

# Disentangling $Z$ decays to light quarks

Krzysztof Mękała  
Uni. of Warsaw & DESY

20.12.2023

The Standard Model is the most successful theory describing particle interactions so far but it is not “self-explanatory”:

- 19 free parameters?
- neutrino masses?
- left-right (a)symmetry?

The Standard Model is the most successful theory describing particle interactions so far but it is not “self-explanatory”:

- 19 free parameters?
- neutrino masses?
- left-right (a)symmetry?

*New Physics* is needed; if it is too heavy to be observed directly, **precision measurements** are crucial to finding any hints of its existence.

- Precision electroweak observables may shed light on the existence of BSM physics, e.g. extra gauge bosons.
- Possible discrepancies in the values of the  $Z$  couplings to *heavy* quarks can be measured by proper identification of the quark flavour (*b*-tagging, *c*-tagging, ...).
- But so far, *u*- and *d*-tagging are not available... how to proceed?

# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W, \quad a_f = 2I_{3,f}.$$

# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W, \quad a_f = 2I_{3,f}.$$

$\Gamma_{had}$  reads:

$$\Gamma_{had} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} \cdot \left(1 + \frac{\alpha_s}{\pi} + \dots\right) \cdot (3c_d + 2c_u)$$

# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W, \quad a_f = 2I_{3,f}.$$

$\Gamma_{had}$  reads:

$$\Gamma_{had} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} \cdot \left(1 + \frac{\alpha_s}{\pi} + \dots\right) \cdot (3c_d + 2c_u)$$

$$\Gamma_{had} = \text{const} \cdot (3c_d + 2c_u) \quad (1)$$



# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W, \quad a_f = 2I_{3,f}.$$

$\Gamma_{had}$  reads:

$$\Gamma_{had} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} \cdot \left(1 + \frac{\alpha_s}{\pi} + \dots\right) \cdot (3c_d + 2c_u)$$

$$\Gamma_{had} = \text{const} \cdot (3c_d + 2c_u) \quad (1)$$

and  $\Gamma_{had+\gamma}$ :

$$\Gamma_{had+\gamma} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} F(y_{cut}) \frac{\alpha}{2\pi} (3q_d^2 c_d + 2q_u^2 c_u)$$

# General idea

We want to measure quark couplings:

$$c_f = v_f^2 + a_f^2.$$

They are given in the SM by:

$$v_f = 2I_{3,f} - 4Q_f \sin^2 \theta_W, \quad a_f = 2I_{3,f}.$$

$\Gamma_{had}$  reads:

$$\Gamma_{had} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} \cdot \left(1 + \frac{\alpha_s}{\pi} + \dots\right) \cdot (3c_d + 2c_u)$$

$$\Gamma_{had} = \text{const} \cdot (3c_d + 2c_u) \quad (1)$$

and  $\Gamma_{had+\gamma}$ :

$$\Gamma_{had+\gamma} = N_c \frac{G_\mu M_Z^3}{24\pi\sqrt{2}} F(y_{cut}) \frac{\alpha}{2\pi} (3q_d^2 c_d + 2q_u^2 c_u)$$

$$\Gamma_{had+\gamma} = \text{const}' \cdot (3c_d + 8c_u) \quad (2)$$

# Measurement at $e^+e^-$ colliders

Separate measurements of  $\Gamma_{had} \sim (3c_d + 2c_u)$  and  $\Gamma_{had+\gamma} \sim (3c_d + 8c_u)$  allow disentangling of the couplings  $c_d$  and  $c_u$ !

# Measurement at $e^+e^-$ colliders

Separate measurements of  $\Gamma_{had} \sim (3c_d + 2c_u)$  and  $\Gamma_{had+\gamma} \sim (3c_d + 8c_u)$  allow disentangling of the couplings  $c_d$  and  $c_u$ !

We want to consider:

$$e^+e^- \rightarrow q\bar{q}(\gamma)$$
$$q \in \{u, d\}$$

# Measurement at $e^+e^-$ colliders

Separate measurements of  $\Gamma_{had} \sim (3c_d + 2c_u)$  and  $\Gamma_{had+\gamma} \sim (3c_d + 8c_u)$  allow disentangling of the couplings  $c_d$  and  $c_u$ !

We want to consider:

$$e^+e^- \rightarrow q\bar{q}(\gamma)$$
$$q \in \{u, d\}$$

However, experimentally measured photons can originate not only from the Final State Radiation but also from the Initial State Radiation, hadronisation and decays...

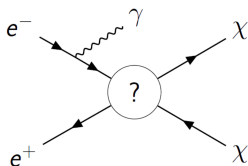
# Monte Carlo generation of radiative events

At the Monte Carlo generation level, it is even more complicated:

- Matrix Element calculations may be either divergent or very slow for small angles,
- ISR structure function can be used for small angles but a proper matching procedure is needed,
- FSR showers are important to account for QCD emissions but they may cause double-counting,
- photons coming from hadronisation and decays have to be included.

# Starting point

Some part of the work has already been done...

