Prospects for light exotic scalar measurements at the e⁺e⁻ Higgs factory

³ Bartłomiej Brudnowski¹, Kamil Zembaczyński¹, and Aleksander Filip Żarnecki^{1,*}

⁴ ¹Faculty of Physics, University of Warsaw, Pasteura 5, 02-093 Warsaw, Poland

Abstract. The physics program of the Higgs factory will focus on measure-5 ments of the 125 GeV Higgs boson, with the Higgs-strahlung process being the 6 dominant production channel at 250 GeV. However, production of extra light 7 scalars is still not excluded by the existing experimental data, provided their coupling to the gauge bosons is sufficiently suppressed. Fermion couplings of 9 such a scalar could also be very different from the SM predictions leading to 10 non-standard decay paterns. Presented in this contribution are results from the 11 ongoing studies on prospects of direct light scalar observation at future Higgs 12 factory experiments in different decay channels. 13

14 1 Motivation

In recent years, a general consensus was reached in the particle physics community about 15 the need for the next-generation large infrastructure to be an electron-positron Higgs fac-16 tory. It was indicated as the highest-priority next collider in the 2020 Update of the European 17 Strategy for Particle Physics [1]. While full exploitation of the Higgs boson physics requires 18 running at collision energies up to the TeV range, see figure 1 (left), most of the precision 19 measurements will be carried out at the collision energy of 240-250 GeV, maximizing the 20 cross section for the Higgs boson production in the so called Higgsstrahlung process, see 21 figure 1 (right). However, existence of additional light scalar particles, with masses of the 22 order of or below the mass of the 125 GeV state observed at the LHC, is by far not excluded 23 experimentally and also well motivated theoretically [3–5]. Higgs factories should be sen-24 sitive to exotic scalar production even for very light scalars and small couplings, thanks to 25 clean environment, precision and hermeticity of the detectors. Still, prospects for light scalar 26 measurements at the e⁺e⁻ Higgs factory were hardly studied in the past. That is why this 27 subject was included as one of the so called focus topics of the ECFA e⁺e⁻ Higgs/EW/top 28 factory study [6]. Two theoretical and phenomenological targets were defined: associated 29 production of the new scalar with the Z boson, $e^+e^- \rightarrow Z S$ (scalar-strahlung process) and 30 light scalar pair-production in 125 GeV Higgs boson decays, $H \rightarrow S S$. 31

2 Decay mode independent search

As for the SM Higgs boson, the production of new scalars in the scalar-strahlung process can be tagged, independent of their decay, based on the recoil mass technique [7]. For best

^{*}e-mail: filip.zarnecki@fuw.edu.pl

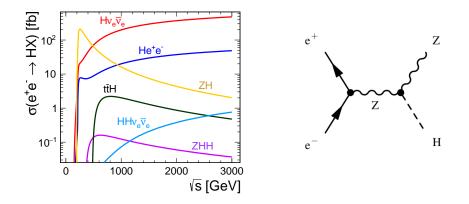


Figure 1. Left: cross section as a function of centre-of-mass energy for the main Higgs production processes at an e^+e^- collider. Right: the leading-order Feynman diagram for the Higgsstrahlung process. Figures taken from [2].

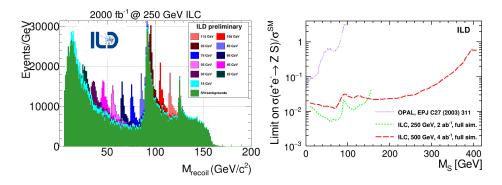


Figure 2. Left: the recoil mass distributions after the selection cuts for signal of light scalar production and SM backgrounds processes for ILC running at 250 GeV. Right: expected 95% C.L. limits on the scalar production cross section, relative to the SM scalar production cross section at given mass, for ILC running at 250 GeV and 500 GeV [8, 9].

recoil mass reconstruction Z decays to muon pair can be used, which were exploited in the full simulation studies performed within ILD [8, 9]. Shown in figure 2 (left) is the recoil mass distribution expected for SM background processed at 250 GeV ILC together with expectations for different signal hypothesis. Expected limits on the scalar production cross section, relative to the SM scalar production cross section at given mass, are compared with LEP limits based on the similar approach in figure 2 (right). Expected sensitivity of the experiment at

⁴⁰ Its based on the similar approach in figure 2 (right). Expected sensitivity of the experiment at ⁴¹ future e^+e^- Higgs factory is an order of magnitude better than the existing LEP limit and the

mass range probed can be significantly extended as well.

