# Decoding Higgs Boson Branching Ratios from Event Shape Variables

based on:

Knobbe, Krauss, DR, Schumann EPJC 84 (2024) 1, 83 [arXiv:2306.03682]

Gehrmann-de Ridder, Preuss, DR, Schumann [arXiv:2403.06929]

#### This Talk...

#### Part I:

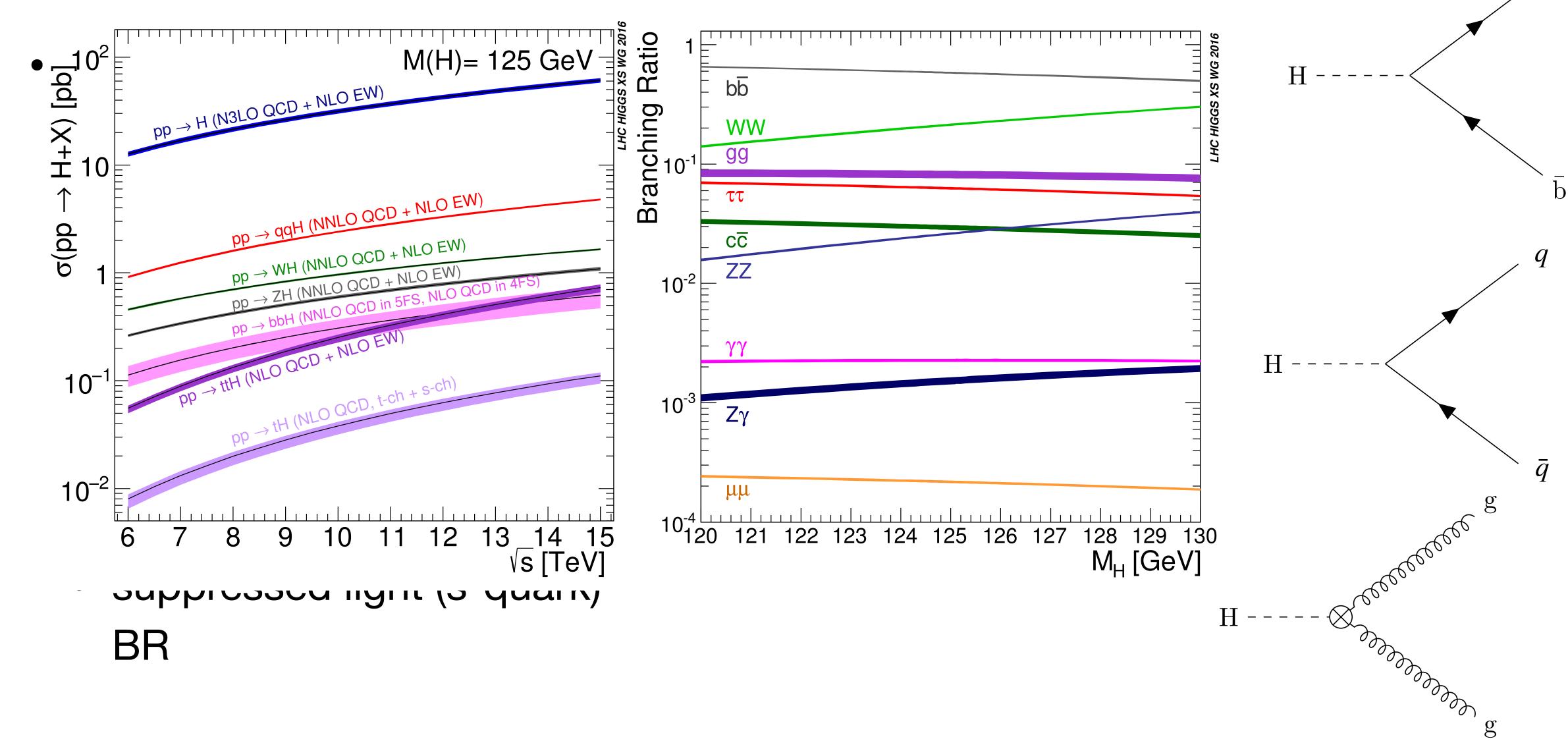
alternative ideas for Measurements of Higgs Boson Branching ratios based on event shapes [Knobbe, Krauss, DR, Schumann '23]

make use of clean environment in FCC-ee setting

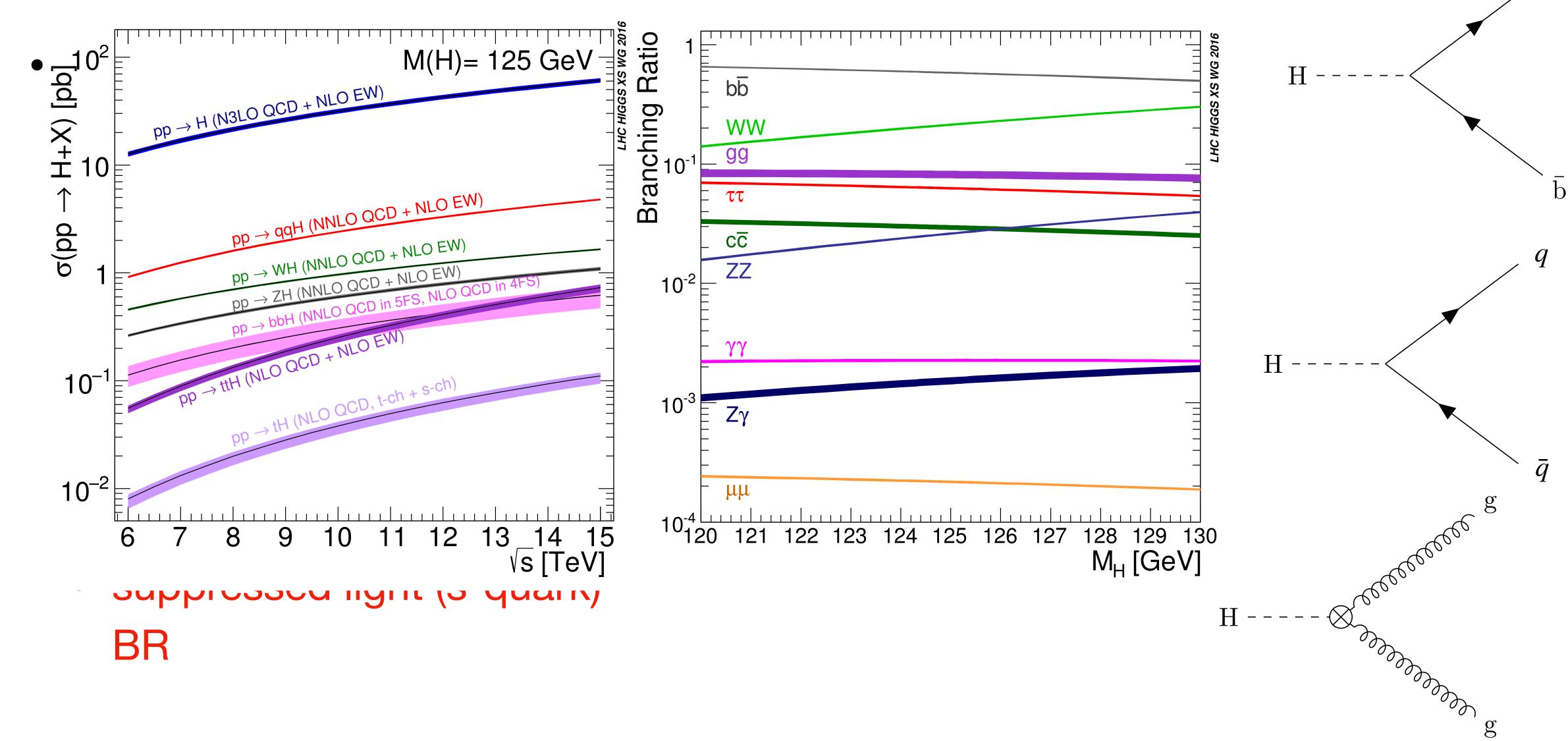
#### Part II:

