

Long-lived particle searches with the ILD experiment

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> D. Jeans⁽¹⁾, <u>J. Klamka⁽²⁾</u>, A. F. Żarnecki⁽²⁾ ⁽¹⁾KEK, ⁽²⁾University of Warsaw

> > jan.klamka@fuw.edu.pl

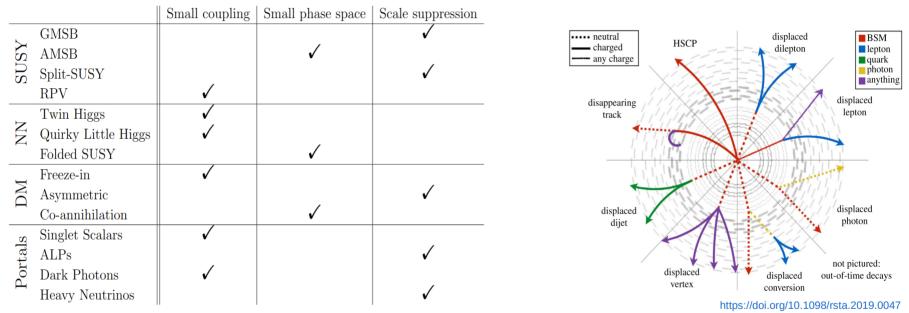


Motivation



Particles with macroscopic lifetimes naturally appear in numerous BSM models

Three main mechanisms are responsible for that... 1810.12602



Multiple LLP searches at the LHC, sensitive to high masses and couplings

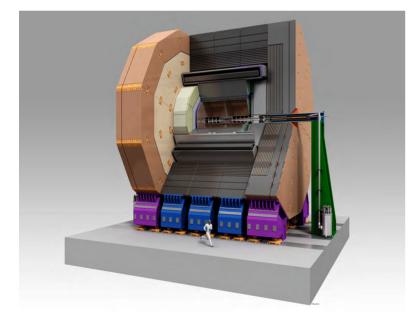
- <u>complementary region</u> could be probed at e⁺e⁻ colliders (small masses, couplings, mass splittings)
- typical properties of feebly interacting massive particles (FIMPs) \rightarrow challenging for hadron colliders

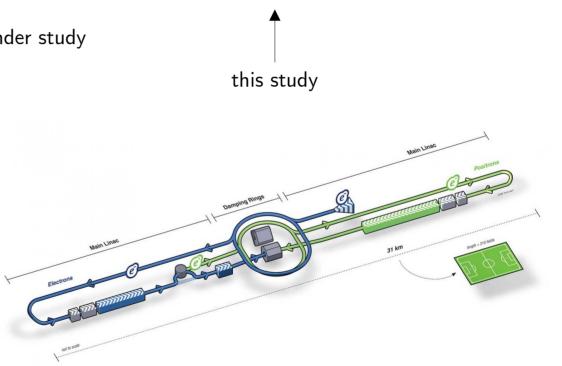


International Large Detector (ILD)



- Multi-purpose detector for an e⁺e⁻ Higgs Factory (HF)
- Example: the International Linear Collider (ILC), with baseline c.m.s. energy **<u>250</u>**-500 GeV
- Possible operation at other HF proposals now under study



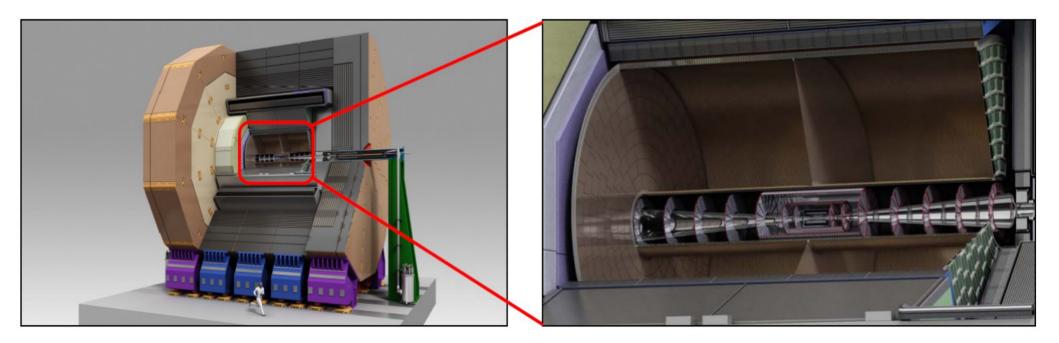




International Large Detector (ILD)



