

# Prospects of Discovering a New Massless Neutral Gauge Boson at the ILC

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In the **Standard Model** based on the  $U_Y(1) \times SU_L(2) \times SU_C(3)$  gauge symmetry,

→ 12 gauge bosons

- 3 electroweak bosons  $W^{+-}$  and  $Z^0$ ,
- 8 massless gluons
- and well known photon  $\gamma$

It is possible to introduce new gauge fields into the SM consistent with basic principles of gauge invariance.

One example of such an extension consists of introducing a new field  $U_P(1)$  which mixes with the SM field  $U_Y(1)$  in a manner that **new particles** come into the game

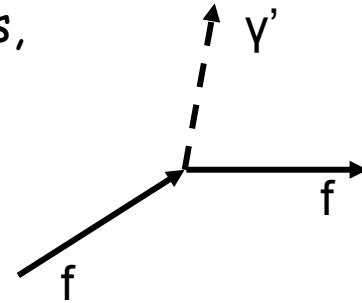
This mechanism provides an elegant way of introducing

- 
- **millicharged particles**
  - the two  $U(1)$  fields can mix such that one linear combination appears as the ordinary photon  $\gamma$  and the other as a **massless spin-1 neutral particle**, the 'paraphoton'  $\gamma'$

Based on the effective **Lagrangian** for the interaction of this new particle with the SM fermions by considering higher dimensional operators, the **Feynman rules for fermion-fermion- $\gamma'$  vertices** are proportional to

*(Bogdan Dobrescu)*

$$\propto C_f \frac{m_f}{M^2}$$



- with  $C_f$  - dimensionless coupling strength, with  $f = u, d, e$
- $m_f$  - mass of fermion
- $M$  - mass scale where the operators are generated
  - for large  $M$ , strong suppression of paraphoton interactions with SM fermions due to  $1/M^2$

The coefficients  **$C_f$  are unknown**, but some **phenomenological constraints** exist:

- e.g. the  $\mu$ - $\mu$ - $\gamma'$  coupling  $M/\sqrt{c_\mu} \geq 1.5 \text{ TeV}$  from expansion rate of universe and of possible paraphoton annihilation to muon pairs, with  $c_\mu$  related to  $C_\mu$  via
 

$c_f = C_f * v_h / (\sqrt{2}m_f)$

 , with  $v_h$  the vacuum expectation value of the Higgs field
- $M/\sqrt{c_e} \geq 3.2 (1.8) \text{ TeV}$  from core of red giant stars
- $M/\sqrt{c_N} \geq 7 \text{ TeV}$  from neutrino signal of the supernova 1987 A

**phenomenological bounds !**

- 
- a massless neutral gauge boson may exist,
  - it has non-zero but very small couplings to ordinary matter,
  - so far, (rather weak) phenomenological bounds make it worthwhile to search for this photon-like state
  - according to Feynman rules,  $\gamma'$  coupling prop. to the fermion mass  
→ favours coupling to the top quark!
  - no bounds on  $c_+$  exist, search for  $\gamma'$  at the next generation of colliders;  
since  $M$  might be very large  
→ at least limits on  $M/\sqrt{c_+}$  can be derived

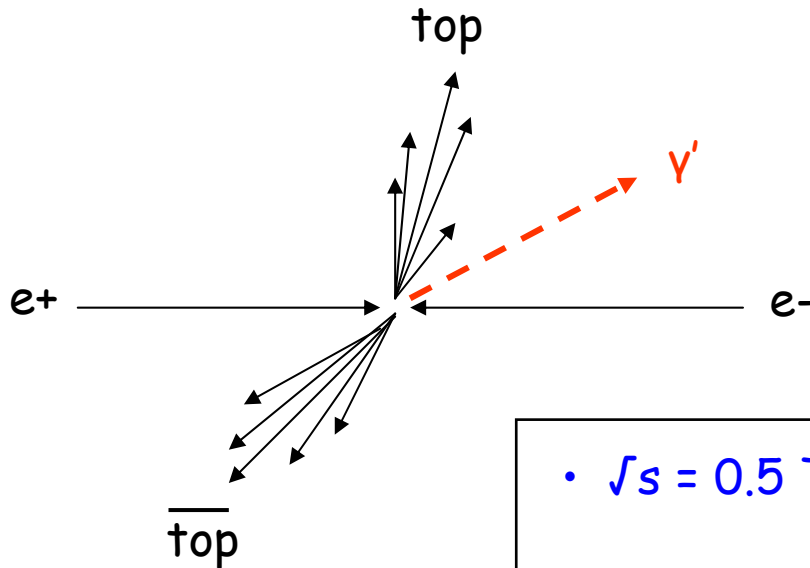
At the **LHC**, huge  $t t$ -bar (+jets) backgrounds  
→ (probably) precludes a sensitive search for  $\gamma'$ .

