

Charge asymmetry in $\gamma\gamma \rightarrow \mu^+\mu^- + \nu$'s $\gamma\gamma \rightarrow W^\pm \mu^\mp + \nu$'s with polarized photons

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- Charge asymmetry in processes like

$$\gamma\gamma \rightarrow \mu^+\mu^-\nu_\mu\bar{\nu}_\mu, \quad \gamma\gamma \rightarrow W^\pm\mu^\mp\nu$$

appears due to P nonconservation in the SM.

- Processes like

$$\gamma\gamma \rightarrow \tau\mu\nu\nu \quad (\gamma\gamma \rightarrow W\tau\nu) \quad \rightarrow \quad \mu^+\mu^-\nu\nu\nu \quad (W\mu\nu\nu)$$

(with $\tau \rightarrow \mu\nu\mu\nu$ decay) produce the same observable final state enhancing total event rate by 37%(17%). We consider such cascade processes.

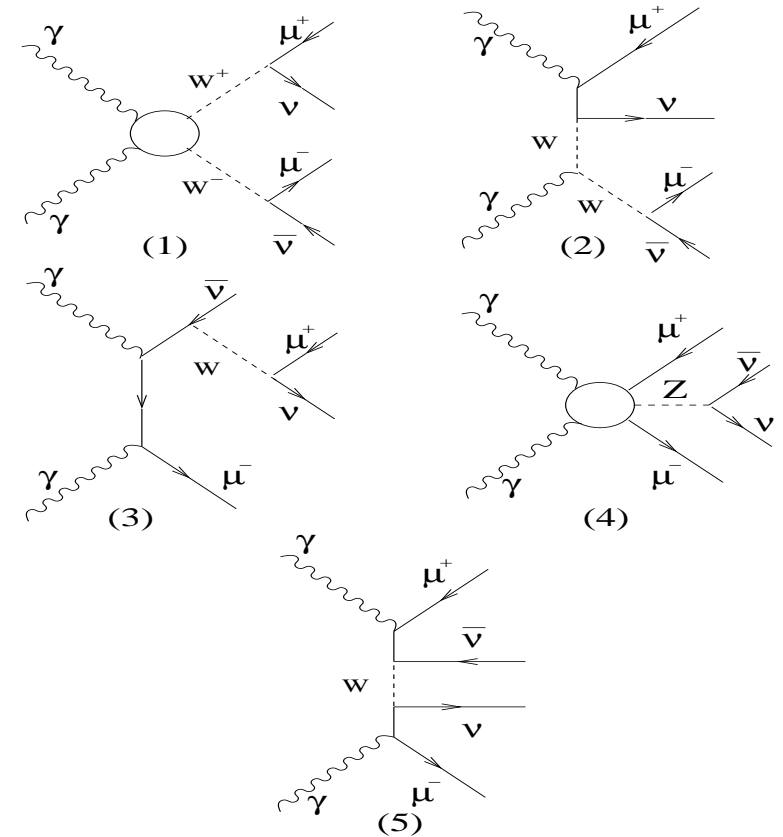
- Photon spectra are non-monochromatic. Photons with energy $E_\gamma < E_\gamma^{max}/\sqrt{2}$ are non-polarized. How this fact reduce asymmetry?

Diagrams for $\gamma\gamma \rightarrow \mu^+\mu^-\nu\mu\bar{\nu}\mu$ ($\gamma\gamma \rightarrow \tau\mu\nu\nu$)

- (1) 3 double-resonant diagrams (DRD)
- (2) 4 single-resonant diagrams (SRD)
- (3) 4 single resonant diagrams with ν exchange
- (5) 2 multi-peripheral non-resonant diagrams

- SRD (2)/DRD (1) is about 5%.
- The interference SRD, DRD is destructive.
- (3),(4) and (5)/DRD (1) $\ll 1$.
- DRD contribution covers almost 98.7 % cross section.

(The $\gamma\gamma \rightarrow W^+\mu^-\bar{\nu}$ is described by only first 3 groups.)



We used CalcHEP for calculations.

For each observed particle:

- Cut in escape angle θ

$$\pi - \theta_0 > \theta > \theta_0 \text{ with } \theta_0 = 10 \text{ mrad},$$

- Cut in transverse momentum p_{\perp} :

$$p_{\perp} > p_{\perp\mu}^c \text{ with } p_{\perp\mu}^c = 10 \text{ GeV}$$

and higher $p_{\perp\mu}^c$ up to 80 GeV.

