

# Probing Gauge-Higgs Unification models with $q\bar{q} A_{FB}$ at ILC

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Gen=T



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# **Introduction & motivation**

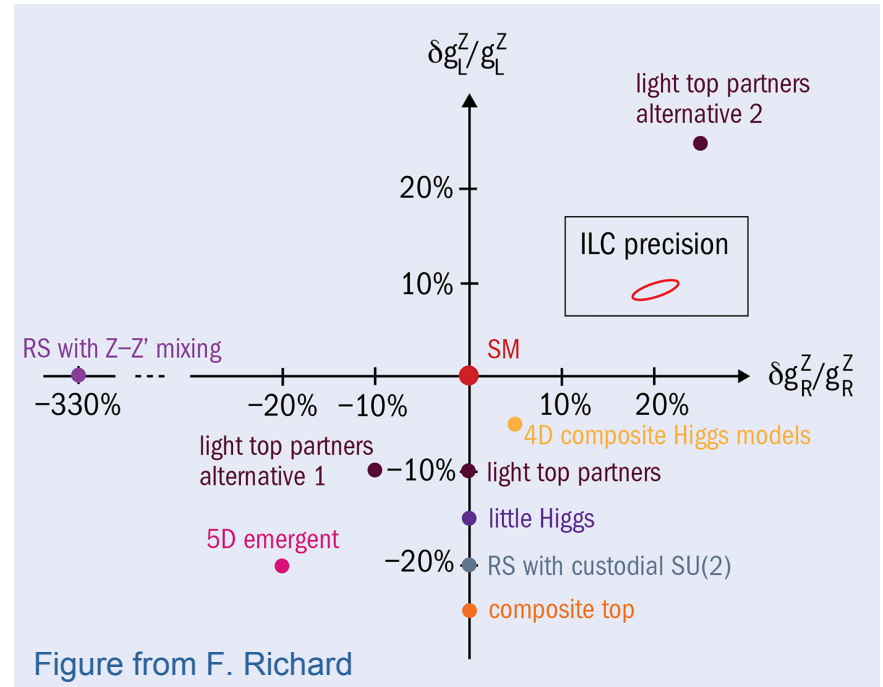


# Motivation: BSM Z' resonances

► Many **BSM scenarios** (i.e. Randal Sundrum, compositeness, Gauge Higgs unification models...) predict heavy resonances coupling to the (t,b) doublet and also lighter fermions (i.e. c/s quarks)

- Only coupling to (t,b) doublet
  - → Peskin, Yoon arxiv:1811.07877
  - → Djouadi et al arxiv:hep-ph/0610173
- Coupling also to lighter fermions [Funatsu, Hatanaka, Hosotani, Orikasa, Yamatsu]

(arxiv:1705.05282) (arxiv:2309.01132) (arxiv:2301.07833)



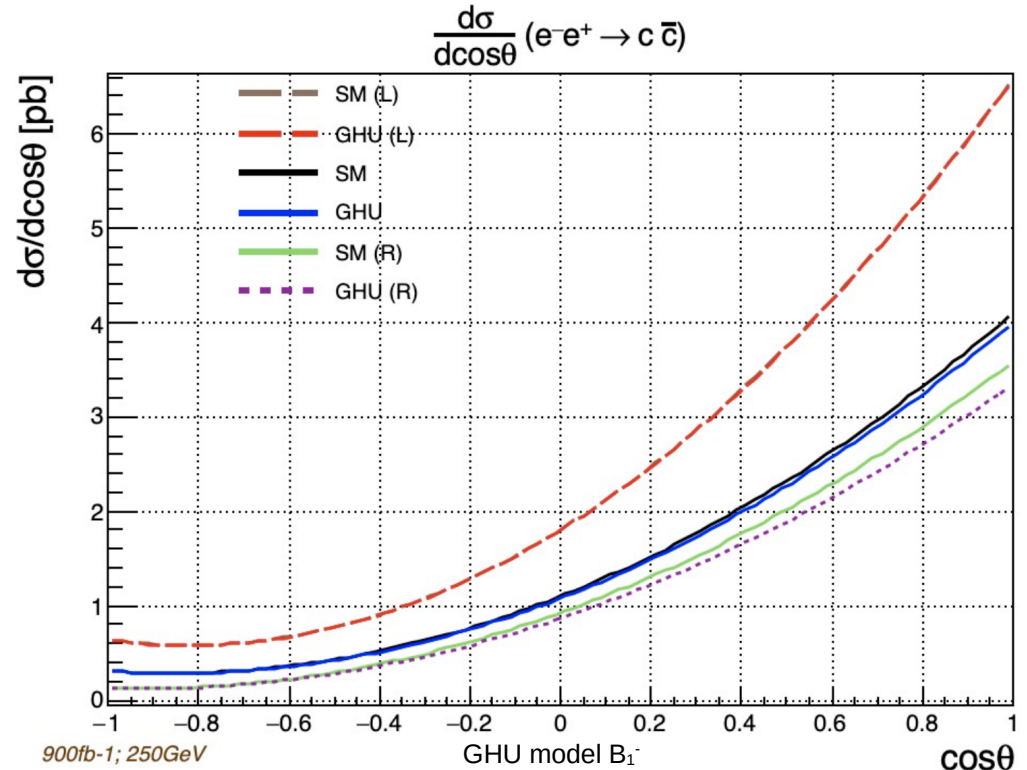
**Probing such scenarios require at least per mil level experimental precision**

tt/**bb/cc**/ss/... **Achievable at future colliders?**

# Gauge-Higgs Unification Models (GHU)

- ▶ Randall-Sundrum metric (5D).
- ▶ The symmetry breaking pattern is different than in the SM and features the *Hosotani mechanism*:
  - Masses are generated dynamically from the extra-dimension properties.
- ▶ Only one parameter, *Hosotani's angle*  $\theta_H$ , determines the projection of the 5D fields, fixing all physical effects:
  - KK resonances of the Z/ $\gamma$  with  $m_{\text{KK}} \sim 10\text{-}25$  TeV.
  - **Modifications and new EW couplings/helicity amplitudes.**
  - Already visible effects at 250GeV.

As **Benchmark**, we will use the [Funatsu, Hatanaka, Hosotani, Orikasa, Yamatsu] models.



▶ A models: ([arxiv:1705.05282](https://arxiv.org/abs/1705.05282))

$$A_1 : \theta_H = 0.0917, m_{KK} = 8.81 \text{ TeV} \rightarrow m_{Z^1} = 7.19 \text{ TeV};$$

$$A_2 : \theta_H = 0.0737, m_{KK} = 10.3 \text{ TeV} \rightarrow m_{Z^1} = 8.52 \text{ TeV},$$

▶ B models: ([arxiv:2309.01132](https://arxiv.org/abs/2309.01132)) ([arxiv:2301.07833](https://arxiv.org/abs/2301.07833))

$$B_1^\pm : \theta_H = 0.10, m_{KK} = 13 \text{ TeV} \rightarrow m_{Z^1} = 10.2 \text{ TeV};$$

$$B_2^\pm : \theta_H = 0.07, m_{KK} = 19 \text{ TeV} \rightarrow m_{Z^1} = 14.9 \text{ TeV};$$

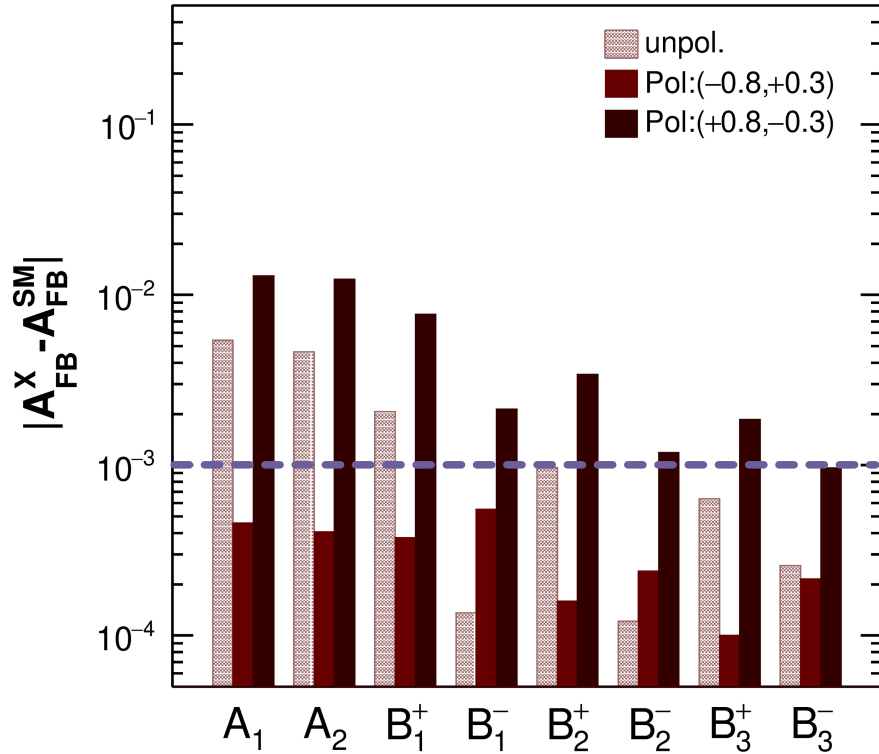
$$B_3^\pm : \theta_H = 0.05, m_{KK} = 25 \text{ TeV} \rightarrow m_{Z^1} = 19.6 \text{ TeV};$$

Resonances of O(10) TeV: Only indirect measurements are possible!