- Nearly 4π angular coverage, optimised for particle flow
- Time projection chamber (TPC) as the main tracker allows for continuous tracking and dE/dx PID
- High granularity calorimeter with minimal material in front of it inside 3.5 T solenoid



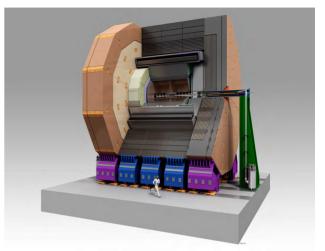


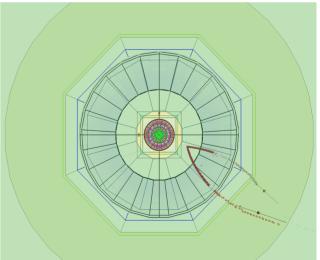




ILD especially promising with a $\underline{\mathsf{TPC}}$ as the main tracker

- \rightarrow we want to investigate experimental aspects
- \rightarrow study based on full simulation
- Study such challenging signatures from the **experimental perspective**
 - → experimental/kinematic properties, not points in a model parameter space
- Focus on a generic (and most challenging) case two tracks from a displaced vertex
- No other assumptions about the final state, approach as general as possible







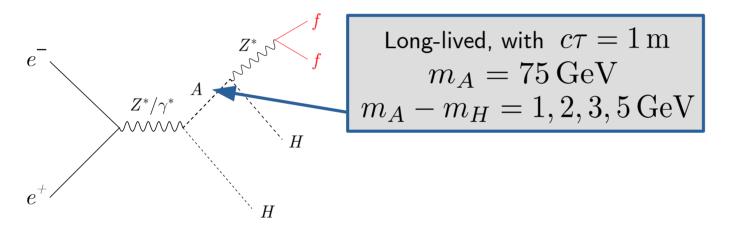
Framework and signatures



 $\sqrt{s} = 250 \,\mathrm{GeV}$

As a challenging case (small boost, low-pT final state) we considered:

ightarrow heavy scalar LLP (A) and DM (H) pair-production with small mass splitting, $Z^*
ightarrow \mu\mu$





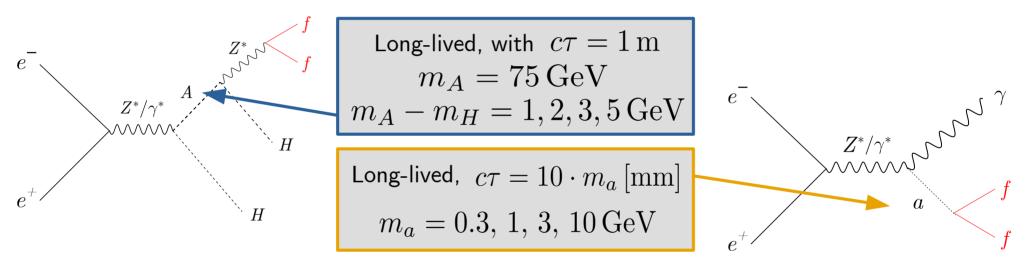
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The opposite extreme case, (<u>large boost, high-pT final state</u>)

 \rightarrow light pseudoscalar LLP $a \rightarrow \mu \mu$

Very simple vertex finding (inside the TPC) based on a distance between track pairs