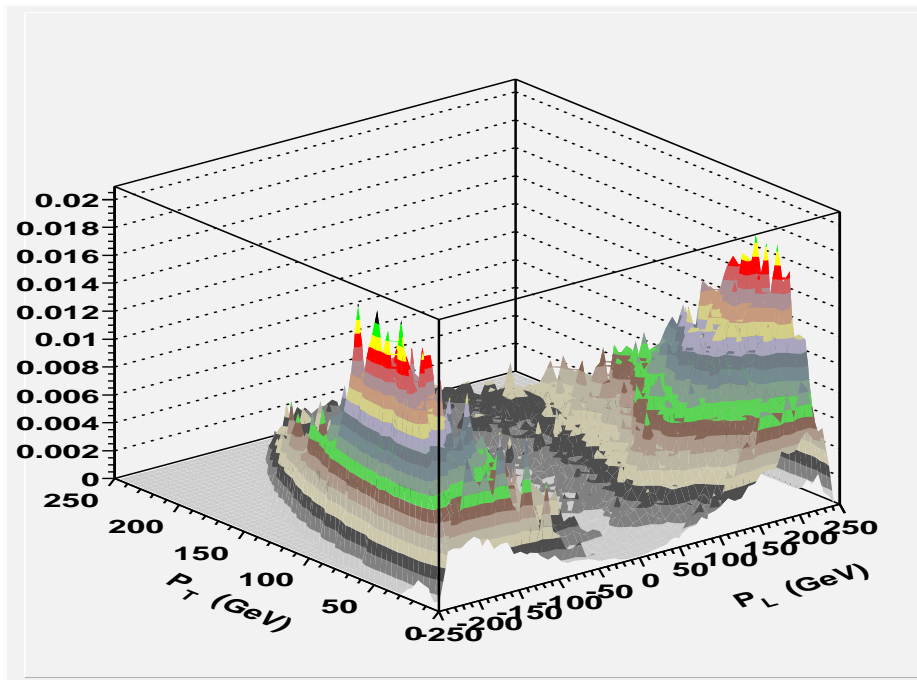
These simultaneous cuts allow many backgrounds to be eliminated.

The number of generated events = anticipated
annual number $\simeq 10^6$ events,

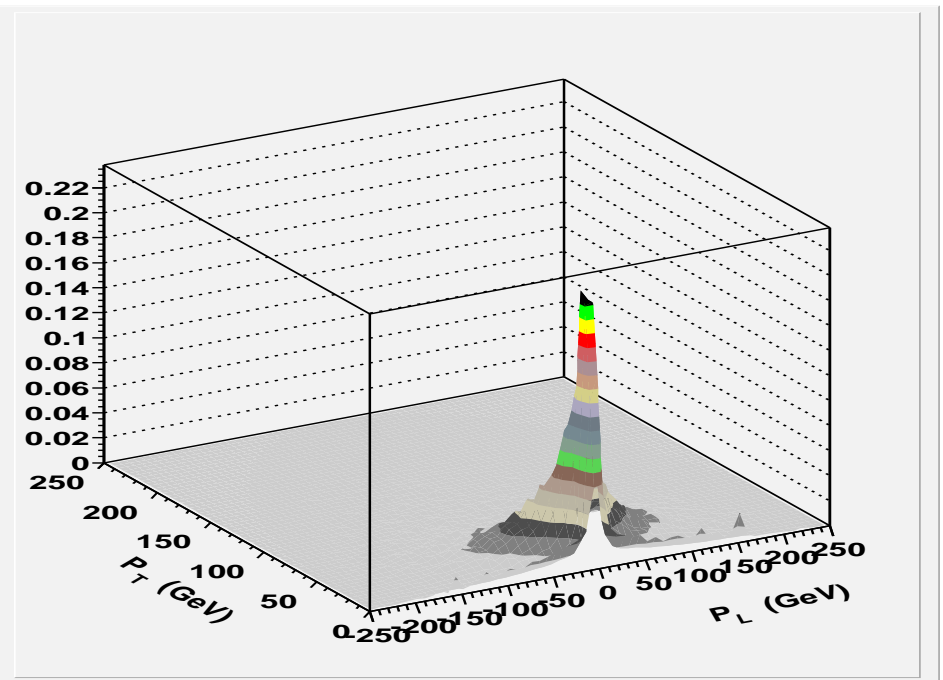
In the non-monochromatic case – that is # of events for fraction
generated by photons with $E_{\gamma} > E_{\gamma}^{max} / \sqrt{2}$.

Difference between distributions of positive and negative muons in $\gamma_{\lambda_1} \gamma_{\lambda_2} \rightarrow W \mu \nu$ (no cuts).