⁴³ 3 Search for S $\rightarrow \tau^+ \tau^-$

44 While limits resulting from the decay independent search based on the recoil mass distribu-

45 tion only are the most general ones, significant improvement of sensitivity is expected when

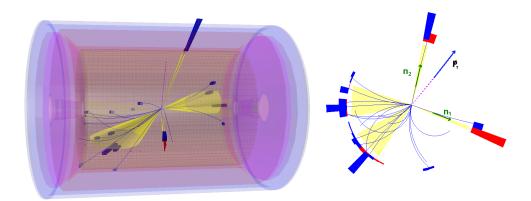


Figure 3. Left: example signal event, with hadronic decays of the two tau leptons produced in the light scalar decay. Right: same event in the transverse plane, missing transverse momentum \vec{p}_T and two unit vectors along tau jet directions (\vec{n}_1 and \vec{n}_2) are indicated.

particular decay channels of the new scalars are addressed, at the price of some model de-46 pendence. As some of the discrepancies from SM predictions observed in LEP and LHC 47 data suggested possible existence of the new scalar with mass of about 95 GeV and enhanced 48 branching ratio to the $\tau^+\tau^-$ final state [4], we decided to consider this decay channel as the 49 possible discovery scenario. Event samples used for the presented study were generated using 50 WHIZARD [10, 11] version 3.1.2. Both background (including 125 GeV Higgs boson produc-51 tion assuming nominal couplings) and light scalar signal production were modeled using the 52 built-in SM_CKM model. As signal samples we generated Higgs boson production with subse-53 quent decays to tau lepton pair, for sets of modified Higgs masses. Total lumionsity of 2 ab⁻¹ 54 was assumed for ILC running at 250 GeV, as expected in the H-20 running scenario [12], with 55 $\pm 80\%$ and $\pm 30\%$ polarisation for electron and positron beams, respectively. The ILC beam 56 energy profile, as simulated with GuineaPig [13], was taken into account based on CIRCE2 57 parametrisation and hadronisation was simulated with the PYTHIA 6 [14]. The fast detector 58 simulation framework DELPHES [15] was used to simulate detector response, with built-in 59 cards for parametrisation of the ILC detector, delphes_card_ILCgen.tcl [16]. Example 60 of signal event with hadronic final state, as simulated by DELPHES, is shown in figure 3 left. 61

Depending on the decays of the two tau leptons, three decay channels can be considered 62 for the signal events: hadronic (with both taus decaying hadronically), semi-leptonic (with 63 one leptonic tau decay) and leptonic (with leptonic decays of both taus). As a tight selection, 64 we require each tau lepton to be identified either as an isolated lepton (and missing p_T) or 65 hadronic jet with τ -tag. In addition to two tau candidates from the decay of the scalar, we 66 also require reconstruction of two (untagged) hadronic jets from the hadronic decay of the Z 67 boson. However, as the efficiency of tau jet tagging implemented in DELPHES is relatively poor 68 (at most 70%), we also consider loose event selection, when we require only one identified 69 tau candidate (isolated lepton or τ -tagged jet) and three untagged hadronic jets, and take the 70 jet with smaller invariant mass as the second tau candidate. At the pre-selection stage we thus 71 select five event categories, as summarised in table 1. 72

One of the challenges in the search for scalar decays into tau leptons is to properly reconstruct the invariant mass of the produced scalar, which can be significantly underestimated due to the escaping neutrinos. To correct for the neutrino energy, we use the so called collinear approximation [17]. For high energy tau leptons, decay products are highly boosted in the ini-

 Table 1. Event categories considered in the search for light scalar production with scalar decay to tau lepton pair.

Event	Isolated	Selection requirements	
category	leptons	tight selection	loose selection
hadronic	zero	4 jets, 2 with τ -tag	4 jets, 1 with τ -tag
semi-leptonic	one	3 jets, 1 with τ -tag	3 jets with no τ -tag
leptonic	two	two jets without τ -tag	

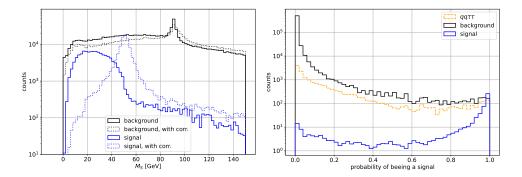


Figure 4. Left: reconstructed invariant mass of the two tau candidates after tight selection for SM background (black) and signal of 50 GeV scalar production (blue) before (solid) and after (dashed) collinear correction. Right: example of the BDT response distribution for 50 GeV scalar signal (blue) and SM background (black) events, for tight semi-leptonic event selection and ILC running with $e_L^-e_R^+$ polarisation combination. The signal cross section is normalized to 1% of the SM Higgs boson production cross section at given scalar mass. Orange dashed lines indicated the contribution of the leading background process, $e^+e^- \rightarrow qq\tau\tau$.

tial lepton direction and one can therefore assume that the initial tau lepton, escaping neutrino
 and the observed tau candidate are collinear. Neutrino energies can be found from transverse
 momentum balance:

80

$$\vec{p}_{T} = E_{\nu_{1}} \cdot \vec{n_{1}} + E_{\nu_{2}} \cdot \vec{n_{2}}$$
(1)

where $\vec{n_1}$ and $\vec{n_2}$ are directions of the two tau candidates in the transverse plane (see the right plot in figure 3). While more advanced reconstruction methods exist, based on the reconstruction of secondary decay vertex position, the advantage of this method is that it can be applied to all events and the solution is unique.

