

# EFT in LHC

Tatsuya Masubuchi

University of Tokyo (ICEPP)



東京大学  
素粒子物理国際研究センター  
International Center for Elementary Particle Physics  
The University of Tokyo

# Higgs EFT interpretation in LHC

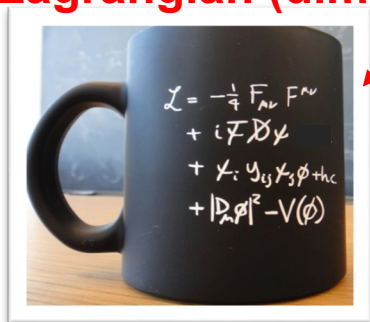
- No any BSM evidence in LHC (so far) from direct searches
- Effective Field Theory (EFT) can set model-independent constraint on BSM physics and indirectly searches beyond LHC reach
- LHC often uses SMEFT (Warsaw basis)

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i^{N_{d=6}} \frac{c_i}{\Lambda^2} O_i^{(6)} + \sum_j^{N_{d=8}} \frac{b_j}{\Lambda^4} O_j^{(8)} + \dots,$$

dim6                      dim8

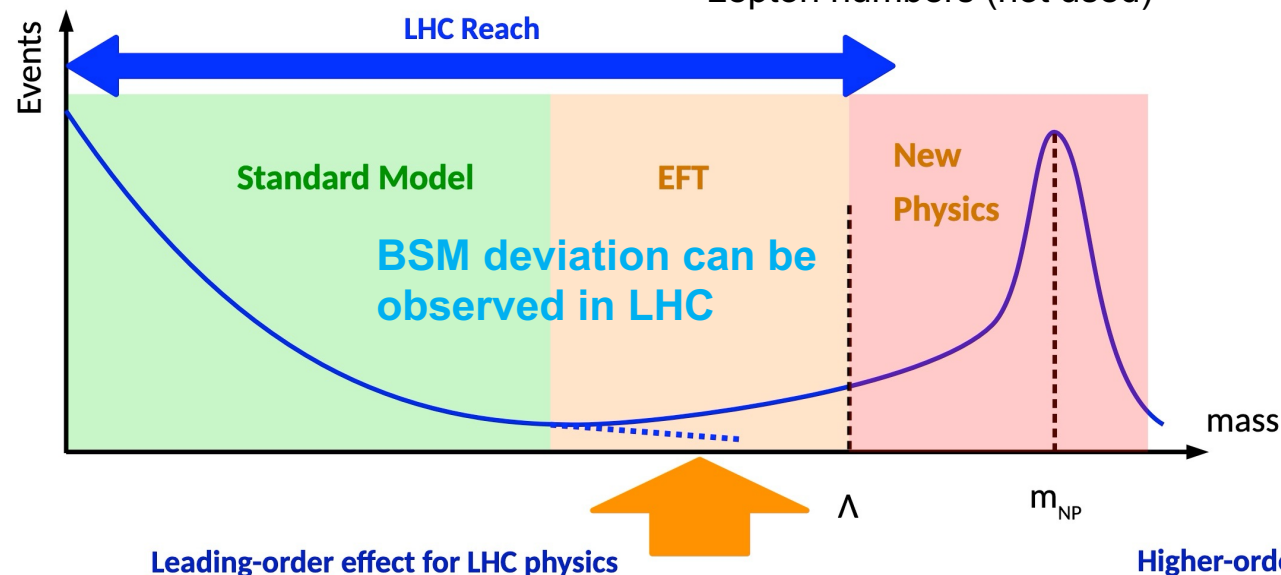
$c_i, b_j$ : Wilson coefficients  
 $O_i^{(6)}, O_j^{(8)}$ : dim6/8 operators

SM Lagrangian (dim4)



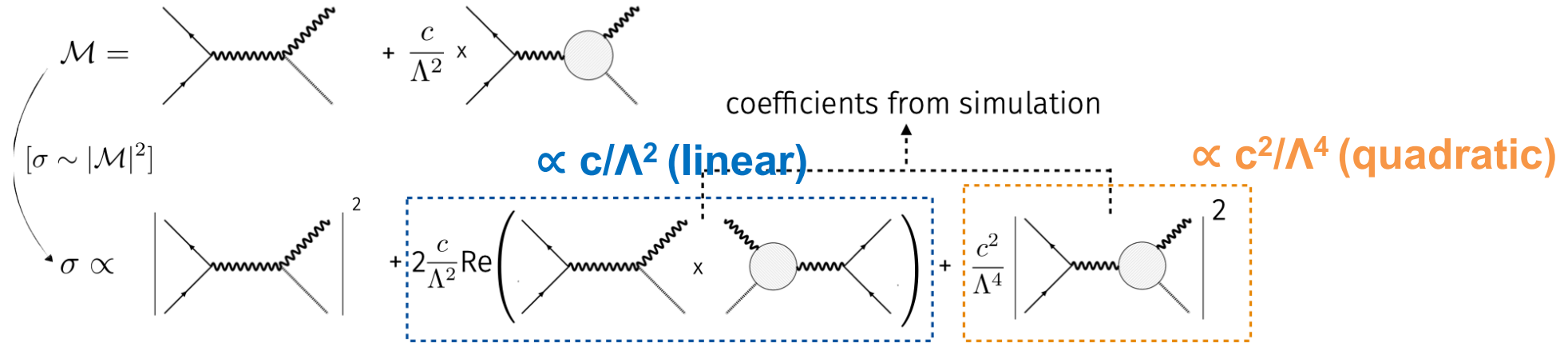
\* dim5/7 violate Baryon and/or Lepton numbers (not used)

- ✓ dim6(8) suppressed by  $1/\Lambda^2(\Lambda^4)$
- ✓ dim8 is not available
- ✓ Common and powerful tool for other measurements → Allow one to combine various measurements



# Higgs EFT formalism

$$\sigma_{\text{SMEFT}} = |\mathcal{L}_{SM} + \mathcal{L}_{BSM}^{\text{dim6}}|^2 = \sigma_{SM} + \sigma_{int} + \sigma_{BSM}$$



Assuming,  $\Lambda=1$  TeV

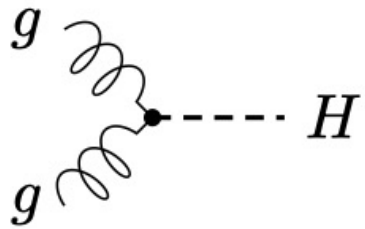
## → Observe cross-section enhancement(kinematic dependence)

- Only CP-even dim-6 operators
- Describe 3<sup>rd</sup> generation from first two generations independently (Top flavor symmetry scheme)
- All lepton generations are modeled independently

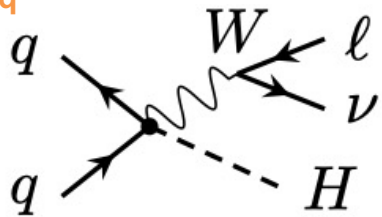
# EFT operators in Higgs

Wilson coefficient	Operator	Wilson coefficient	Operator
$c_H$	$(H^\dagger H)^3$	$c_{Qq}^{(1,1)}$	$(\bar{Q}\gamma_\mu Q)(\bar{q}\gamma^\mu q)$
$c_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	$c_{Qq}^{(1,8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{q}T^a\gamma^\mu q)$
$c_G$	$f^{abc}G_\mu^{ab}G_\nu^{bc}G_\rho^{ca}$	$c_{Qq}^{(3,1)}$	$(\bar{Q}\sigma^i\gamma_\mu Q)(\bar{q}\sigma^i\gamma^\mu q)$
$c_W$	$\epsilon^{IJK}W_\mu^I W_\nu^J W_\rho^K$	$c_{Qq}^{(3,8)}$	$(\bar{Q}\sigma^i T^a\gamma_\mu Q)(\bar{q}\sigma^i T^a\gamma^\mu q)$
$c_{HDD}$	$(H^\dagger D^\mu H)^*(H^\dagger D_\mu H)$	$c_{qq}^{(3,1)}$	$(\bar{q}\sigma^i\gamma_\mu q)(\bar{q}\sigma^i\gamma^\mu q)$
$c_{HG}$	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	$c_{tu}^{(1)}$	$(\bar{t}\gamma_\mu t)(\bar{u}\gamma^\mu u)$
$c_{HB}$	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	$c_{tu}^{(8)}$	$(\bar{t}T^a\gamma_\mu t)(\bar{u}T^a\gamma^\mu u)$
$c_{HW}$	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	$c_{td}^{(1)}$	$(\bar{t}\gamma_\mu t)(\bar{d}\gamma^\mu d)$
$c_{HWB}$	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	$c_{td}^{(8)}$	$(\bar{t}T^a\gamma_\mu t)(\bar{d}T^a\gamma^\mu d)$
$c_{HL,11}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{l}_1\gamma^\mu l_1)$	$c_{Qu}^{(1)}$	$(\bar{Q}\gamma_\mu Q)(\bar{u}\gamma^\mu u)$
$c_{HL,22}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{l}_2\gamma^\mu l_2)$	$c_{Qu}^{(8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{u}T^a\gamma^\mu u)$
$c_{HL,33}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{l}_3\gamma^\mu l_3)$	$c_{Qd}^{(1)}$	$(\bar{Q}\gamma_\mu Q)(\bar{d}\gamma^\mu d)$
$c_{HL,11}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}_\mu^I H)(\bar{l}_1\tau^I\gamma^\mu l_1)$	$c_{Qd}^{(8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{d}T^a\gamma^\mu d)$
$c_{HL,22}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}_\mu^I H)(\bar{l}_2\tau^I\gamma^\mu l_2)$	$c_{tq}^{(1)}$	$(\bar{q}\gamma_\mu q)(\bar{t}\gamma^\mu t)$
$c_{HL,33}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}_\mu^I H)(\bar{l}_3\tau^I\gamma^\mu l_3)$	$c_{tq}^{(8)}$	$(\bar{q}T^a\gamma_\mu q)(\bar{t}T^a\gamma^\mu t)$
$c_{He,11}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{e}_1\gamma^\mu e_1)$	$c_{eH,22}$	$(H^\dagger H)(\bar{l}_2 e_2 H)$
$c_{He,22}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{e}_2\gamma^\mu e_2)$	$c_{eH,33}$	$(H^\dagger H)(\bar{l}_3 e_3 H)$
$c_{He,33}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{e}_3\gamma^\mu e_3)$	$c_{uH}$	$(H^\dagger H)(\bar{q}Y_u^\dagger u \tilde{H})$
$c_{Hq}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{q}\gamma^\mu q)$	$c_{tH}$	$(H^\dagger H)(\bar{Q}\tilde{H}t)$
$c_{Hq}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}_\mu^I H)(\bar{q}\tau^I\gamma^\mu q)$	$c_{bH}$	$(H^\dagger H)(\bar{Q}\tilde{H}b)$
$c_{Hu}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{u}_p\gamma^\mu u_r)$	$c_{tG}$	$(\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{H}G_{\mu\nu}^A$
$c_{Hd}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{d}_p\gamma^\mu d_r)$	$c_{tW}$	$(\bar{Q}\sigma^{\mu\nu}t)\tau^I\tilde{H}W_{\mu\nu}^I$
$c_{HQ}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{Q}\gamma^\mu Q)$	$c_{tB}$	$(\bar{Q}\sigma^{\mu\nu}t)\tilde{H}B_{\mu\nu}$
$c_{HQ}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}_\mu^I H)(\bar{Q}\tau^I\gamma^\mu Q)$	$c_{ll,1221}$	$(\bar{l}_1\gamma_\mu l_2)(\bar{l}_2\gamma^\mu l_1)$
$c_{Ht}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{t}\gamma^\mu t)$		
$c_{Hb}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar{b}\gamma^\mu b)$		