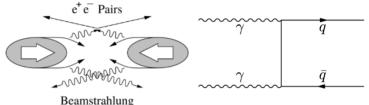


Overlay events background



At linear e⁺e⁻ colliders beams are strongly focused and radiate photons, so $\gamma\gamma$ interactions also occur in detector. On average, in each bunch-crossing (BX) at ILC250, produced are:

- 1.55 γγ → low-p_T hadrons events
- **O(10⁵) incoherent e⁺e⁻ pairs**, only a small fraction enters detector



These events are soft, usually important because they **overlay** on physical events

...but can also look like signal on their own



Overlay events background



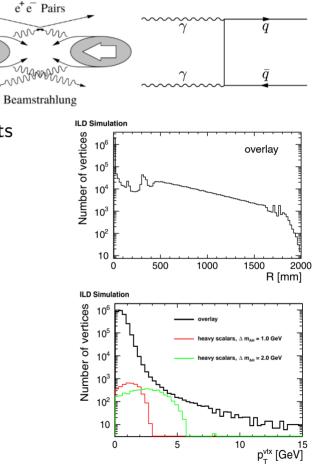
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- ~10¹¹ BXs per year at ILC \rightarrow overwhelming number of overlay events
- Similar kinematics to the signal considered and can be busy
 - \rightarrow many secondary vertices (mostly fake, also V⁰s and photon conversions)
 - \rightarrow significant background





Overlay events background



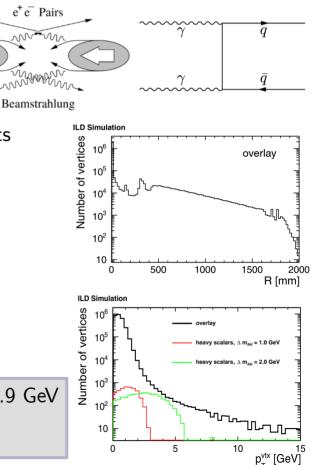
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- 1.55 γγ → low-p_T hadrons events
- $O(10^5)$ incoherent e^+e^- pairs, only a small fraction enters detector

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- ~10¹¹ BXs per year at ILC \rightarrow overwhelming number of overlay events
- Similar kinematics to the signal considered and can be busy
 - \rightarrow many secondary vertices (mostly fake, also V^os and photon conversions)
 - \rightarrow significant background
 - Can be suppressed using cuts on the track pair geometry and $p_{_{\rm T}}^{_{_{\rm Vtx}}} > 1.9~\text{GeV}$
 - Total expected reduction factor at the level of $\sim 10^{-10}$





Background from high-p_T events



Backg. sources occur mainly inside jets, so we consider (hard) e^+e^- and $\gamma\gamma$ processes with jets in final state

The following survive overlay selection in the hard e^+e^- processes:

- Decays of kaons, lambdas, photon conversions
- Secondary tracks from interactions with detector material

Selection eff. depends on number of jets, so:

Estimate efficiency based on full simulation

Use qq efficiency for the remaining processes

	$\operatorname{sgn}(\operatorname{P}(\operatorname{e}^{-}), \operatorname{P}(\operatorname{e}^{+}))$	(-,+)	(+, -)	(-, -)	(+,+)	
	$\operatorname{channel}$	σ [fb]				
	qq	$127,\!966$	$70,\!417$	0	0	
	qqqq	$28,\!660$	970	0	0	
	$\mathrm{q}\mathrm{q}\ell u$	29,043	261	191	191	
	${ m ZZ} ightarrow { m qq} \ell \ell, { m qq} u u$	838	467	0	0	
	$Z\nu_e\nu_e o qq\nu_e\nu_e$	454	131	0	0	
	$\text{Zee} \rightarrow \text{qqee}$	$1,\!423$	$1,\!219$	$1,\!156$	$1,\!157$	
	process	BB	BW	WB	WW	
	hard $\gamma^{B/W}\gamma^{B/W}$	$42,\!150$	90,338	90,120	71,506	