# GHU vs SM (250 GeV)

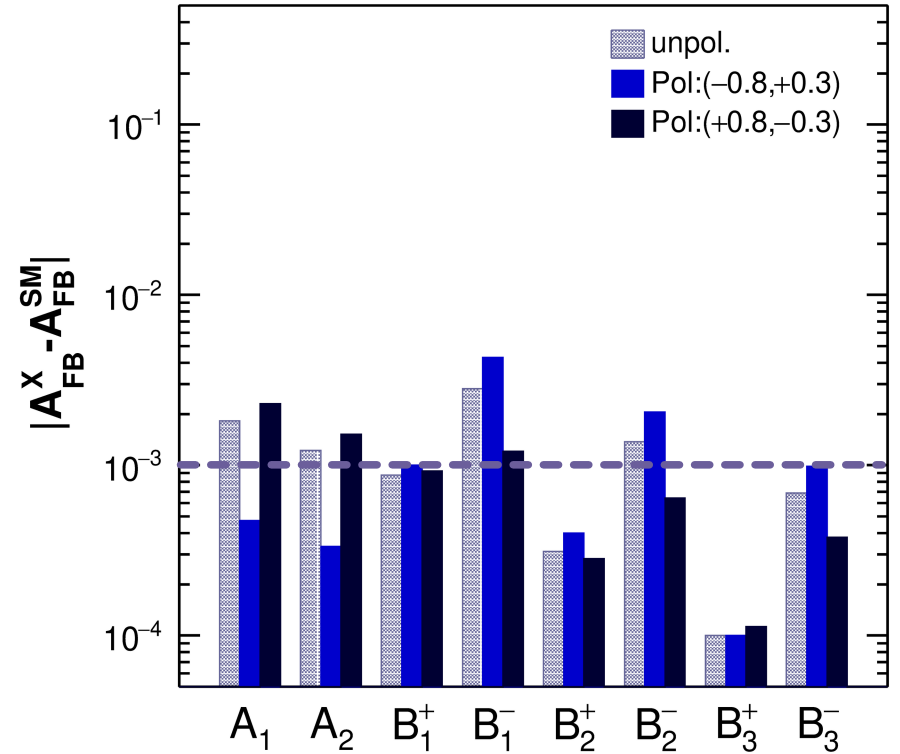
$\sqrt{s_{e^+e^-}} = 250 \text{ GeV}$

**b-quark**



$\sqrt{s_{e^+e^-}} = 250 \text{ GeV}$

**c-quark**



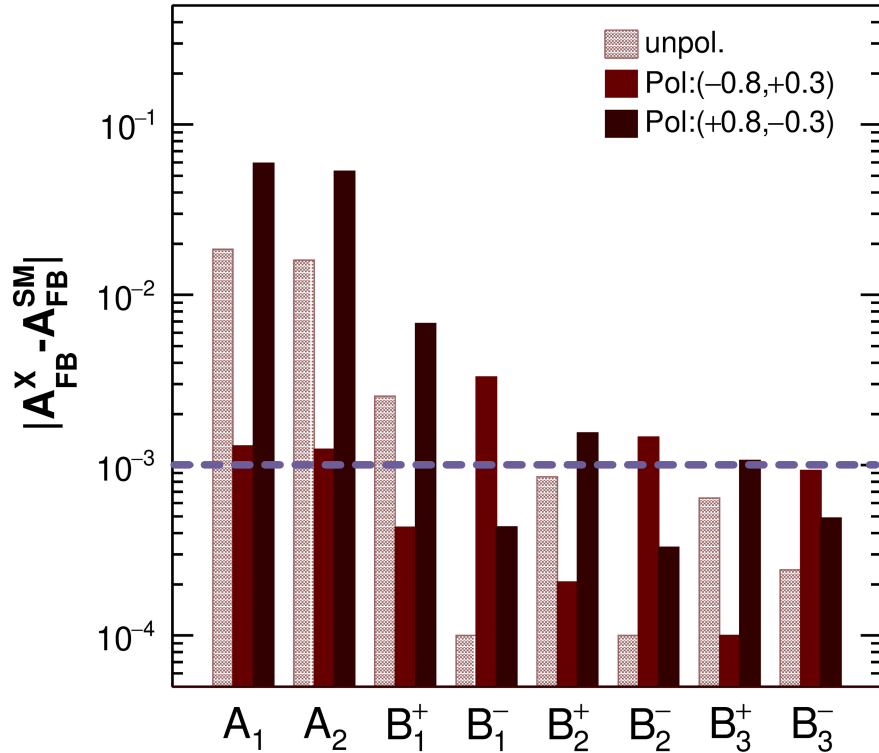
Deviations at the **per mil** level

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

# GHU vs SM (500 GeV)

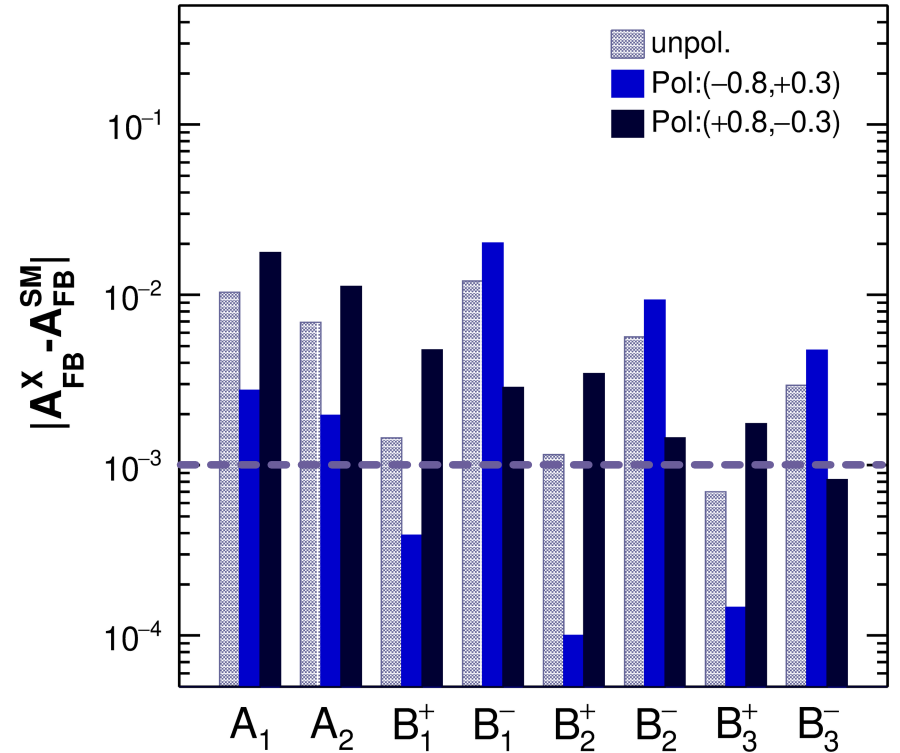
$\sqrt{s_{e^+e^-}} = 500 \text{ GeV}$

**b-quark**



$\sqrt{s_{e^+e^-}} = 500 \text{ GeV}$

**c-quark**



Deviations at the **per mil** level

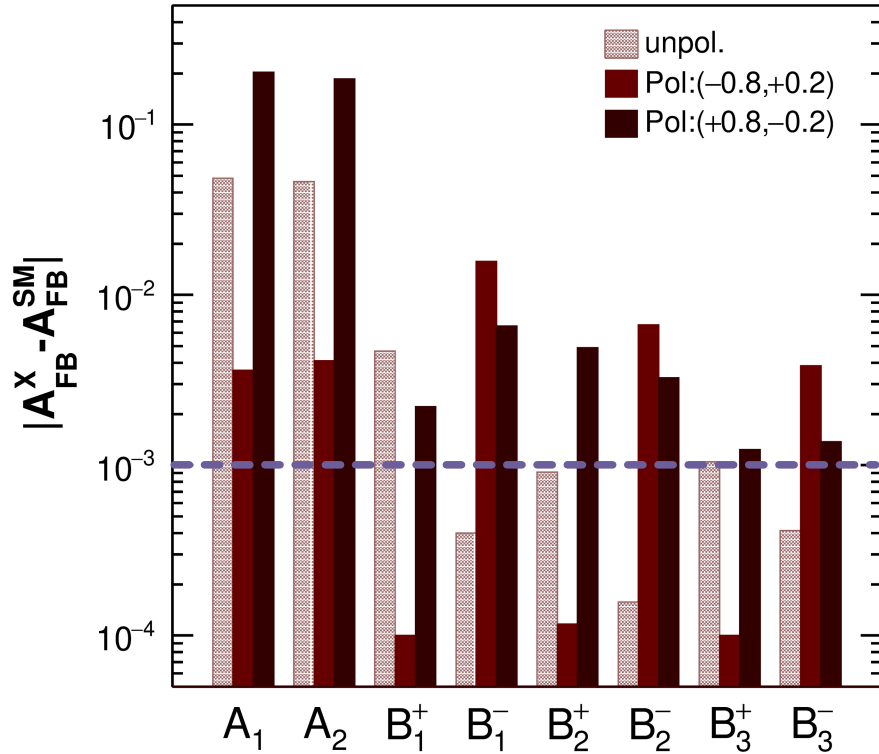
$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$



# GHU vs SM (1 TeV)

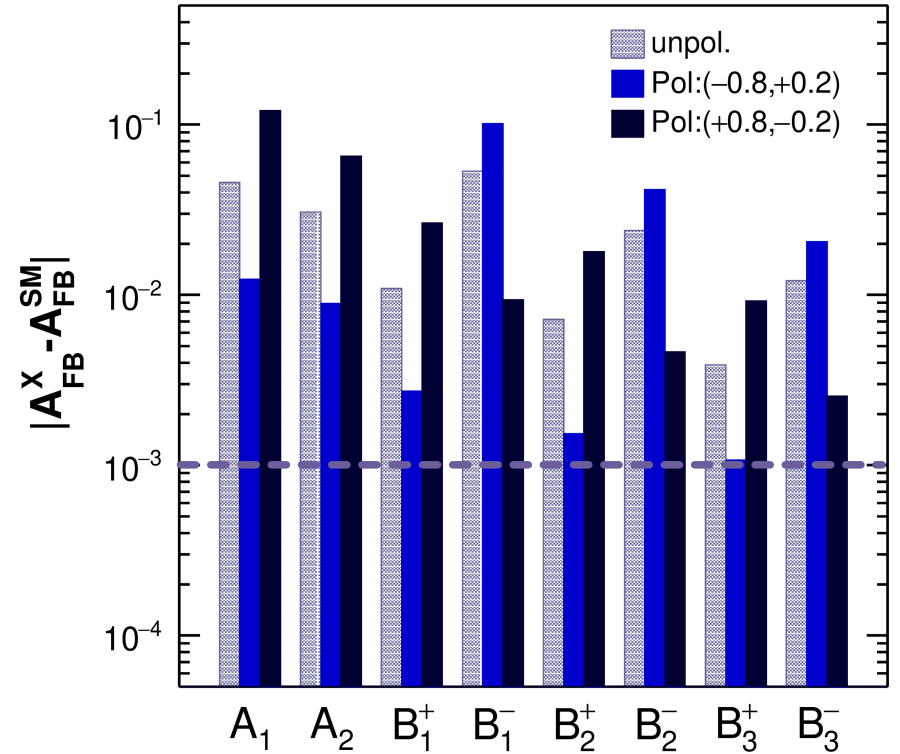
$\sqrt{s_{e^+e^-}} = 1000 \text{ GeV}$

**b-quark**



$\sqrt{s_{e^+e^-}} = 1000 \text{ GeV}$

**c-quark**



Deviations at the **per mil** level

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$



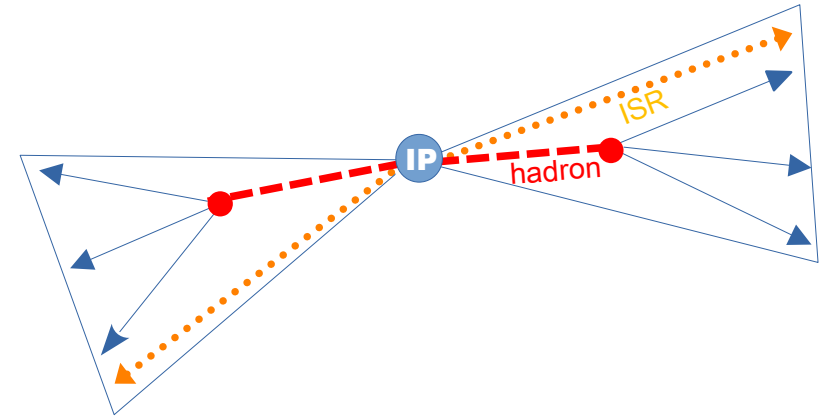
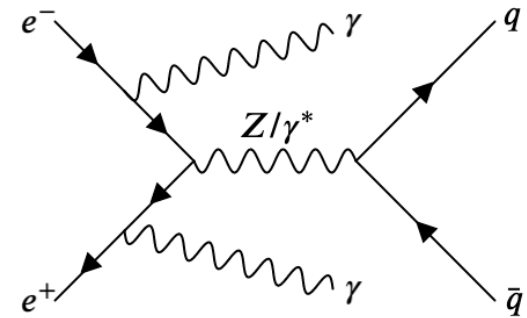
# **Experimental study with full simulation**



- ▶ ILD note (2022) [2306.11413](#)
  - ILC250, b and c studies. (A. Irles, F. Richard, R. Pöschl).
  
- ▶ Work presented in different conferences (J.P. Márquez):
  - Proceeding LCWS (2023) [2307.14888](#)
    - ▶ Optimization of flavor tagging, use of dNdx PID and extension to 500 GeV.
  - Proceeding EPS-HEP (2023) [2310.17617](#)
    - ▶ First phenomenology prospects.
  
- ▶ A paper has been submitted for revision to EPJ-C.
  - First draft (2024) [2403.09144](#)

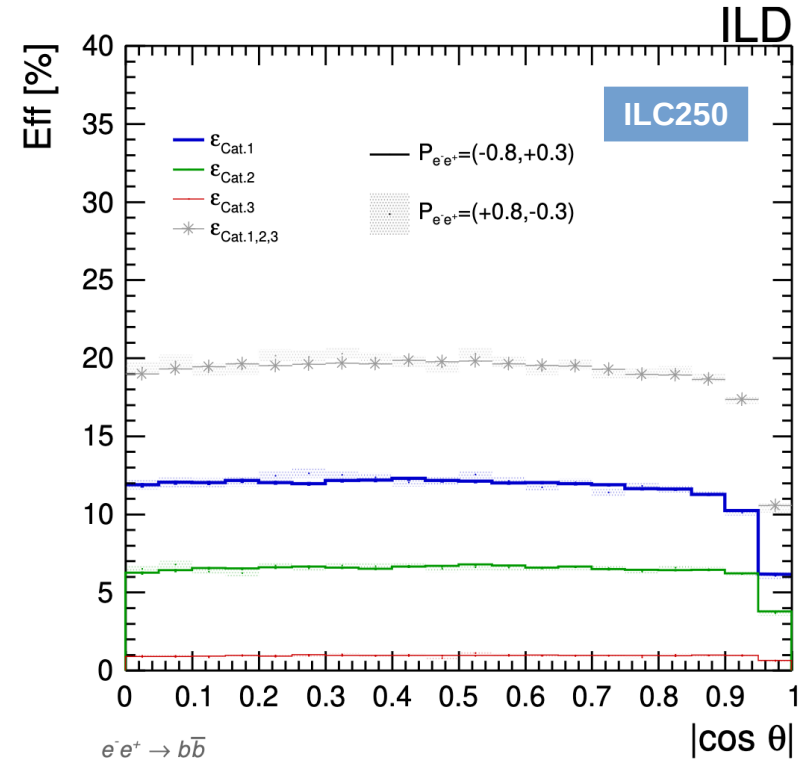
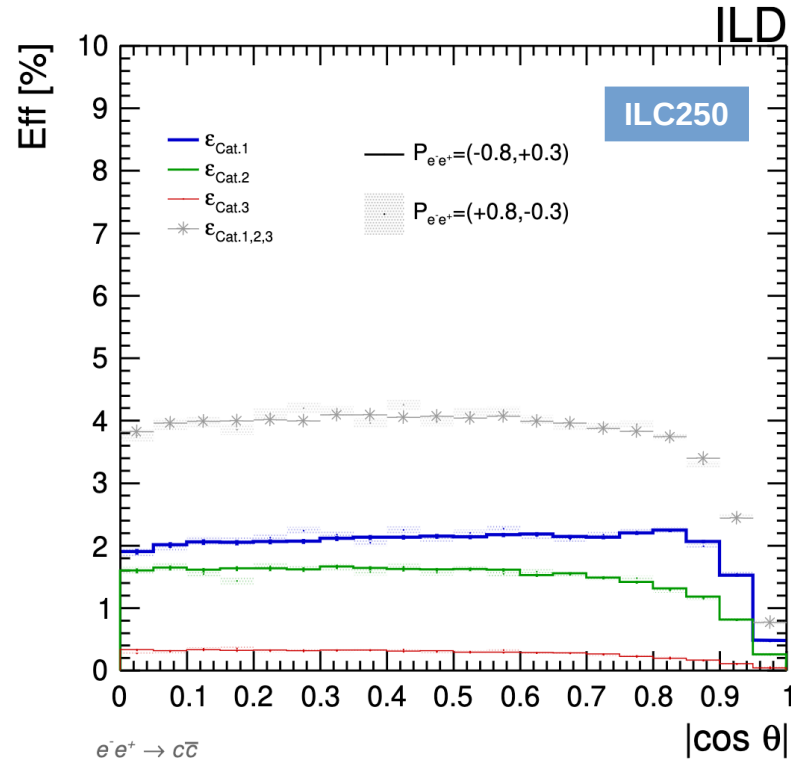


- ▶ We work with  $A_{FB}$  for b and c quarks.
  - MC simulations at 250 and 500 GeV.
    - ▶ International Linear Collider (ILC) run plan.
  - Full simulation of the International Large Detector (ILD).
- ▶ Topology: Two back-to-back jets.
- ▶ Procedure:
  - Background suppression → Selection of  $q\bar{q}$  events.
  - Flavor tagging → Selection of  $b\bar{b}$  &  $c\bar{c}$  events.
    - ▶ Double tagging.
  - Charge measurement → Quark-Antiquark identification.
    - ▶ Double charge.