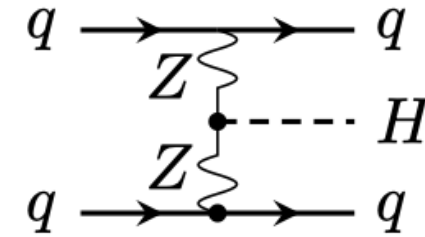
$C_{HG}$



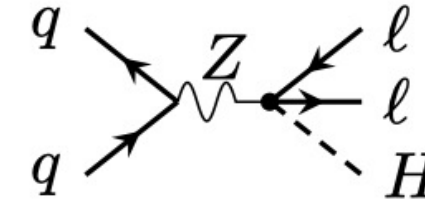
$C^3_{Hq}$



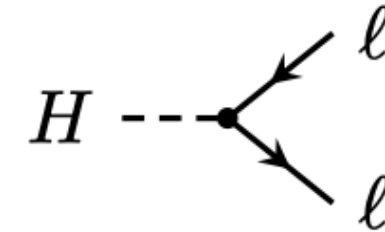
$C_{HDD}$



$C_{HI}$



$C_{eH}$

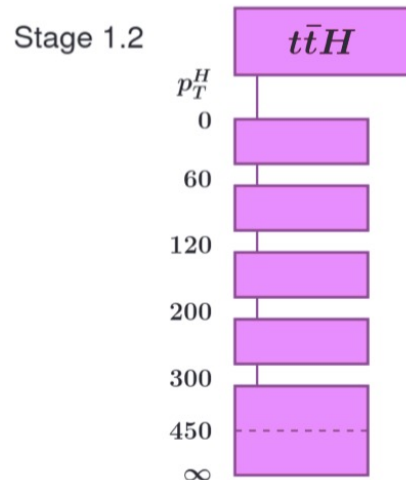
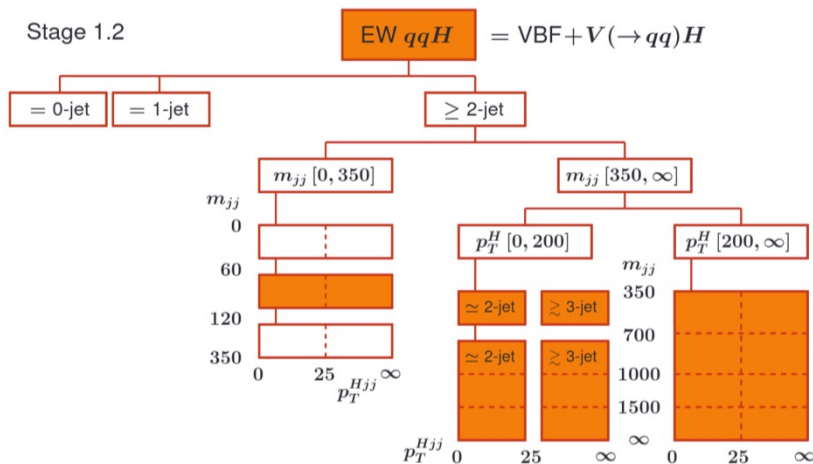
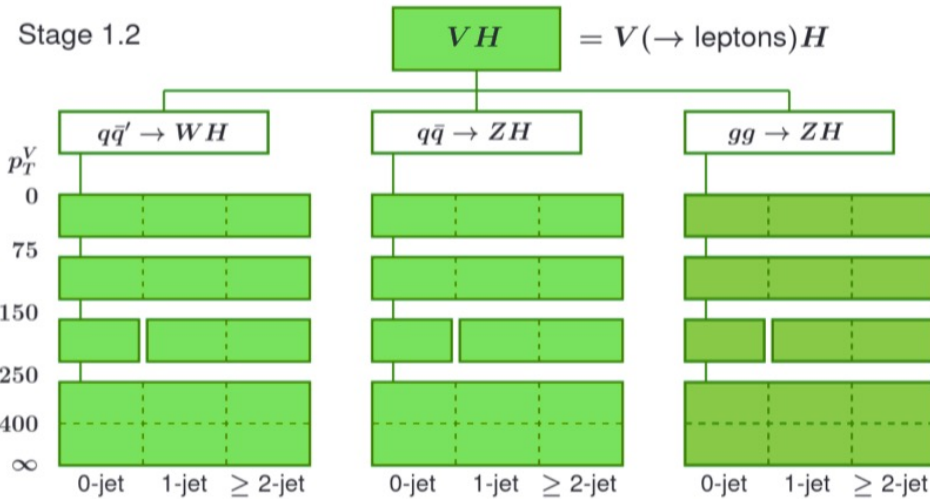
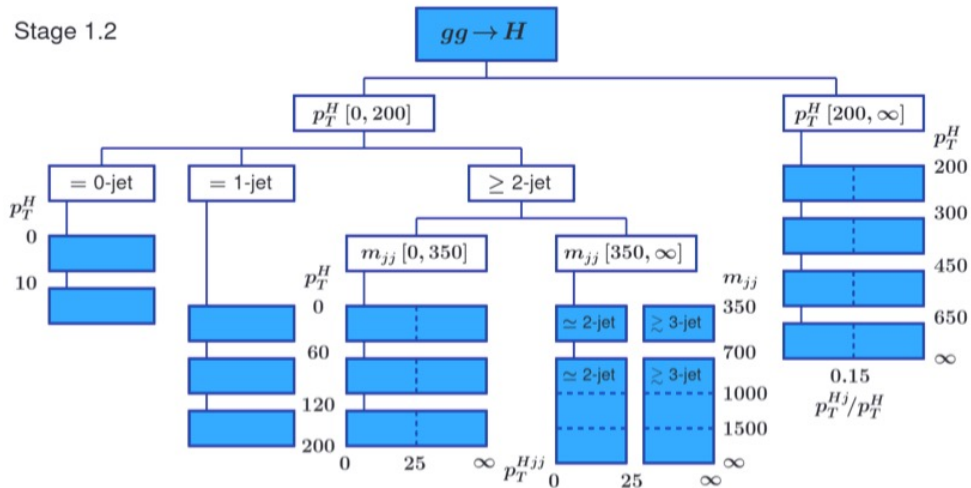


~50 related to Higgs measurements considered



# Higgs Combination and STXS

- How to constrain EFT parameters? What measurements are powerful?  
 → Simplified template cross-section (STXS)



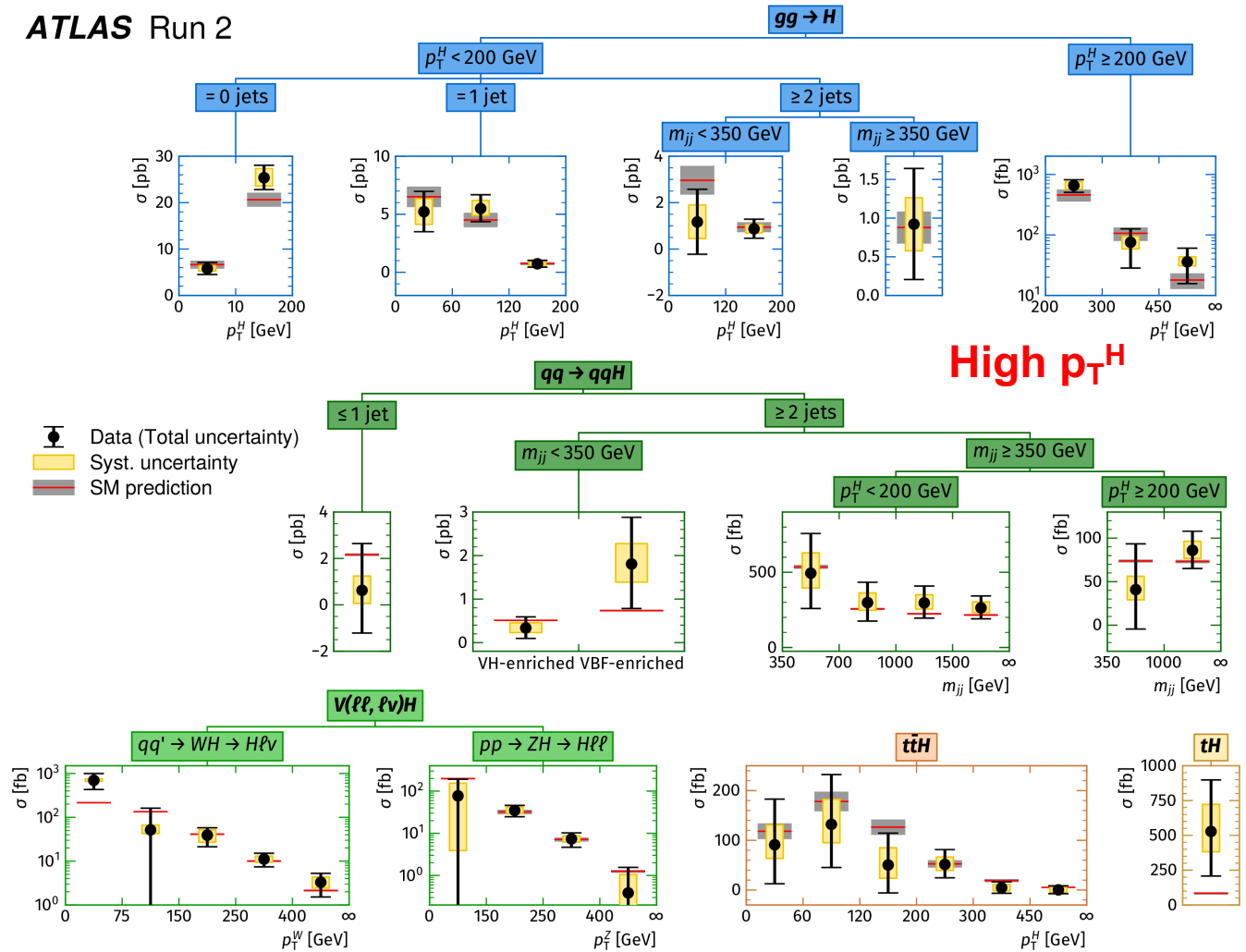
- STXS defines “bins” which are sensitive to the physics
- Bin definitions are tunable to maximize sensitivity based on the experimental precision
- Minimize theory dependences  
 → STXS measurements are sensitive to the EFT parameters

# Combined STXS measurement in LHC

Decay channel	ggF	VBF	VH	ttH/tH
$H \rightarrow \gamma\gamma$	✓	✓	✓	✓
$H \rightarrow ZZ$	✓	✓	✓	✓
$H \rightarrow WW$	✓	✓	✓*	✓*
$H \rightarrow \tau\tau$	✓	✓	✓	✓
$H \rightarrow bb$	✓	✓	✓	✓
$H \rightarrow Z\gamma$	Inclusive			
$H \rightarrow \mu\mu$	✓	✓	✓	✓

Most of channels use full Run2 dataset (~140fb<sup>-1</sup>)

ATLAS Run 2

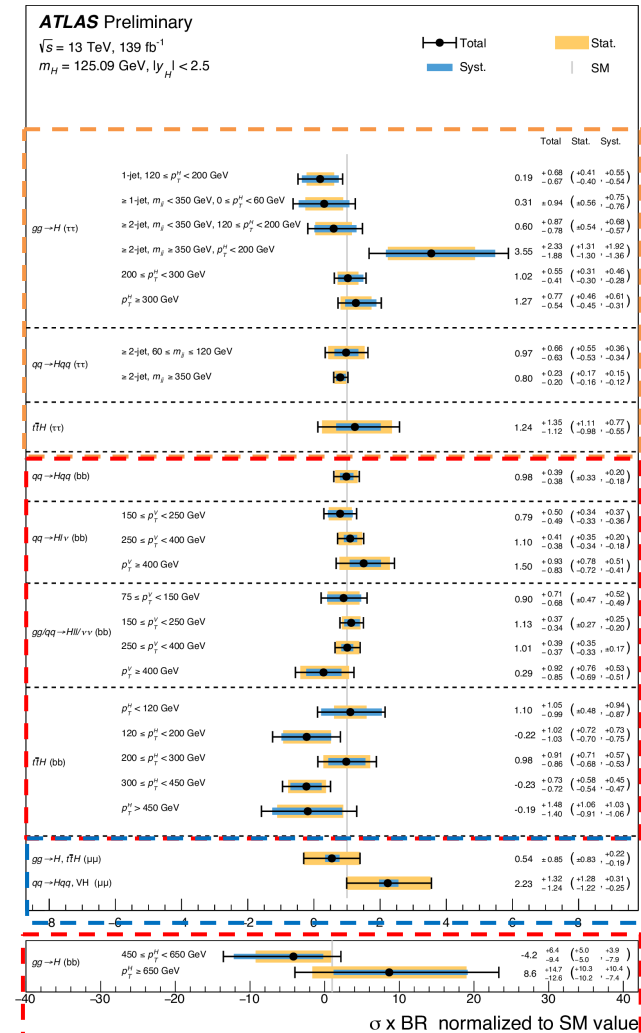
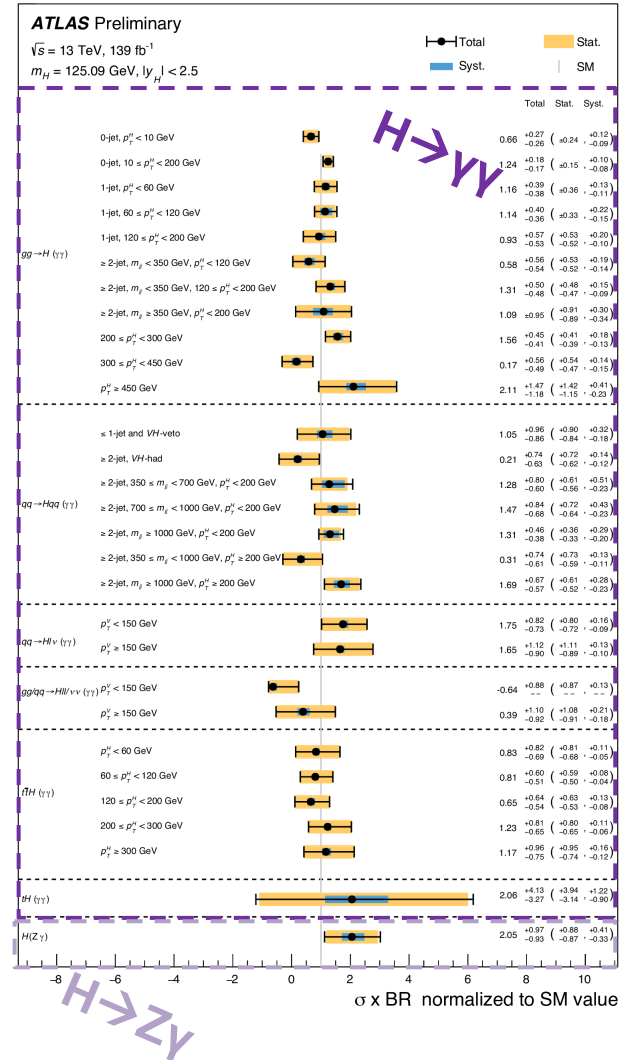
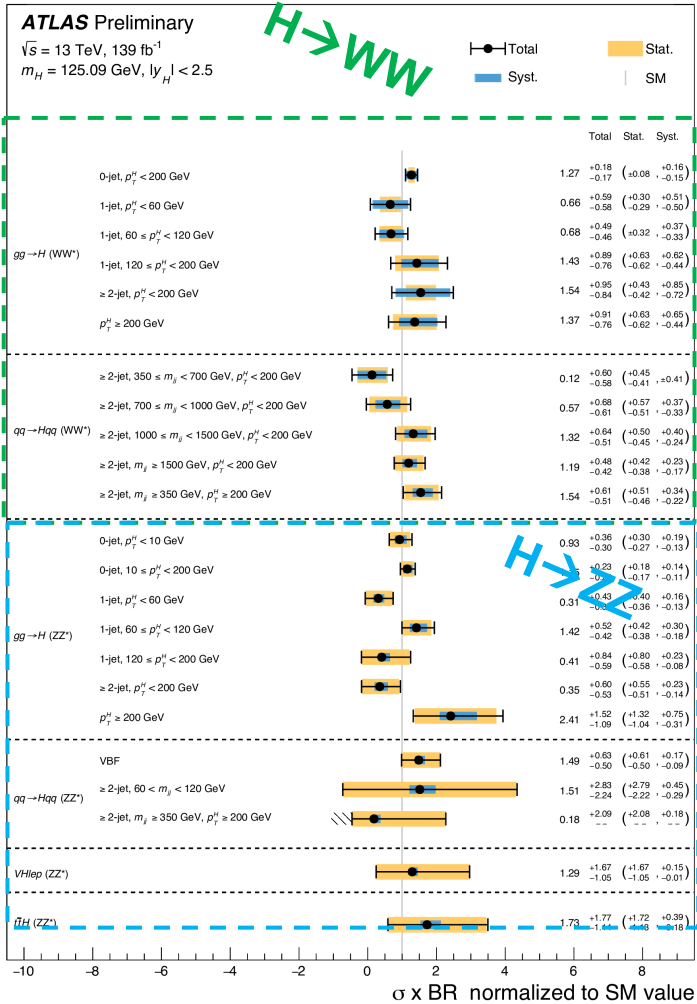