The **ILC** ideally suited to evaluate prospects of discovering (excluding) the  $\gamma'$

$$e^+ e^- \rightarrow t t\text{-bar } \gamma'$$

## Search strategy of the $\gamma'$ at ILC

→ select a pair of **acoplanar top quark jets** with **missing transverse energy**  $E_T^{\text{miss}}$  carried away by the paraphoton



- $\sqrt{s} = 0.5 \text{ TeV}$  with  $\int L dt = 500 \text{ fb}^{-1}$
- $\sqrt{s} = 1.0 \text{ TeV}$  with  $\int L dt = 1000 \text{ fb}^{-1}$

All **signal** event simulations with '**coupling**'  $M/\sqrt{c_+} = 0.2 \text{ TeV}$ , resulting to reasonable number of events  
→ differentiation w.r.t. background

## # of $t\bar{t}\gamma'$ events:

$M/\sqrt{c_t}$ , TeV	$\sqrt{s} = 0.5$ TeV	$\sqrt{s} = 1.0$ TeV
0.1	90 000	680 000
0.2	5 700	42 500
0.3	1 100	8 200
0.5	40	1 100
1.0	10	70

all W's  
decay modes  
included

- Event generation:
- **CompHEP** (extended to implement  $\gamma'$ )  $\rightarrow$  parton level; interfaced to PYTHIA
  - **PYTHIA** (ISR / FSR; jet hadronisation, resonance decays); beamstrahlung
  - **SIMDET** (detector response  $\rightarrow$  energy flow objects)

### Background:

*other backgrounds are negligible after selection procedure*

	$\sqrt{s} = 0.5$ TeV	$\sqrt{s} = 1.0$ TeV
$t\bar{t}\gamma$	276 675	200 310
$t\bar{t}\nu\bar{\nu}$	75	930

$$M/\sqrt{c_+} < 0.1 \text{ TeV} ?$$

- for such values, mixing between  $U_y(1)$  and  $U_p(1)$  fields would lead to **millicharged particles with charges  $> 10^{-4}$**