Both photons are left polarized, $\gamma_- \gamma_-$

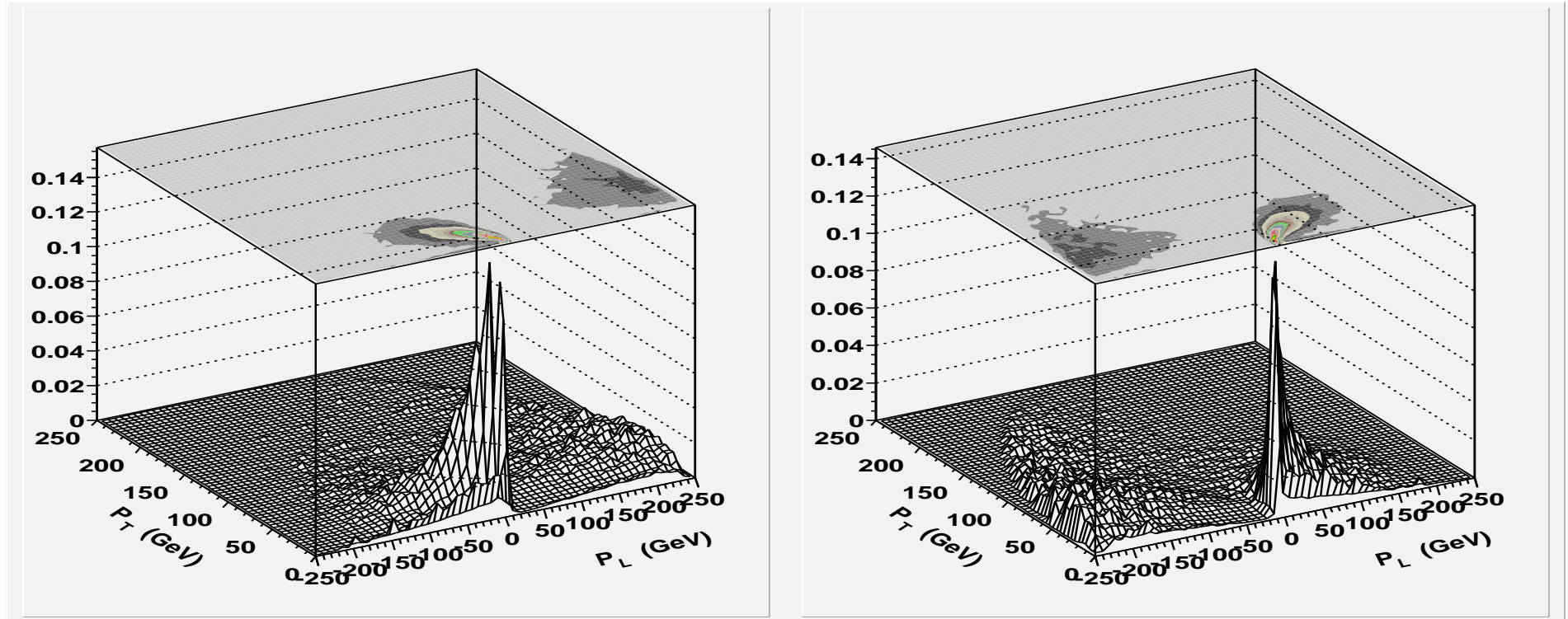


Negative μ distribution.



Positive μ distribution.

First photon is left polarized, second is right polarized, $\gamma_-\gamma_+$



Negative μ distribution.

Positive μ distribution.

Note: the distributions are mirror-symmetric.

For $\gamma\gamma \rightarrow W^\pm \mu^\pm + \nu$'s processes we considered normalized mean values of longitudinal p_{\parallel}^\mp and transverse p_{\perp}^\mp momenta of muons:

$$P_L^\pm = \frac{\int p_{\parallel}^\pm d\sigma}{E_\gamma^{max} \int d\sigma}, \quad P_T^\pm = \frac{\int p_{\perp}^\pm d\sigma}{E_\gamma^{max} \int d\sigma},$$

and taken their relative difference as a measure of charge asymmetry:

$$\Delta_L = \frac{P_{L+}^- - P_{L+}^+}{P_{L+}^- + P_{L+}^+}, \quad \Delta_T = \frac{P_{T+}^- - P_{T+}^+}{P_{T+}^- + P_{T+}^+}.$$

- Quantities for $\gamma_+\gamma_+$ and $\gamma_-\gamma_+$ can be obtained with $\mu^+ \leftrightarrow \mu^-$ exchange for P_N and with sign change for Δ_N .