High  $p_T^H$

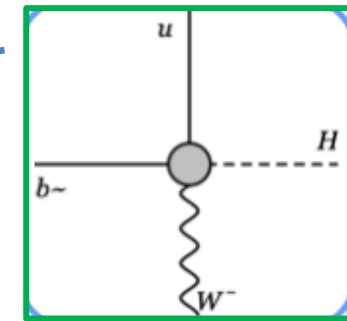
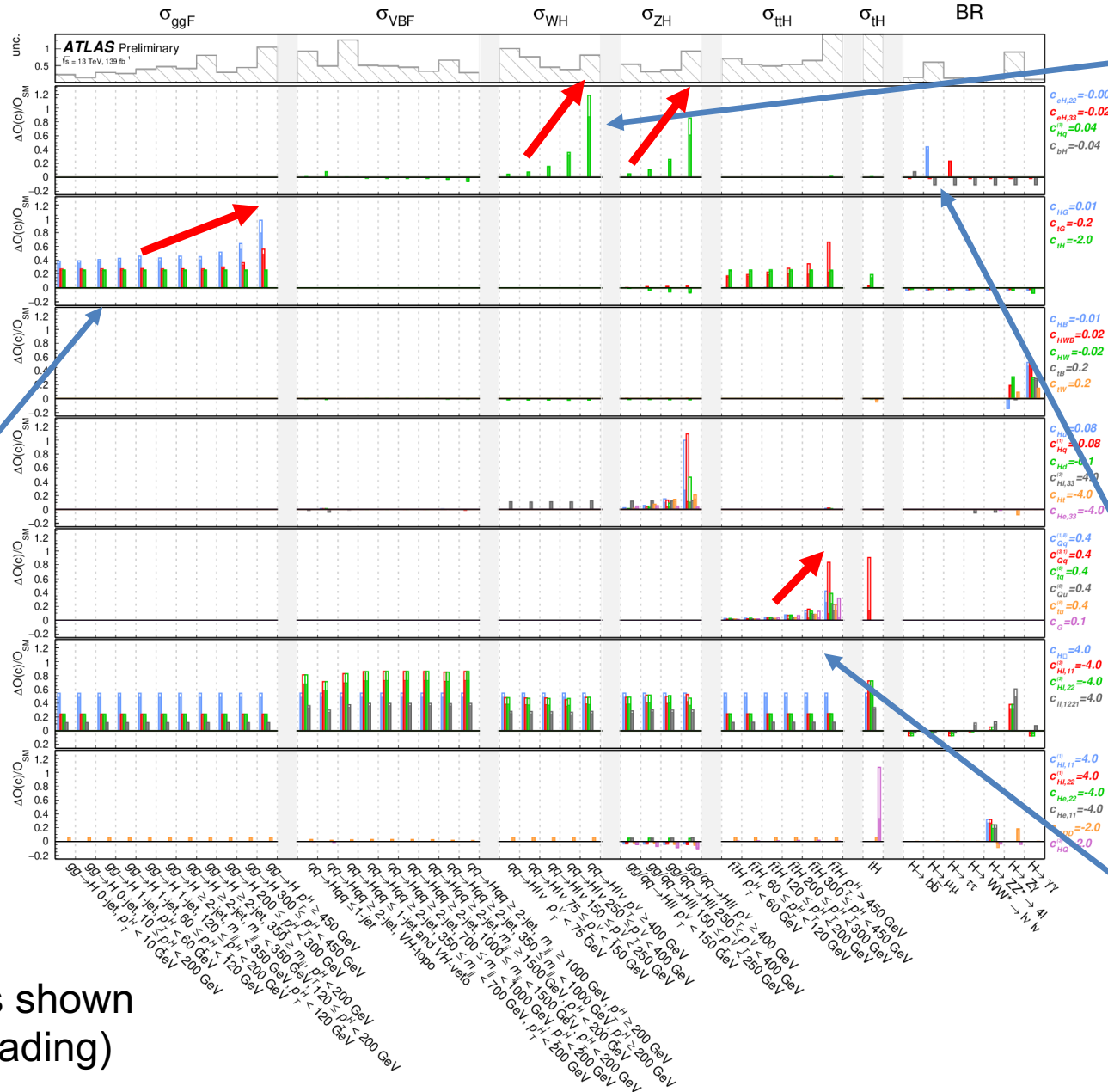
High  $p_T^H$

# STXS measurements

- Decay branching ratios are also considered in EFT interpretation



# Relative impact of EFT operators

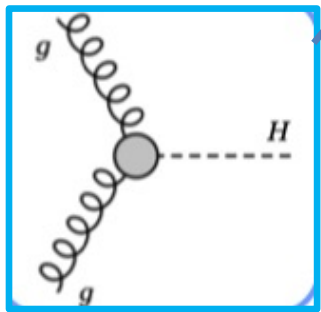


$C_{Hq}^3$ : Large BSM effects on cross-section in high  $p_T^H$  ( $p_T^H$ )

$C_{eH,22}, C_{eH,33}$   
Unique impact on  $H \rightarrow \tau\tau/\mu\mu$

$C_{Qq}$ : Large BSM effects on cross-section in high  $p_T^H$  (Huge effect in quadratic term)