precision calculations for event shapes in Higgs decays with EERAD [Coloretti, Gehrmann-de Ridder, Preuss '22] [Gehrmann-de Ridder, Preuss, Williams '23] and resummation in Sherpa-CAESAR framework [Gehrmann-de Ridder, Preuss, DR, Schumann '24]

# Hadronic Higgs Decays



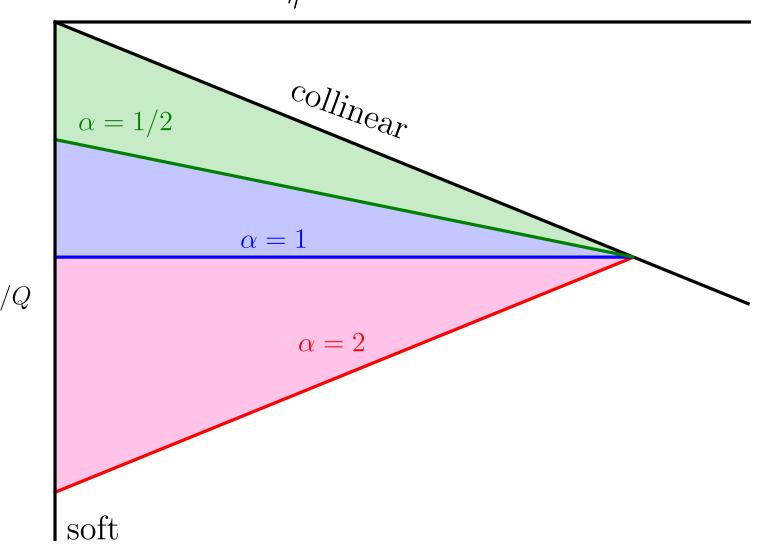
# Hadronic Higgs Decays



uisuivuuviis. Numerical inputs: All predictions are provided with the Fiducial cross section: In Fig. bljet the cros **EVEN INNPOF3.1 NNLO PDF set 1573** with  $\alpha_s(M_Z) = 0.118$  and S predictions for the process  $pp \to Z + b$ -jet at  $n_f^{\rm max} = 5$ , where both the PDF and  $\alpha_s$  values are accessed within the fiducial region defined accord tagging in the specifical theorists personal and  $M_{\ell\bar{\ell}} \in [71, 111]$  the b-jets are reconstruction. and  $M_{e\bar{e}} \in [71, 111]$  GeV. The b-jets are reco • looknears displaced wertex, videntify = with 52 GeV, with the flavor- $k_T$  algorithm with R=0.5, with  $M_W^{os} = 80.385 \text{ GeV}, \quad \Gamma_W^{os} = 2.085 \text{ GeV}, \quad \text{and} \quad G_{\mu} = 0.085 \text{ GeV}$ tional constraint of  $\Delta R(b,\ell) > 0.5$ . As discuss decaying hadron (~ heaving parton) universal corthis matches the fiducial region of the data [8] improvable with various techniques from machine exception of the choice of the jet clustering algorithms and  $\alpha$  and  $\alpha$  and  $\alpha$  are  $\alpha$  as in [59], leads to  $\alpha_{\rm eff} = 0.007779$  . The cross section defined according to Eq. (1) learnith  $\dot{\phi}$   $\dot{\phi}$   $\dot{\phi}$   $\dot{\phi}$  = 0.2293. An uncertainty due to the impact as FONLL, and predictions are shown at both C  $\mathcal{O}(\alpha_s^3)$  as a function of  $m_h$  [as it arises explici of missing higher-order corrections is assessed in the apphreation to varyobecays lues dust acounty tagged prets hesis on the rhs of Eq. (1). The filled band the uncertainty due to scale variation above, the state variation alone, the of 2 around the central scale  $\mu_0 \equiv E_{T,Z}$ , with the additional the upper particular particular performance, but of ten hardes of again, the original particular performance in the performance of the particular particular performance in the upper particular performance in the upper performance in the performance of the performance in the upper performance in the upp under stärldighetween the coefficients appearing in Eq. (1). We **NNLOJET**  $pb \rightarrow Z+b-jet+X$ follow the specific setup of the flavor- $k_T$  algorithm adopted flavor- $k_T$ , R = 0.5,  $\alpha = 2$ NLO 5fs • theory stycky with event or jet shapes distance measure that includes a sum over both QCD partons as well  $m_h^{phys.} = 4.92 \text{ GeV}$ NNLO 5fs necessarily most degrees but accurate 4 understanding possible (i.e. well defined FO diction for a flavored-jet cross section as defined in Eq. (1) restimphation) LHA  $\lambda_{0.5}^1$ can we apply this with landerays a (challeinge! without any vertex info)? flavor. However, there are no data available for the process  $7 \perp b_{iet}$  [8 60\_65] (or in fact any process) that uses

# Event Shapes - Fractional Energy Correlations

- class of observables, typically normalised such that
  - $FC_x \to 0 \Rightarrow$  pencil-like event (little radiation),
  - $FC_x \to FC_x^{\max} < 1 \Rightarrow$  spherical event (a lot of radiation)