Shown in figure 4 (left) are the mass distributions of the tau candidate pairs in signal and
 background events, before (solid) and after (dashed) collinear correction. After the correction, the scalar mass can be reconstructed with about 5 GeV precision, also for semi-leptonic
 and leptonic events.

For best event classification, we consider each event category (see table 1) and each beam polarisation combination separately, resulting in 20 independent BDTs trained for event classification, for each scalar mass considered. Example of BDT response distribution is presented in figure 4.

Expected exclusion limits, assuming no deviation from SM predictions are observed, are calculated from the Hessian matrix of the template fit to the BDT response distributions. Final results of the study are presented in figure 5. As expected, the semi-leptonic event se-

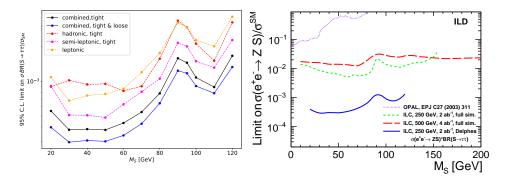


Figure 5. Expected 95% C.L. cross section limits on the light scalar production cross section times di-tau branching ration for ILC running at 250 GeV. Left: comparison of combined limits with limits obtained for different event categories. Right: combined limits compared with limits resulting from decay independent study.

⁹⁶ lection results in the strongest limits, combining high event statistics (about 47% of decays) ⁹⁷ with background lower than for the hadronic channel. Including loose selection categories ⁹⁸ improves the expected limits by 20-30% for the whole considered mass range. When com-⁹⁹ paring to the decay independent limits, we can state that the targeted analysis results in over ¹⁰⁰ order of magnitude increase in sensitivity. However, as the limit includes the branching ratio, ¹⁰¹ we can conclude that the di-tau channel is more sensitive only if the branching ratio is of the ¹⁰² order of 10% or above.

¹⁰³ 4 Search for $S \rightarrow b\bar{b}$

If the structure of new light scalar couplings to SM particles is similar to that of the SM 104 Higgs boson then the decay to $b\bar{b}$ is expected to dominate down to the masses of the order 105 of 10 GeV. As huge hadronic background is expected from pair production of W bosons, we 106 focus on leptonic Z boson decays in this search channel. Again, invariant mass of the new 107 scalar, as reconstructed directly from the two b-quark jets, is poorly measured due to neutri-108 nos escaping in semi-leptonic heavy meson decays. However, as leptons from Z decays can 109 be very precisely measured, we can use conservation of transverse momentum to reconstruct 110 jet energies from leptonic final state and jet angles. This approach was first proposed for 111 Higgs mass measurement at the ILC [18], but we verified that it works very well also for light 112 scalar production. This is shown in figure 6, where reconstructed di-jet mass distributions are 113 compared before and after correction for 125 GeV Higgs boson and new scalar of 50 GeV. 114 Mass reconstructed with the recoil method (based on the Z decay measurement only) is also 115 included for comparison. One can clearly see that jet energy correction based on the trans-116 verse momentum conservation significantly improves scalar mass measurement and allows 117 to obtain precision much better than with uncorrected jet energies or the recoil method. The 118 study of light scalar production in bb decay channel is ongoing and the expected exclusion 119 limits will be presented at the 3rd ECFA workshop in Paris. 120

121 5 Conclusions

BSM scenarios involving light scalars, with masses accessible e^+e^- Higgs factories, are still not excluded by existing data. Sizable production cross sections for new scalars can also

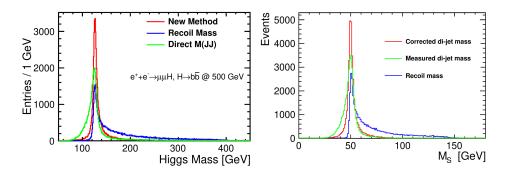


Figure 6. Invariant mass of two jets reconstructed using raw jet energies, corrected jet energies and recoil method. Left: full simulation results for the 125 GeV Higgs boson production at 500 GeV ILC [18]. Right: DELPHES simulation results for the signal of 50 GeV scalar production at 250 GeV ILC.

coincide with non-standard decay patterns, so different decay channels should be considered.
 Strong limits are already expected from decay independent searches, based on recoil mass
 reconstruction. These are expected to be improved further with new analysis methods and
 the corresponding study is ongoing. Compared to the decay independent search, more than
 an order of magnitude limit improvement in search sensitivity is expected for light scalar
 decays to tau pairs, if this is a dominant decay channel. Studies of other decay channels are
 ongoing and the results are being prepared for the ECFA study report.