$C_{HG}$ : Large BSM effects on cross-section in high  $p_T^H$

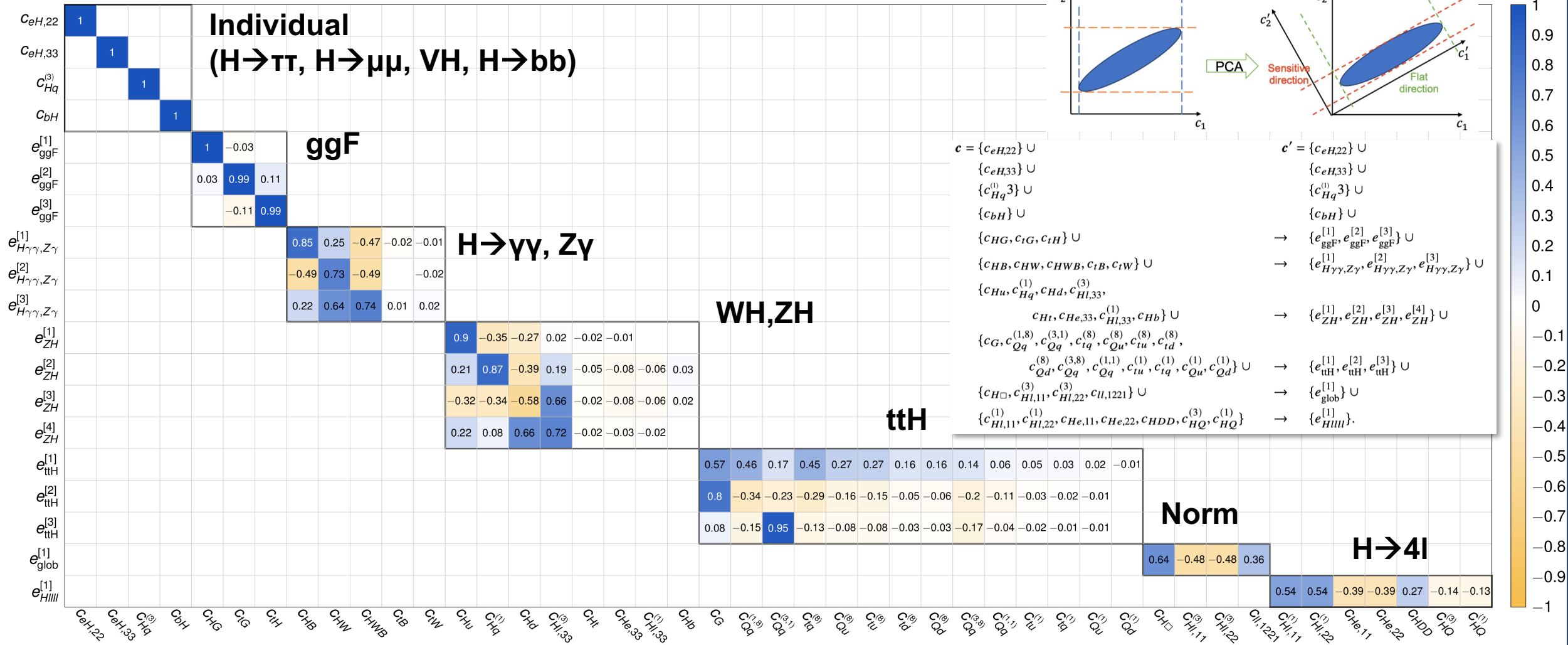


33 Wilson coefficients shown (Remaining are subleading)

# Modified basis with linear combinations

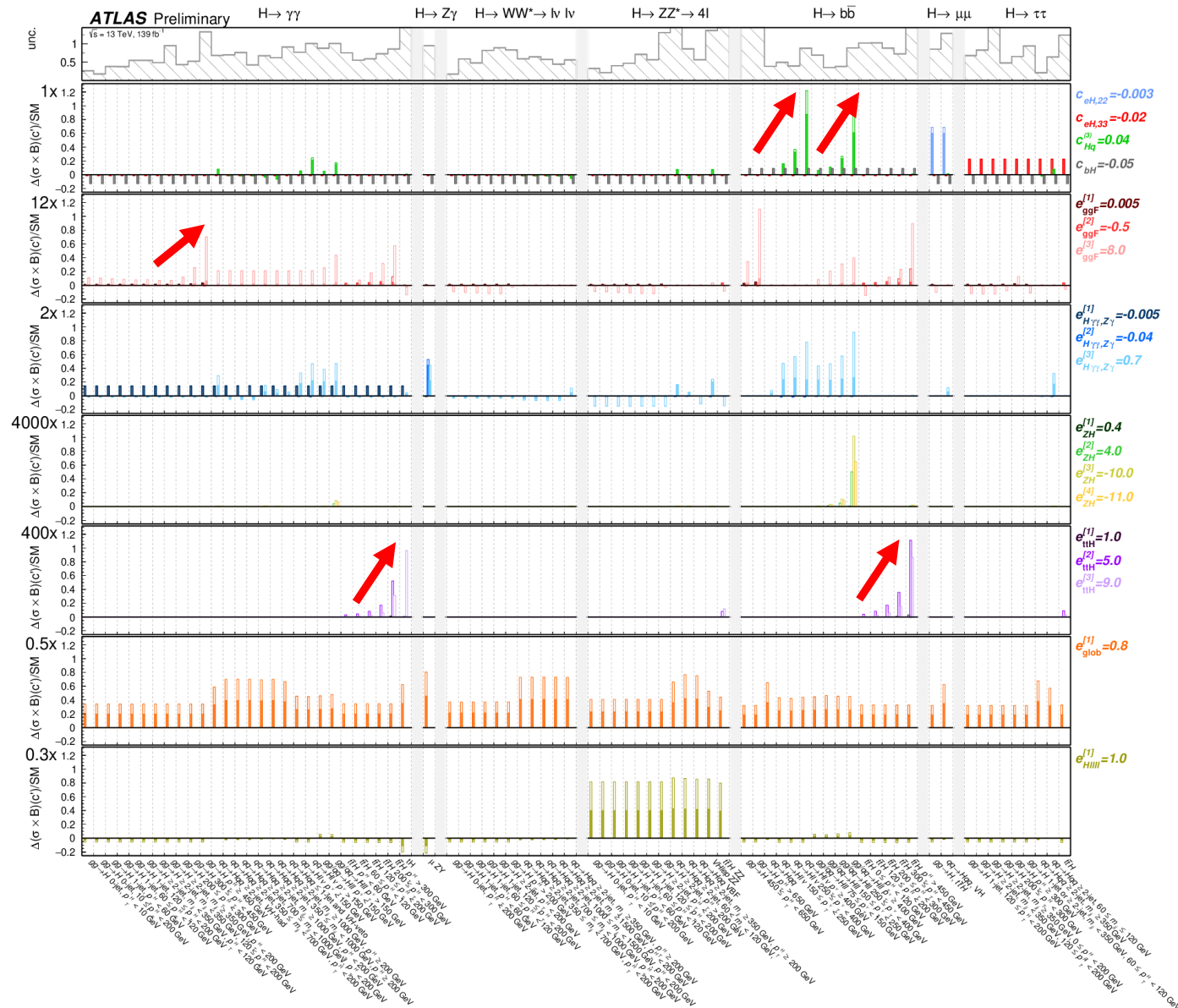
- Not able to constrain all Wilson coefficients in current sensitivity
  - ➔ Huge correlation among measurements
  - ➔ 19 parameters are reparametrized by the linear combinations

ATLAS Preliminary  $\sqrt{s}=13$  TeV,  $139 \text{ fb}^{-1}$





# Impact on reparametrized coefficients



- Extract the feature of each coefficient more clearly
- Mitigated correlations among coefficients

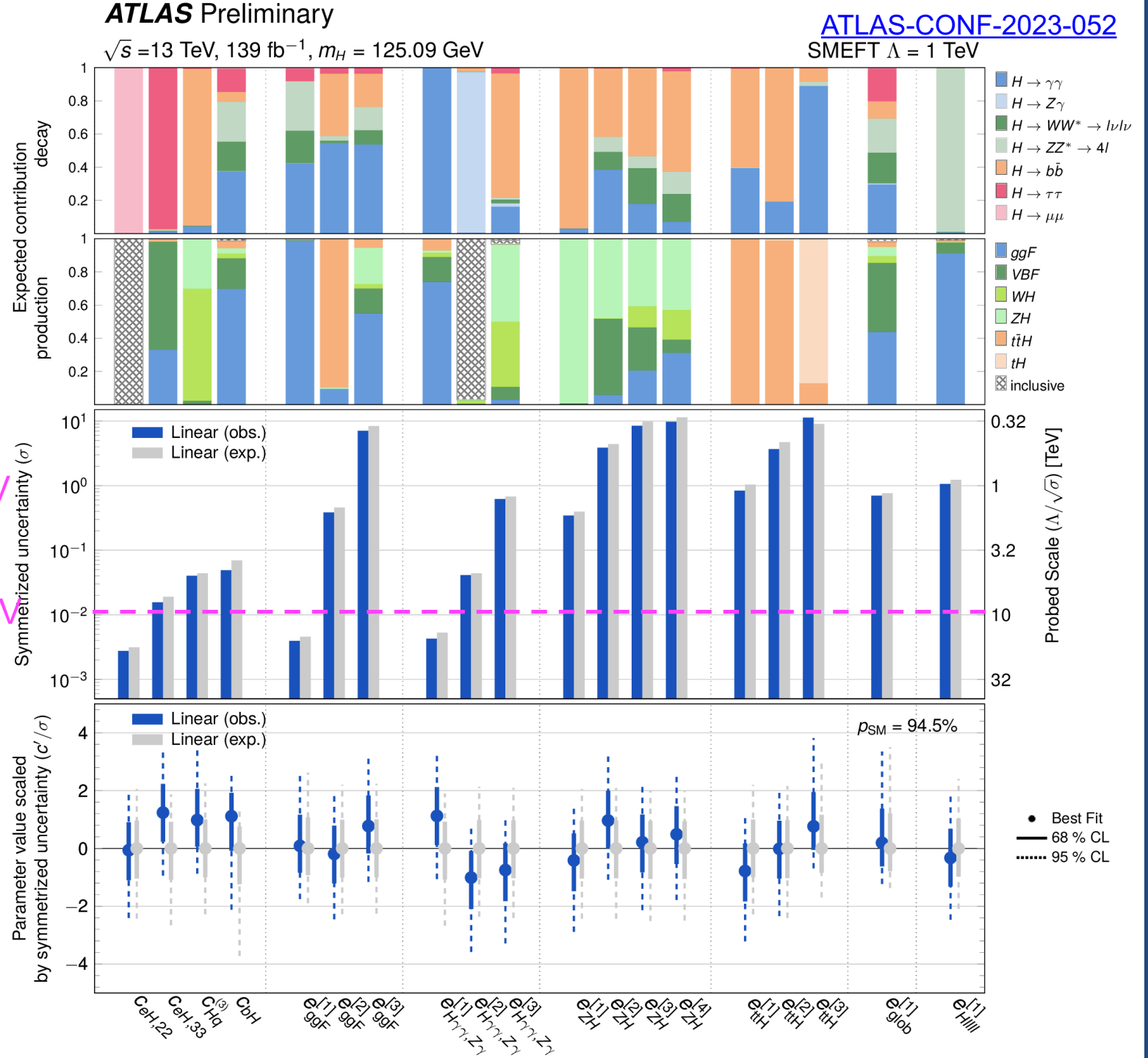


# Results

- Only linear term considered
- No strong deviation from SM (p-value 94.5%)
- O(1 TeV)-O(10 TeV) scale can be probed in the current sensitivity
- Various production/decay modes contributes EFT parameter constraint

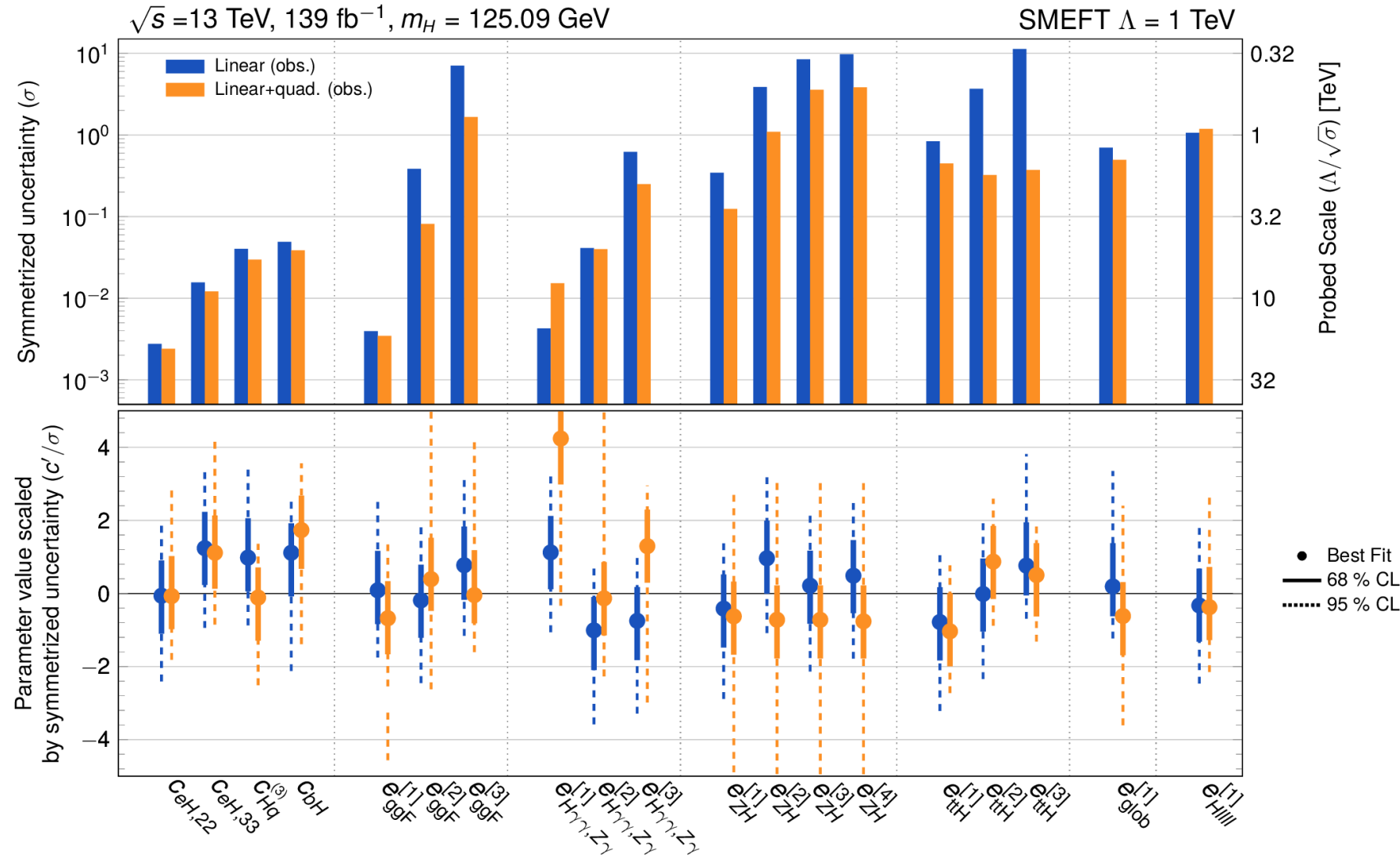
1 TeV

10 TeV



# Linear vs Linear+Quad

- In general, linear+quad provides stronger constraint (dim8 terms are important) **ATLAS Preliminary**



# Constrain UV complete BSM models

# EFT to 2HDM constraints

[Phy. Rev. D 102, 055012 \(2020\)](#)

S.Dawson, S.Homiller, & S.D. Lane

- EFT measurement can be mapped to the UV-complete models
- SMEFT constraints are interpreted 2HDM using theory paper

SMEFT parameters	Type I	Type II	Lepton-specific	Flipped
$\frac{v^2 c_{tH}}{\Lambda^2}$	$-Y_t c_{\beta-\alpha} / \tan \beta$	$-Y_t c_{\beta-\alpha} / \tan \beta$	$-Y_t c_{\beta-\alpha} / \tan \beta$	$-Y_t c_{\beta-\alpha} / \tan \beta$
$\frac{v^2 c_{bH}}{\Lambda^2}$	$-Y_b c_{\beta-\alpha} / \tan \beta$	$Y_b c_{\beta-\alpha} \tan \beta$	$-Y_b c_{\beta-\alpha} / \tan \beta$	$Y_b c_{\beta-\alpha} \tan \beta$
$\frac{v^2 c_{eH,22}}{\Lambda^2}$	$-Y_\mu c_{\beta-\alpha} / \tan \beta$	$Y_\mu c_{\beta-\alpha} \tan \beta$	$Y_\mu c_{\beta-\alpha} \tan \beta$	$-Y_\mu c_{\beta-\alpha} / \tan \beta$
$\frac{v^2 c_{eH,33}}{\Lambda^2}$	$-Y_\tau c_{\beta-\alpha} / \tan \beta$	$-Y_\tau c_{\beta-\alpha} \tan \beta$	$Y_\tau c_{\beta-\alpha} \tan \beta$	$-Y_\tau c_{\beta-\alpha} / \tan \beta$
$\frac{v^2 c_H}{\Lambda^2}$	$c_{\beta-\alpha}^2 M_A^2 / v^2$	$c_{\beta-\alpha}^2 M_A^2 / v^2$	$c_{\beta-\alpha}^2 M_A^2 / v^2$	$c_{\beta-\alpha}^2 M_A^2 / v^2$