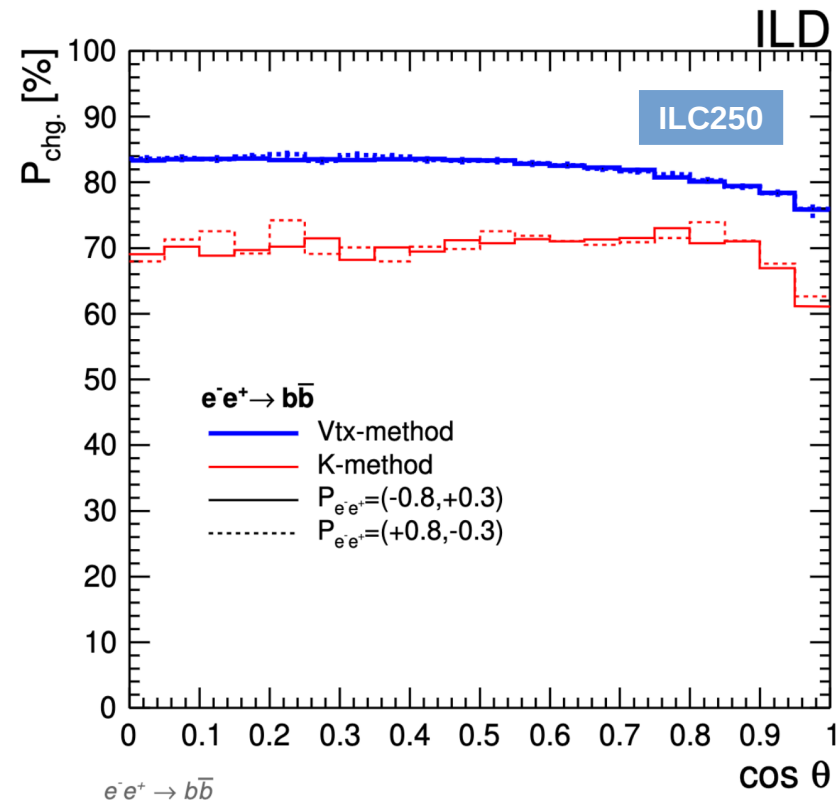
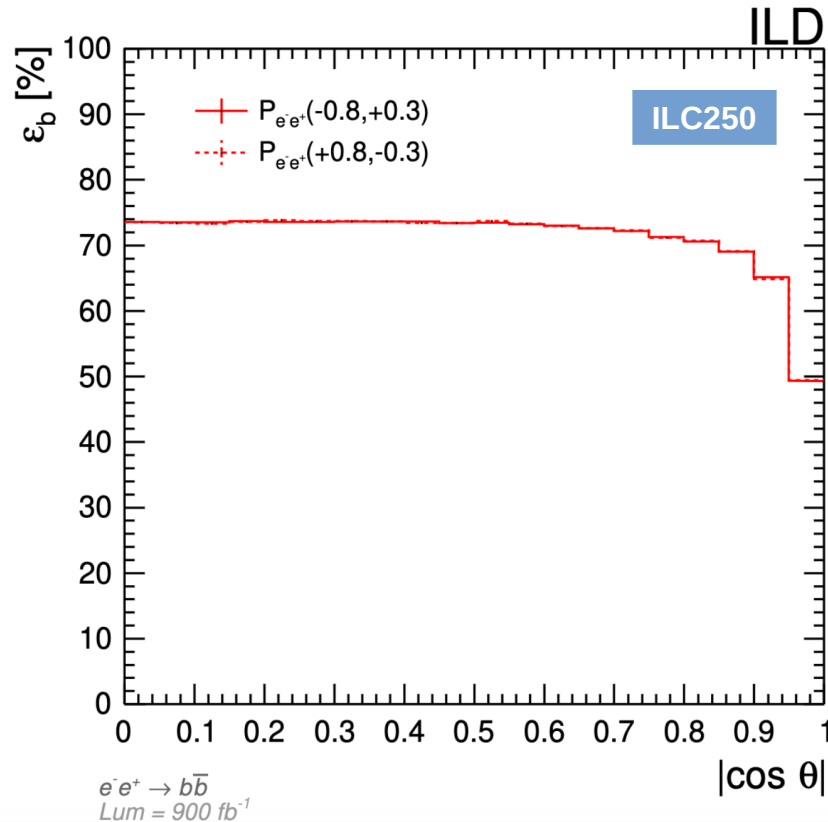
**High-purity & independent samples for each quark flavour.**

- Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )



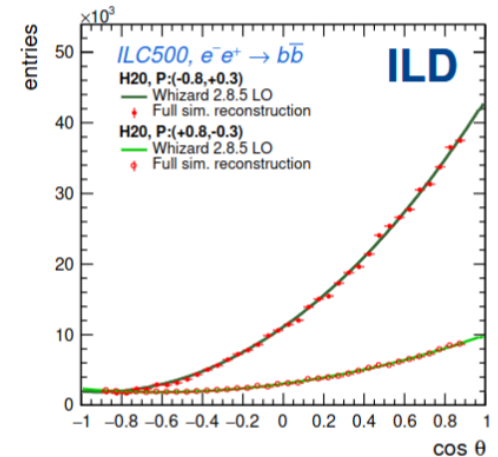
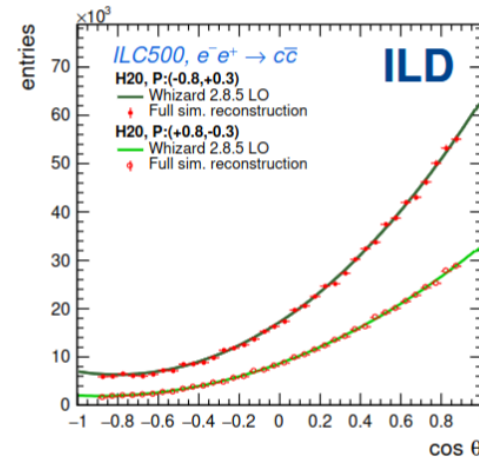
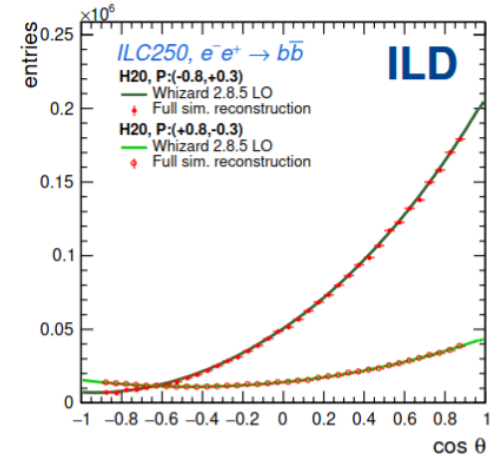
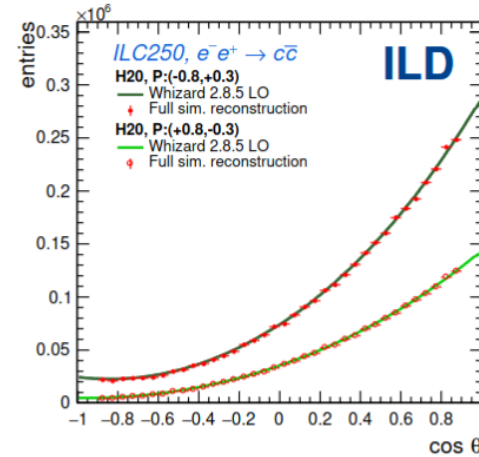
High-purity & independent samples for each quark flavour.

- ▶ Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )
  - To maximally reduce the usage of MC tools (control of fragmentation, QCD correlations... uncertainties)

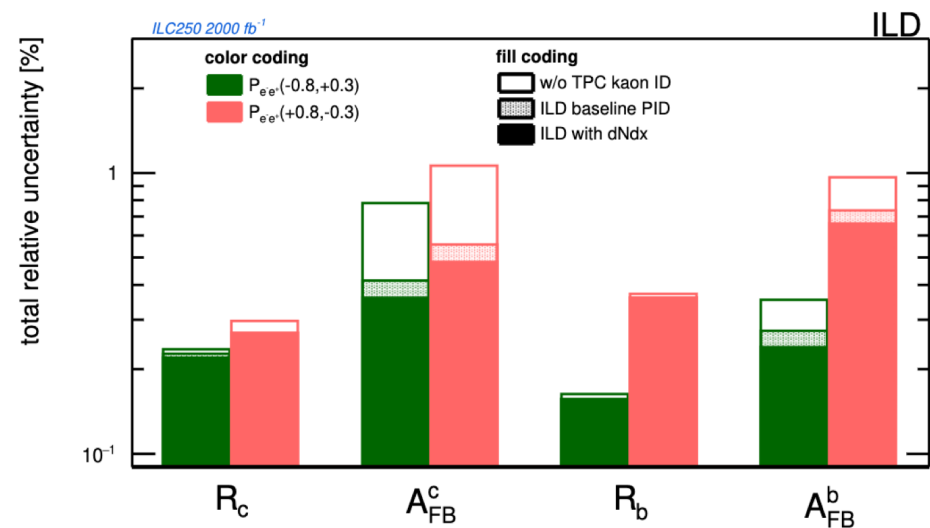
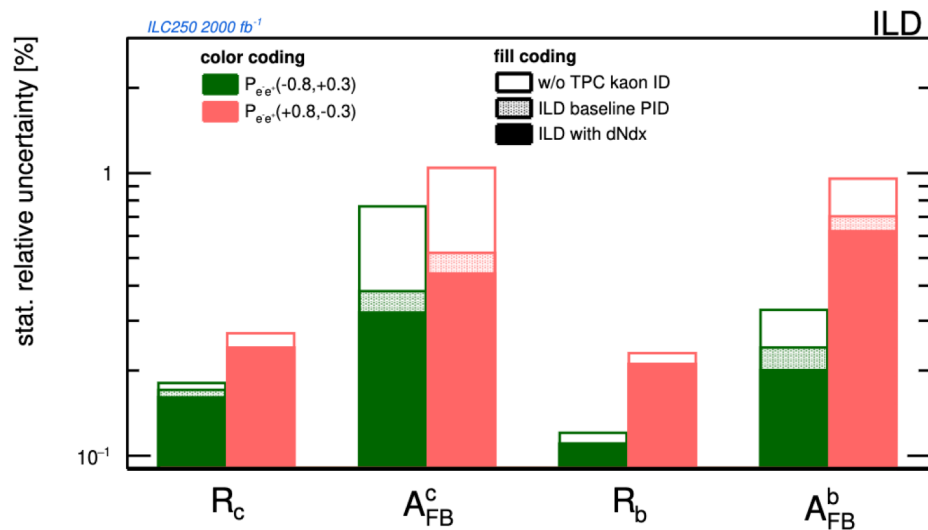


# Fit and results

- ▶ **Results for ILC250 and ILC500:**
  - Presented in LCWS (2023) [2307.14888](https://arxiv.org/abs/2307.14888)
- ▶ **At least 4 observables for AFB at ILC per energy point**
  - 2 quarks (b and c).
  - 2 polarizations ( $e_{LpR}$ ,  $e_{RpL}$ ).
- ▶ **Per mil level statistical uncertainties** reachable for the nominal **ILC program**
  - **Smaller exp syst. Uncertainties.**
    - Fragmentation, angular correlations, preselection efficiency, mistag, etc.



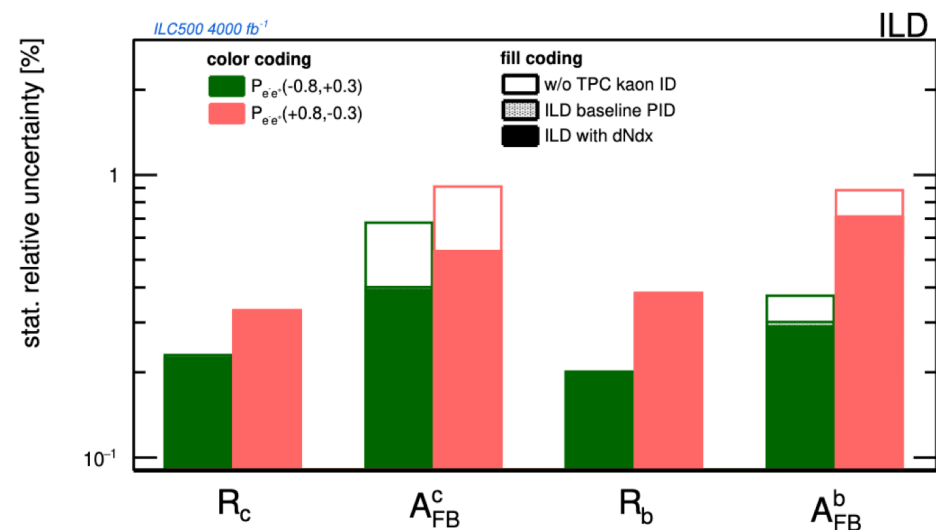
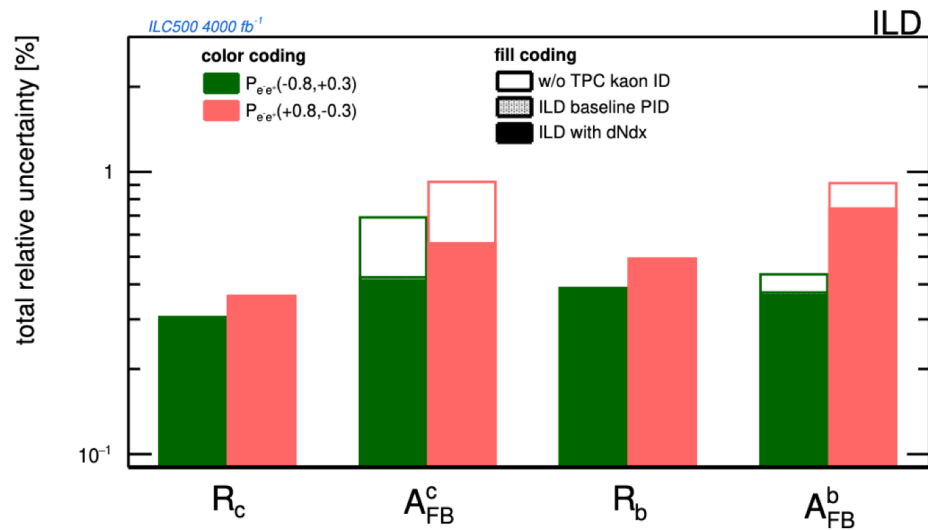
► Presented in LCWS (2023) [2307.14888](#)



Statistical uncertainties dominate over systematic uncertainties

► Presented in LCWS (2023) [2307.14888](#)

- Less benefit from the use of PID, but the  $A_{FB}$  uncertainties are in the same level.



Statistical uncertainties dominate over systematic uncertainties



# Results for ILC250 & ILC500

## ▶ $A_{FB}$ definition:

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

## ▶ At least 4 observables for $A_{FB}$ at ILC per energy point

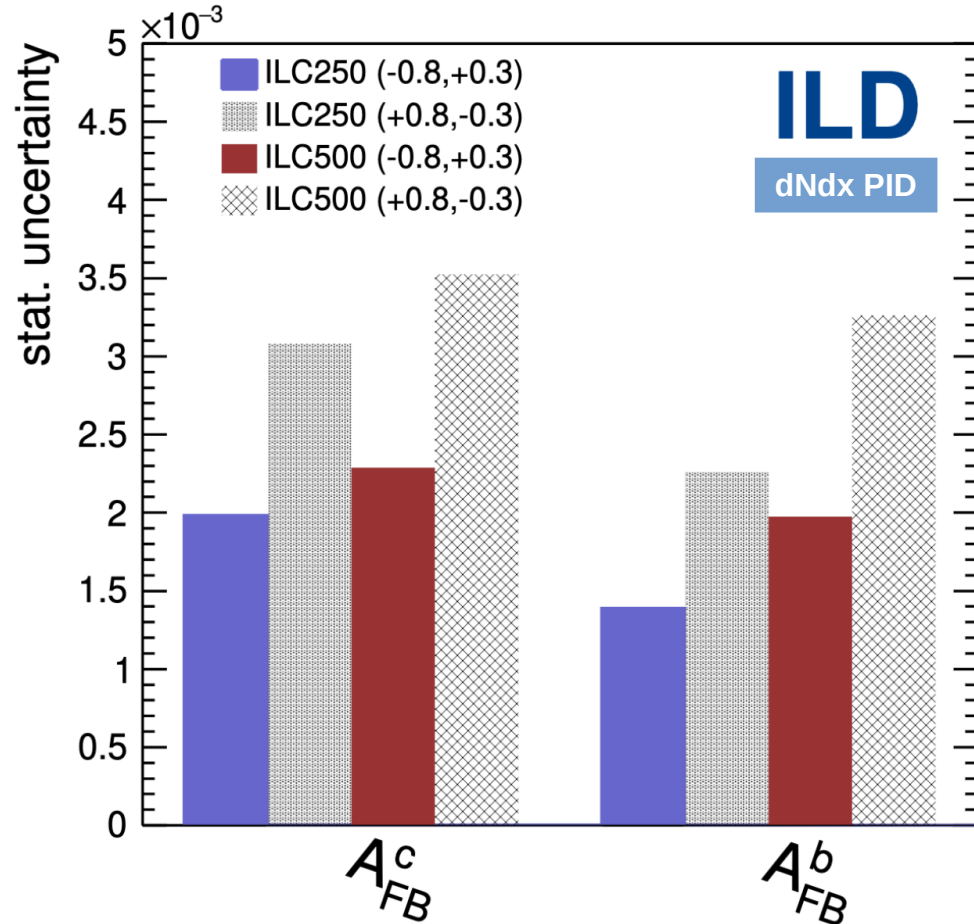
- 2 quarks (b and c).
- 2 polarizations ( $e_L p_R$ ,  $e_R p_L$ ).