This work was carried out in the framework of the ILD concept group as a contribution to the ECFA e^+e^- Higgs/EW/top factory study. This work was supported by the National Science Centre, Poland, under the OPUS research project no. 2021/43/B/ST2/01778.

134 References

- [1] E.S. Group, 2020 Update of the European Strategy for Particle Physics (CERN Council, Geneva, 2020), ISBN 978-92-9083-575-2
- [2] H. Abramowicz et al., Higgs physics at the CLIC electron–positron linear collider, Eur.
 Phys. J. C 77, 475 (2017), 1608.07538. 10.1140/epjc/s10052-017-4968-5
- [3] S. Heinemeyer, C. Li, F. Lika, G. Moortgat-Pick, S. Paasch, Phenomenology of a
 96 GeV Higgs boson in the 2HDM with an additional singlet, Phys. Rev. D 106, 075003
 (2022), 2112.11958. 10.1103/PhysRevD.106.075003
- ¹⁴² [4] T. Biekötter, S. Heinemeyer, G. Weiglein, Mounting evidence for a 95 GeV Higgs bo-¹⁴³ son, JHEP **08**, 201 (2022), 22**03**.13180. 10.1007/JHEP08(2022)201
- [5] T. Robens, A Short Overview on Low Mass Scalars at Future Lepton Colliders, Universe
 8, 286 (2022), 2205.09687. 10.3390/universe8050286
- [6] J. de Blas et al., Focus topics for the ECFA study on Higgs / Top / EW factories (2024),
 2401.07564.
- [7] J. Yan, S. Watanuki, K. Fujii, A. Ishikawa, D. Jeans, J. Strube, J. Tian, H. Yamamoto, Measurement of the Higgs boson mass and $e^+e^- \rightarrow ZH$ cross section using $Z \rightarrow \mu^+\mu^-$
- and $Z \to e^+e^-$ at the ILC, Phys. Rev. D **94**, 113002 (2016), [Erratum: Phys.Rev.D 103, 000002 (2021)] 1694, 97524, 10, 1102 (Phys.Rev.D 04, 112002)
- ¹⁵¹ 099903 (2021)], 1604.07524. 10.1103/PhysRevD.94.113002

- [8] Y. Wang (International Large Detector Concept Group), Search for Light Scalars Produced in Association with a Z boson at the 250 GeV stage of the ILC, PoS ICHEP2018, 630 (2019). 10.22323/1.340.0630
- [9] Y. Wang, M. Berggren, J. List, ILD Benchmark: Search for Extra Scalars Produced in Association with a Z boson at $\sqrt{s} = 500 \text{ GeV} (2020)$, 2005.06265.
- [10] M. Moretti, T. Ohl, J. Reuter, O'Mega: An Optimizing matrix element generator, pp. 1981–2009 (2001), hep-ph/0102195.
- ¹⁵⁹ [11] W. Kilian, T. Ohl, J. Reuter, WHIZARD: Simulating Multi-Particle Processes at LHC
- and ILC, Eur. Phys. J. C71, 1742 (2011), 0708.4233. 10.1140/epjc/s10052-011-1742-y
- [12] P. Bambade et al., The International Linear Collider: A Global Project (2019),
 1903.01629.
- [13] D. Schulte, Beam-beam simulations with GUINEA-PIG, CERN-PS-99-014-LP,
 https://cds.cern.ch/record/382453/files/ps-99-014.pdf (1999)
- [14] T. Sjostrand, S. Mrenna, P.Z. Skands, PYTHIA 6.4 Physics and Manual, JHEP 05, 026
 (2006), hep-ph/0603175. 10.1088/1126-6708/2006/05/026
- [15] J. de Favereau et al., DELPHES 3, A modular framework for fast simulation of a generic collider experiment, JHEP 02, 057 (2014), 1307.6346. 10.1007/JHEP02(2014)057
- [16] A.F. Zarnecki, *The ILC DELPHES card* (IDT-WG3 Software Tutorial), https://agenda.
 linearcollider.org/event/9264/ (2021)
- [17] [17] S. Kawada, K. Fujii, T. Suehara, T. Takahashi, T. Tanabe, A study of the measurement precision of the Higgs boson decaying into tau pairs at the ILC, Eur. Phys. J. C 75, 617 (2015), 1509.01885. 10.1140/epjc/s10052-015-3854-2
- ¹⁷⁴ [18] Junping Tian, A new method for measuring the Higgs mass at the ILC, ILD-PHYS-¹⁷⁵ PUB-2019-001 (2020)