- Monte Carlo simulations have statistical uncertainty $\delta P_{L,T}, \delta \Delta_{L,T}$ similar to experimental.

| $\gamma_{\lambda_1}\gamma_{\lambda_2}$ | N | P_N^- δP_N^- | P_N^+ δP_N^+ | Δ_N $\delta \Delta_N$ |
|--|---|---------------------------|---------------------------|---------------------------------|
| $\gamma_-\gamma_-$ | L | 0.606 0.29% | 0.201 0.55% | 0.501 0.57% |
| | T | 0.333 0.61% | 0.159 0.28% | 0.335 0.44% |
| | L | 0.223 0.82% | 0.609 0.19% | -0.463 0.47% |
| $\gamma_+\gamma_-$ | T | 0.164 0.08% | 0.262 0.31% | -0.231 0.76% |

Charge asymmetry quantities and statistical uncertainties for $\gamma_{\lambda_1}\gamma_{\lambda_2} \rightarrow W\mu\nu$. Monochromatic case.

We also studied inaccuracy of DRD approximation for various asymmetries.

- Inaccuracy of DRD approximation in $\gamma\gamma \rightarrow W\tau\nu$ for $P_{L,T}$ and $\Delta_{L,T}$ quantities $\lesssim 5\%$.

- **Inaccuracy for $\gamma\gamma \rightarrow W\tau\nu \rightarrow W\mu\nu\nu$**

is $\lesssim 0.17 \cdot 5\% = 0.85\%$.

- **within statistical MC inaccuracy for 10^6 events.**

Cascade process

Muons with missing transverse momentum can appear via processes

$$\gamma\gamma \rightarrow \tau^+ \mu^- \nu_\tau \bar{\nu}_\mu \quad (\gamma\gamma \rightarrow W \tau \nu)$$

followed by $\tau \rightarrow \mu \nu_\mu \nu_\tau$.

Total event rate enhancement:

- for $\gamma\gamma \rightarrow W \mu + \nu's$: $B \equiv Br(\tau \rightarrow \mu \nu \nu) = 17\%$
- for $\gamma\gamma \rightarrow \mu^+ \mu^- + \nu's$: $2B + B^2 \approx 37\%$.

Calculation of such processes (6 or more final particles) is a computationally challenging task. Reasonable (DRD) approximations provides high enough accuracy for our purposes.

In the frame of DRD each τ is produced from W decay \Rightarrow τ polarisation is known and we are allowed to *convolute* generated distribution of τ in $\gamma\gamma \rightarrow W\tau\nu$ with distribution of μ in τ decay:

$$f = \frac{4}{\pi E_\tau m_\tau^4} \left[(3m_\tau^2 - 4pk)pk + ks \cdot m_\tau(4pk - m_\tau^2) \right] d\Gamma$$

Here k and p are 4-momenta of μ and τ .

$$\text{Spin of } \tau: \pm s/2, \quad s = \frac{1}{\sqrt{2}} \left(\frac{p_\nu m_\tau}{(pp_\nu)} - \frac{p}{m_\tau} \right) \begin{cases} + & \text{for } \tau^+, \\ - & \text{for } \tau^-. \end{cases}$$

Essential feature

Decay $\tau \rightarrow \mu\nu\tau\nu\mu$ involves 3 particles,

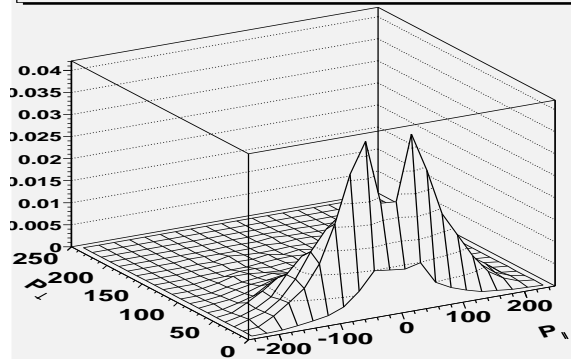
the effective mass of the $\nu\bar{\nu}$ system $m_{\nu\nu}$ varies from 0 to m_τ

\Rightarrow the μ distribution is *contracted* in comparison with τ distribution:

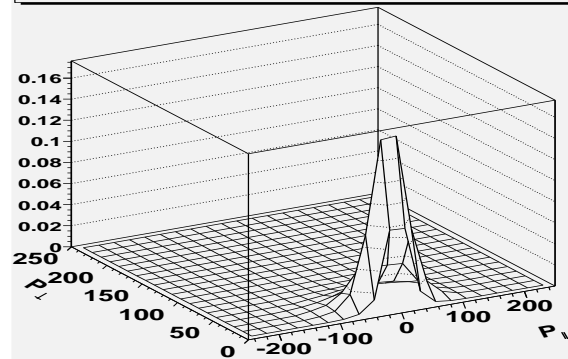
$$E_\mu \leq E_\tau (1 - m_{\nu\nu}^2/m_\tau^2).$$

