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## Z-boson couplings to quarks at Higgs factories

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Electroweak Precision Measurements are stringent tests of the Standard Model and sensitive probes to New Physics. Accurate studies of the Z-boson couplings to the first-generation quarks could reveal potential discrepancies between the fundamental theory and experimental data. Future  $e^+e^-$  colliders running at the Z pole would be an excellent tool to perform such a measurement, unlike the LHC where hadronic Z decays are only available in boosted topologies. The measurement is based on comparison of radiative and non-radiative hadronic decays. Due to the difference in quark charge, the relative contribution of the events with final-state radiation (FSR) directly reflects the ratio of decays involving up- and down-type quarks. Such an analysis requires proper modeling and statistical discrimination between photons coming from different sources, including initial-state radiation (ISR), FSR, parton showers and hadronisation. In our contribution, we show how to extract the values of the Z couplings to light quarks and present the estimated measurement uncertainties.

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## 28 1 Introduction

29 The Standard Model (SM) of particle physics is the most successful theory describing interactions at the  
30 fundamental level. Despite its ability to explain experimental data from various experiments, several cosmolo-  
31 gical observations suggest that it is not the ultimate theory. Furthermore, the SM is a theory consistent with  
32 data only when many internal parameters, such as particle masses and couplings, are set to some particular  
33 values. They have to be measured experimentally and thus, precision studies are needed to complete our  
34 knowledge of the SM and possibly find any discrepancies hinting towards more general models explaining the  
35 observations not included in the SM.

36 Future Higgs factories operating at the Z-pole would produce  $10^9 - 10^{12}$  Z bosons, allowing for precise  
37 measurement of the Z-boson couplings to fermions. In this contribution, we present prospects for measuring  
38 electroweak couplings of quarks, especially highlighting the measurement for light quarks.

## 39 2 Measurement idea

40 A similar measurement has already been performed at LEP [1–5]. The main idea relies on the fact that up-  
41 and down-type quarks differ in electric charge and thus, their electromagnetic couplings are distinguishable.  
42 The coupling strength of the Z boson to a given fermion,  $c_f$ , is conventionally defined as a sum of its squared  
43 vector and axial couplings. The total width of the Z boson to hadrons,  $\Gamma_{had}$ , is proportional to the sum of the  
44 couplings for down- and up-type quarks:

$$45 \Gamma_{had} \sim (3c_d + 2c_u). \quad (1)$$

46 Analogously, the total width to radiative hadronic decays for exactly one photon emission,  $\Gamma_{had+\gamma}$ , scales as

$$47 \Gamma_{had+\gamma} \sim \frac{\alpha}{2\pi} f(y_{cut}) \left( 3Q_d^2 c_d + 2Q_u^2 c_u \right), \quad (2)$$

48 where  $\alpha$  is the electromagnetic coupling constant,  $f(y_{cut})$  is a form factor depending on an arbitrary parameter  
49  $y_{cut}$  incorporating the isolation criteria for photons and  $Q_d$  ( $Q_u$ ) is the electric charge of down-type (up-type)  
50 quarks. In the following, we will assume  $y_{cut}$  is the photon transverse momentum with respect to the jet  
51 direction,  $q^T$ . Since  $Q_d \neq Q_u$ , the expressions for the radiative and the total hadronic widths include different  
52 coupling combinations, and the couplings of the up- and down-type quarks can potentially be disentangled.

53 In our approach, we collect two-jet events inclusively and classify them into 10 categories, corresponding  
54 to all the combinations of 4 possible jet tags ("light", s, c, b). Additionally, we tag events with exactly one  
55 isolated photon. To estimate the statistical precision, we fit cross sections to minimise the  $\chi^2$ -test statistic  
56 by comparing the numbers of expected and "measured" events, where the second number is drawn from the  
57 Poisson distribution.

## 58 3 Monte Carlo event generation

59 A realistic and precise Monte Carlo simulation is crucial for the measurement. The process of  $e^+e^- \rightarrow q\bar{q}$  is  
60 often perceived as a benchmark point for Monte Carlo generators, but the reconstruction of isolated photons

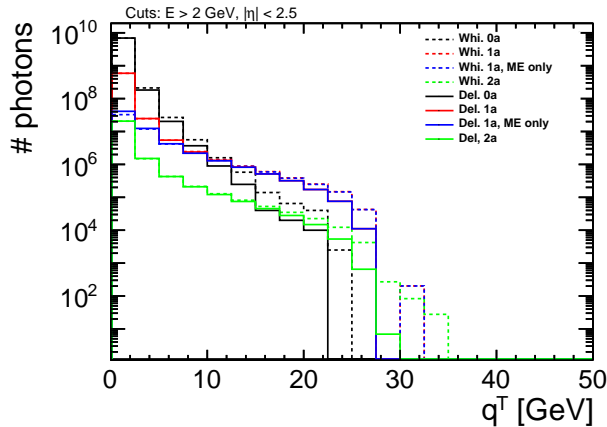


Figure 1: The number of photons visible in the detector as a function of  $q^T$ . See text for details.