→ which would be detected in various SM reactions

contrary to experimental findings

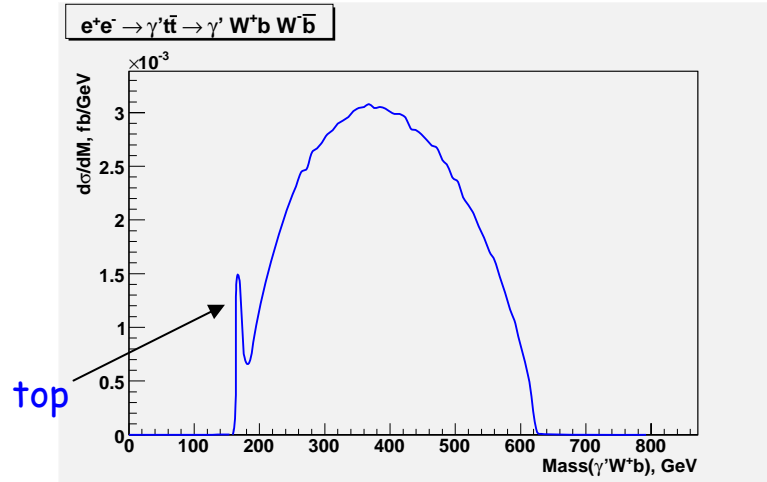
- eff. Lagrangian approach not valid

for small ( $< \sim 0.1 \text{ TeV}$ )  $M/\sqrt{c_+}$  values

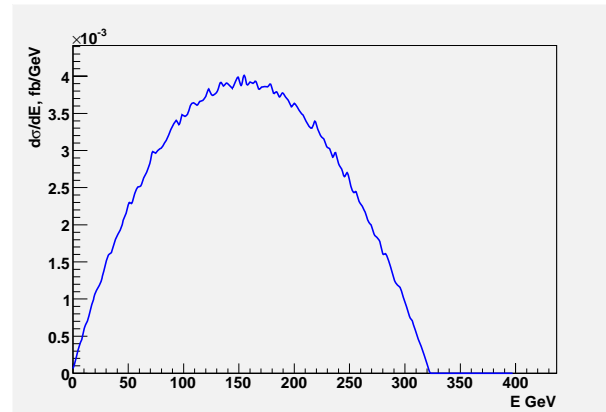
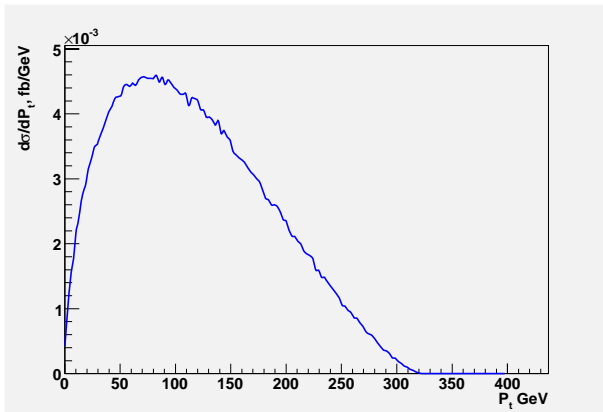
# Few **properties of the signal** reaction relevant for selection procedure

**invariant mass of  $Wb\gamma'$**

- in most of events, the top which radiates the  $\gamma'$  is off-shell,
 
$$t \rightarrow (Wb)\gamma',$$
 with  $M(Wb) = M_{\text{top}}$



**missing transverse momentum and missing energy (= trans.momentum / energy of parafoton)**



**at 1 TeV**

plus some other variables → **selection procedure**



- Selection of only hadronic W decays,  $W \rightarrow q \bar{q}$ , since leptonic decays  $W \rightarrow l \nu$  with neutrino(s) mimic  $\gamma'$  signal events
  - jet finder selects 6- (7-) jet events from  $t\bar{t} \rightarrow (Wb)(Wb) \rightarrow (qq)b (qq)b$
- 

**First attempt** to select the signal -

consecutive cuts on kinematical variables

**18 variables of 3 categories** were considered:

- global event kinematics (aplanarity, thrust, missing energy, ...)
- variables based on jet properties ( $\theta$ -jet, ...)
- variables based on jet correlations (angle between top quarks, ...)

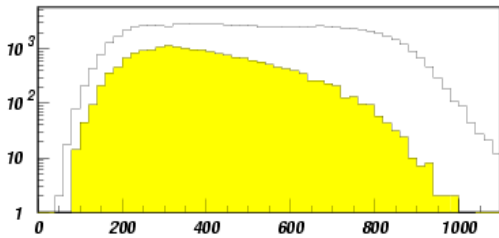


**This method was found to be inefficient to select signal from large bkg, mainly due to the failure of distinct properties between signal and bkg events**