Distributions of μ in cascade process

μ^- from $\tau \Delta\sigma/\Delta P_{\parallel}^- \Delta P_{\perp}^-$, (--), $\sqrt{s}=500$ GeV

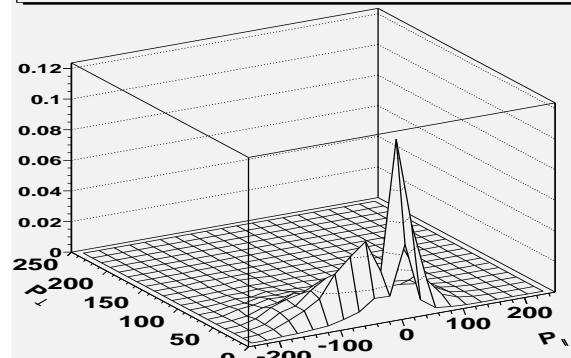


μ^+ from $\tau \Delta\sigma/\Delta P_{\parallel}^+ \Delta P_{\perp}^+$, (--), $\sqrt{s}=500$ GeV

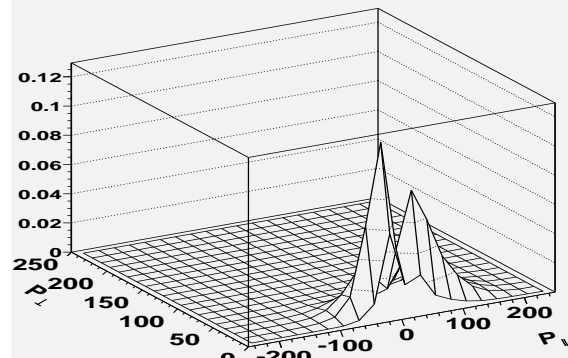


Muon distribution in $\gamma_-\gamma_- \rightarrow W\mu\nu\nu$

μ^- from $\tau \Delta\sigma/\Delta P_{\parallel}^- \Delta P_{\perp}^-$, (+-), $\sqrt{s}=500$ GeV



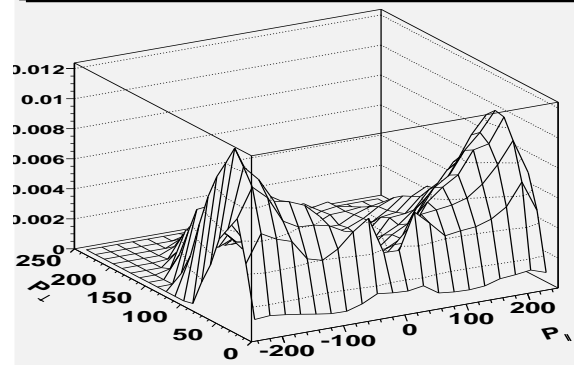
μ^+ from $\tau \Delta\sigma/\Delta P_{\parallel}^+ \Delta P_{\perp}^+$, (+-), $\sqrt{s}=500$ GeV



Muon distribution in $\gamma_+\gamma_- \rightarrow W\mu\nu\nu$ left - μ^- , right - μ^+

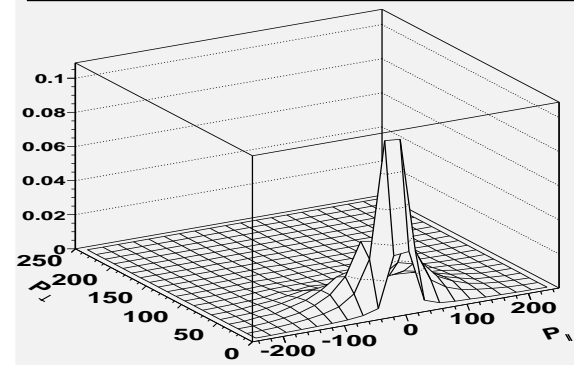
Entire distributions of μ

Full $\mu^- \Delta\sigma/\Delta P_{\parallel}^- \Delta P_{\perp}^-$,(--), $\sqrt{s}=500$ GeV