- $\Lambda \gg v$  in EFT  $\rightarrow \cos(\beta-\alpha) \rightarrow 0$  (valid near alignment limit)

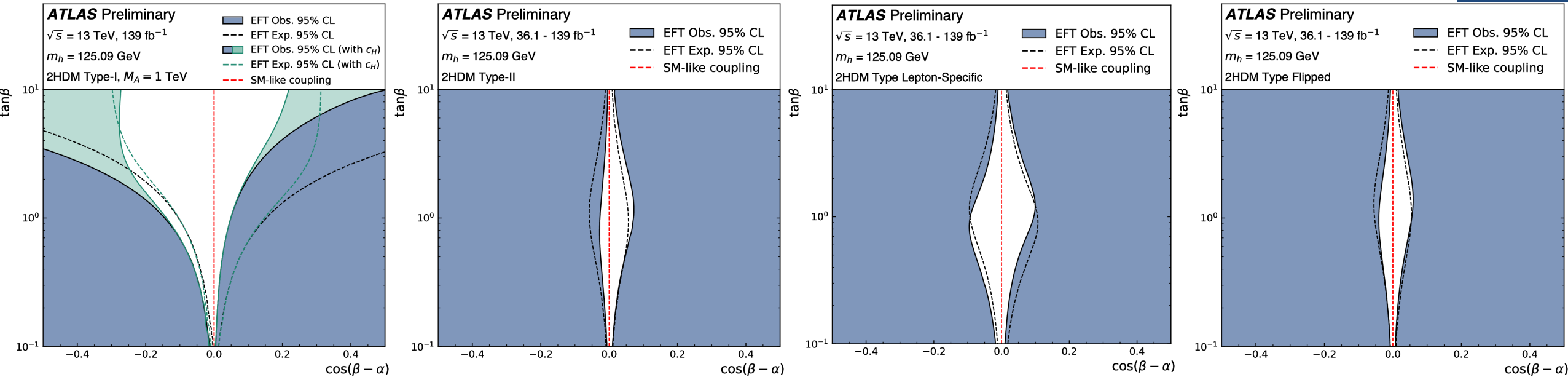
$$Y_{b,t,\mu,\tau} = \frac{\sqrt{2}m_i}{v} (\sim \text{SM})$$

- Constrain in  $\tan\beta$  vs  $\cos(\beta-\alpha)$

- use only single Higgs for  $c_H$  constraint (no HH direct measurement included)

# 2HDM constraints

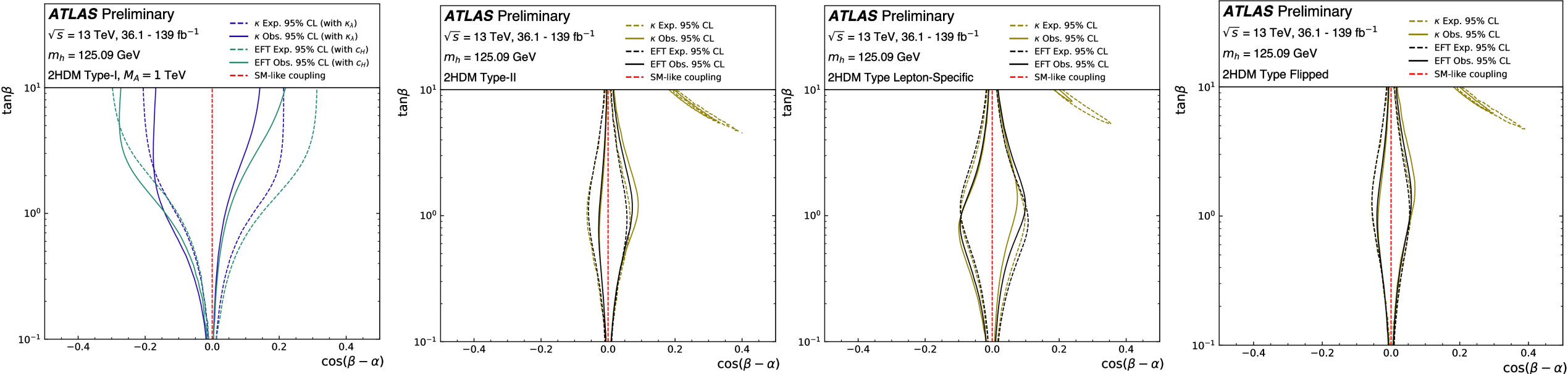
- Only linear expansion for EFT



- No surprise and large 2HDM parameter spaces are excluded
- $c_H$  constraint is included in Type-I

# 2HDM constraint

- Compare with interpretation with coupling measurement(k-framework)



- Coupling measurement gives similar (slightly better) constraint
  - Missing dim-8 operators
  - No petal structure in EFT (no 2<sup>nd</sup> minimum)

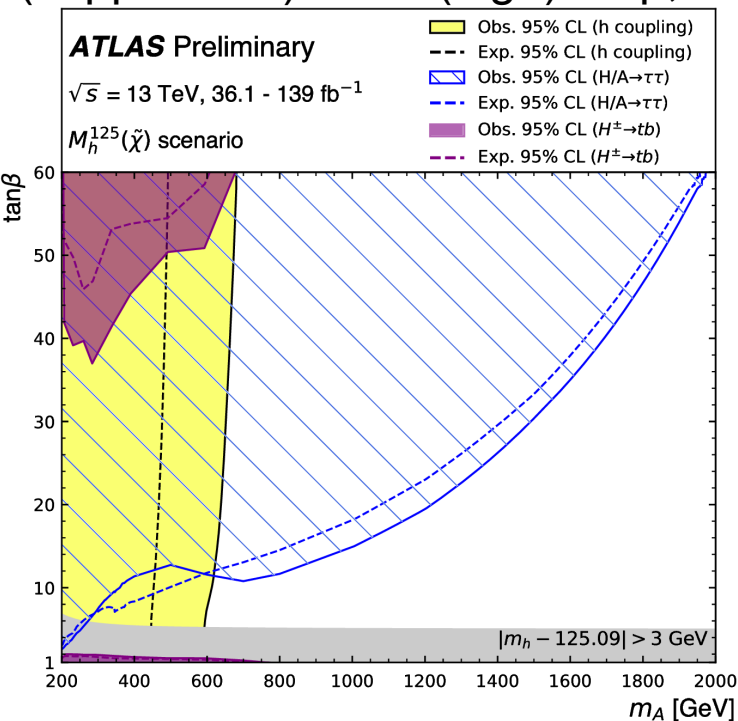


# Constraints on MSSM

- Not use EFT interpretation results but Constraint from STXS measurements
- Several MSSM benchmark scenarios are considered (more in backup)

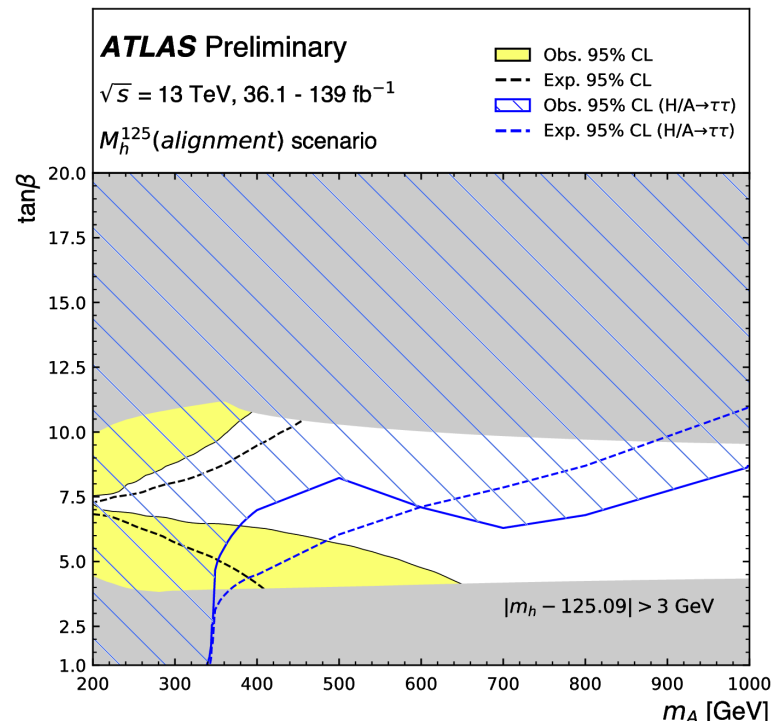
## $M_H^{125}(\bar{\chi})$ scenario:

All charginos and neutralinos are relatively light with significant higgsino-gaugio mixing  
 $H \rightarrow bb$ ,  $H \rightarrow \gamma\gamma$  decay enhanced (suppressed) at low(high)  $\tan\beta$ ,



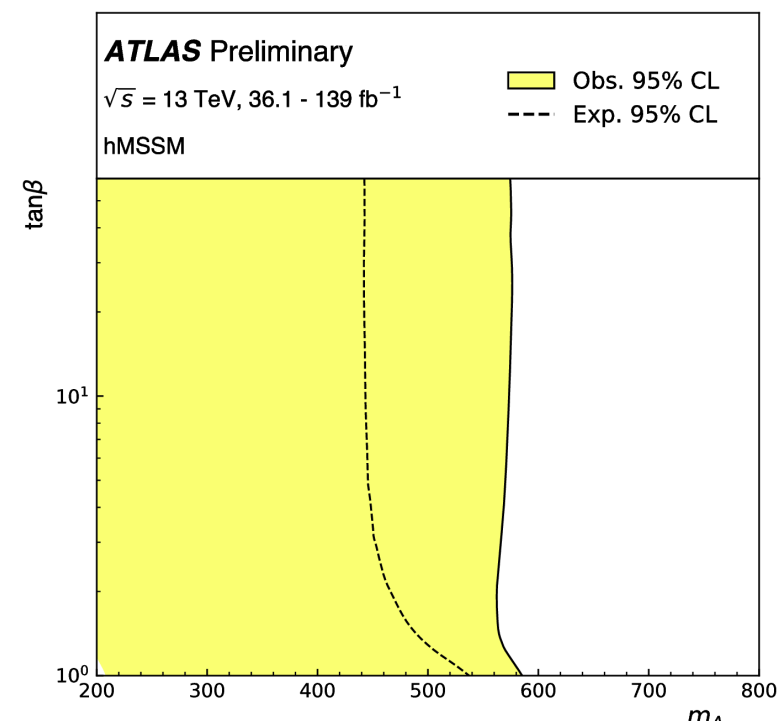
## $M_H^{125}(\text{alignment})$ scenario:

Alignment without decoupling scenario one of CP even scalars have SM-like couplings  
 $\tan\beta \sim 7 \rightarrow$  nearly SM-like



## hMSSM scenario:

$m_H$  is 125.09 GeV with radiative correction from stop-top sector, determine  $\alpha$ ,  $m_H$ ,  $m_{H\pm}$ , couplings

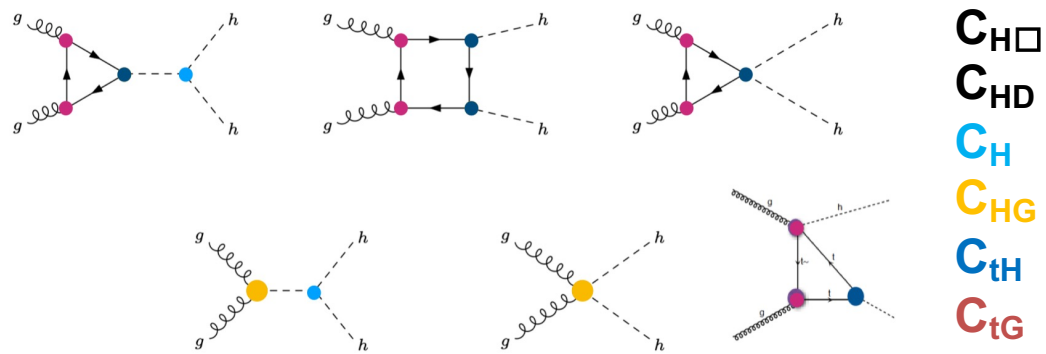


# EFT in DiHiggs

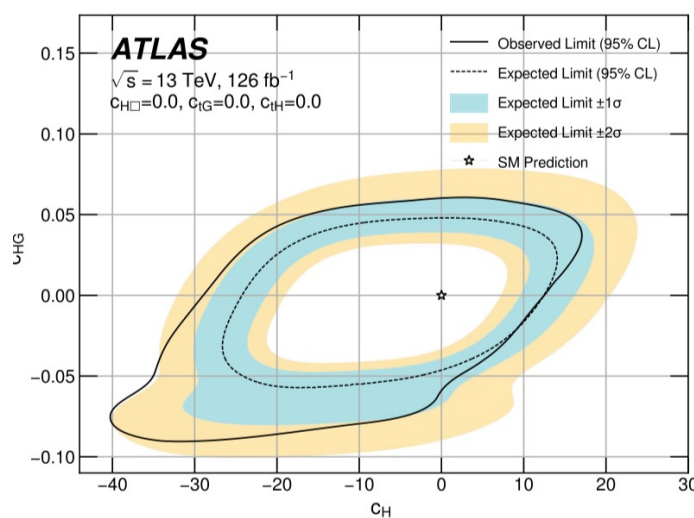
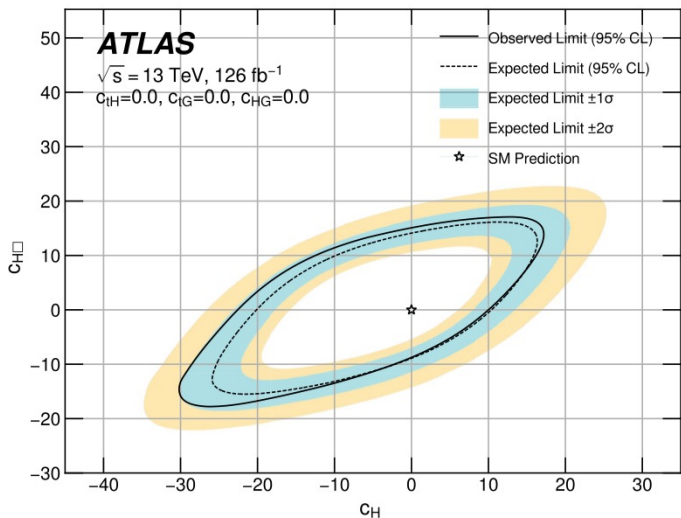
arXiv:2310.12301