61 poses several challenges. Not only does the final-state photon radiation (FSR) from Matrix Elements (ME)  
 62 have to be included but also the initial-state radiation (ISR), as well as their matching to parton showers  
 63 and hadronisation modeling. As a recipe, one can generate data samples using fixed-order ME calculations,  
 64 with exclusive emissions of hard photons, and match them with ISR structure function and FSR showers  
 65 accounting for collinear and soft emissions. Building upon [6], we have developed a dedicated matching  
 66 procedure for simulating photons using both ME calculations (for hard emissions) and ISR structure function  
 67 and FSR showers (for soft emissions) in WHIZARD [7, 8].

## 68 4 Experimental prospects

69 The difference in quark charges affects only the photons emitted from the final state and hence, experimental  
 70 selection should be optimised to enhance the numbers of measured photons  $q$  of this kind and suppress the  
 71 other contributions (ISR and decays of hadronisation products). To simulate detector effects, we used DEL-  
 72 PHES 3.5.0 [9] with the *ILCgen* cards. The steering cards were modified to cluster all the photons into jets. It  
 73 allowed for testing photon isolation criteria independently of the current assumptions.

74 In Figure 1, we compare the distributions of photons at the generator and detector level for different gener-  
 75 ated photon multiplicities. Besides the distributions for 0, 1 and 2 ME-photon samples, we plot the distribution  
 76 of the 1 ME-photon sample when only "proper" photons (those generated in the full ME picture, not coming  
 77 from showers or hadronisation) are tagged. The plot shows that above 10 GeV, the 1-ME-photon sample  
 78 becomes dominant.

79 To optimise the cut, we studied 7 sources of uncertainties: luminosity of 0.01%, acceptance of radiative  
 80 (non-radiative) events of 5% (50%), and tagging of b jets of 1%, c jets of 2%, s jets of 5%, light quarks of 10%.  
 81 The assumptions are preliminary and were used as a "proof of concept" of our statistical framework. Figure 2  
 82 shows preliminary results for the uncertainty of the measurement of the couplings and the ratio of light-quark  
 83 couplings, depending on the cut on  $q^T$ . The precision reaches the sub-percent level for light quarks (u and  
 84 d) and the sub-permille level for heavier quarks. The study will be continued to estimate the uncertainties in  
 85 greater detail.

## 86 5 Conclusions

87 Future  $e^+e^-$  colliders operating at the Z-pole will further constrain the Standard Model parameters. Precision  
 88 measurements of the Z-boson couplings to quarks will be possible thanks to very high data statistics. In  
 89 this study, we established a dedicated event generation procedure including photons coming from different  
 90 sources – ISR, FSR, hadronisation and hadron decays. We studied photon isolation criteria and estimated the  
 91 statistical and systematic uncertainties of the measurement. Preliminary results suggest that a sub-percent  
 92 precision will be achievable for light quarks while for heavier quarks, a sub-permille precision is envisioned.

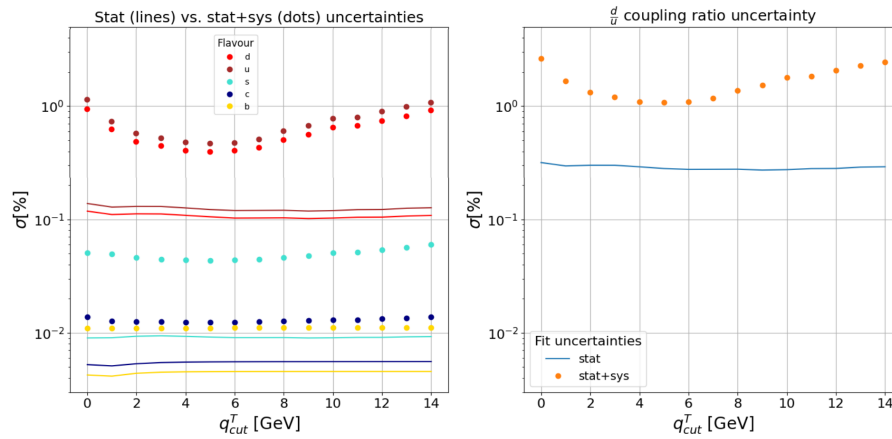


Figure 2: Preliminary results for the uncertainty of the Z-boson couplings to quarks (left) and the ratio of the u and d couplings (right) as a function of the cut on  $q^T$ . Solid lines indicate statistical precision only, while dots stand for a sum of statistical and systematic uncertainties.

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