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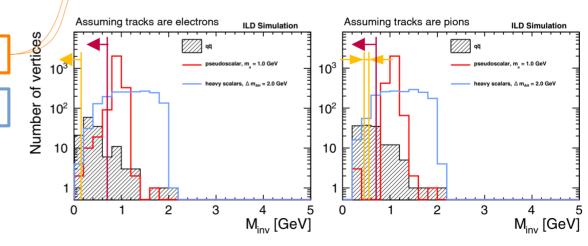
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→ Additional cuts on invariant mass are applied, with two working points: **standard** and **tight** (tight involving also **isolation** criterium)

	$\operatorname{sgn}(\operatorname{P}(\operatorname{e}^{-}), \operatorname{P}(\operatorname{e}^{+}))$	(-,+)	(+, -)	(-, -)	(+,+)	
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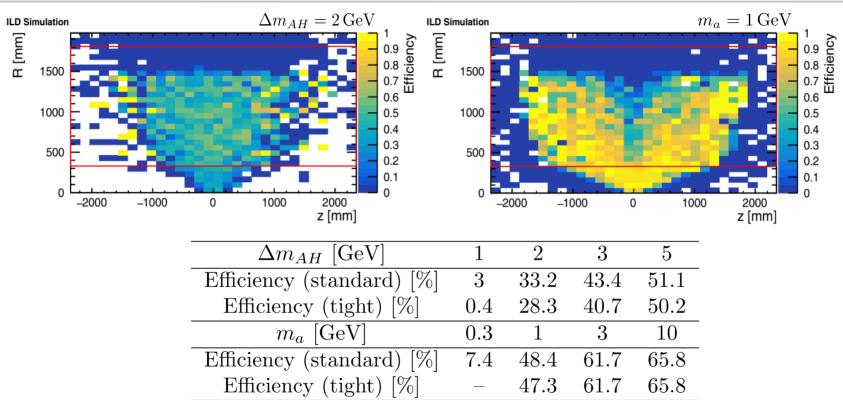


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Vertex finding results





- Efficiency = (correct / decays within TPC acceptance), "correct" if distance to the true vtx < 30 mm
- Signal selection depends strongly on the mass splitting (Z* virtuality) and mass of a (final state boost)
- A dedicated approach could enhance sensitivity for $\Delta m_{\text{AH}}=1$ GeV and $m_{\text{a}}=300$ MeV scenarios

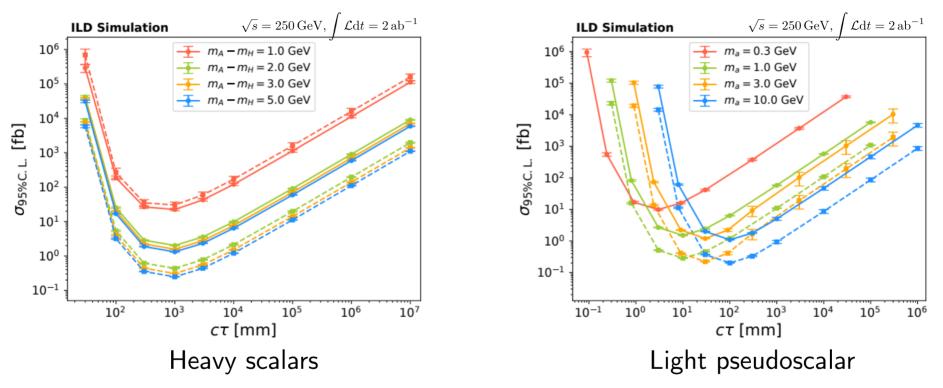
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Cross section limits





- Tight selection: dashed line, standard selection: solid line
- A wide range of models with heavy scalars with small mass splittings, or light pseudo scalar particles, can be excluded down to 0.1 fb