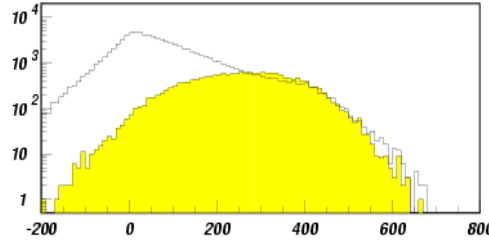
□ background events

■ signal events

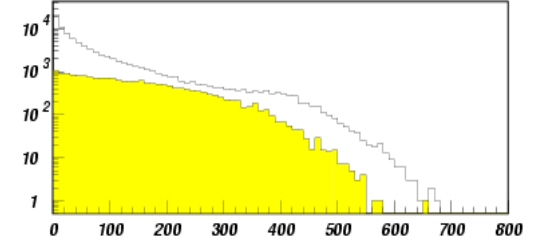
**Signal vs. Background at 1 TeV**



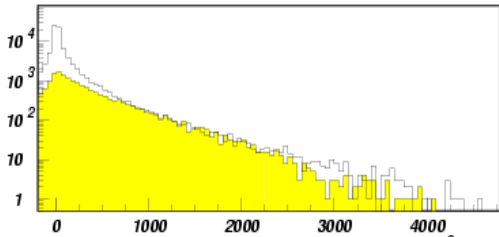
*Sum of trans. energies*



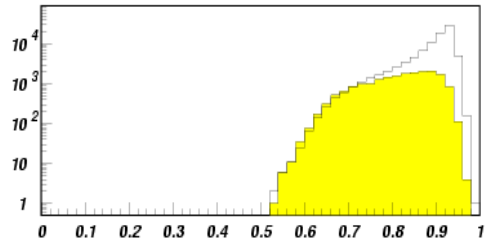
*Total missing energy*



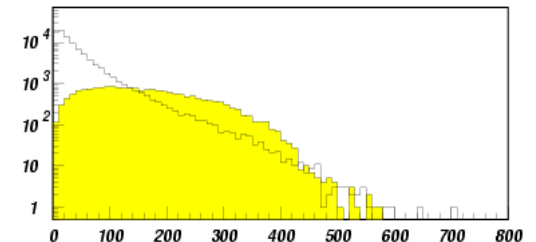
*Total missing trans. energy*



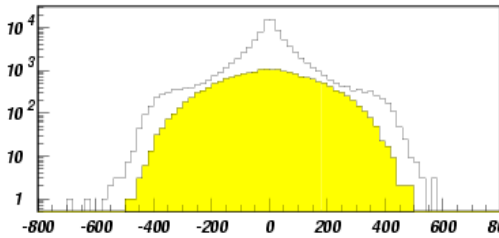
*Miss. mass squared*



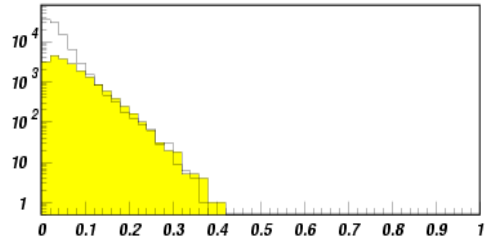
*Thrust*



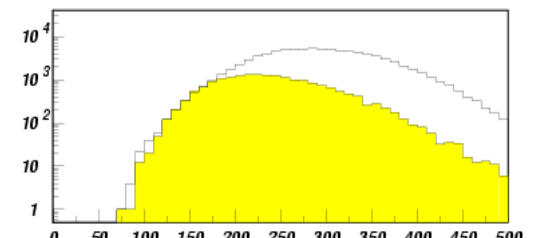
*Total missing trans. momentum*



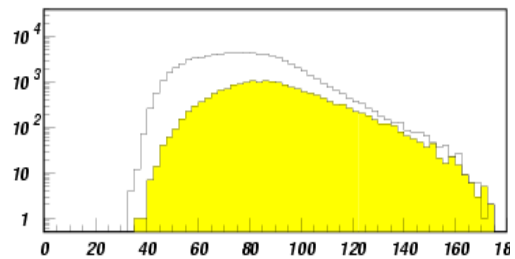
*Total long. momentum*



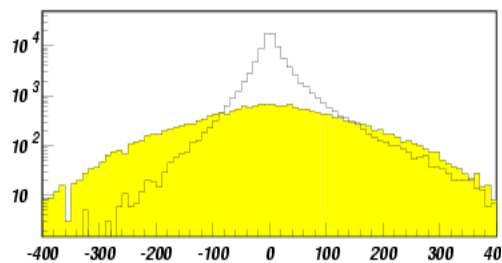
*Aplanarity*



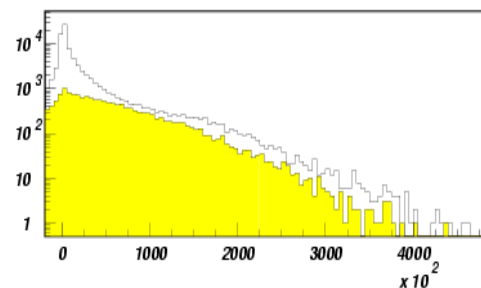
*Max. jet energy per event*



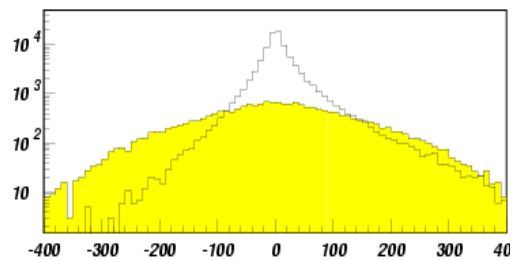
*Max. angle between two jets in given hemisphere*



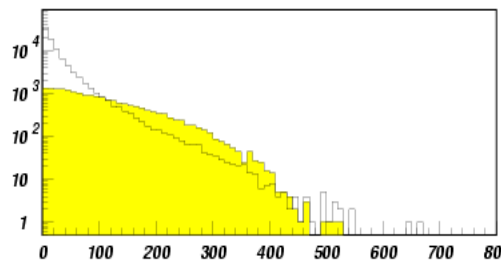
*Missing z-momentum, from jet information*