2023/12/20

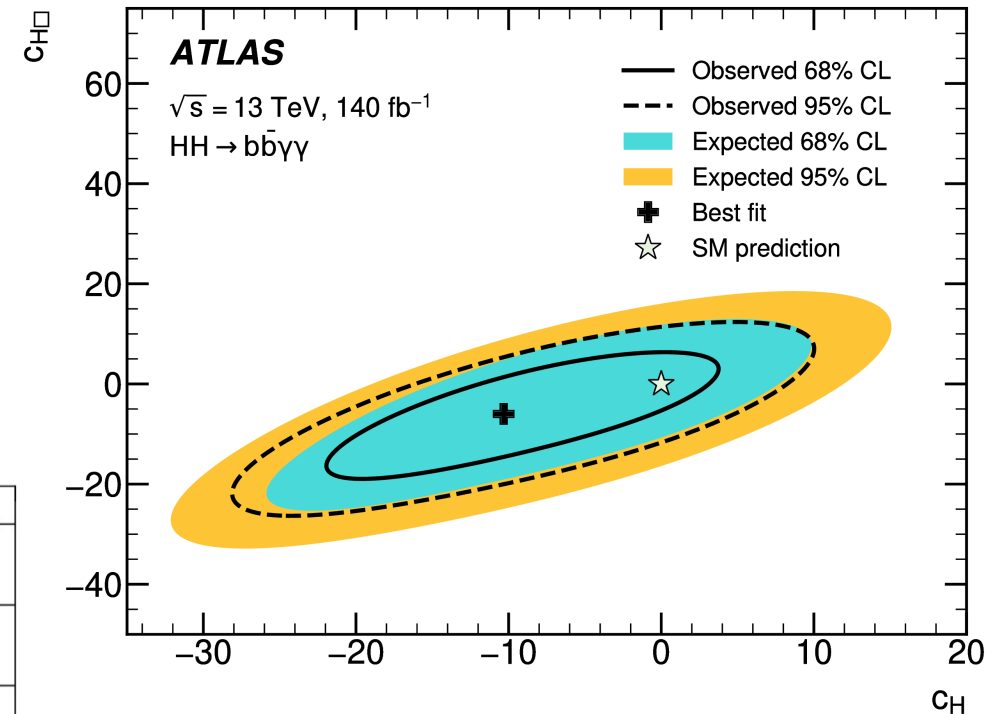
- DiHiggs production is sensitive to the EFT parameters
  - Both interpretations for HEFT and SMEFT
  - 1D or 2D parameter scan (other parameters fixed to zero)



**HH→4b**



**HH→bbγγ**



Wilson coefficient	95% CL Observed	95% CL Expected
$C_H$	[-14.4, 6.2]	[-16.8, 9.7]
$C_{H\Box}$	[-9.4, 10.2]	[-12.4, 13.7]

ILC meeting

# Constraint on CP-odd operators

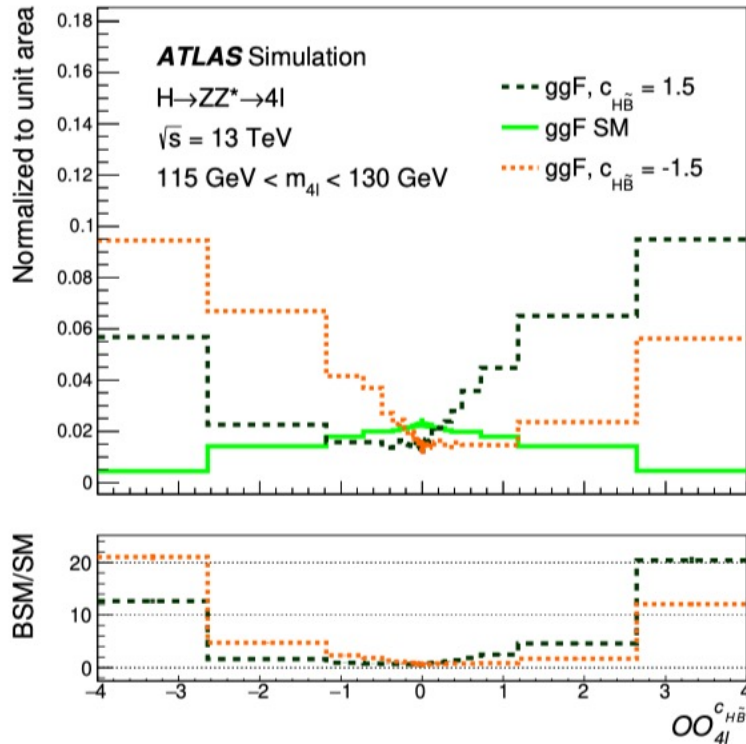
- $H \rightarrow ZZ \rightarrow 4l$  constraint with CP-odd EFT parameters
- “Optimal observable” using the interference term of SM and BSM

$$OO = \frac{2\Re(\mathcal{M}_{SM}^* \mathcal{M}_{BSM})}{|\mathcal{M}_{SM}|^2}$$

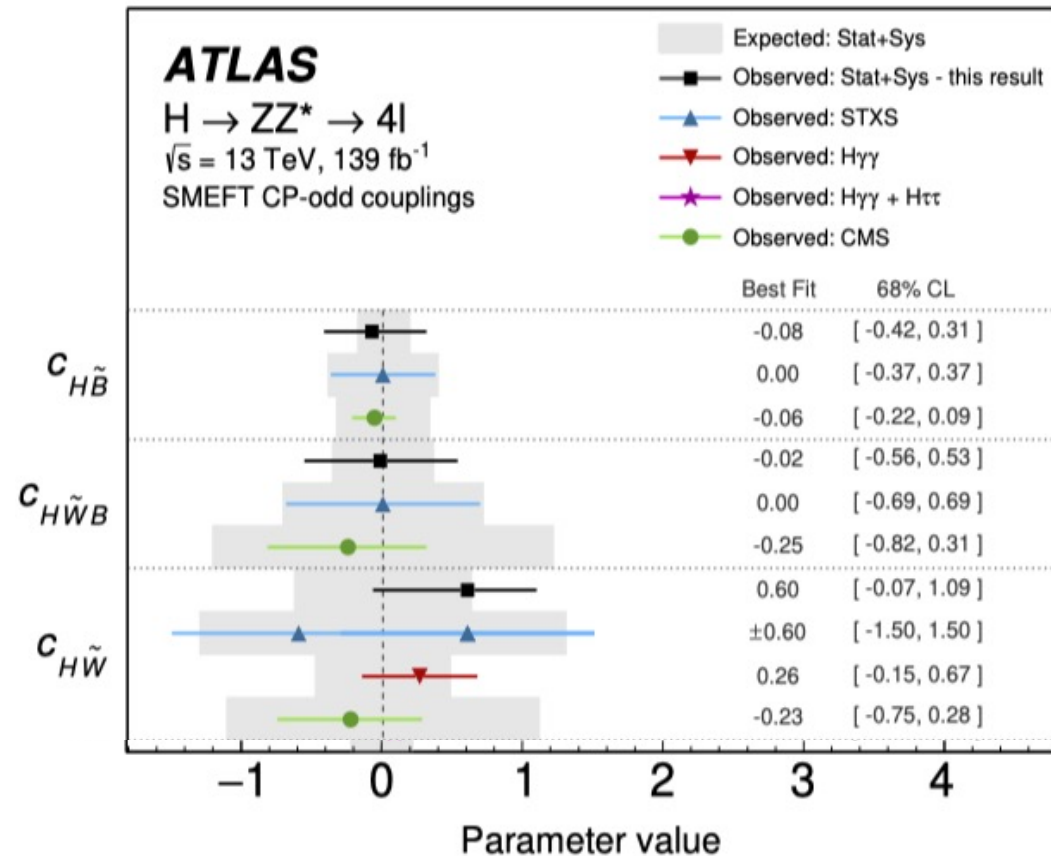
Operator	Structure	Coupling
Warsaw Basis		
$O_{\Phi\tilde{W}}$	$\Phi^\dagger \Phi \tilde{W}_{\mu\nu}^I W^{\mu\nu I}$	$c_{H\tilde{W}}$
$O_{\Phi\tilde{W}B}$	$\Phi^\dagger \tau^I \Phi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$c_{H\tilde{W}B}$
$O_{\Phi\tilde{B}}$	$\Phi^\dagger \Phi \tilde{B}_{\mu\nu} B^{\mu\nu}$	$c_{H\tilde{B}}$

Asymmetric  
CP-odd

Symmetric  
CP-even



(b)  $OO_{4l}^{c_{H\tilde{B}}}$



# Other EFT interpretations

- Combined EFT interpretations with other EW measurements

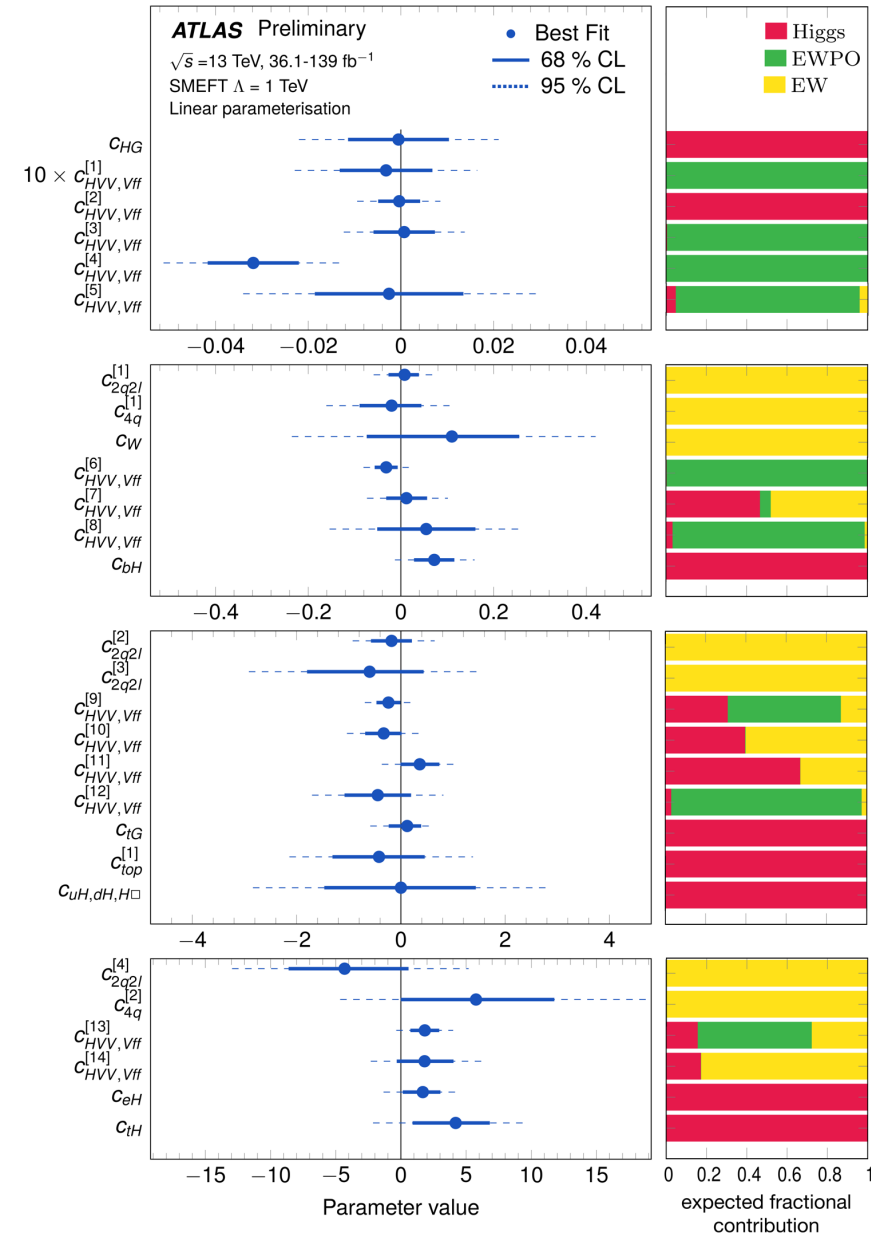
- Differential cross section measurements with diboson and VBF Z production