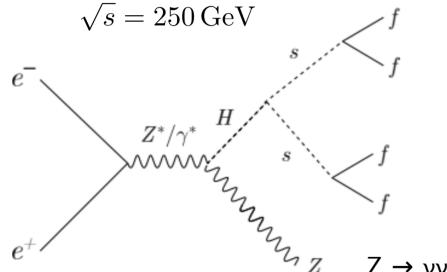


Higgs decays to LLPs



Higgsstrahlung with H(125) decay to two long-lived scalars

Generated using the Triple Real Singlet Higgs model with fixed lifetimes of s



Generated scenarios:

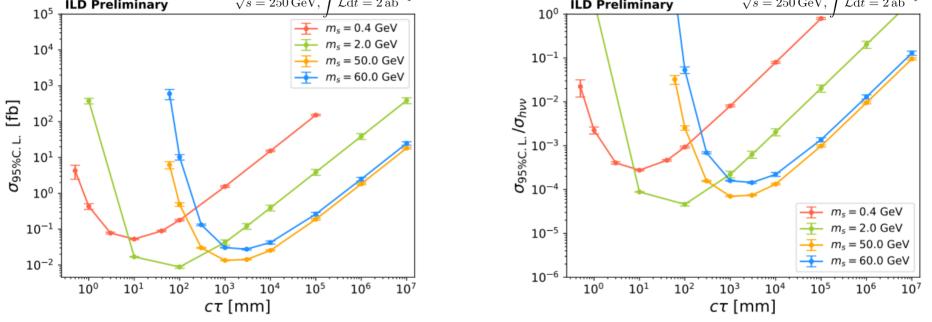
$$m_s = 400 \text{ MeV}, c\tau = 10 \text{ mm}$$
$$m_s = 2 \text{ GeV}, c\tau = 10 \text{ mm}$$
$$m_s = 50 \text{ GeV}, c\tau = 1 \text{ m}$$
$$m_s = 60 \text{ GeV}, c\tau = 1 \text{ m}$$

 $Z \rightarrow \nu\nu$, s $\rightarrow \mu\mu$ decays used to simplify the simulation

Use the <u>same analysis procedure</u>, but further <u>optimise for this channel</u> by requiring: \rightarrow no additional prompt tracks with pT > 2 GeV

 \rightarrow total pT > 10 GeV of tracks forming a vertex (to neglect the overlay)





- ILD can improve the current constraints and probe higher lifetimes already @ ILC250 thanks to higher TPC acceptance
- The limits could be further improved by dedicated searches using vertex detector and by more data at higher energy stages

SHIV



Conclusions



- ILD has a good potential to study long-lived particles, considering the model-independent approach and extreme signatures tested
- TPC plays the key role by enhancing the acceptance, allowing to probe very high lifetimes
- Additional selection utilizing features of a given signature can greatly improve sensitivity
- Presented expected limits on SM-like Higgs decays to LLPs would improve current constraints by order of magnitude or probe longer lifetimes





BACKUP

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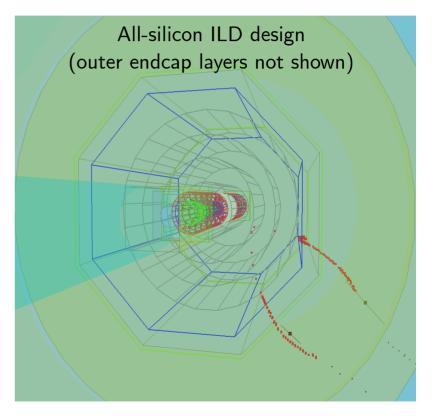


Alternative all-silicon ILD design



<u>Alternative ILD design</u> implemented for tests

- **TPC replaced** by the **silicon Outer Tracker**, modified from the CLICdet
- One **barrel layer** added and **endcap layers spacing** increased w.r.t. CLICdet
- Conformal tracking algorithm (designed for CLICdet) used for reconstruction at all-silicon ILD



