

**Our analytic results in the FD gauge provide a new insight into gauge theories !!**

[1] J. Chen, K. Hagiwara, J. Kanzaki and K. Mawatari, “Helicity amplitudes without gauge cancellation for electroweak processes”, arXiv:2203.10440, Eur.Phys.J.C 83 (2023) 922

[2] J. Chen, K. Hagiwara, J. Kanzaki, K. Mawatari and Y. Zheng, “Helicity amplitudes in light-cone and Feynman-diagram gauges”, arXiv:2211.14562, (2022), Eur.Phys.J.Plus 139 (2024) 332

[3] K. Hagiwara, J. Kanzaki, O. Mattelaer, K. Mawatari and Y. Zheng, “Automatic generation of helicity amplitudes in Feynman-Diagram gauge”, arXiv:2405.01256