## ▶ **Per mil level statistical uncertainties** reachable for the nominal ILC program

- **Smaller exp. syst. Uncertainties**

## ▶ Running at ILC500

- Similar uncertainties but bigger deviations.
- Possibility of combining with the ILC250 results.

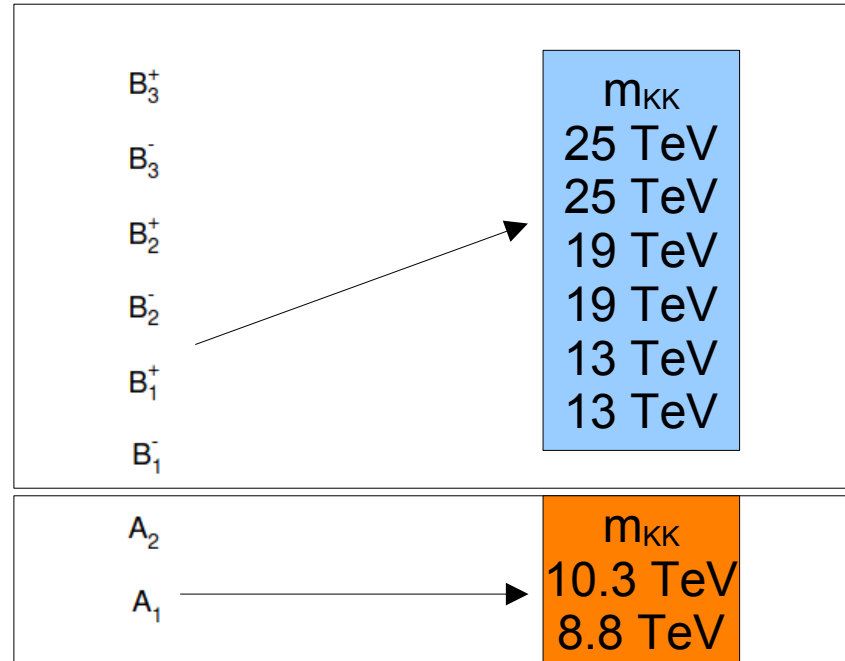


**Discrimination  
power between  
GHU & SM**



- ▶ Procedure: Testing the statistical significance of model  $AFB_{\text{test}}$  vs a reference model  $AFB_{\text{ref}}$  assuming that one of them is measured.
- ▶ The uncertainties are considered normally distributed:
  - Significance in  $\sigma$ .
  - P-value: Gaussian at  $d_\sigma$ .
$$d_\sigma = \frac{\|AFB_{\text{test}} - AFB_{\text{ref}}\|}{\Delta_{AFB_{\text{ref}}}}$$
- ▶ Combination of multiple measurements is done with a *multivariate gaussian*.
  - Assuming no correlations for AFB.
- ▶ We also assumed different precisions for the SM Z boson couplings:
  - Current precision, ILC250 and Giga-Z (ILC run at the Z-Pole).

# GHU vs SM: GHU energy scale

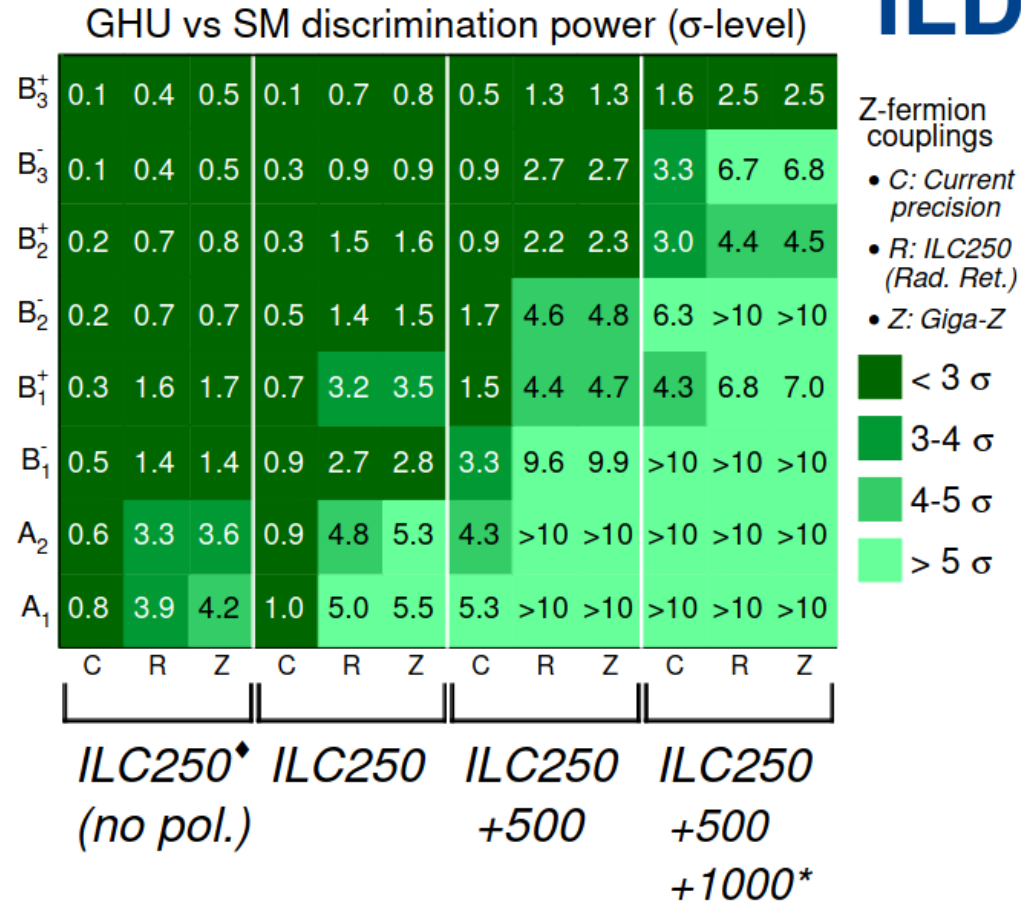


## Similar structure for all plots:

- More massive resonances, i.e., hardest to detect models as we move up.
  - Higher energy accessed by the ILC runs as we move to the right.

# GHU vs SM: discrimination power plots

ILD

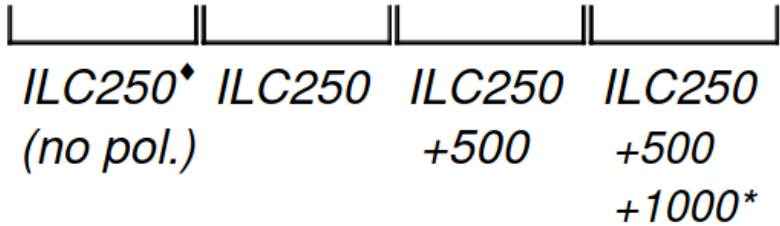


# GHU vs SM: Beam scenarios

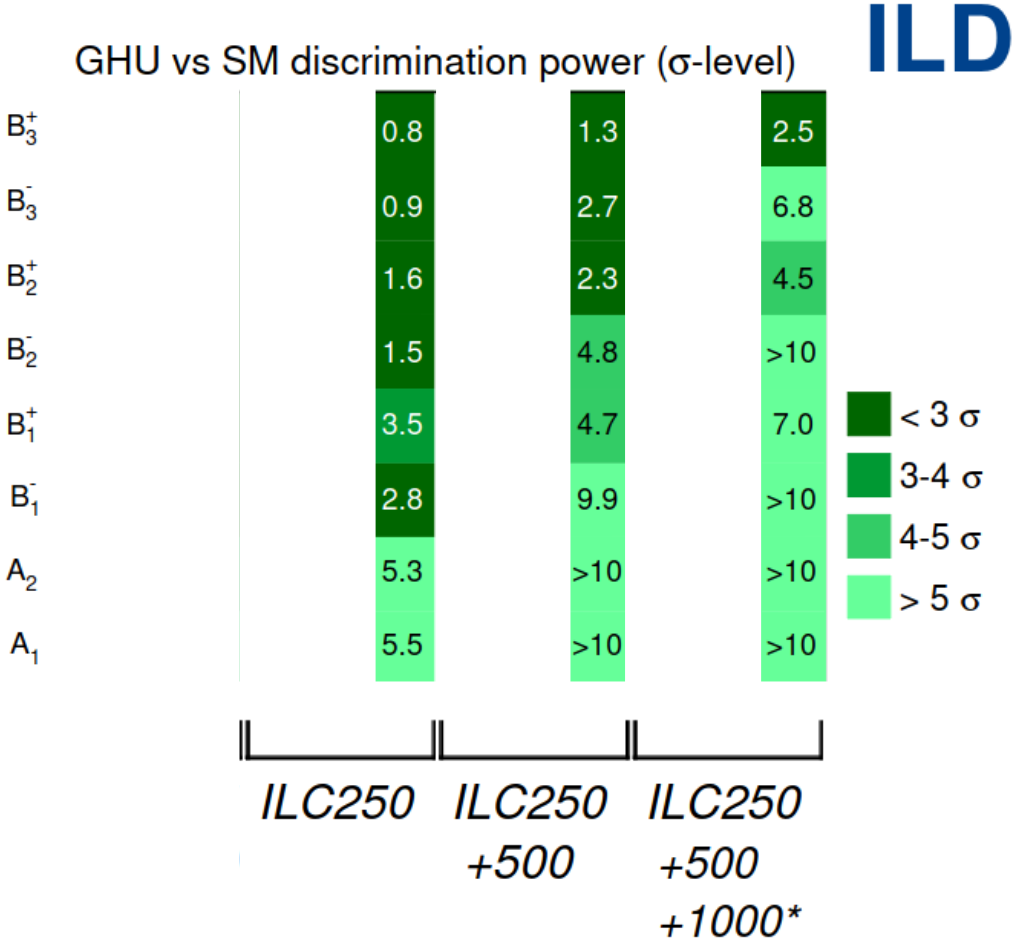
*Hypothetical case*  
**ILC250\* no pol**  
 $\int L = 2000 \text{ fb}^{-1}$   
  
**Full ILD simulation**  
**assuming**  
**no beam pol.**

**H20 nominal program**  
**ILC250**  
(  $P_{e-}=0.8, P_{e+}=0.3$  )  
 $\int L = 2000 \text{ fb}^{-1}$   
  
**ILC500**  
(  $P_{e-}=0.8, P_{e+}=0.3$  )  
 $\int L = 4000 \text{ fb}^{-1}$   
  
**Full ILD simulation**  
**assuming beam pol.**

**H20 nominal program**  
**ILC1000**  
(  $P_{e-}=0.8, P_{e+}=0.2$  )  
 $\int L = 8000 \text{ fb}^{-1}$   
  
*Not full simulation studies*  
*but extrapolations from ILC500*

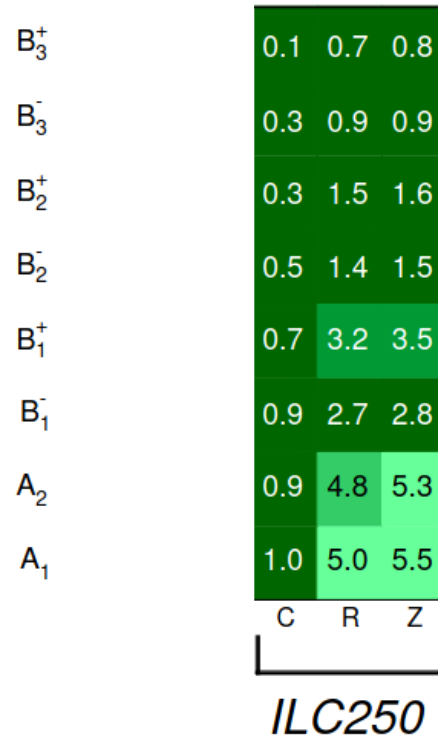


# GHU vs SM: center of mass energy



# GHU vs SM: Precision on Z-couplings

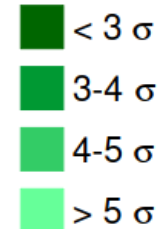
GHU vs SM discrimination power ( $\sigma$ -level)