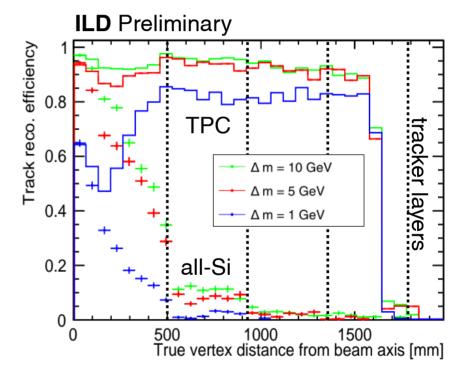
 \rightarrow Check how the **results** for <u>heavy scalars</u> are influenced by a **change of tracker** design



Heavy scalars at all-silicon ILD



- <u>Vertex reconstruction</u> driven by **track reconstruction efficiency**
- Performance similar to baseline design (TPC) <u>near</u> <u>the beam axis</u>
- Smaller number of hits available → efficiency drops faster with vertex displacement
- At least 4 hits required for track reconstruction
 → limited reach
- For large decay lengths, efficiency significantly higher for "standard" ILD with TPC



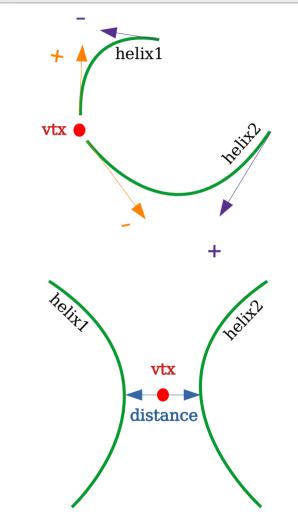


Vertex finding strategy



Approach as simple and general as possible:

- Consider tracks in pairs
- As the TPC is not sensitive to track direction:
 - \rightarrow use both track direction (charge) hypothesis for vertex finding
 - \rightarrow consider opposite-charge track pairs only
 - \rightarrow select pair with closest starting points
- Reconstruct vertex in between points of closest approach of helices
 - \rightarrow Require distance < 25 mm

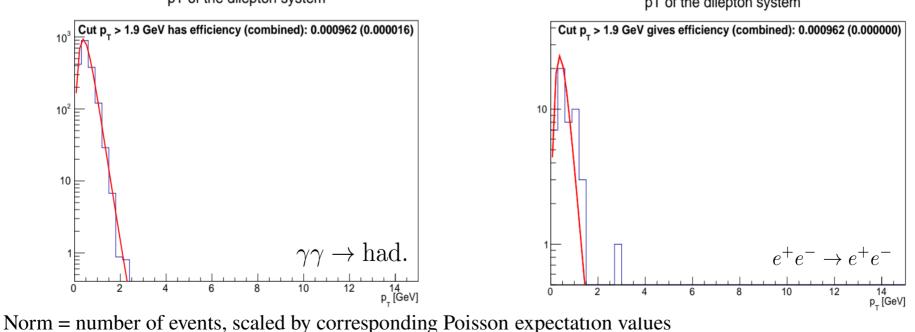




Final selection – pT



- We consider $yy \rightarrow had$. and e^+e^- samples separately
- Estimated background eff. from fitted distributions ~10⁻³ (~10⁻⁵–10⁻⁷ with preselection)
- Very small statistics in e^+e^- sample after preselection \rightarrow fit shape from $\gamma\gamma \rightarrow$ had. with floating normalisations



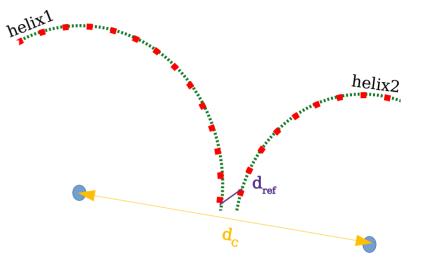
pT of the dilepton system

pT of the dilepton system

Final selection – other variables



- At least one more (independent) variable needed to achieve the assumed reduction
- We expect that signal tracks should come out of a single point → reference points should be close
- In busier backgound events, still many tracks evade the cuts e.g. curlers, secondary decays
- \rightarrow either far reference points or close centres of helices