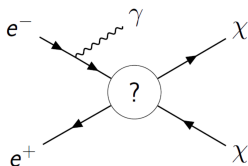


*Simulating hard photon production with WHIZARD*

J. Kalinowski *et al.*, arXiv:2004.14486

# Starting point

Some part of the work has already been done...



*Simulating hard photon production with WHIZARD*

J. Kalinowski *et al.*, arXiv:2004.14486

General idea:

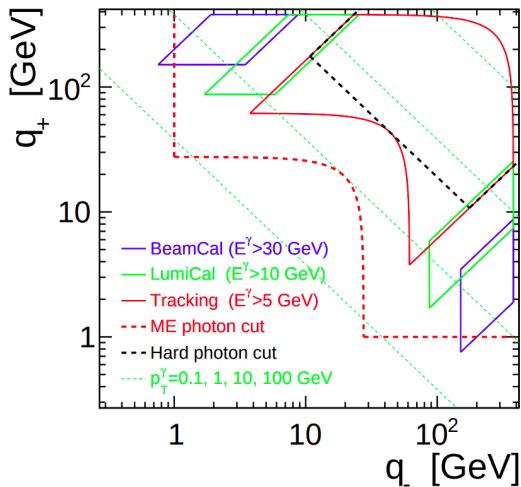
- soft ISR photons simulated using the built-in structure function
- hard ISR photons simulated at the ME level
- matching in  $q_{\pm}$  and  $E_{\gamma}$  ('ISR rejection'):

$$q_- = \sqrt{4E_0 E_{\gamma}} \sin \frac{\theta_{\gamma}}{2}$$

$$q_+ = \sqrt{4E_0 E_{\gamma}} \cos \frac{\theta_{\gamma}}{2}$$



# $(q_+, q_-)$ plane



CLIC 380 GeV, arXiv:2004.14486

# Extension of the procedure

Simulating events with Whizard and Pythia6 (shower and hadronisation)

① ME cuts:

- all  $\gamma$ 's:  $q_{\pm} > 1 \text{ GeV}$  and  $E > 1 \text{ GeV}$  and  $M(\gamma, q_i) > 1 \text{ GeV}$
- any  $\gamma$ :  $p^T > 2 \text{ GeV}$  and  $5^\circ < \theta < 175^\circ$  [useful for efficiency]

# Extension of the procedure

Simulating events with Whizard and Pythia6 (shower and hadronisation)

① ME cuts:

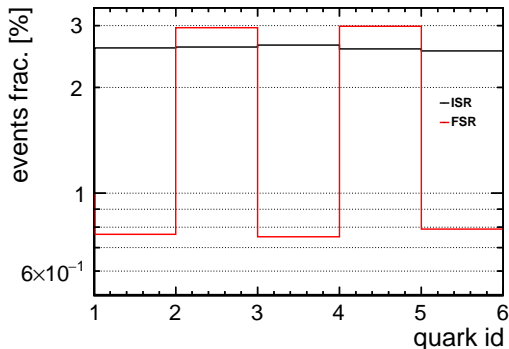
- all  $\gamma$ 's:  $q_{\pm} > 1 \text{ GeV}$  and  $E > 1 \text{ GeV}$  and  $M(\gamma, q_i) > 1 \text{ GeV}$
- any  $\gamma$ :  $p^T > 2 \text{ GeV}$  and  $5^\circ < \theta < 175^\circ$  [useful for efficiency]

② event selection:

- all ISR SF  $\gamma$ 's:  $q_{\pm} < 1 \text{ GeV}$  or  $E < 1 \text{ GeV}$  or  $M(\gamma, q_i) < 1 \text{ GeV}$
- all FSR show.  $\gamma$ 's whose parents are initial  $q/\bar{q}$ :  
 $q_{\pm} < 1 \text{ GeV}$  or  $E < 1 \text{ GeV}$  or  $M(\gamma, q_i) < 1 \text{ GeV}$

## Efficiency of the procedure

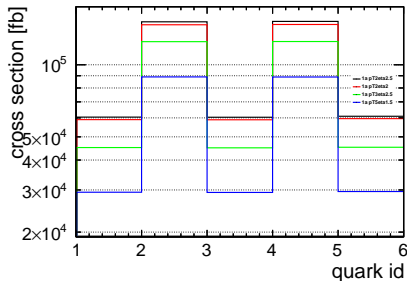
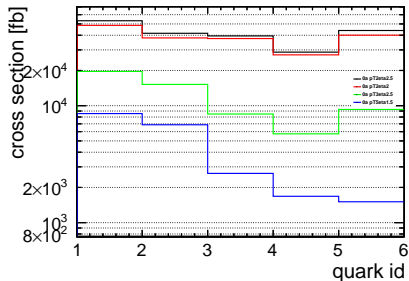
- At the Z-pole, the ISR is reduced so the dominant contribution should come from the FSR.
- Only 4% of Whizard events are rejected.