## ILD

Z-fermion couplings

- C: Current precision
- R: ILC250 (Rad. Ret.)
- Z: Giga-Z

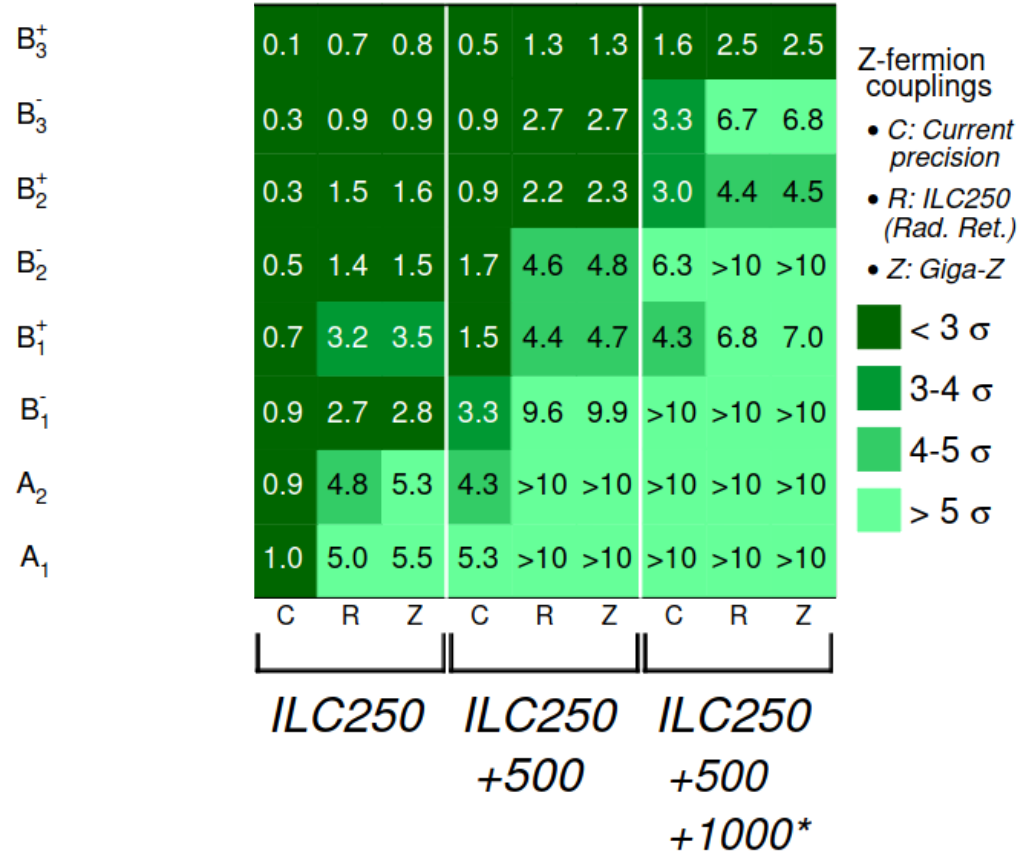




# GHU vs SM: Precision on Z-couplings

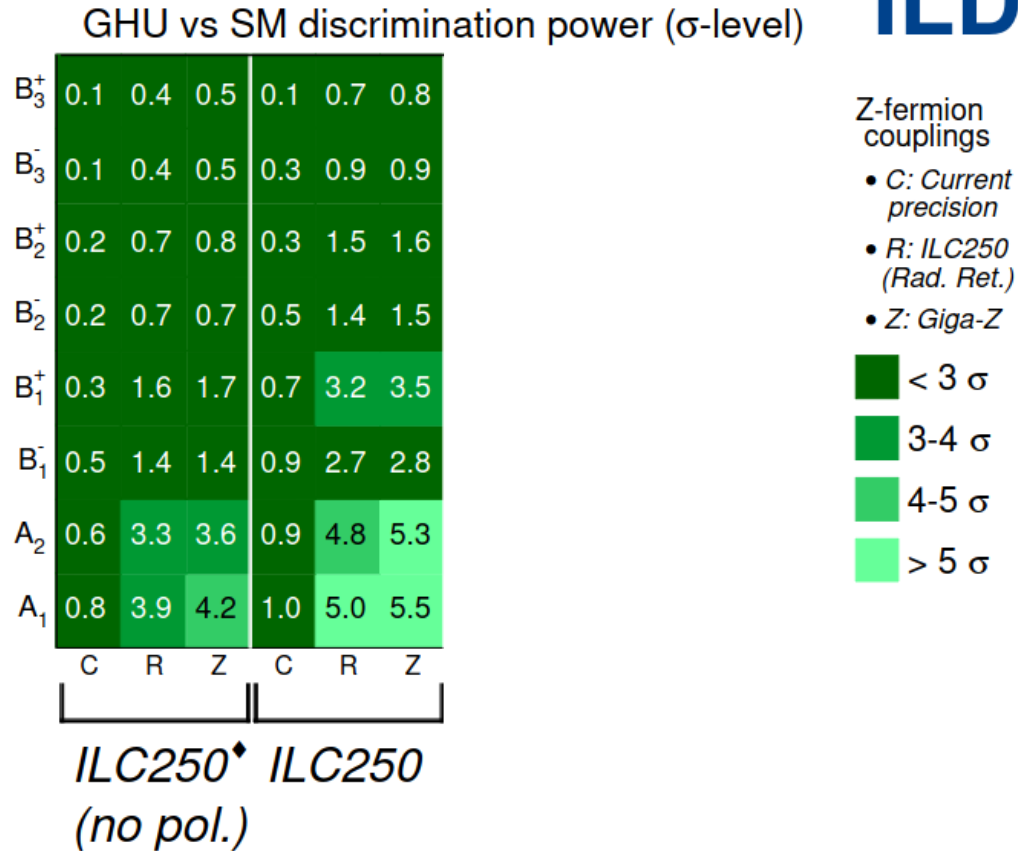
ILD

GHU vs SM discrimination power ( $\sigma$ -level)



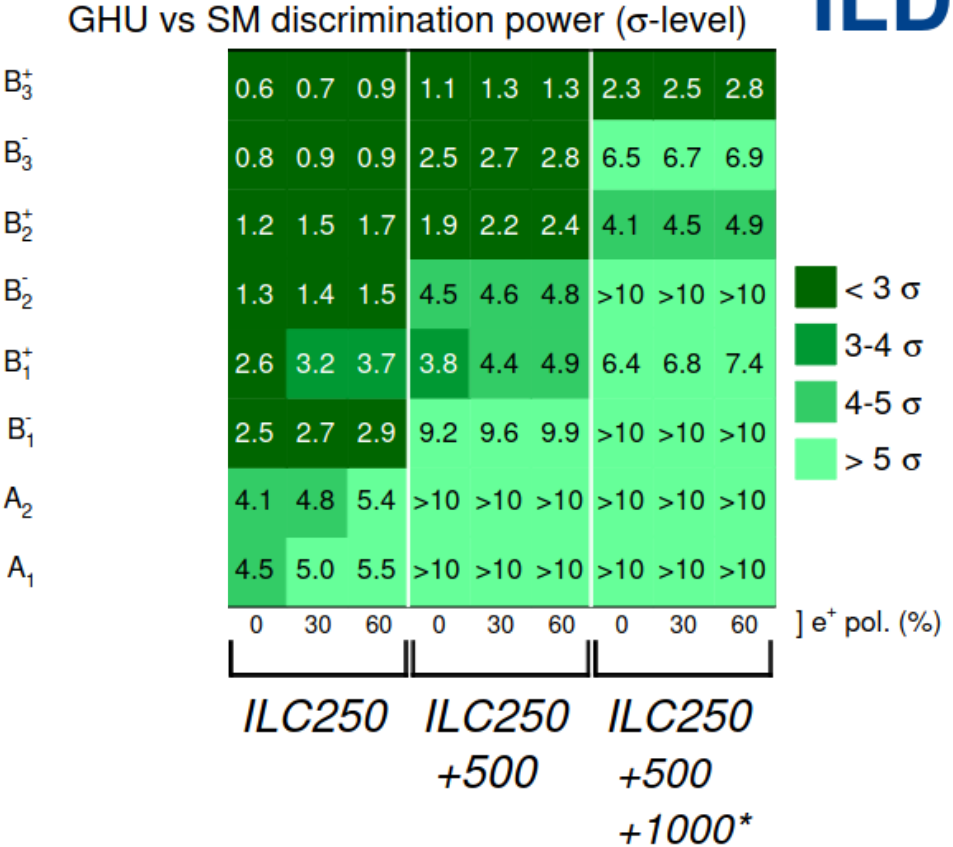
# GHU vs SM: Beam(s) polarization

ILD



# GHU vs SM: Positron beam polarization

ILD

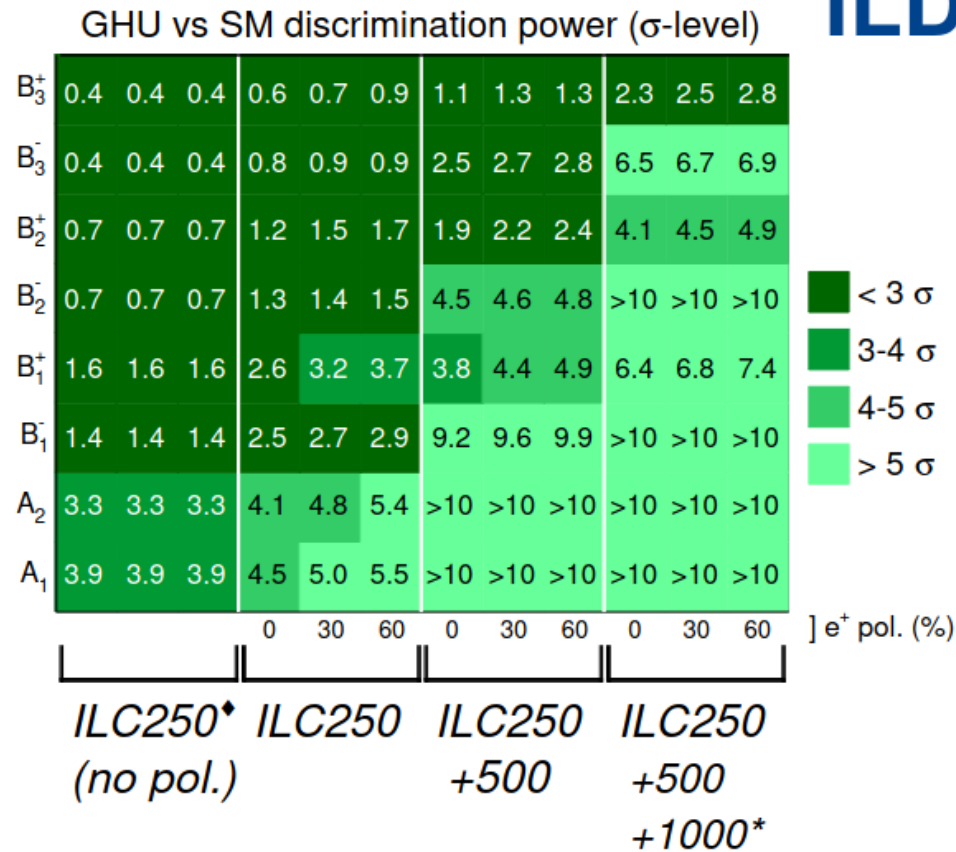


**e- polarization at 80% in all these scenarios**



# GHU vs SM: Positron beam polarization

ILD



# GHU vs SM: Particle ID dependence

GHU vs SM discrimination power ( $\sigma$ -level)

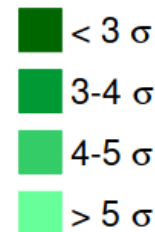
$B_3^+$	0.5	0.7	0.7
$B_3^-$	0.5	0.8	0.9
$B_2^+$	0.9	1.4	1.5
$B_2^-$	0.8	1.3	1.4
$B_1^+$	2.2	3.1	3.2
$B_1^-$	1.4	2.4	2.7
$A_2$	3.3	4.7	4.8
$A_1$	3.5	4.9	5.0
	O	E	N

ILC250

ILD

Ch. had. PID

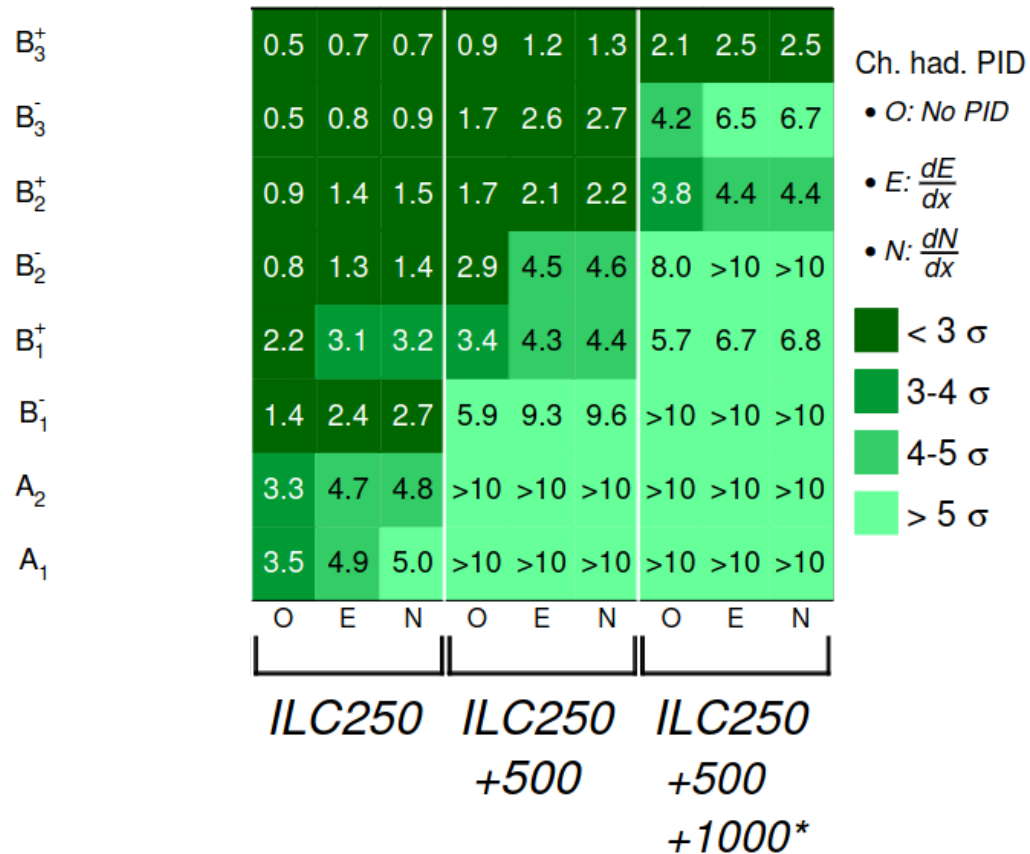
- O: No PID
- E:  $\frac{dE}{dx}$
- N:  $\frac{dN}{dx}$



# GHU vs SM: Particle ID dependence

GHU vs SM discrimination power ( $\sigma$ -level)