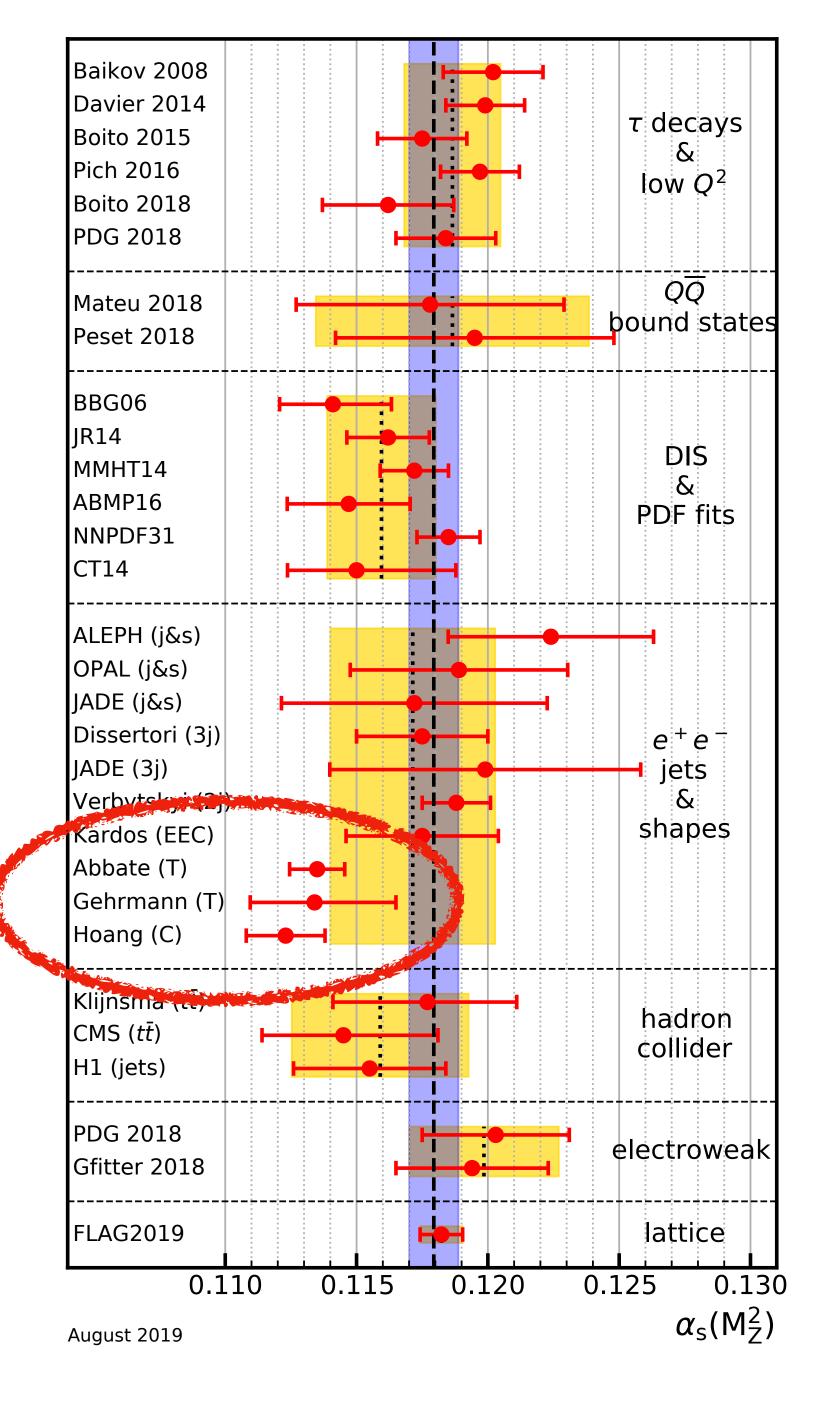


$$FC_x \equiv \sum_{i \neq j} \frac{E_i E_j |\sin \theta_{ij}|^x (1 - |\cos \theta_{ij}|)^{1-x}}{(\sum_i E_i)^2} \Theta\left[ (\vec{q}_i \cdot \vec{n}_T)(\vec{q}_j \cdot \vec{n}_T) \right]$$

- parameter x determines weight of collinear emissions
- analogous to angularities at the LHC ( $\alpha \sim 2 x$ )

#### Event Shapes and $\alpha_s$

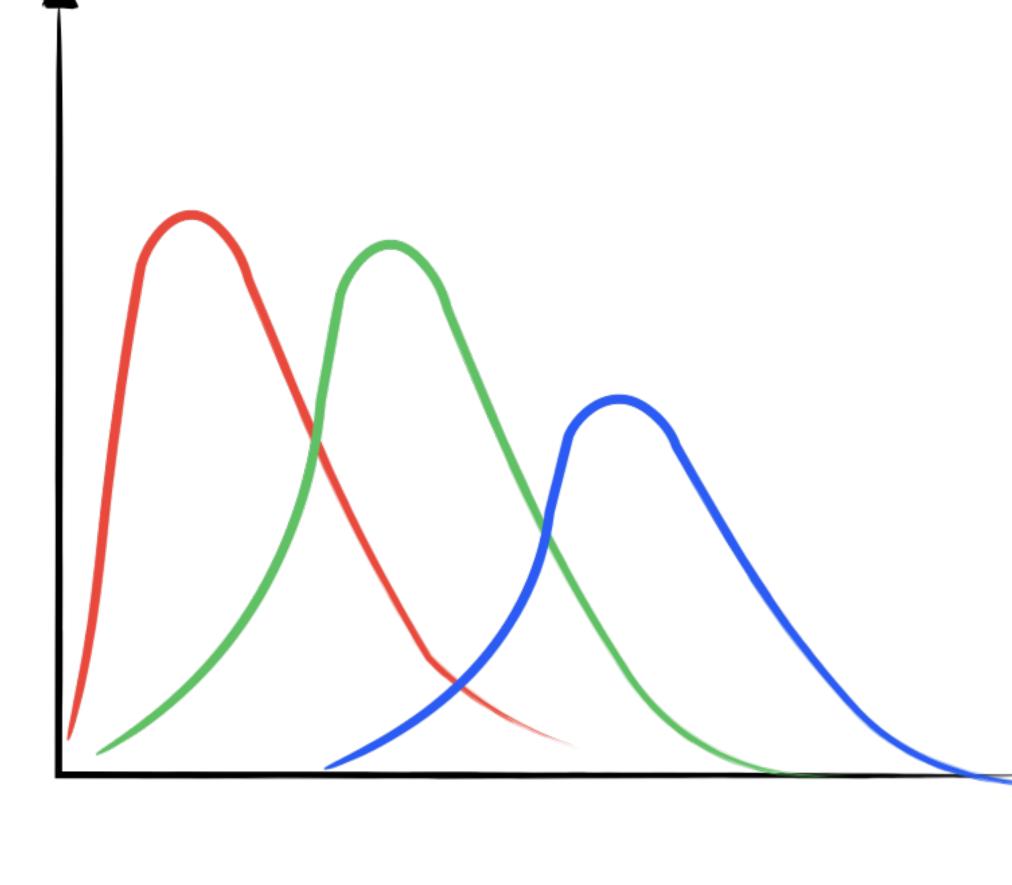
- one traditional way to extract strong coupling constant:
  - high accuracy (NNLO+NNLL) of event shapes (Thrust, C-Parameter etc.) fitted to LEP data at the Z-pole
  - simple 1 or 2 parameter fit, can be pushed by theorists long after experiments are done