normalised per flavour

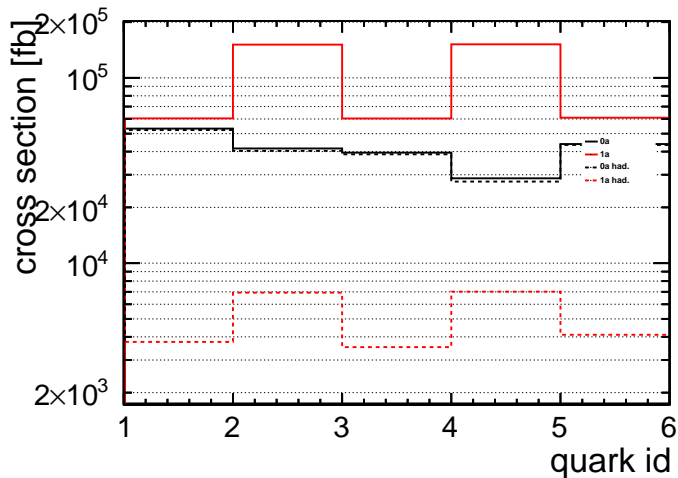
# Reconstruction

- detector simulation in Delphes with default *ILCgen* cards
- analysis cuts:
  - 2 jets
  - exactly 1 reconstructed photon with specific  $(p^T, \eta)$
  - no other activity in the detector

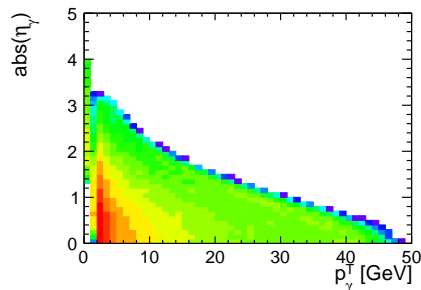
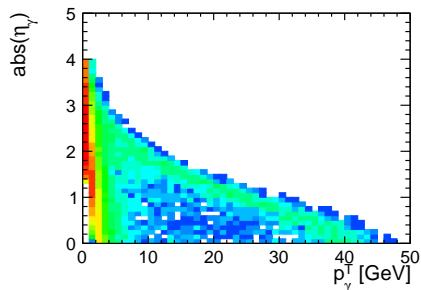


$0\gamma$  vs.  $1\gamma$

# Hadronisation photons



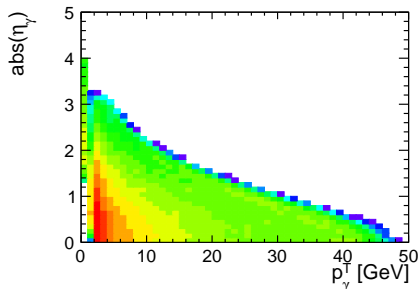
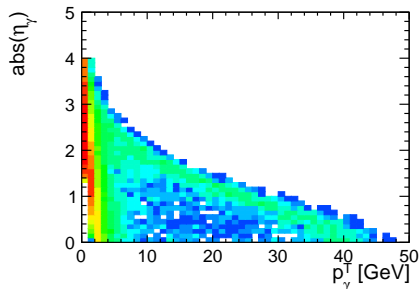
# $(p^T, \eta)$ distribution



$0\gamma$  vs.  $1\gamma$

By a wise choice of the cuts, one may obtain a ratio of  $N_{1\gamma}/N_{0\gamma} > 20$ , keeping the event statistics at a reasonable level.

# $(p^T, \eta)$ distribution

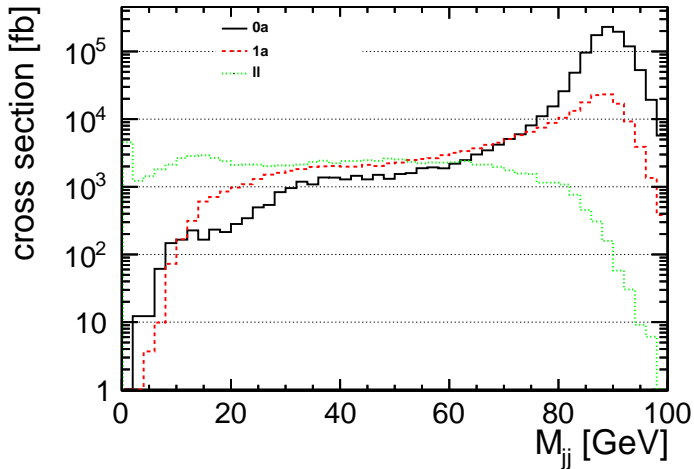


0 $\gamma$  vs. 1 $\gamma$

By a wise choice of the cuts, one may obtain a ratio of  $N_{1\gamma}/N_{0\gamma} > 20$ , keeping the event statistics at a reasonable level.

*Disclaimer:*  $Z \rightarrow \tau^+ \tau^-$  has to be kept under control.





# Prospects and conclusions

- We wanted to estimate how well couplings of the  $Z$  boson to light quarks can be measured.
- We have established a dedicated generation procedure accounting for photons coming from different sources.
- We have performed preliminary studies on the experimental cuts and their efficiency.
- The next step is to study photon isolation criteria which are crucial for reducing the contribution originating from hadronisation and decay.
- Our ultimate goal is to estimate the uncertainty of the measurement at ILC depending on the reconstruction criteria and experimental cuts.
- Work in progress...