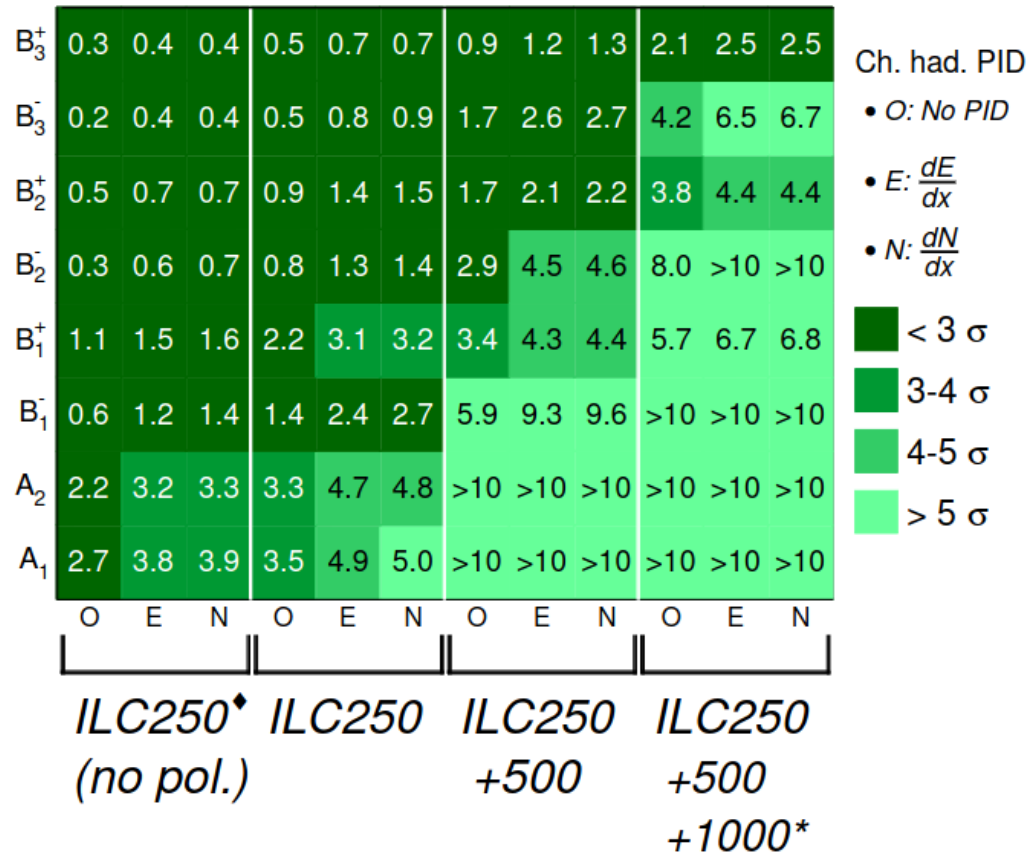
ILD



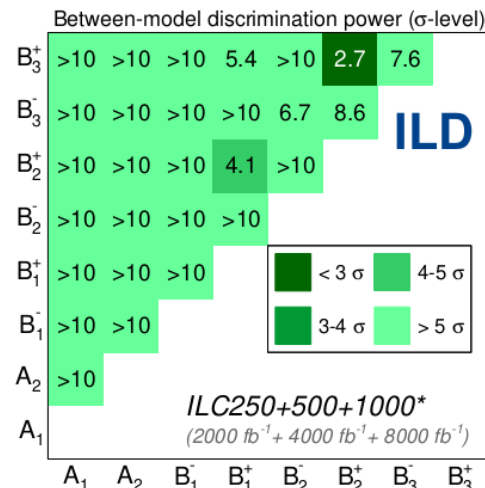
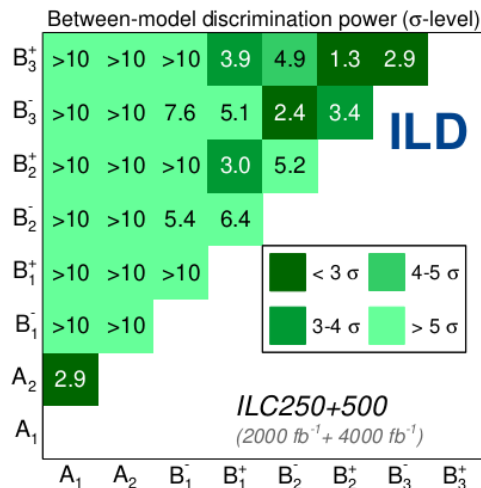
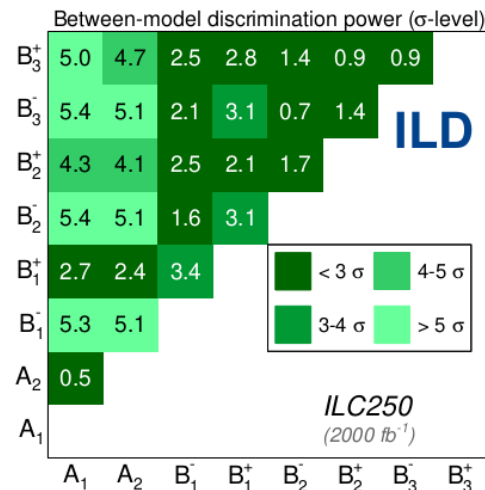
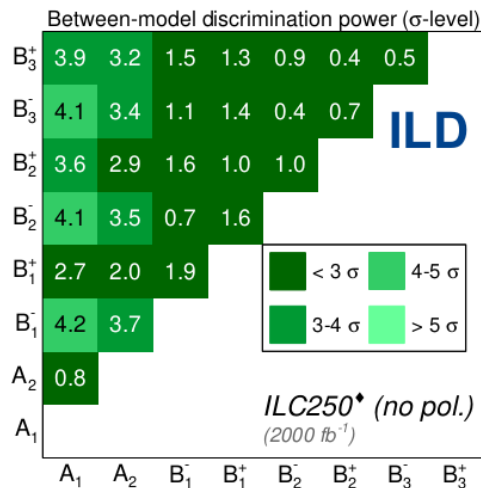
# GHU vs SM: Particle ID dependence

ILD

GHU vs SM discrimination power ( $\sigma$ -level)



# GHU between model discrimination





# **Conclusion/ summary**



- ▶ ILC offers unique capabilities to explore these signatures and discriminate GHU vs SM:
  - **High energy reach.**
  - Electron and positron **beam polarization** → enhancing the sensitivity but also allowing combination of measurements with different BSM sensitivity (for control of systematics).
- ▶ Comprehensive study done at ILC250/ILC500 with ILD simulations:
  - Backgrounds, beam features, polarization, realistic reconstruction tools.
  - Uncertainties dominated by statistics, above the Z-pole.
  - Room for improvement (modern algorithms for flavour tagging, event selection, etc.
- ▶ Full discrimination of almost all of the proposed models (and within models) is possible with the H20 nominal run plan for ILC!



## ► Studies and paper reviewed in the ILD collaboration:

- By the ILD editorial board: Mikael Berggren and Daniel Jeans.
- Circulated on the ILD mailing list.

## ► It has been submitted for review to EPJ-C.

Eur. Phys. J. C manuscript No.  
(will be inserted by the editor)

### Probing Gauge-Higgs Unification models at the ILC with di-quark forward-backward asymmetry at center-of-mass energies above the $Z$ mass. \*

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Received: date / Accepted: date

**Abstract** The International Linear Collider (ILC) will allow the precise study of  $e^+e^- \rightarrow q\bar{q}$  interactions at different center-of-mass energies from the  $Z$ -pole to 1 TeV. In this paper we discuss the experimental prospects for measuring differential observables in  $e^+e^- \rightarrow b\bar{b}$  and  $e^+e^- \rightarrow c\bar{c}$  at the ILC baseline energies, 250 and 500 GeV. The studies are based on full detector simulation samples and reconstruction of the International Large Detector (ILD) concept. Two gauge-Higgs unification models predicting new high-mass resonances beyond the Standard Model are discussed. These models predict sizable deviations of the forward-backward observables at the ILC running above the  $Z$  mass and with longitudinally polarized electron and positron beams. The ability of the ILC to probe these models via high-precision forward-backward asymmetry measurements is discussed. Alternative scenarios with other energy points or different beam polarisation schemes are also discussed, extrapolating the estimated uncertainties from the two baseline scenarios.

**Keywords** First keyword · Second keyword · More

#### 1 Introduction

The Standard Model (SM) is a successful theory, well-established experimentally and theoretically. With the discovery of the Higgs boson [1, 2], the structure of the SM seems to be confirmed. However, some inconsistencies in the SM still need to be answered. For instance,

the striking mass hierarchy in the fermion sector. Moreover, while the dynamic of the SM gauge bosons, the photon,  $W$  and  $Z$  bosons, and gluons are governed by the gauge principle, the dynamic of the Higgs boson is different and unique in the SM. The SM does not predict the values of the Higgs couplings of quarks and leptons, nor the Higgs self-couplings. Large quantum corrections have to be canceled by fine-tuning the parameters to calculate the Higgs boson mass matching the measured value. One possible solution to this issue, achieving stabilization of the Higgs mass against quantum corrections, appears when the Higgs boson is associated with the zero mode of a dimension-five component of extensions of the SM gauge group. These models are referred to as gauge-Higgs unification (GHU) models.

The two most precise determinations of  $\sin^2 \theta_{eff}$  by the LEP and SLC differ in  $3.7\sigma$ , and none of them agrees with the SM prediction [3, 3]. In particular, the LEP value was extracted from the forward-backward asymmetry measurement for  $b$ -quarks with LEP1 data, and it is nearly three standard deviations away from the predicted value in the SM. Clarifying the  $A_{FB}^b$  value as well as exploring the possibility of BSM physics motivate the study of quark pair production in high energy  $e^-e^+$  collisions at future colliders not only at the  $Z$ -mass energy but also at higher energies. In the SM, these interactions are produced and mediated by a photon, a  $Z$ -boson, and the interference between them. Some BSM theories predict deviation of such couplings or even new sizable contributions to these processes from new mediators (such as heavy  $Z'$  resonances). These deviations would be accessible experimentally by probing high precision measurements of  $e^-e^+ \rightarrow q\bar{q}$  observables at different center-of-mass energies ( $\sqrt{s}$ ).

\*This work was carried out in the framework of the ILD concept group.

<sup>a</sup>Corresponding author: adrian.irles@ific.uv.es

<sup>b</sup>On leave from Tohoku University, Sendai, Japan



**Thanks for your attention!**

**back-up**



► <https://arxiv.org/pdf/2203.07622.pdf>

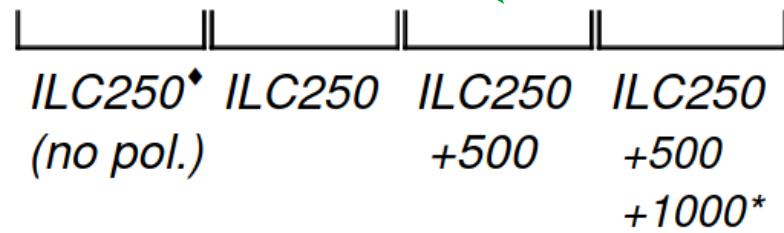
Quantity	Value	current $\delta[10^{-4}]$	Z pole		ILC250	
			$\delta_{stat}[10^{-4}]$	$\delta_{sys}[10^{-4}]$	$\delta_{stat}[10^{-4}]$	$\delta_{sys}[10^{-4}]$
boson properties						
$m_W$	80.379	1.5	-	-	0.08	0.3
$m_Z$	91.1876	0.23		0.022		-
$\Gamma_Z$	2.4952	9.4	0.5	-	6	-
$\Gamma_Z(had)$	1.7444	11.5		4.	-	-
Z-e couplings						
$1/R_e$	0.0482	24.	2.	5	5.5	10
$A_e$	0.1513	139.	1.5	1.2	12.	9.
$g_L^e$	-0.632	16.	1.0	3.2	2.8	7.6
$g_R^e$	0.551	18.	1.0	3.2	2.9	7.6
Z- $\ell$ couplings						
$1/R_\mu$	0.0482	16.	2.	2.	5.5	10
$1/R_\tau$	0.0482	22.	2.	2.	5.7	10
$A_\mu$	0.1515	991.	2.	5	54.	3.
$A_\tau$	0.1515	271.	2.	5.	57.	3
$g_L^\mu$	-0.632	66.	1.0	2.3	4.5	7.6
$g_R^\mu$	0.551	89.	1.0	2.3	5.5	7.6
$g_L^\tau$	-0.632	22.	1.0	2.8	4.7	7.6
$g_R^\tau$	0.551	27.	1.0	3.2	5.8	7.6
Z-b couplings						
$R_b$	0.2163	31.	0.4	7.	3.5	10
$A_b$	0.935	214.	1.	5.	5.7	3
$g_L^b$	-0.999	54.	0.32	4.2	2.2	7.6
$g_R^b$	0.184	1540	7.2	36.	41.	23.
Z-c couplings						
$R_c$	0.1721	174.	2.	30	5.8	50
$A_c$	0.668	404.	3.	5	21.	3
$g_L^c$	0.816	119.	1.2	15.	5.1	26.
$g_R^c$	-0.367	416.	3.1	17.	21.	26.

# GHU vs SM: Beam scenarios

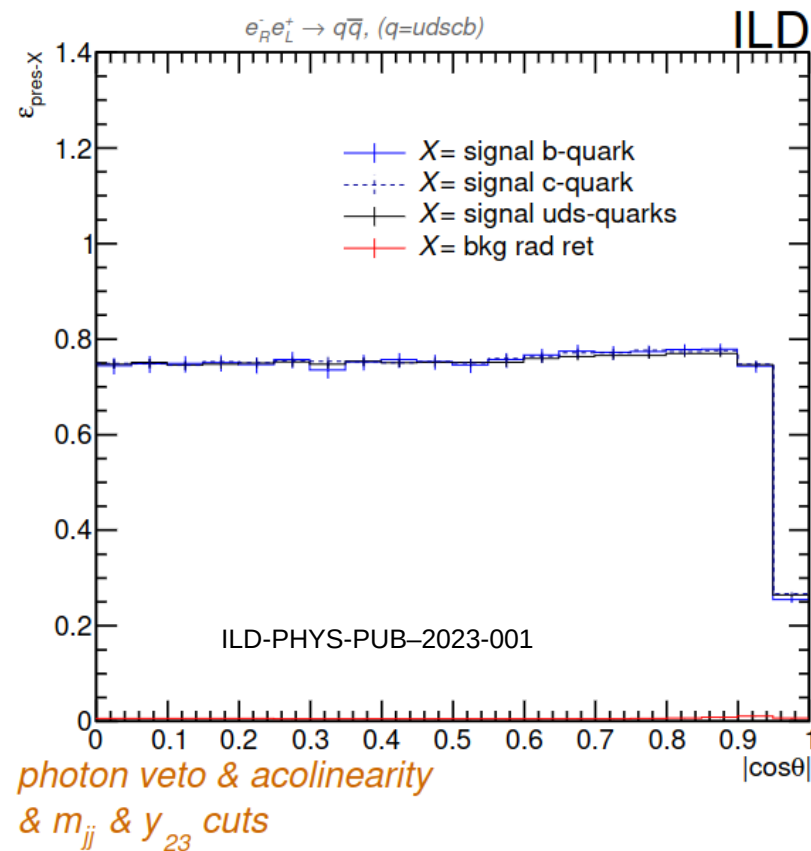
*Hypothetical case*  
**ILC250\* no pol**  
 $\int L = 2000\text{fb}^{-1}$   
OSP|SSP [%] = 45 | 5  
**Full ILD simulation**  
**assuming**  
**no beam pol**

**H20 nominal program**  
**ILC250**  
( $P_{e-}=0.8, P_{e+}=0.3$ )  
 $\int L = 2000\text{fb}^{-1}$   
OSP|SSP [%] = 45 | 5  
**ILC500**  
( $P_{e-}=0.8, P_{e+}=0.3$ )  
 $\int L = 4000\text{fb}^{-1}$   
OSP|SSP [%] = 40 | 10  
**Full ILD simulation**  
**assuming beam pol.**

**H20 nominal program**  
**ILC1000**  
( $P_{e-}=0.8, P_{e+}=0.2$ )  
 $\int L = 8000\text{fb}^{-1}$   
OSP|SSP [%] = 40 | 10  
  
*Not full simulation studies*  
*but extrapolations from ILC500*



- ▶ **Topology: 2 back-to-back jets (pencil-like topology)**
- ▶ **Preselection aiming for high background rejection and high efficiency.**
- ▶ Main bkg  $ee \rightarrow Z\gamma$  (radiative return through ISR)
  - $\sim x10$  larger than signal
  - **$\sim 90\%$  of such ISR photons are lost in the beam pipe**  $\rightarrow$  events filtered by energy & angular mom. conservation arguments
  - The **remaining  $\sim 10\%$  are filtered by identifying photons** in the detector (efficiency of  $>90\%$ )
  - PFA detector!!
- ▶ Other backgrounds from diboson production decaying hadronically are removed with extra topological cuts.





