New Developments in Physics of Electroweak Symmetry Breaking

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What do we expect to see at TeV? → Physics of electroweak symmetry breaking Is there New Physics at TeV? → We don't know

Possible hints

- Problem of Naturalness
- Existence of the Dark Matter

Good motivations for Physics beyond the SM May also be responsible for weak-Planck hierarchy ... e.g. theory of radiative EWSB → Target of Next Experiments (LHC, LC, ...)

Naturalness Problem of the SM

The SM Lagrangian is not stable under quantum corrections:



"Experimentally" $m_h^2 < O(100 \text{ GeV})^2$ (v = $2m_h^2/\lambda$ = 174 GeV) Naturalness: $m_h^2 < O(100 \text{ GeV})^2 \rightarrow m < O(\text{TeV})$ What do we know about New Physics?
→ The SM describes physics at < O(100 GeV) extremely well</p>

Contributions from New Physics:





Weak Scale Supersymmetry

Superparticle at the TeV scale

Bosons	Fermions
A_{μ}	λ
q	q
Ĩ	1
h	ĥ

After SUSY breaking, λ , \tilde{q} , \tilde{I} and \tilde{h} obtain masses of O(TeV)

→ Beautiful cancellation of δm_h^2 between contributions from the SM and superparticles

R parity \rightarrow the existence of Dark Matter

Gauge coupling unification at a high scale:

nonSUSY



• SUSY



 $M_{unif} \sim 10^{16} \text{ GeV}$

Superparticle spectrum provides a window for a deeper level of physical theory

• "Gravity" mediation



Gauge mediation



- Anomaly mediation
- ...
- → Distinct spectra

Combination of LHC and LC important

LSP ~ DM: weakly interacting --- Exploration at LC



Supersymmetry after the LEPII
"Minimal" Supersymmetry is fine-tuned

... Supersymmetric fine-tuning problem

- Minimization condition

$$\frac{M_{\rm Higgs}^2}{2} \simeq -m_h^2 - |\mu|^2$$

Natural EWSB requires

$$\frac{M_{\rm Higgs}^2}{2} \sim |m_h^2| \sim |\mu|^2$$

In the MSSM,

$$M_{\text{Higgs}} \lesssim 130 \text{ GeV} \xrightarrow{\Delta^{-1} \ge 20\%} |m_h^2|, \ |\mu|^2 \lesssim (200 \text{ GeV})^2$$

m_h² receives contribution from top-stop loop

$$\delta m_h^2 \simeq -\frac{3y_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln\left(\frac{M_{\text{mess}}}{m_{\tilde{t}}}\right)$$

M_{mess}: the scale where suerparticle masses are generated

What's wrong?

- $M_{Higgs} < M_Z$ at tree level
 - need radiative corrections from top-stop loop $M_{\text{Higgs}} \gtrsim 114 \text{ GeV} \implies m_{\tilde{t}} \gtrsim 1 \text{ TeV} \quad (\text{for small } A_t)$
- Tension with the other superparticle mass bounds
 - e.g. mediation by the SM gauge interactions

 $\frac{m_{\tilde{t}}}{m_{\tilde{e}}} \simeq \frac{(4/3)g_3^4 + \delta}{(3/5)g_1^4} \simeq (7 \sim 8) \quad \Rightarrow \quad m_{\tilde{t}} \gtrsim 700 \text{ GeV for } m_{\tilde{e}} \gtrsim 100 \text{ GeV}$



Possible approaches

- We don't care
- Higgs boson may be lighter Dermisek, Gunion; Chang, Fox, Weiner; ...
 - → LEPII may have missed the Higgs h⁰ because h⁰ decays into final states that are hard to detect
- Higgs boson may be heavier Barbieri, Hall, Y.N., Rychkov; ...
 ... alleviates fine-tuning (in the most straightforward way)
 → We may have been misled in interpreting EWPT
- Problem of SUSY breaking mechanism? Kitano, Y.N.; ...
 Some mechanism may be preferred over others
- Environmental
 - Split SUSY Arkani-Hamed, Dimopoulos; Giudice, Romanino; ...
 - Living dangerously Giudice, Rattazzi; ...

SUSY without a light Higgs boson

Heavier Higgs boson alleviates tuning

$$V_{\text{Higgs}} = m_h^2 |h|^2 + \lambda |h|^4 / 4$$

$$\rightarrow v^2 = |h|^2 = -2m_h^2 / \lambda, \ M_{\text{Higgs}}^2 = \lambda v^2$$

In SUSY

$$\frac{M_{\rm Higgs}^2}{2} \simeq -m_h^2 - |\mu|^2$$

Allowed by EWPT?

(We imagine e.g. $M_{Higgs} \sim 200-300 \text{ GeV}$)



Constraint on M_{Higgs} on the S-T plane



→ easy to have large M_{Higgs} if additional ΔT exist Maybe we are being fooled by $\Delta T|_{New Physics}$

λ SUSY framework

Barbieri, Hall, Y.N., Rychkov, hep-ph/0607332

Higgs boson can be made heavy in SUSY by

 $W = \lambda S H_u H_d$

with $\lambda \approx (1 \sim 2)$ (λ perturbative up to ~10TeV)

• Large λ makes M_{Higgs} heavy and **at the same time** induces sizable ΔT from singlet-doublet mixings!