# **Event Shapes in Higgs Decays**

- naive picture
  - b-quarks: radiation suppressed due to masses (dead-cone)
  - light quarks: massless radiators, collinear enhancement  $\propto C_F = 4/3$
  - gluons: massless radiators, collinear enhancement

$$\propto C_A > C_F$$



less/softer particles

more/harder particles

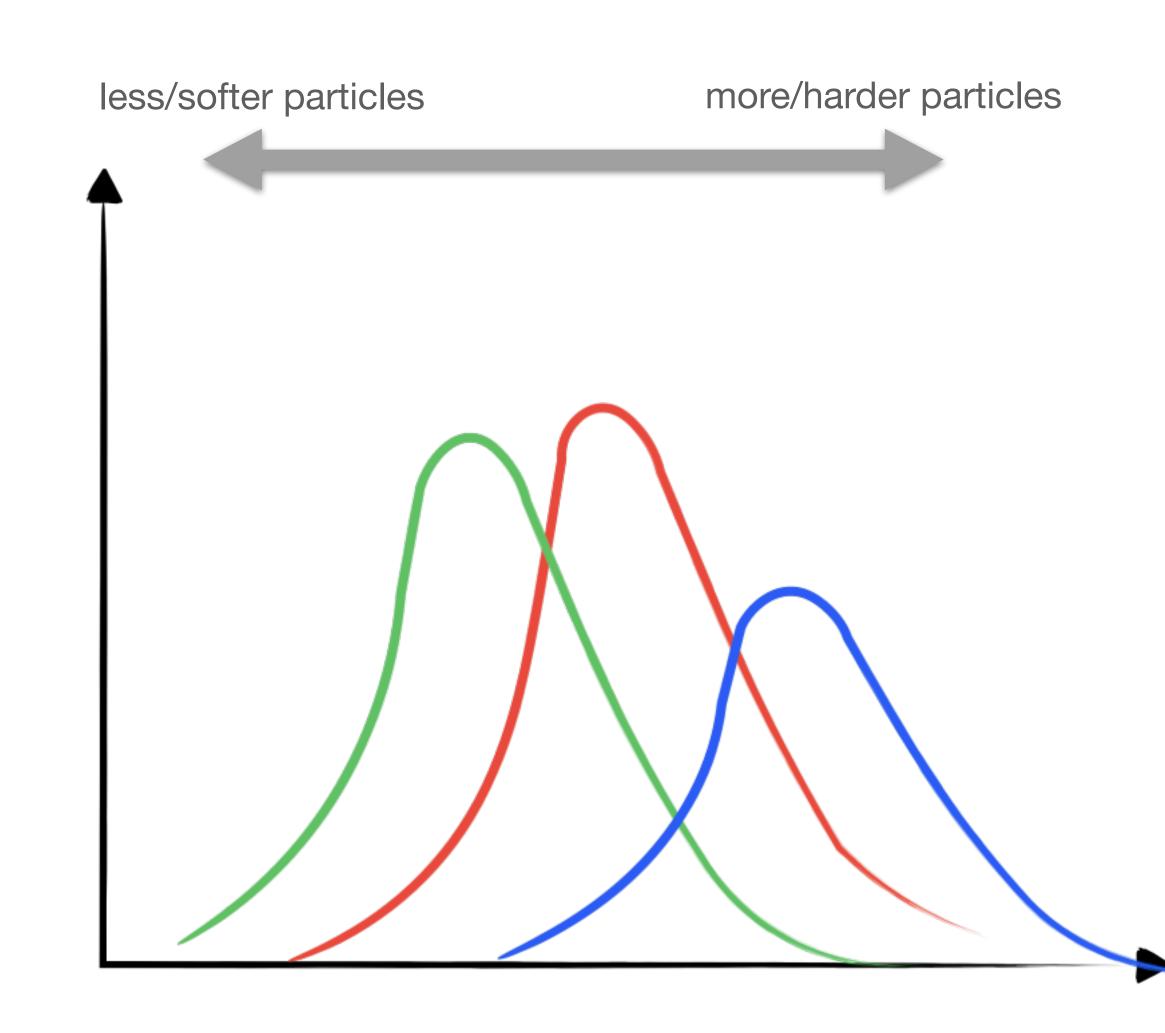
## **Event Shapes in Higgs Decays**

- naive picture, at hadron level
  - b-quarks: radiation suppressed due to masses (dead-cone), high mass decay
  - light quarks: massless radiators, collinear enhancement