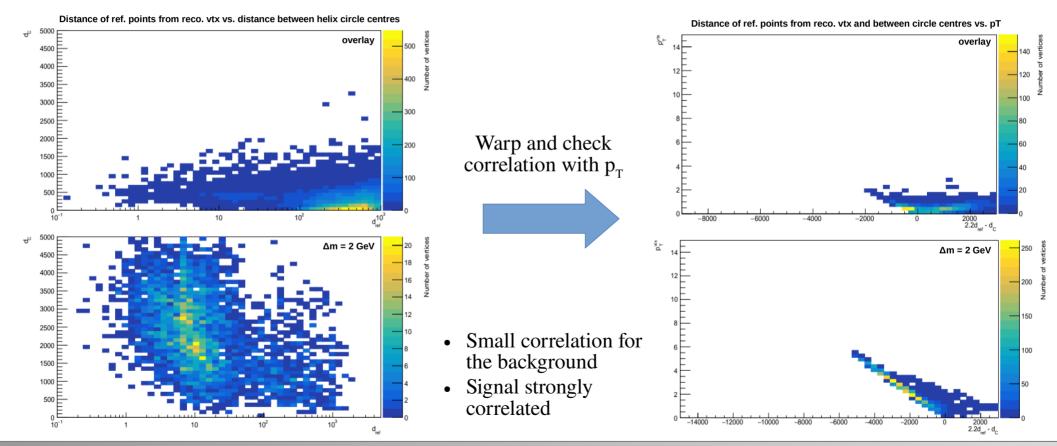
- **d**_{ref} distance between reference points (TrackStates / first hits)
- d_c distance between centres of helices projections into XY plane



Final selection – second variable



- New variable(s) should be uncorrelated with pT to make the cuts independent
- $2.2d_{ref} d_C$ good for optimal signal-background separation \rightarrow use it to look for correlation



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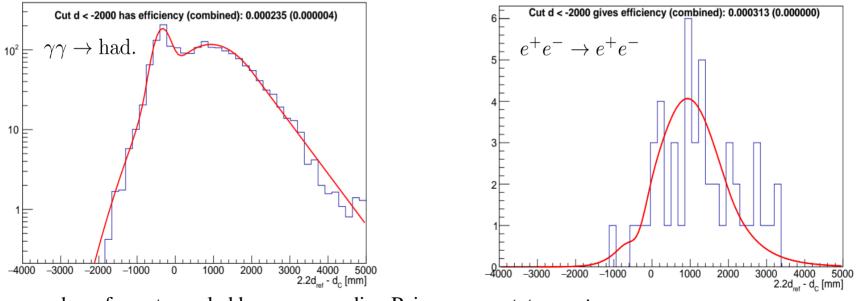
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Final selection – second variable



- Same approach as for the pT
- For $2.2d_{ref} d_{C} \le -2000 \text{ mm}$, signal eff. $\sim 37\%$ ($\Delta m = 2 \text{ GeV}$)
- Estimated background eff. from fitted distributions ~10⁻⁴ (~10⁻⁶–10⁻⁷ with preselection)
- Total expected efficiency at the level of $\sim 10^{-9}$ ($\sim 10^{-10}$) for $\gamma\gamma \rightarrow had.$ (e^+e^- pairs)



Norm = number of events, scaled by corresponding Poisson expectation values

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For small correlations r between x and y, total selection efficiency can be described as

$$\epsilon_{xy} = \epsilon_y^{(1-r)} \epsilon_x, \ \epsilon_x > \epsilon_y$$

For cuts on \mathbf{p}_{T} and $\mathbf{2.2d}_{ref} - \mathbf{d}_{C}$ (slide 5), assuming **30%** correlation, for $\gamma\gamma \rightarrow$ had. (e⁺e⁻ pairs) that gives:

• 2.8·10⁻⁶ (3.4·10⁻⁶)

• $4.6 \cdot 10^{-8} (1.7 \cdot 10^{-9}) \leftarrow$ combined with preselection

Combined cut efficiency $x > 2 \cap y > 3$