Parameter of the model:

Fermion sector Scalar sector μ, Μ

 $\mu_1^2, \mu_2^2, \mu_3^2 \rightarrow \tan\beta, m_{H^+}, v$

in the limit of decoupling the S scalar and gauginos

Contribution from the scalar Higgs sector



λ=2

m_{H+}=350,500,750 GeV

- $\Delta T \rightarrow 0$ for tanβ→1 (custodial symmetry)
- $\Delta T>0$ can make large M_{Higgs} consistent

\rightarrow tan β cannot be large

also reinforced by the stop-sbottom contributions



Contribution from the Higgsino sector



tan β < 3 preferred in λ SUSY \rightarrow rich Higgs physics at ~O(200-700 GeV)



Gauge Coupling Unification

Compositeness of S, H_u , H_d and/or top can induce large λ , keeping the desert Harnik, Kribs, Larson, Murayama; Chang, Kilic, Mahbubani; Birkedal, Chacko, Y.N.; Delgado, Tait; ...

- 5D realization/modeling Birkedal, Chacko, Y.N.



 \rightarrow Gauge coupling unification can be preserved

Higgs sector physics in supersymmetry can be quite rich – potential window for the DM sector Exploration at LC very useful

Higgs as a pseudo Nambu-Goldstone boson

- Why $m_h << M_{Pl}$?
 - → Higgs is a pseudo Nambu-Goldstone boson
 - Old composite models Georgi, Kaplan; ...
 - Little Higgs models Arkani-Hamed, Cohen, Georgi; ...
 - Holographic Higgs models Contino, Y.N., Pomarol; ...
 - Twin Higgs models Chacko, Goh, Harnik; ...

— . . .

Cancellation of δm_h^2 is between the same statistics fields – e.g. we have t' instead of t

Higgs as a holographic pseudo Nambu-Goldstone boson

Contino, Y.N., Pomarol, hep-ph/0306259 Agashe, Contino, Pomarol, hep-ph/0412089

Technicolor:

 $(\Lambda_{QCD} \sim \langle q\bar{q} \rangle^{1/3} \sim 1 \text{ GeV gives } f_{\pi} \sim m_W \sim 100 \text{ MeV})$

 $m_W \sim 100 \text{ GeV} \rightarrow \Lambda_{TC} \sim 1 \text{ TeV}$

The scale of New Physics too close v --- contradict with the precision electroweak data

Pions in massless QCD:

 $(m_{\pi^{\pm}} \sim 10 \text{ MeV with } \Lambda_{QCD} \sim 1 \text{ GeV})$ $m_{h} \sim 100 \text{ GeV} \rightarrow \Lambda_{NEW} \sim 10 \text{ TeV}$

Safer in terms of the precision electroweak data

 $\Lambda_{\sf NFW}$ —

V_{FW} -

 Λ_{TC} _____

V_{EW} ------

- Basic structure (omitting details, e.g. U(1)_Y assignment) $SU(3)_{global} \supset SU(2)_{L}$ $SSB: SU(3)_{global} \rightarrow SU(2)$ produces PNGB which is $SU(2)_{L}$ doublet $\begin{pmatrix} SU(2) \\ * \\ * \end{pmatrix}$
- Higgs compositeness is not enough



Analogy with QCD:



 ν $\mathcal{L} \approx \frac{1}{M^2} u dd \nu$ "Yukawa" couplings suppressed because
dim[O_n] (O_n =udd) is "large" = 9/2
Need 1.1

Need interactions strong for a wide energy range to reduce dim[O] \rightarrow CFT

• How to realize such theories?

--- Gauge theory/gravity correspondence



Realization in 5D



$$A_{\mu}: \begin{pmatrix} (+,+) & (+,+) & (-,-) \\ (+,+) & (+,+) & (-,-) \\ \hline (-,-) & (-,-) & (+,+) \end{pmatrix} \qquad A_{5}: \begin{pmatrix} (-,-) & (-,-) & (+,+) \\ (-,-) & (-,-) & (+,+) \\ \hline (+,+) & (+,+) & (-,-) \end{pmatrix}$$

Higgs arises as an extra dimensional component of the gauge boson, A_5

Realistic Yukawa couplings obtained



- Higgs potential is finite and calculable → EWSB (due to higher dimensional gauge invariance)
- Existence of resonances (KK towers) for the gauge fields as well as quarks and leptons ... Physics at the LHC (and LC)
- Nontrivial wavefunctions for W and Z as well as for matter → couplings deviate from the SM ... Physics at LC

Summary

- We are about to explore physics of EWSB
- New Physics at the TeV scale is well motivated
 - Problem of Naturalness, DM, ...
 - Weak scale supersymmetry
 - Higgs as a pseudo Nambu-Goldstone boson
 - ...
- Exploration of this New Physics is the prime target of experiments in the next decades
 A variety of possibilities for how New Physics shows up
 → Combination of the LHC and LC very useful
- We hope to understand Nature at a deeper level