(All EW measurements are not included yet)

Process	Important phase space requirements	Observable	$\mathcal{L}$ [fb <sup>-1</sup> ]	Ref.
$pp \rightarrow e^\pm \nu \mu^\mp \nu$	$m_{\ell\ell} > 55 \text{ GeV}, p_T^{\text{jet}} < 35 \text{ GeV}$	$p_T^{\text{lead. lep.}}$	36	[19]
$pp \rightarrow \ell^\pm \nu \ell^+ \ell^-$	$m_{\ell\ell} \in (81, 101) \text{ GeV}$	$m_T^{WZ}$	36	[20]
$pp \rightarrow \ell^+ \ell^- \ell^+ \ell^-$	$m_{4\ell} > 180 \text{ GeV}$	$m_{Z2}$	139	[21]
$pp \rightarrow \ell^+ \ell^- jj$	$m_{jj} > 1000 \text{ GeV}, m_{\ell\ell} \in (81, 101) \text{ GeV}$	$\Delta\phi_{jj}$	139	[22]

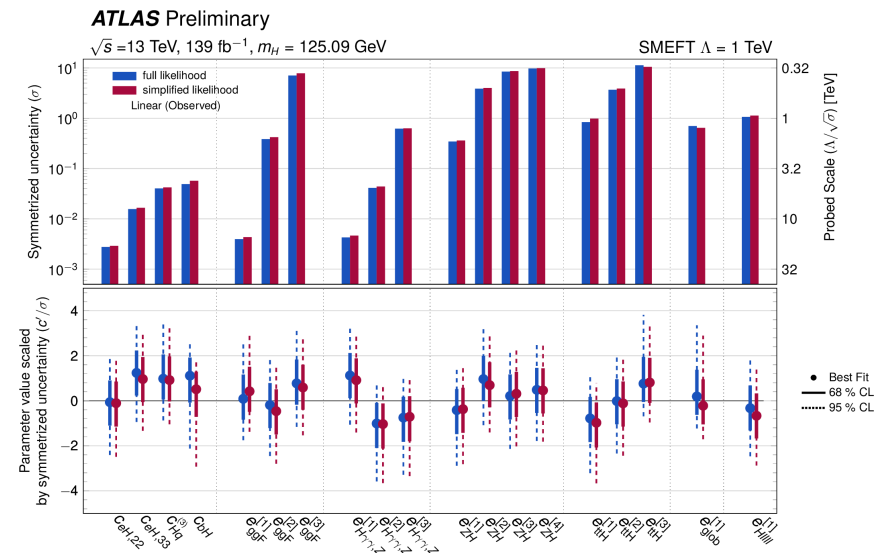
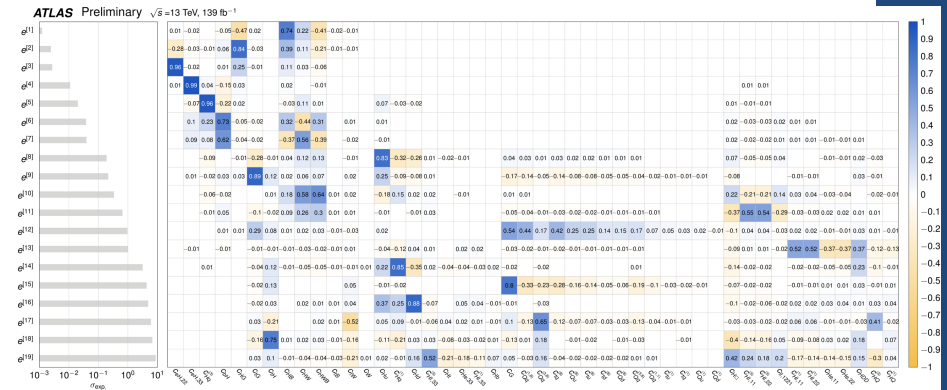
- Electroweak precision observable (mainly from SLC and LEP)

Observable	Measurement	Prediction	Ratio
$\Gamma_Z$ [MeV]	$2495.2 \pm 2.3$	$2495.7 \pm 1$	$0.9998 \pm 0.0010$
$R_\ell^0$	$20.767 \pm 0.025$	$20.758 \pm 0.008$	$1.0004 \pm 0.0013$
$R_c^0$	$0.1721 \pm 0.0030$	$0.17223 \pm 0.00003$	$0.999 \pm 0.017$
$R_b^0$	$0.21629 \pm 0.00066$	$0.21586 \pm 0.00003$	$1.0020 \pm 0.0031$
$A_{\text{FB}}^{0,\ell}$	$0.0171 \pm 0.0010$	$0.01718 \pm 0.00037$	$0.995 \pm 0.062$
$A_{\text{FB}}^{0,c}$	$0.0707 \pm 0.0035$	$0.0758 \pm 0.0012$	$0.932 \pm 0.048$
$A_{\text{FB}}^{0,b}$	$0.0992 \pm 0.0016$	$0.1062 \pm 0.0016$	$0.935 \pm 0.021$
$\sigma_{\text{had}}^0$ [pb]	$41488 \pm 6$	$41489 \pm 5$	$0.99998 \pm 0.00019$



# Summary

- No BSM particle found (so far) at LHC 🙄
- Higgs, EW, top precise measurements provide valuable EFT interpretations and stronger constraints
- Still can improve EFT formalism for both theory and experimental sides
  - Principal component analysis
  - Global EFT interpretation
  - Provide simplified likelihood in HEPdata
  - Dim-8 calculation
- Showed EFT interpretation is usable to constrain UV complete model  
 → Beyond 2HDM scenarios?



# Backup



# EFT formulation

- SMEFT cross sections are calculated by LO diagrams (NLO QCD for loop, NLO QED for  $H \rightarrow \gamma\gamma, Z\gamma$ )
- Higher order calculation computed by scaling SM cross section assuming the same relative effect on  $\sigma_{\text{int}}$  and  $\sigma_{\text{BSM}}$  as on  $\sigma_{\text{SM}}$

$$\sigma_{\text{SMEFT}} = \sigma_{\text{SM}}^{((\text{N})\text{N})\text{NLO}} \times \left( 1 + \frac{\sigma_{\text{int}}^{(\text{N})\text{LO}}}{\sigma_{\text{SM}}^{(\text{N})\text{LO}}} + \frac{\sigma_{\text{BSM}}^{(\text{N})\text{LO}}}{\sigma_{\text{SM}}^{(\text{N})\text{LO}}} \right)$$

- Branching ratio effects are taken into account

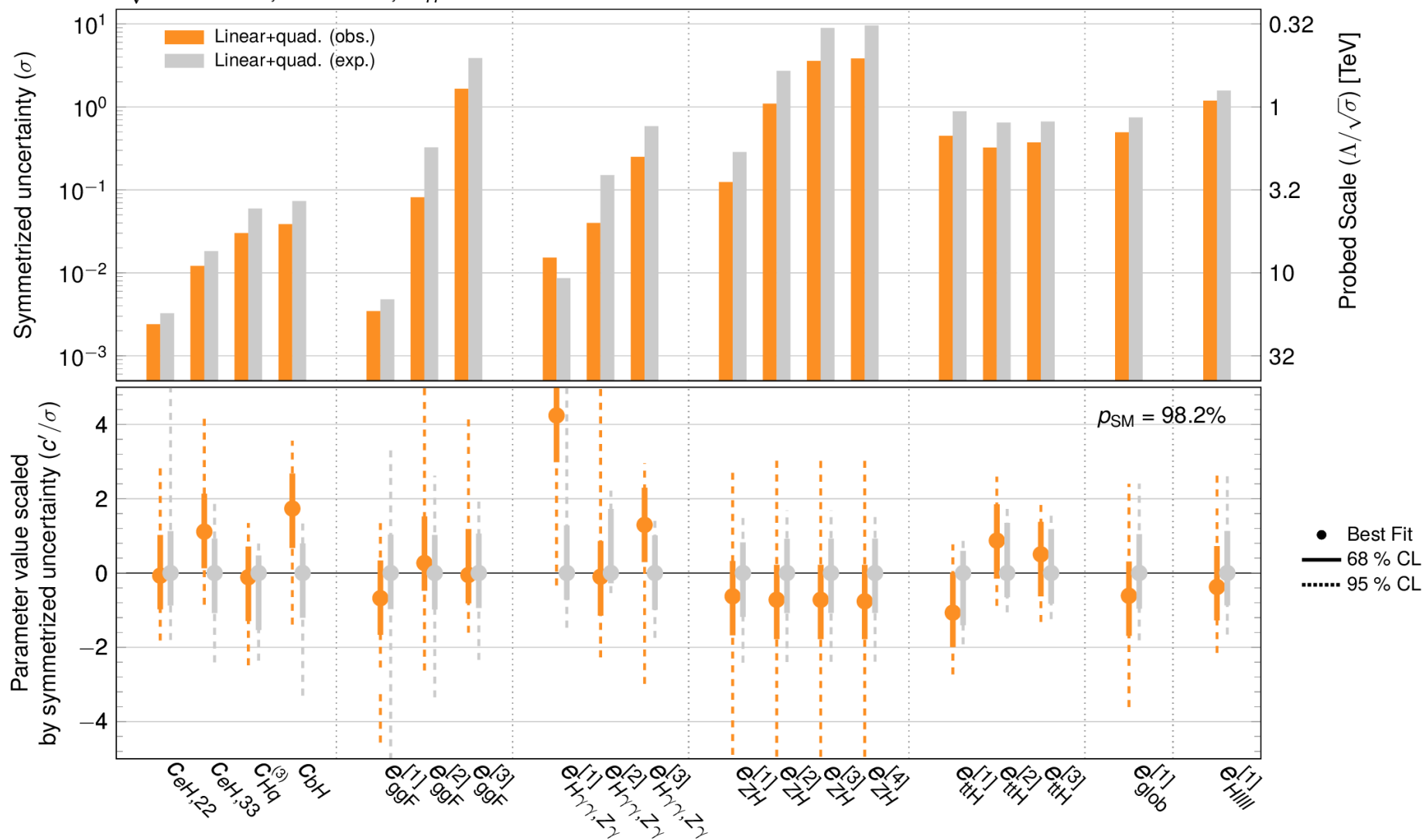
$$(\sigma \times B)_{\text{SMEFT}}^{i,k',H \rightarrow X} = (\sigma \times B)_{\text{SM},(\text{N}(\text{N}))\text{NLO}}^{i,k',H \rightarrow X} \left( 1 + \frac{\sigma_{\text{int},(\text{N})\text{LO}}^{i,k'}}{\sigma_{\text{SM},(\text{N})\text{LO}}^{i,k'}} + \frac{\sigma_{\text{BSM},(\text{N})\text{LO}}^{i,k'}}{\sigma_{\text{SM},(\text{N})\text{LO}}^{i,k'}} \right) \left( \frac{1 + \frac{\Gamma_{\text{int}}^{H \rightarrow X}}{\Gamma_{\text{SM}}^{H \rightarrow X}} + \frac{\Gamma_{\text{BSM}}^{H \rightarrow X}}{\Gamma_{\text{SM}}^{H \rightarrow X}}}{1 + \frac{\Gamma_{\text{int}}^H}{\Gamma_{\text{SM}}^H} + \frac{\Gamma_{\text{BSM}}^H}{\Gamma_{\text{SM}}^H}} \right)$$

# Linear+Quad

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}, m_H = 125.09 \text{ GeV}$