$$\propto C_F = 4/3$$

• gluons: massless radiators, collinear enhancement

$$\propto C_A > C_F$$



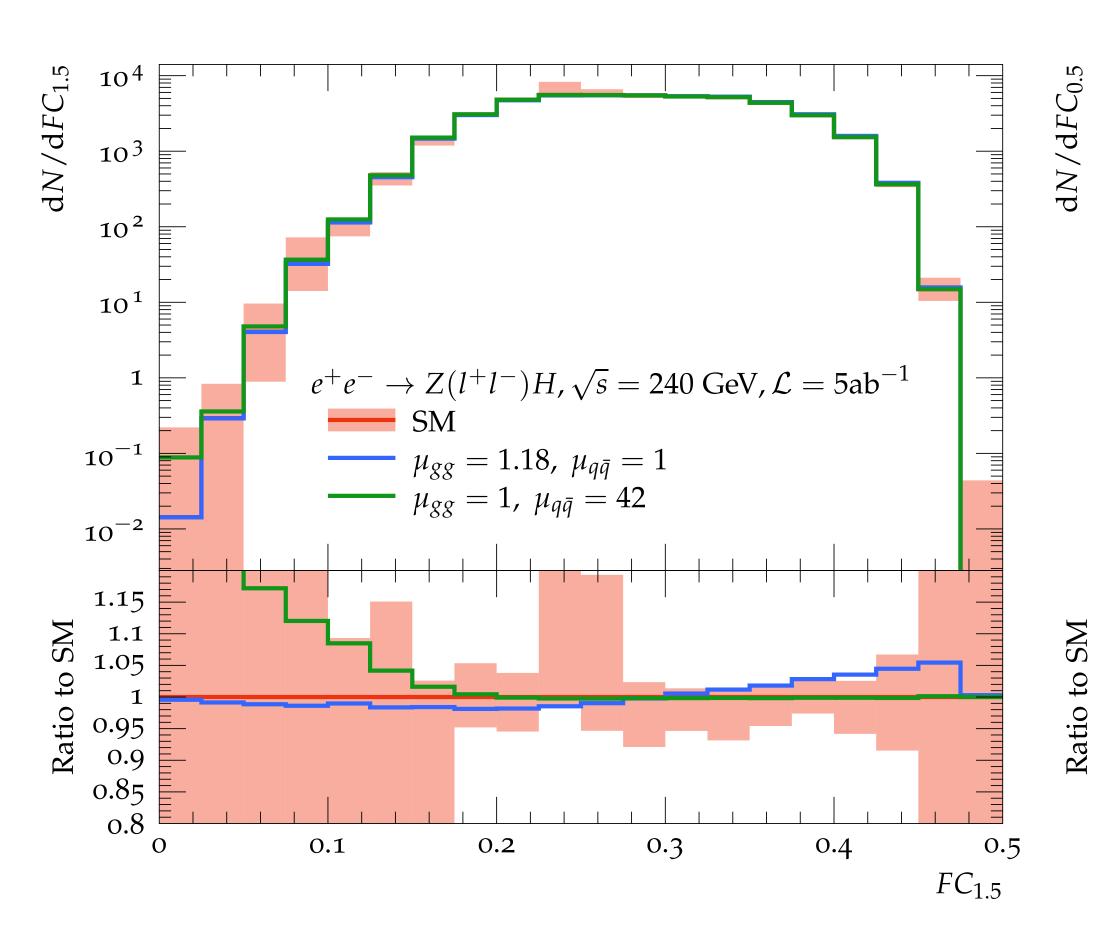
### Effect of Yukawa couplings

overall distribution is sum over hadronic decay channels

$$\frac{\mathrm{d}\sigma}{\mathrm{d}v} = \sum_{i \in \{q\bar{q}, c\bar{c}, b\bar{b}, gg, WW, ZZ\}} \mu_i \frac{\mathrm{d}\sigma_i}{\mathrm{d}v} + \frac{\mathrm{d}\sigma_{ZZ}}{\mathrm{d}v}$$

- can determine relative contribution of each channel (here 2 parameters,  $\mu_{gg}, \mu_{q\bar{q}}$ )
- boundary condition

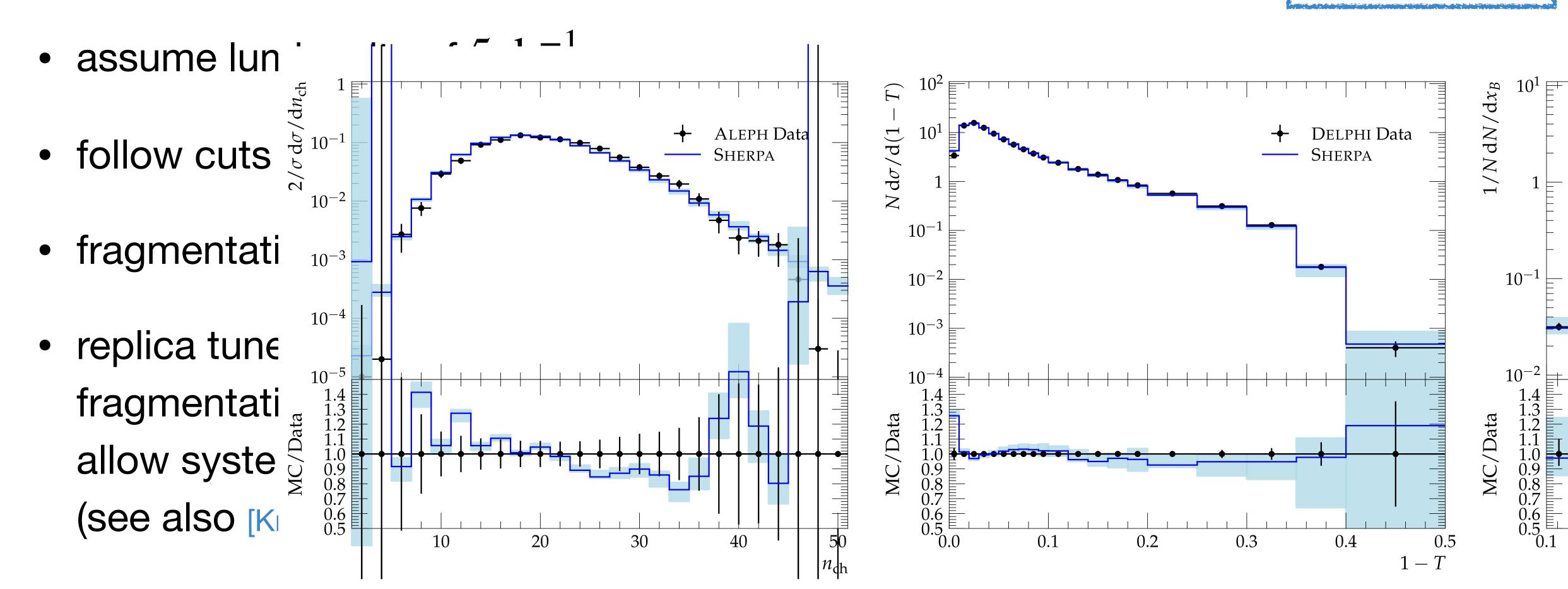
$$\mu_{b\bar{b}} = 1 - (\mu_{gg} - 1)\frac{\sigma_{gg}}{\sigma_{b\bar{b}}} - (\mu_{q\bar{q}} - 1)\frac{\sigma_{q\bar{q}}}{\sigma_{b\bar{b}}}$$



# Study Setup

• run  $e^+e^- \to HZ(\to l^+l^-)$  at  $\sqrt{s}=240~{\rm GeV}$  with Sherpa  $3.0.\alpha$ 