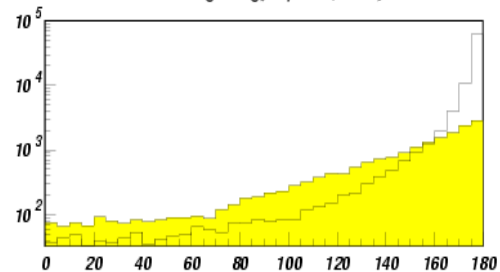
*Missing energy squared, from jet information*



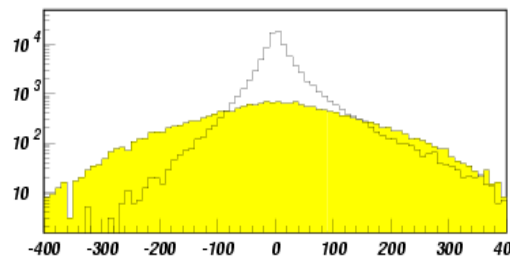
*Missing x-momentum, from jet information*



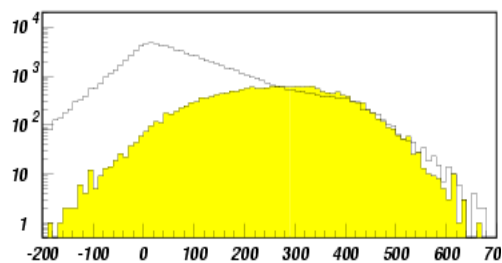
*Missing trans. momentum, from jet information*



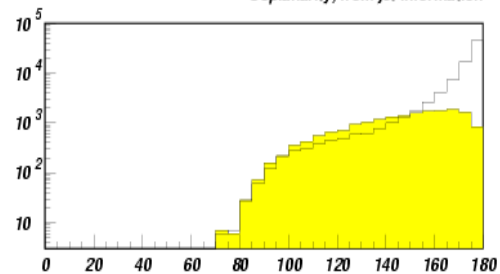
*Coplanarity, from jet information*



*Missing y-momentum, from jet information*



*Missing trans. energy, from jet information*



*Angle between top/anti-top, from jet information*

Thus, selection method was applied:

**multivariate selection procedure**

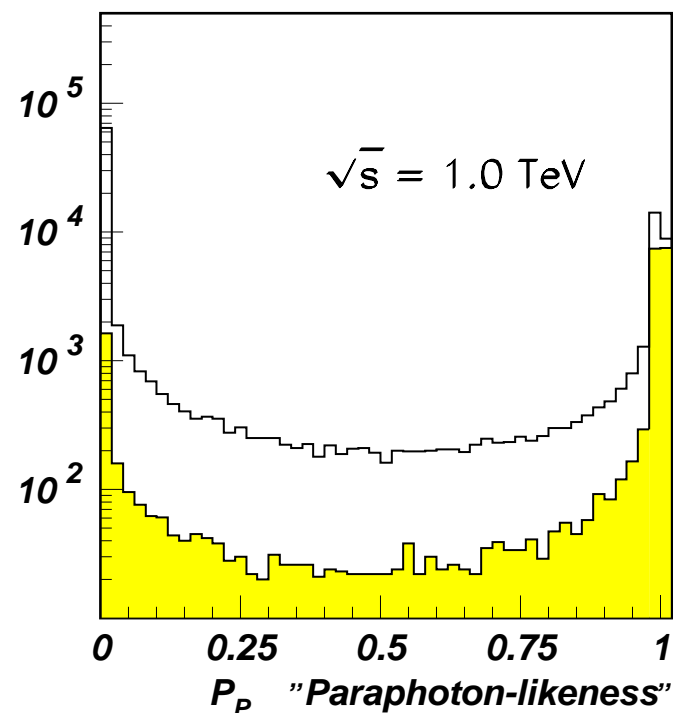
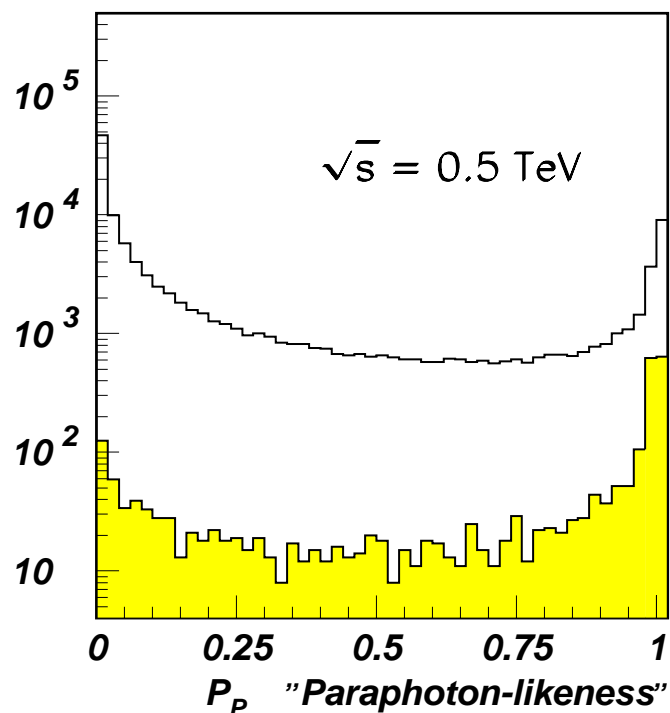
Many variables (18) were considered simultaneously;