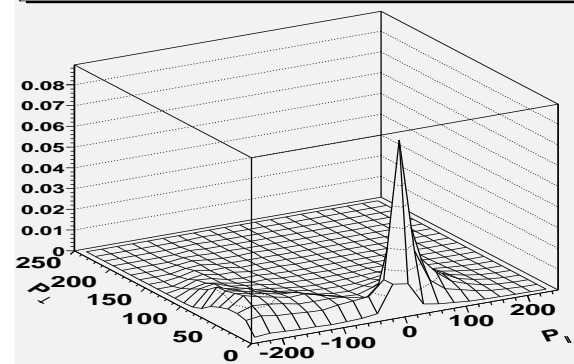


Total muon distribution in $\gamma_- \gamma_- \rightarrow W \mu + \nu$'s

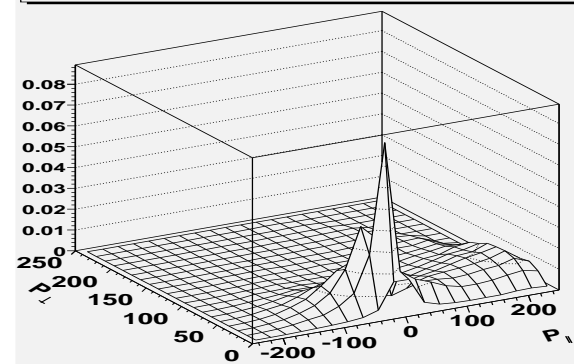
Full $\mu^+ \Delta\sigma/\Delta P_{\parallel}^+ \Delta P_{\perp}^+$,(--), $\sqrt{s}=500$ GeV



Full $\mu^- \Delta\sigma/\Delta P_{\parallel}^- \Delta P_{\perp}^-$,(--), $\sqrt{s}=500$ GeV



Full $\mu^+ \Delta\sigma/\Delta P_{\parallel}^+ \Delta P_{\perp}^+$,(--), $\sqrt{s}=500$ GeV



Total muon distribution in $\gamma_+ \gamma_- \rightarrow W \mu + \nu$'s left $-\mu^-$, right $-\mu^+$

- Cascade process changes μ distribution only at small momenta.
- Asymmetry parameters decrease by $\lesssim 3\%$

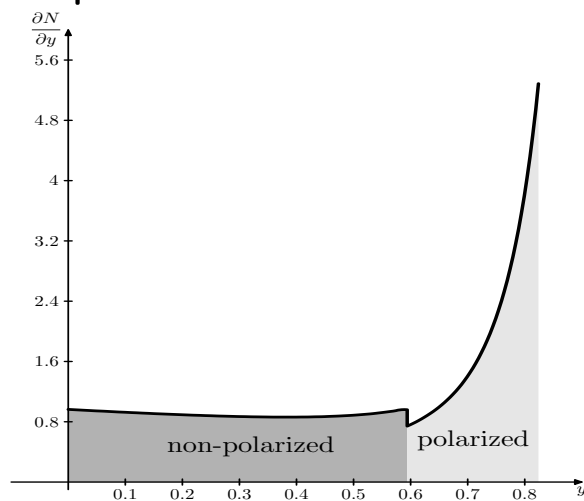
| $\gamma_{\lambda_1} \gamma_{\lambda_2}$ | N | P_N^- | P_N^+ | Δ_N |
|---|---|---------|---------|------------|
| $\gamma_- \gamma_-$ | L | 0.548 | 0.164 | 0.539 |
| | T | 0.311 | 0.142 | 0.374 |
| $\gamma_+ \gamma_-$ | L | 0.199 | 0.513 | -0.440 |
| | T | 0.152 | 0.232 | -0.207 |

Total charge asymmetry quantities.

- Applied cuts reduce the contribution of cascade process stronger than the main contribution \Rightarrow *reduce* inaccuracy of DRD approximation in the description of charge asymmetry with growth of $p_{\perp\mu}^c$.

Effect of photon non-monochromaticity

High energy part $E_\gamma > E_\gamma^{max}/\sqrt{2}$ is obtained from ideal one (*Compton spectrum*) with known factor dependent on photon energy and distance **conversion point – collision point**. The polarization distribution on energy is the same as in an ideal case (**Ginzburg, Kotkin**). At lower energies factorization is broken, polarization disappears. Details depend on real collision scheme strong.



To **IMITATE**, we used factorized spectra – for $E_\gamma > E_\gamma^{max}/\sqrt{2}$ mentioned **Ginzburg, Kotkin**, for $E_\gamma < E_\gamma^{max}/\sqrt{2}$ – ideal Compton spectrum without geometrical factors, no polarization.