SMEFT  $\Lambda = 1 \text{ TeV}$



# Constraints on 2HDM

2HDM: one of natural extension of Higgs sector

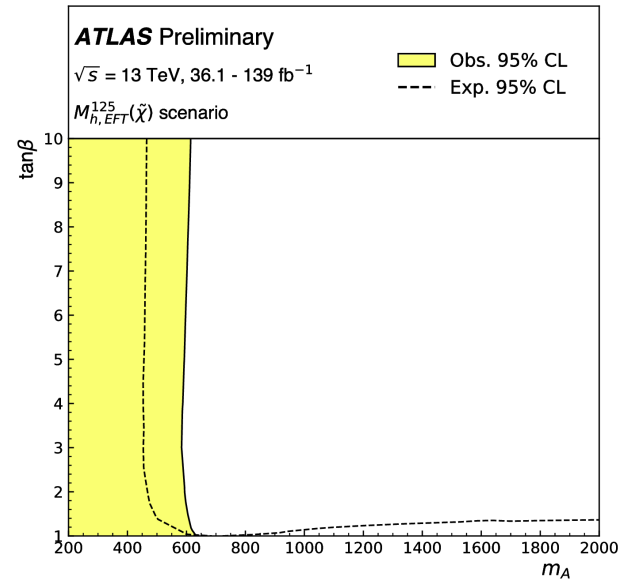
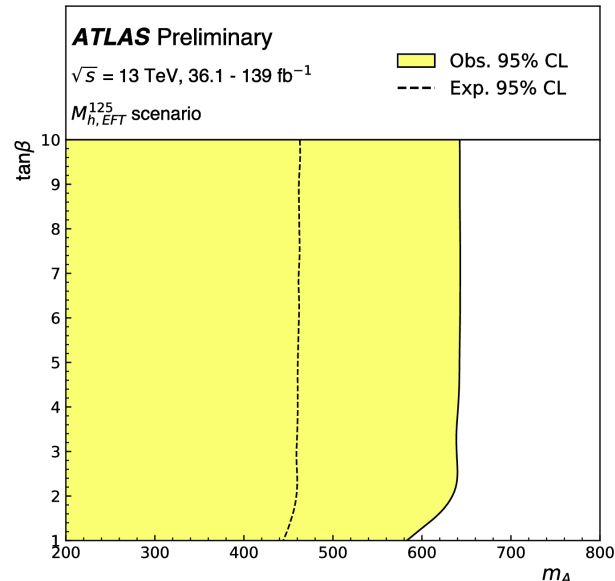
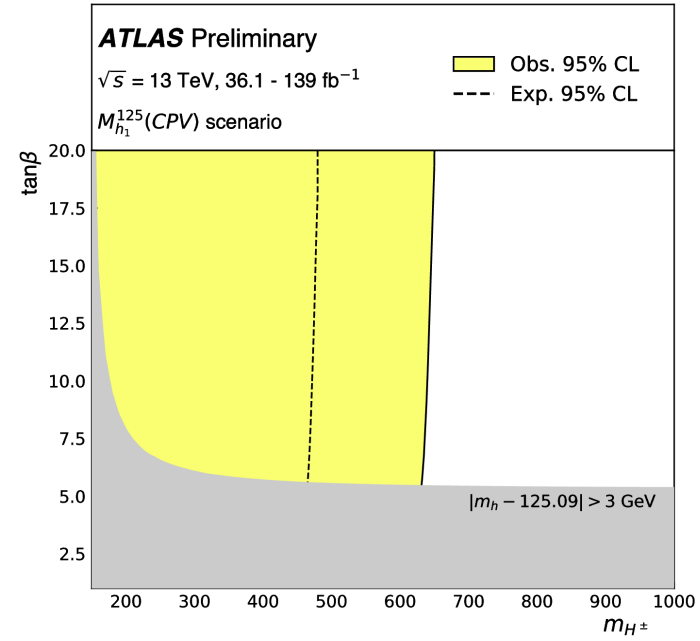
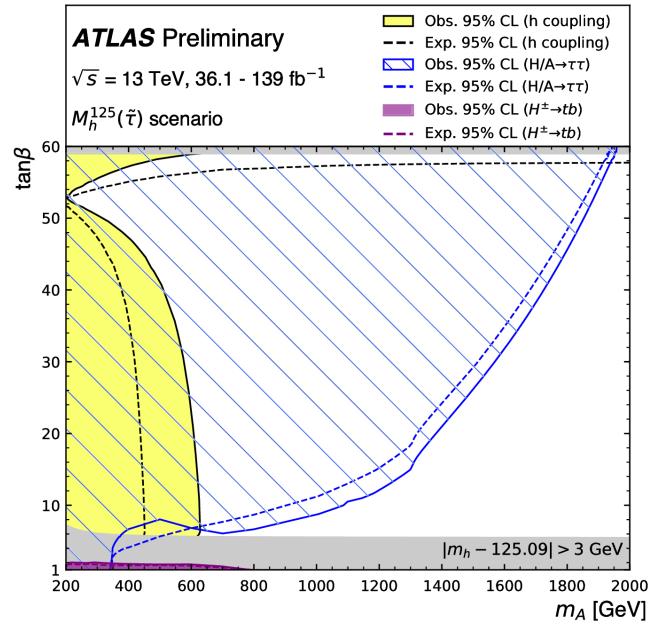
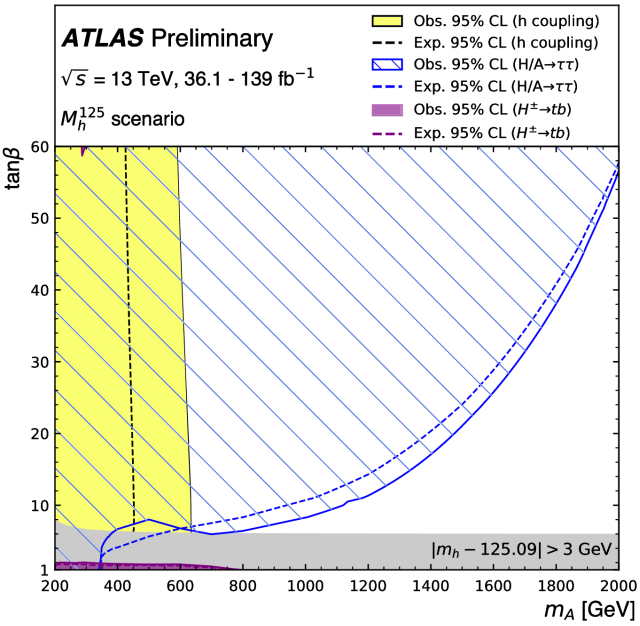
- parameter:  $h, H, A, H^\pm$ , CP-even mixing angle  $\alpha$ , two doublet mixing  $\beta$ ,  $m_{12}$

Coupling	Type I	Type II	Lepton-specific	Flipped
$u, c, t$		$s_{\beta-\alpha} + c_{\beta-\alpha}/\tan\beta$		
$d, s, b$	$s_{\beta-\alpha} + c_{\beta-\alpha}/\tan\beta$	$s_{\beta-\alpha} - c_{\beta-\alpha} \times \tan\beta$	$s_{\beta-\alpha} + c_{\beta-\alpha}/\tan\beta$	$s_{\beta-\alpha} - c_{\beta-\alpha} \times \tan\beta$
$e, \mu, \tau$	$s_{\beta-\alpha} + c_{\beta-\alpha}/\tan\beta$	$s_{\beta-\alpha} - c_{\beta-\alpha} \times \tan\beta$	$s_{\beta-\alpha} - c_{\beta-\alpha} \times \tan\beta$	$s_{\beta-\alpha} + c_{\beta-\alpha}/\tan\beta$
$W, Z$		$s_{\beta-\alpha}$		
$H$		$s_{\beta-\alpha}^3 + \left(3 - 2\frac{\bar{m}^2}{m_h^2}\right) c_{\beta-\alpha}^2 s_{\beta-\alpha} + 2 \cot(2\beta) \left(1 - \frac{\bar{m}^2}{m_h^2}\right) c_{\beta-\alpha}^3$		

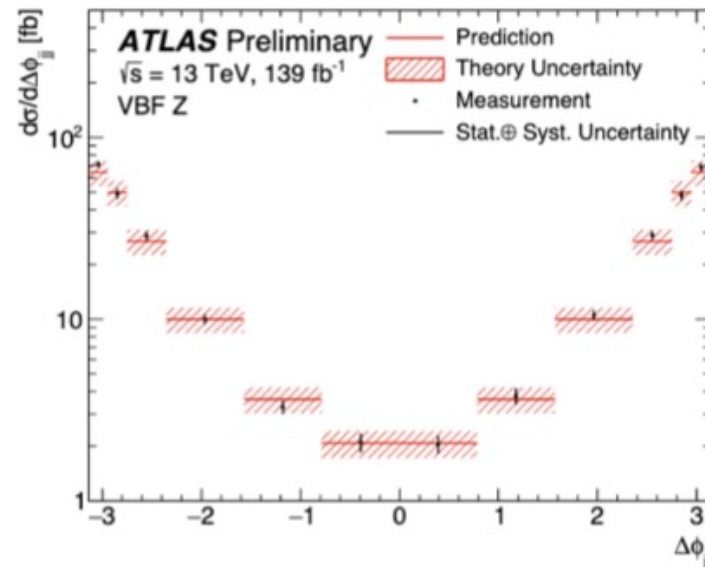
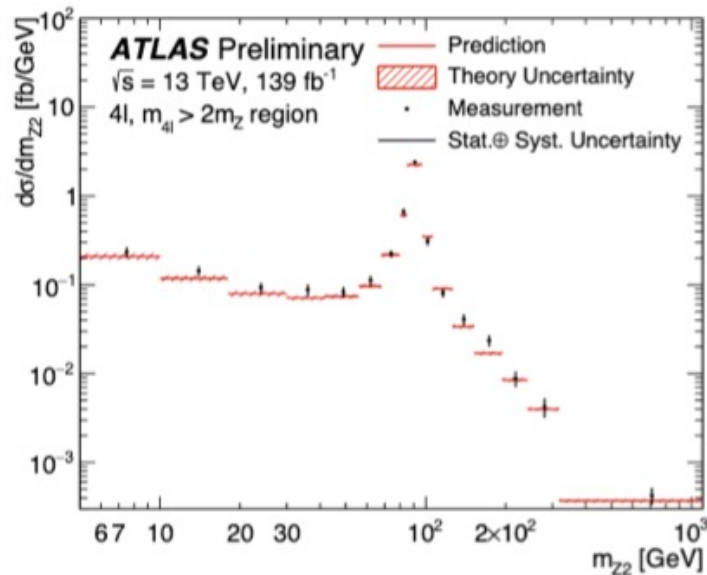
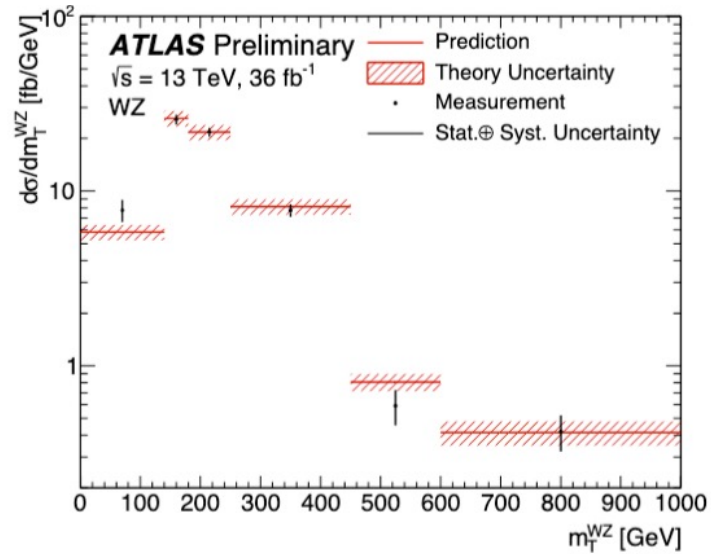
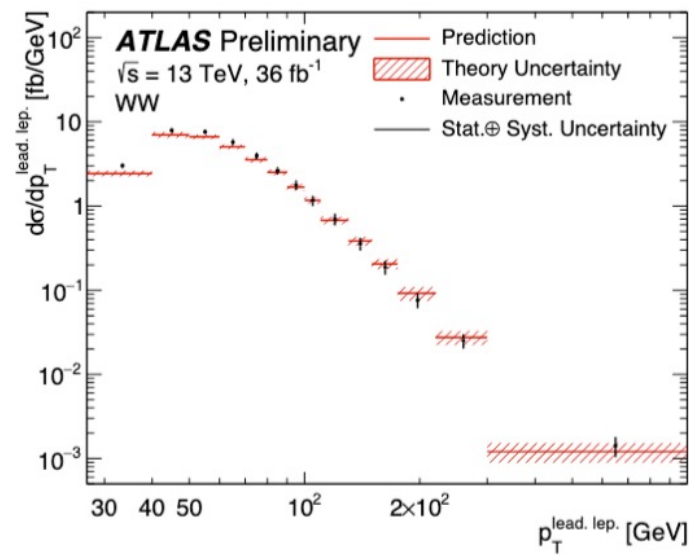
$$v_1^2 + v_2^2 = v^2, \tan\beta = \frac{v_2}{v_1}, \bar{m}^2 = \frac{m_{12}^2}{\sin\beta \cos\beta} = m_A^2 + \lambda_5 v^2$$

$$\text{Assuming, } \lambda_5 v^2 \ll m_A^2, \bar{m} = m_A = 1\text{TeV} (\lambda_5 = 0)$$

# Constraints on MSSM



# Global EFT EW inputs



# Ref

- <https://indico.cern.ch/event/1296757/timetable/>
- <https://indico.cern.ch/event/1276727/timetable/?view=standard>