Sherpa 3.0.0 now released, see talk by F. Siegert

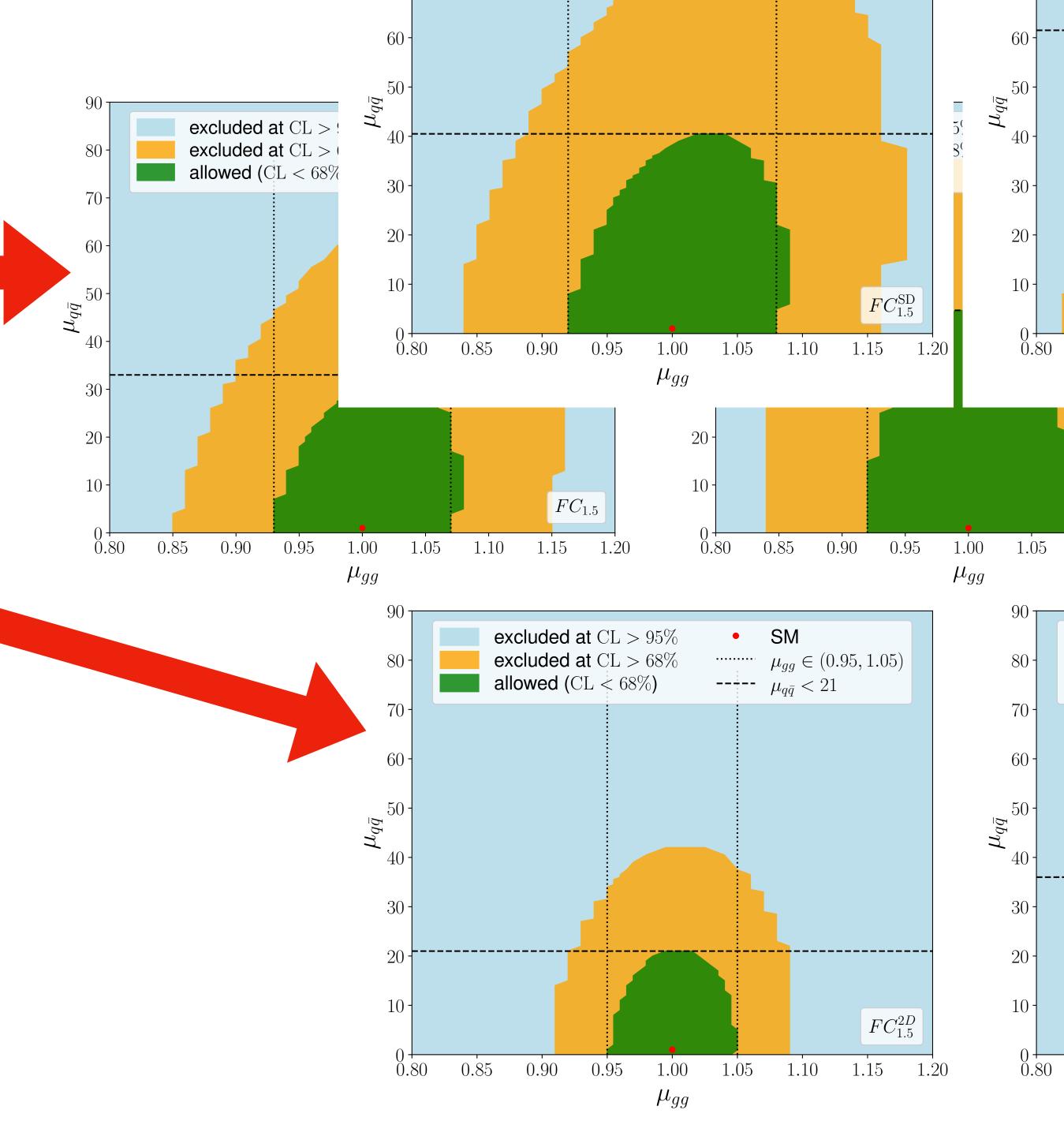


#### Results

• generally best limit from  $FC_{1.5}$  (higher weight on collinear emissions)

 slight improvement from taking into account correlated distribution in each jet

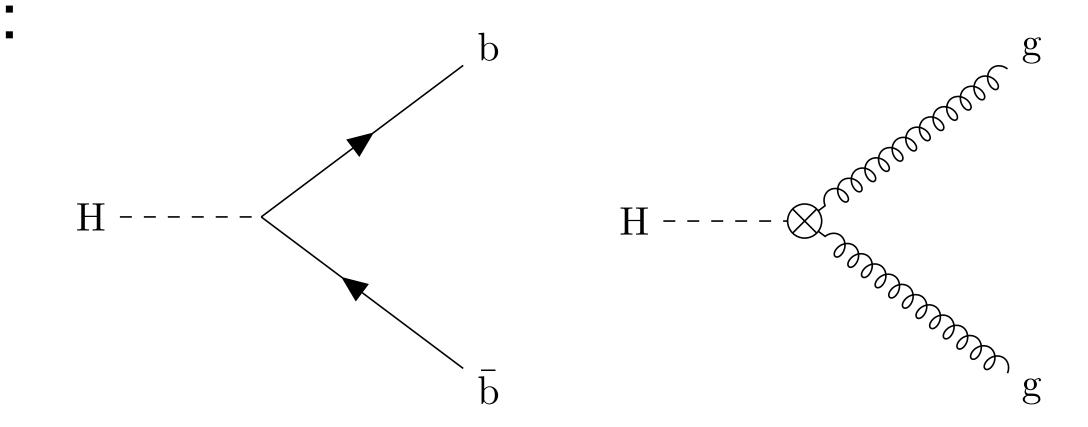
 limit of the same orders as tagging techniques, though not quite competitive (but provides independent methodology)



# Part II - precision calculations for event shapes

#### Precision calculations - Fixed Order

From [Coloretti, Gehrmann-de Ridder, Preuss '22]:
two types of Higgs decays, to
(massless) quarks and gluons via
effective vertex, implemented in
EERAD3



• produces coefficients A,B for IR safe event shape O:

$$\frac{1}{\Gamma^{n}(s,\mu_{\rm R})} \frac{\mathrm{d}\Gamma(s,\mu_{\rm R},O)}{\mathrm{d}O} = \frac{\Gamma^{0}(\mu_{\rm R})}{\Gamma^{n}(s,\mu_{\rm R})} \left(\frac{\alpha_{\rm s}(\mu_{\rm R})}{2\pi}\right) \frac{\mathrm{d}A(s)}{\mathrm{d}O} + \frac{\Gamma^{0}(\mu_{\rm R})}{\Gamma^{n}(s,\mu_{\rm R})} \left(\frac{\alpha_{\rm s}(\mu_{\rm R})}{2\pi}\right)^{2} \frac{\mathrm{d}B(s,\mu_{\rm R})}{\mathrm{d}O}$$

 needs addition of all orders (resummed) calculation

## Resummation - CAESAR in Sherpa

CAESAR formalism for soft gluon resummation at NLL

[Banfi, Salam, Zanderighi '04]