Luminosity was normalized for product of high energy photon fluxes.

| $\gamma_{\lambda_1} \gamma_{\lambda_2}$ | N | P_N^- δP_N^- | P_N^+ δP_N^+ | Δ_N $\delta \Delta_N$ | P_N^- δP_N^- | P_N^+ δP_N^+ | Δ_N $\delta \Delta_N$ |
|---|---|---------------------------|---------------------------|---------------------------------|---------------------------|---------------------------|---------------------------------|
| $\gamma-\gamma-$ | L | 0.606 0.29% | 0.201 0.55% | 0.501 0.57% | 0.365 0.31% | 0.157 0.22% | 0.398 0.18% |
| | T | 0.333 0.61% | 0.159 0.28% | 0.335 0.44% | 0.284 0.38% | 0.179 0.11% | 0.228 0.81% |
| $\gamma+\gamma-$ | L | 0.223 0.82% | 0.609 0.19% | -0.463 0.47% | 0.174 0.24% | 0.338 0.28% | -0.321 0.43% |
| | T | 0.164 0.08% | 0.262 0.31% | -0.231 0.76% | 0.200 0.09% | 0.236 0.16% | -0.082 0.42% |

Monochromatic case

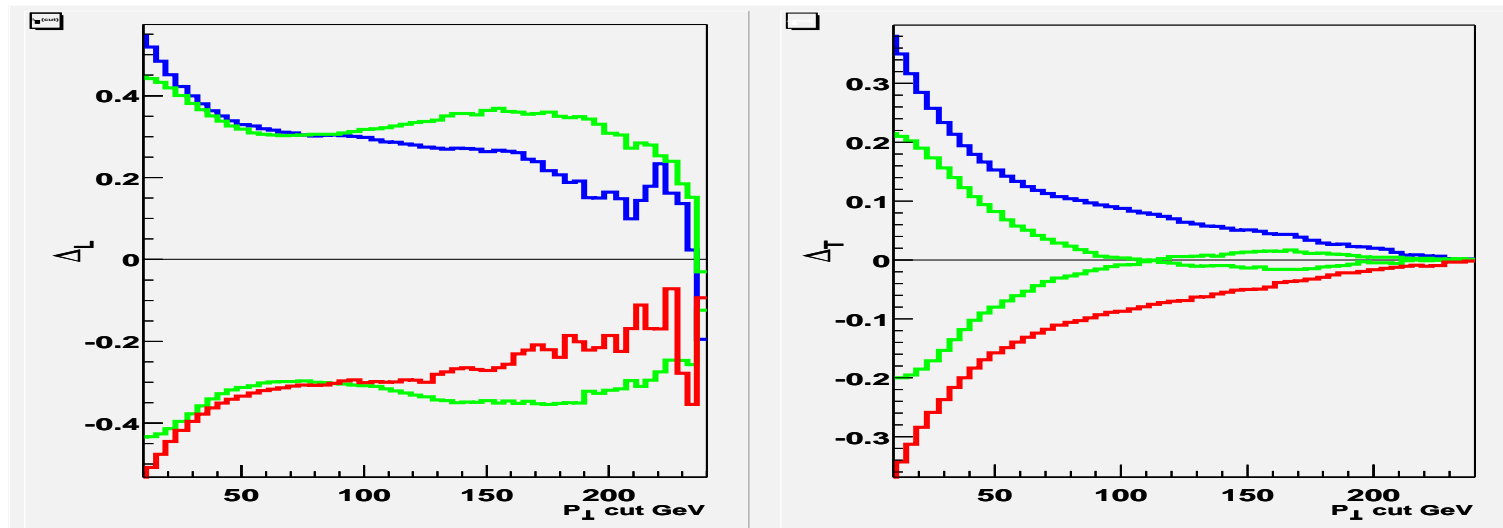
Non-monochromatic case

*Charge asymmetry quantities and statistical uncertainties for
 $\gamma_{\lambda_1} \gamma_{\lambda_2} \rightarrow W_{\mu\nu}, \text{cut} = 10 \text{ GeV}.$*

In the non-monochromatic case charge asymmetry quantities are reduced typically by a factor 1.3-1.5, their statistical uncertainties vary not so strong.