- Compare samples with 1 tag vs 2 tags (after preselection)

$$f_{1b} = \overline{\varepsilon_c} \overline{R_b} + \overline{\tilde{\varepsilon}_c} \overline{R_c} + \overline{\tilde{\varepsilon}_{uds}} (1 - \overline{R_b} - \overline{R_c})$$
$$f_{2b} = \overline{\varepsilon_b^2} (1 + \overline{\rho}) \overline{R_b} + \overline{\varepsilon_c^2} \overline{R_c} + \overline{\varepsilon_{uds}^2} (1 - \overline{R_b} - \overline{R_c})$$

Measured observables

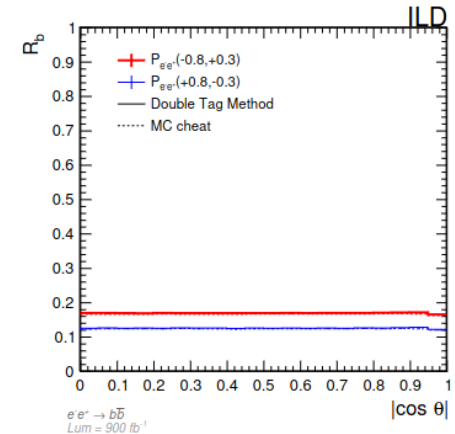
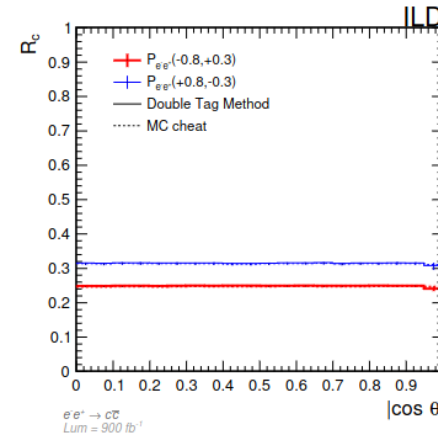
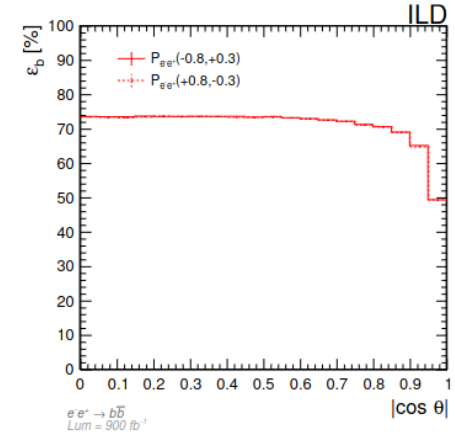
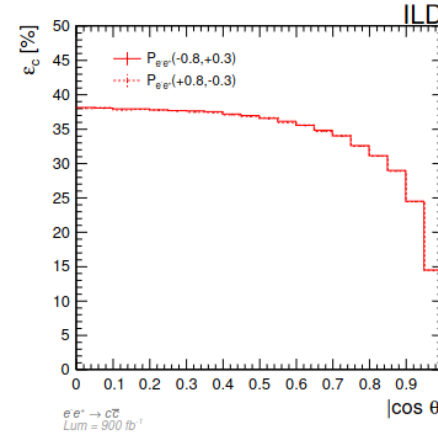
PHYSICS!  
Indirect observables

Inputs (MC or independent measurements)

**Similar set of equations  
for the c-quark  
solved simultaneously**

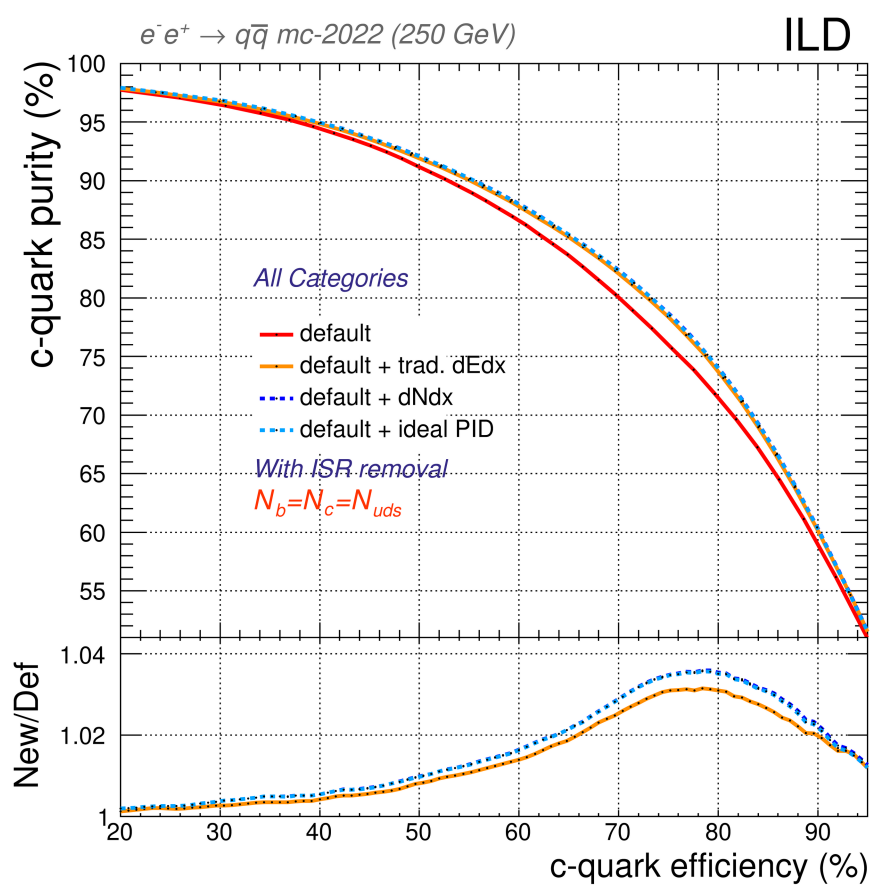
# Double flavour tagging – control of systematics

- ▶ Flavour tagging efficiency will be measured (double tagging)
  - Not estimated with MC
  - Per mil level reachable because the contamination from lighter quarks is minimal and the tight IP constraint
- ▶ Fully differential analysis !!
- ▶ Rb and Rc measured at the same time
  - than the tagging efficiencies
  - No assumption needed in  $R_{uds}$
- ▶ Per mil level stat. Uncertainty
- ▶ Comparable/lower exp syst. uncertainty
  - Dominated by flavour tagging and followed by angular correlations

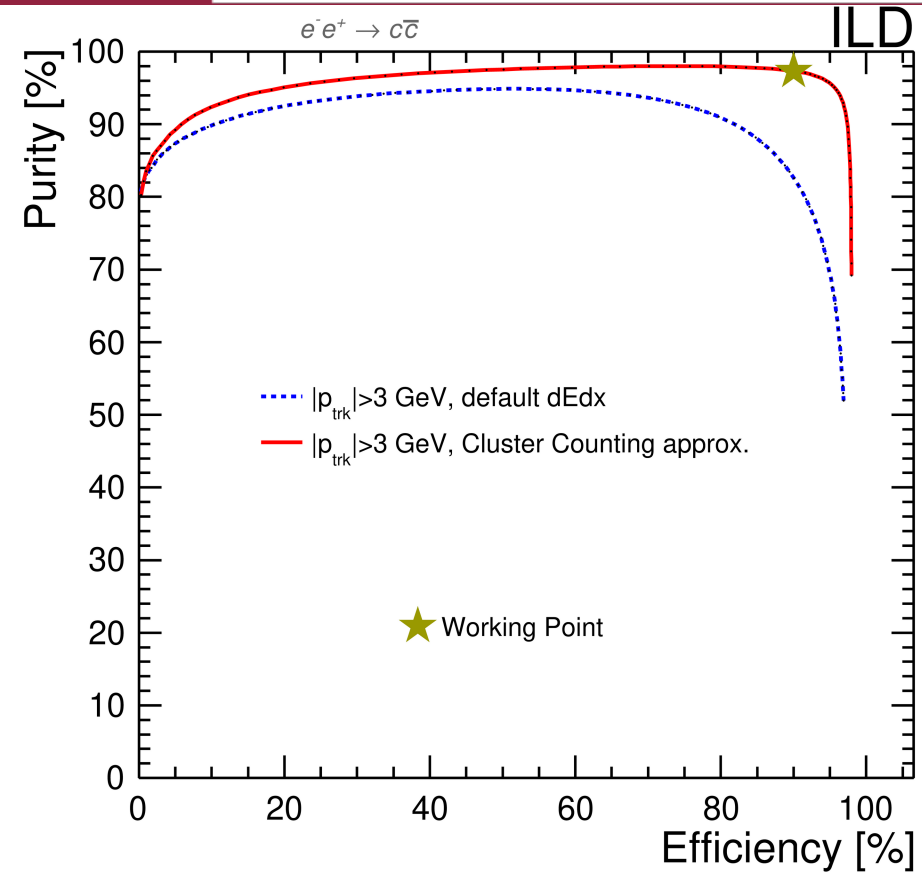


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# PID: From dEdx to dNdx



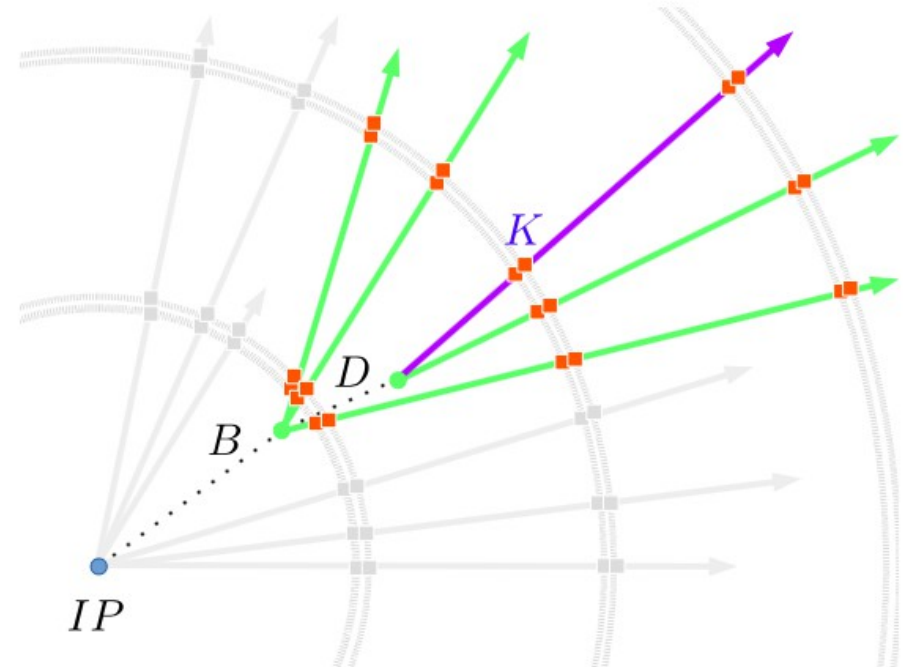
Effects in Flavour Tagging

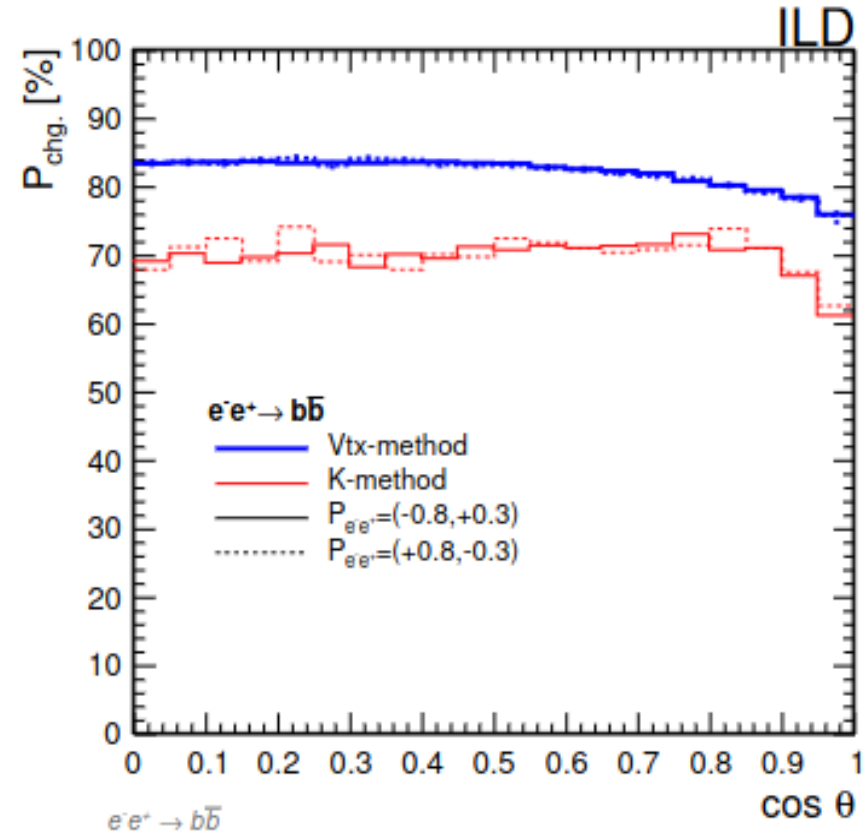
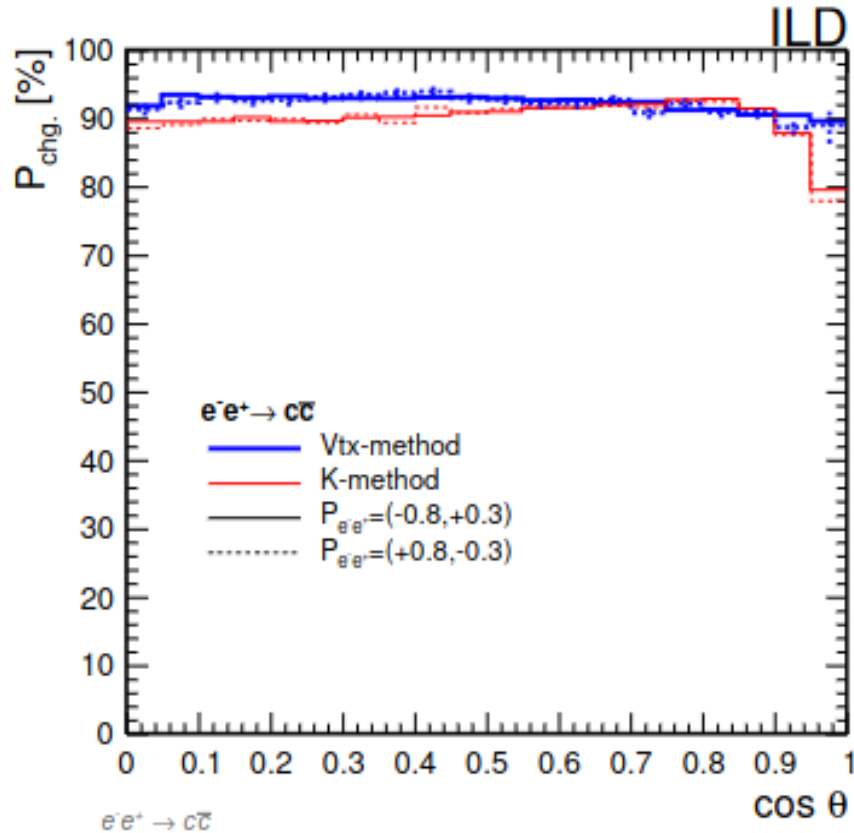


Effects in Kaon ID for charge reco.