• available as implementation in Sherpa

[Gerwick, Höche, Marzani, Schumann '15] [Baberuxki, Preuss, DR, Schumann '19]

- multiplicative matching (⇒ NLL' accurate)
- necessary extensions for jet observables...:
- [Dasgupta, Khelifa-Kerfa, Marzani, Spannowski '12]
- modified wide angle behaviour [Caletti, Fedkevych, Marzani, DR, Schumann, Soyez, Theeuwes '21]
  - [DR, Caletti, Fedkevych, Marzani, Schumann, Soyez '22]

non-global logs

[Dasgupta, Salam '01]

• ... and for soft drop grooming

[Larkoski, Marzani, Soyez, Thaler '14]

CEASAR style formulas available

[Baron, DR, Schumann, Schwanemann, Theeuwes '20]

#### Precision Calculations - Resummation with CAESAR

• master formula for rIRC save observable: [Banfi, Salam, Zanderighi '04]

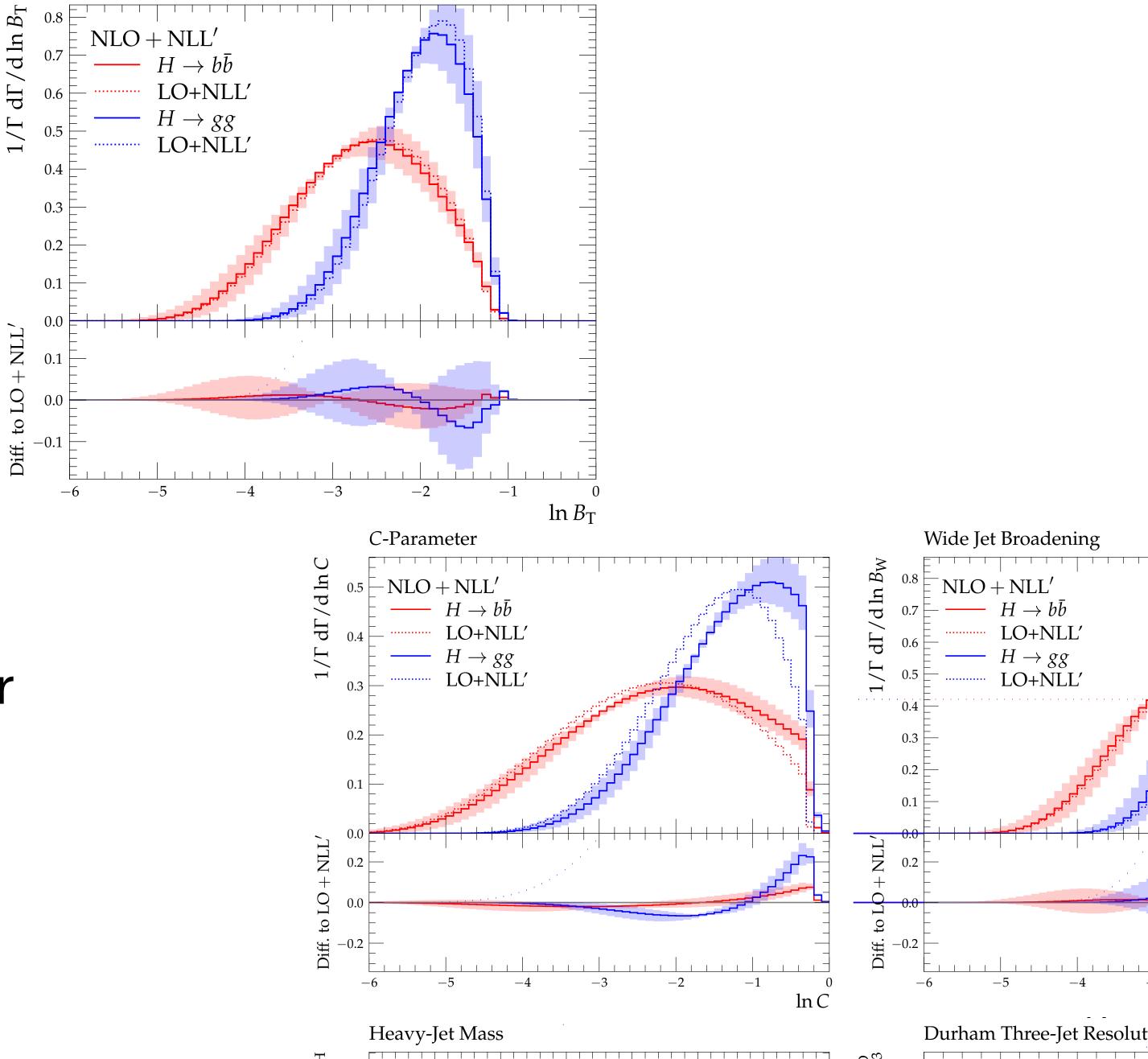
$$\Sigma_{\rm res}^{\delta}(v) = \int d\mathcal{B}_{\delta} \frac{d\sigma_{\delta}}{d\mathcal{B}_{\delta}} \exp \left[ -\sum_{l \in \delta} R_{l}^{\mathcal{B}_{\delta}}(L) \right] \mathcal{P}^{\mathcal{B}_{\delta}}(L) \mathcal{S}^{\mathcal{B}_{\delta}}(L) \mathcal{F}^{\mathcal{B}_{\delta}}(L) \mathcal{H}^{\delta}(\mathcal{B}_{\delta})$$

- ingredients known analytically in our cases
- matching:  $\Sigma_{\mathrm{matched}} = \Sigma_{\mathrm{res}} \left( 1 + \frac{\Sigma_{\mathrm{fo}}^{(1)} \Sigma_{\mathrm{res}}^{(1)}}{\sigma^{(0)}} + \frac{\Sigma_{\mathrm{fo}}^{(2)} \Sigma_{\mathrm{res}}^{(2)}}{\sigma^{(0)}} \frac{\Sigma_{\mathrm{res}}^{(1)}}{\sigma^{(0)}} \frac{\Sigma_{\mathrm{fo}}^{(1)} \Sigma_{\mathrm{res}}^{(1)}}{\sigma^{(0)}} \right)$
- here for the first time, handle external (to Sherpa) fixed order calculation, given in terms of A,B distributions for matching:

$$\frac{1}{\Gamma^n(s,\mu_{\rm R})} \frac{\mathrm{d}\Gamma(s,\mu_{\rm R},O)}{\mathrm{d}O} = \frac{\Gamma^0(\mu_{\rm R})}{\Gamma^n(s,\mu_{\rm R})} \left(\frac{\alpha_{\rm s}(\mu_{\rm R})}{2\pi}\right) \frac{\mathrm{d}A(s)}{\mathrm{d}O} + \frac{\Gamma^0(\mu_{\rm R})}{\Gamma^n(s,\mu_{\rm R})} \left(\frac{\alpha_{\rm s}(\mu_{\rm R})}{2\pi}\right)^2 \frac{\mathrm{d}B(s,\mu_{\rm R})}{\mathrm{d}O} \quad 16$$