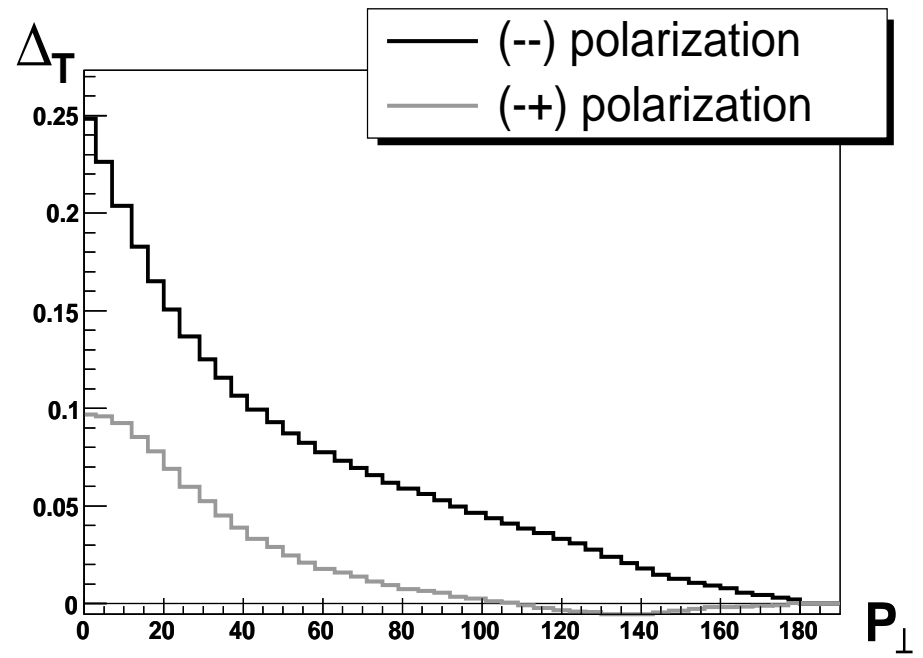
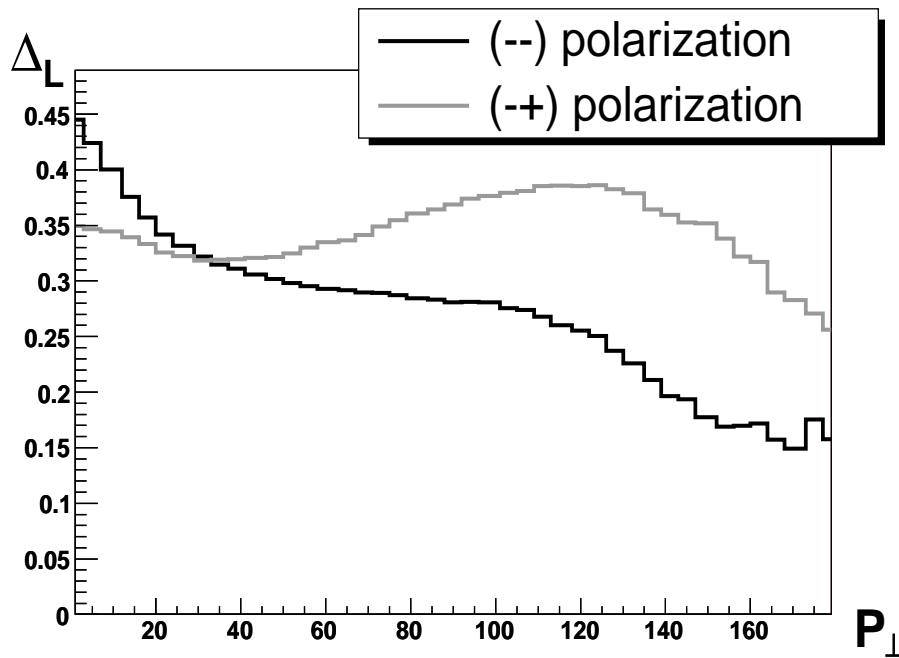
Dependence on cut $p_{\perp\mu}^c$

New Physics is expected to be switched on at large transverse momenta. We study the dependence of asymmetry on the cut $p_{\perp\mu}^c$.



The $p_{\perp\mu}^c$ dependence of asymmetry. Left – Δ_L , right – Δ_T ,

blue – $\gamma_-\gamma_-$, green – $\gamma_-\gamma_+$ and $\gamma_+\gamma_-$, red – $\gamma_+\gamma_+$



The mentioned decreasing of asymmetry become not so strong with growth of cut momentum for longitudinal asymmetry.

CONCLUSIONS AND PLANS

- Huge and easily observable effect.
- Cascade process weakly affect the asymmetry.
- Introduced quantities (especially Δ_L) large even with large $p_{\perp\mu}^c$ cuts.
- Taking into account same effects for e^+e^- , $e^+\mu^-$, μ^+e^- enhance statistics by 4 times (it is taken into account).
- Non-monochromaticity of photon spectra decreases the considered asymmetries but retain them large enough.

We plan to consider charge asymmetry for discovery of New Physics effects (e.g. MSSM).