- for each event a probability, the **PARAPHOTON-LIKENESS**, was estimated

so that **background events**  $\rightarrow P_p \sim 0$

while **signal events**  $\rightarrow P_p \sim 1$

□ signal plus background  
■ only signal



With  $P_p > 0.98 \rightarrow$  efficient signal selection with small bkg !

$\sqrt{s}$ , TeV	signal events	bkg events	$S/\sqrt{B}$
0.5	1264	10310	12.4
1.0	14882	8116	165.2



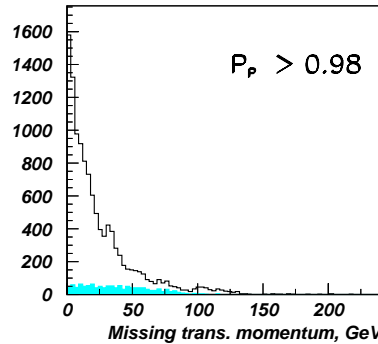
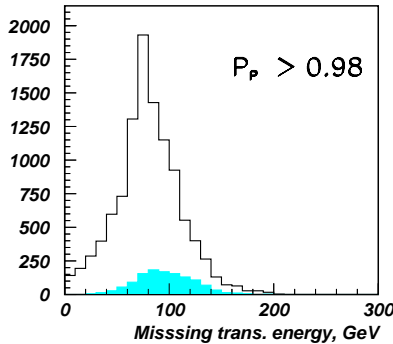
Probability for bkg to fluctuate up to produce signal events

0.5  $10^{-12}$  at 0.5 TeV

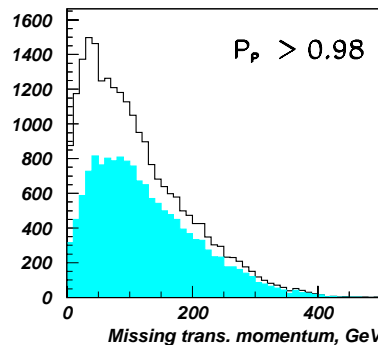
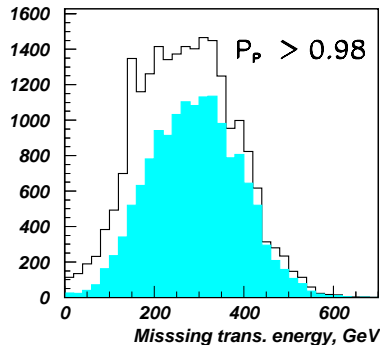
$< 10^{-15}$  at 1.0 TeV  
(assuming Gaussian distribution)

□ signal plus bkg

■ only signal



← 0.5 TeV



← 1.0 TeV



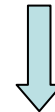
cut on the missing transverse momentum/energy improves the signal/background ratio and provides an almost background-free event sample

e.g. miss.  $p_{\text{trans}} > 100 \text{ GeV}$

→ what value of  $M/\sqrt{s}_+$  is needed for a  $5\sigma$  effect?

• at 0.5 TeV, the limit is  $M/\sqrt{s}_+ = 0.32$  TeV

• at 1.0 TeV, the limit is  $M/\sqrt{s}_+ = 0.61$  TeV



seems to be the most stringent limit  
accessible from the next generation  
of colliders

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