#### Results

- example: jet/hemisphere broadening
- expected separation between quark and gluon final states
  - gluons dominate much "harde configurations than quarks
- expected "Sudakov shoulder" for observables like C-parameter

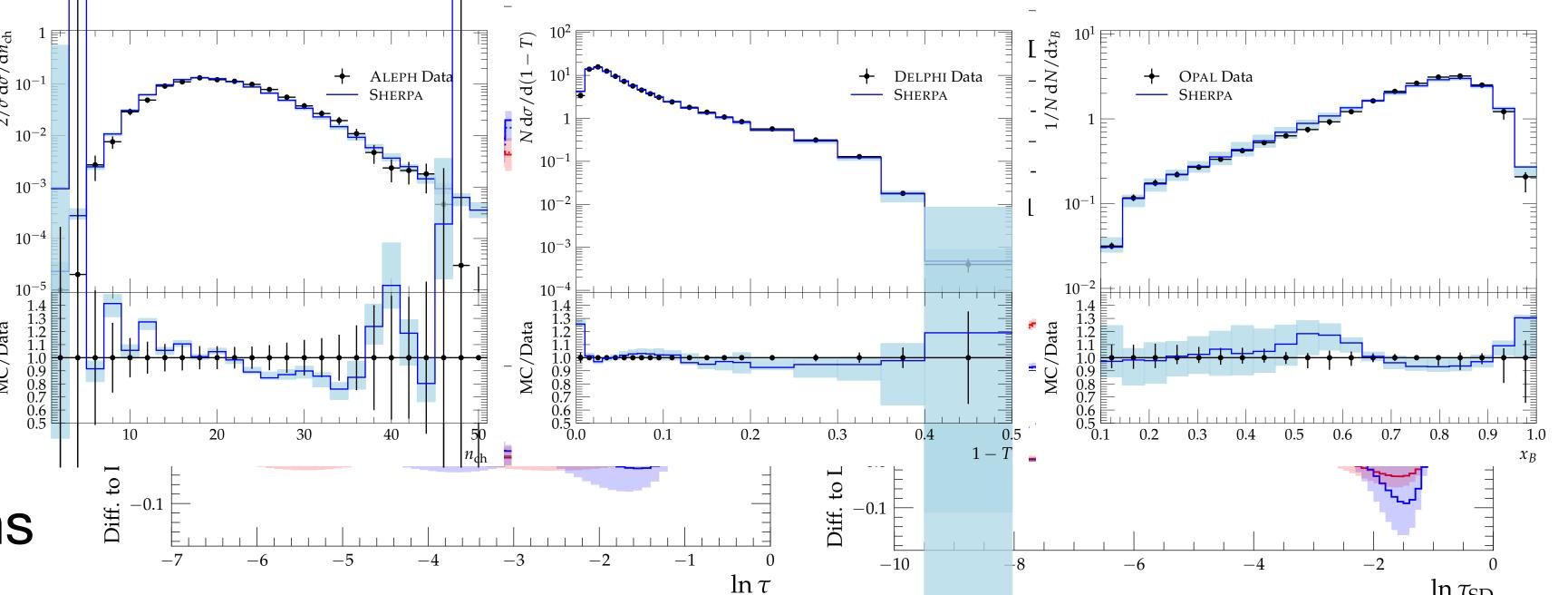


Total Jet Broadening

#### Results

 example: Thrust and Soft-Drop groomed tl

 no grooming: well separated distributions

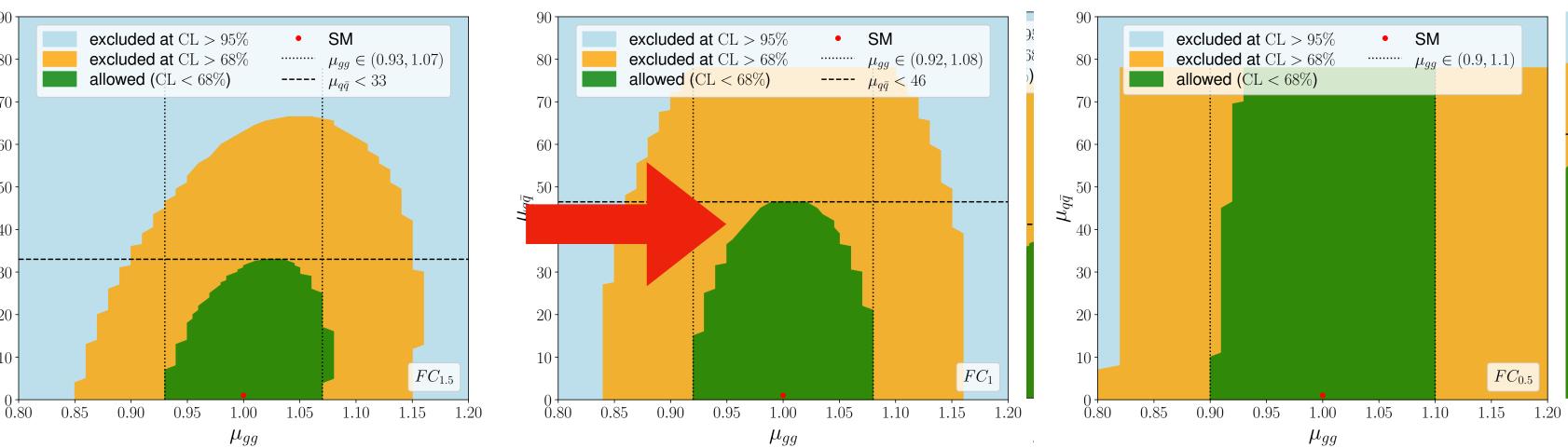


Soft-Drop Thrust

 with grooming: overlap of Sudakov peak for gluons with transition region of quark distribution, possible explanation for limited additional separation

Thrust

power



# Summary

- Event shapes as theoretically well controllable taggers
  - applicable to hadronic Higgs decays
  - enables measurement with minimal (possibly without) modelling input
  - strongest limits for  $FC_{1.5}$  observable, in particular 2D version
- accompanied by precision calculation of event shapes
  - NLO+NLL' almost trivially available
  - NNLO+NNLL in principle available
- outlook: studies with next generation of NLL dipole showers (see for example Alaric talk on Wednesday)