- ▶ We start from a very pure & background-free **double tagged** sample
- ▶ We are required to **measure the jet charge**
  - Using K-ID and/or full Vtx charge measurement
  - K-ID is better suited for the C-quark (Vtx is better suited for b-quark)
- ▶ We use the **double charge** measurements
  - To control / reduce the systematic uncertainties





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# Double charge method

## ► Double Tag + Double Charge

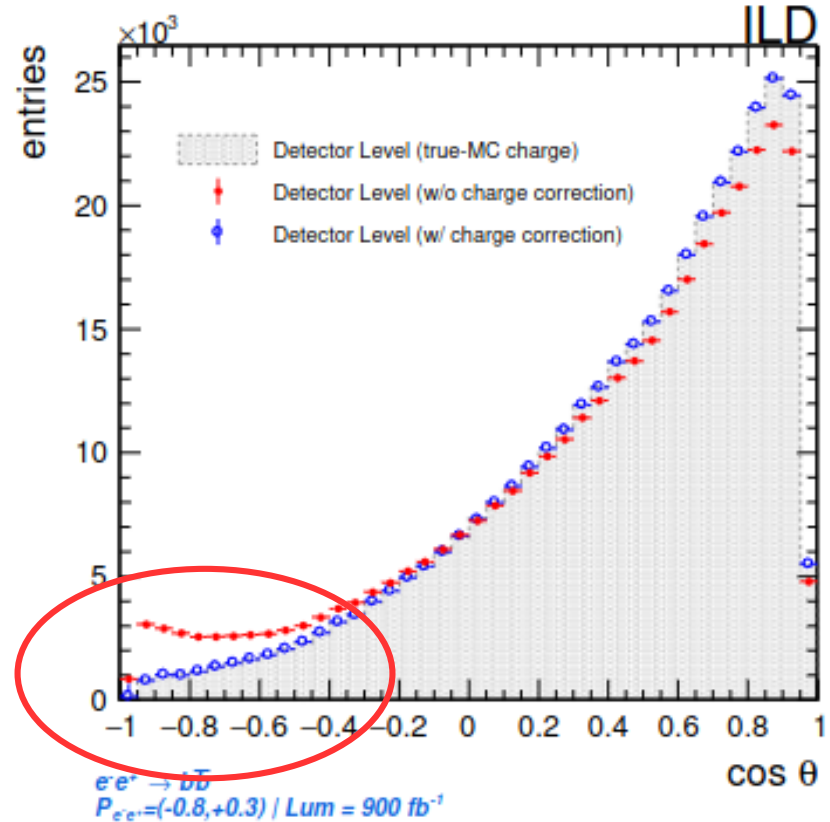
- Both jets need to have a charge measurement compatible with the 2 quarks back to back scenario

## ► Double mistakes are unlikely but still not negligible and lead to “sign flip” → migrations

BSM or simple migrations?

Red shows the distribution without sign correction.

Gray is the parton level distribution

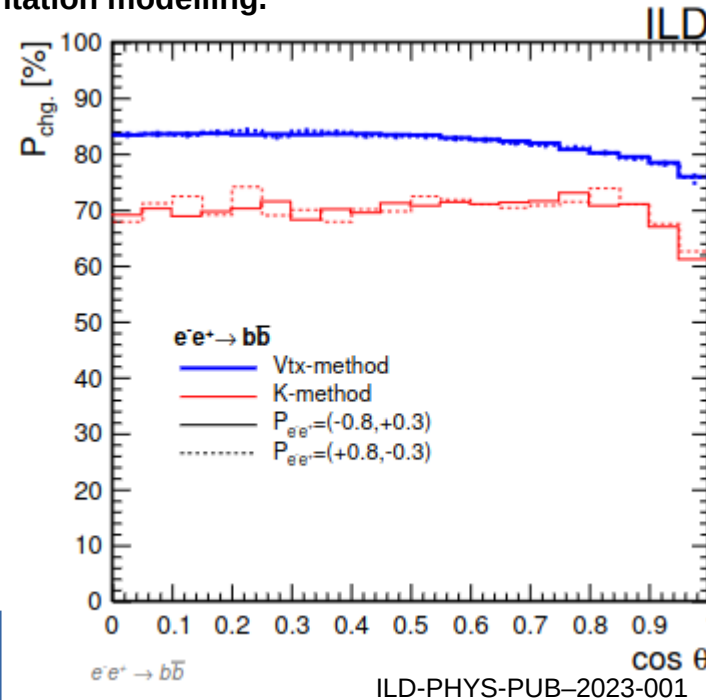
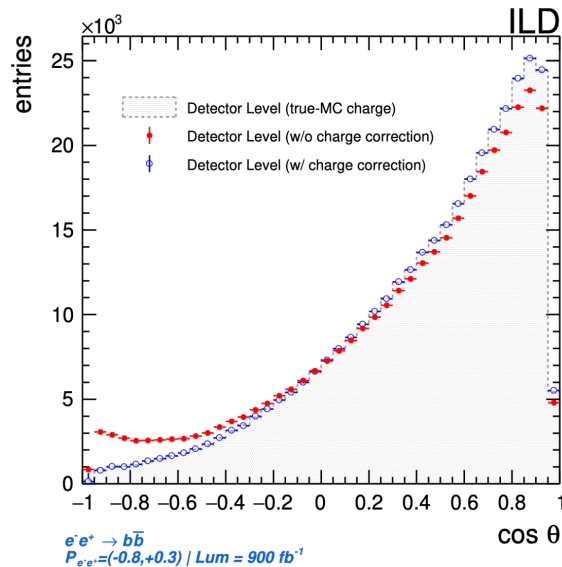


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# Migration correction

- ▶ Migrations look as “new physics” → we need to correct them
  - Using data: double charge measurements with same and opposite charges (see back-up slides)
  - We measure the probability to reconstruct correctly the charge ( $P_{\text{ch}}$ ) and use it for correction
  - DATA DRIVEN METHOD → non sensitive to fragmentation modelling.

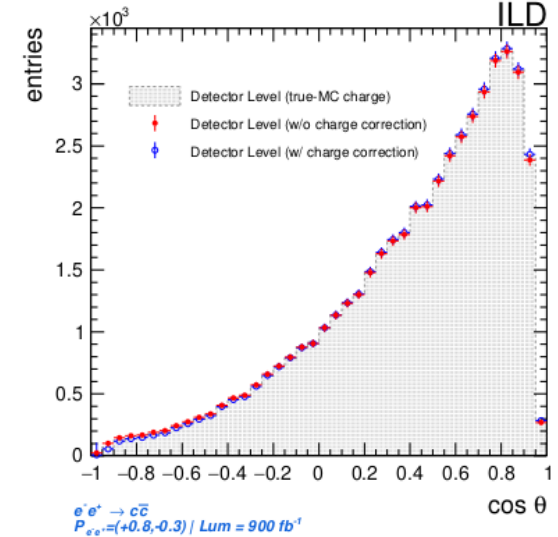
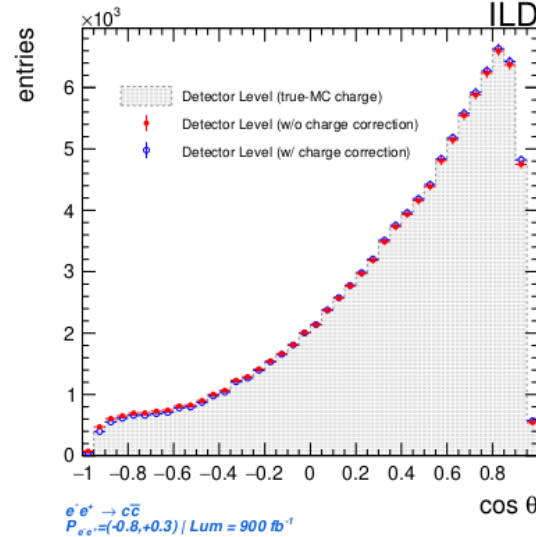
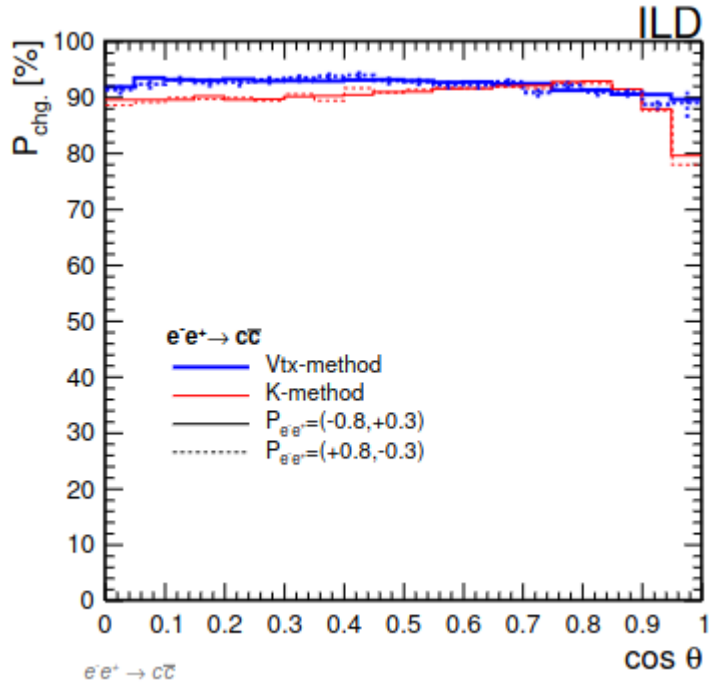


▶ **P<sub>chg</sub> limited by vertex reconstruction efficiency, Particle ID efficiency and B<sub>0</sub> oscillations (b-quark case).**

blue shows the distribution after sign correction.

Gray is the parton level distribution

# Migration correction – c quark case

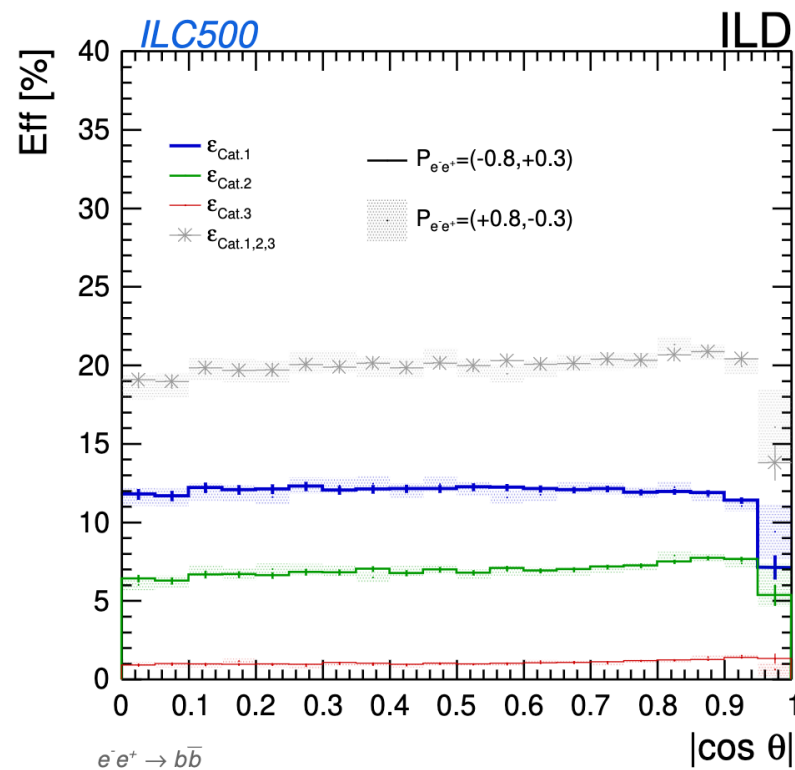
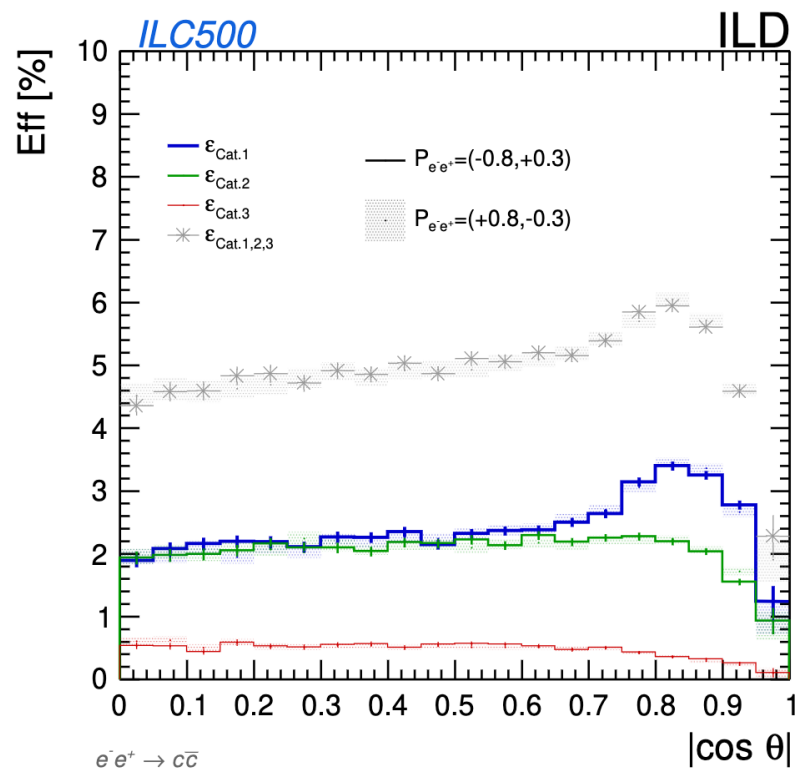


Minimal migration effects  
(and corrections!)



# Jet flavour tagging & charge measurement

► Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )



High-purity & independent samples for each quark flavour.

# Jet flavour tagging & charge measurement

- ▶ Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )
  - To maximally reduce the usage of MC tools (control of fragmentation, QCD correlations... uncertainties)

