

Precision Measurement of the Stop Mass at the Linear Collider

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*In Collaboration with
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Publication in Preparation*

Introduction

- We have previously studied the light stop, with a small mass difference to the neutralino, in an attempt to understand EW baryo-genesis the asymmetry matter anti-matter and the role of the stop in dark matter annihilation.

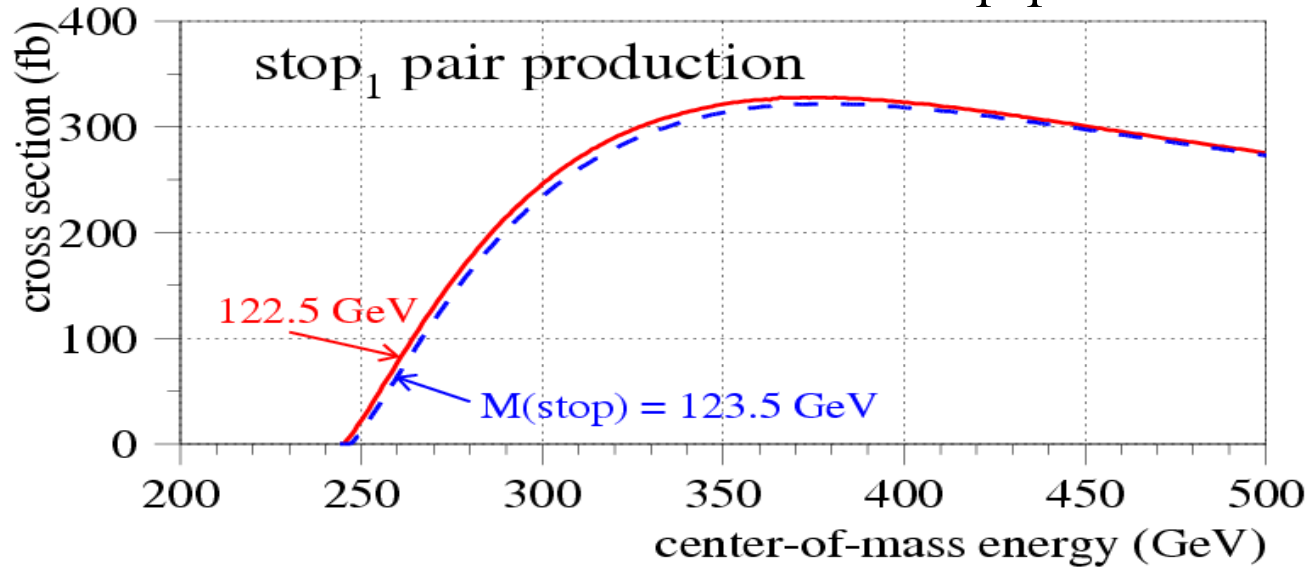
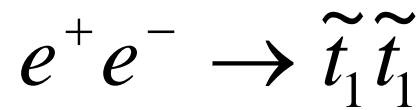
Phys. rev. D 72, 115008(2005)

M. Carena, A. Finch, A. Freitas, C. Milstene, H. Nowak, A. Sopczak

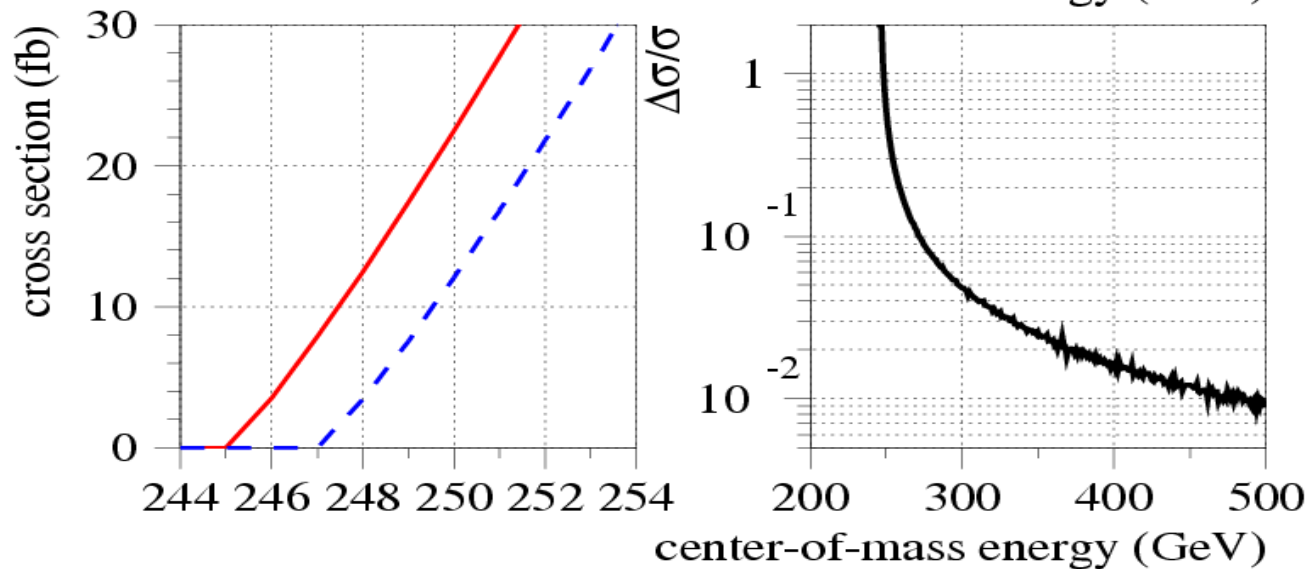
The mass precision measurement reached was $\Delta m \sim 1.2 \text{ GeV}$ including theoretical errors

- This analysis aims at the minimization of the systematics while using more realistic data, stop hadronization/fragmentation included. We will show that:
- The precision is improved in two ways:
 - a/ The systematic uncertainties are minimized by measuring the production cross-section at two energies \rightarrow cancellations .
 - b/ The 2nd energy point chosen at or close to the production energy threshold \rightarrow increased sensitivity to mass changes.
- The stop hadronization is included at production of the data \rightarrow the c quark energy is spread out in the process of hadronization. As a result:
 - the final number of jets increases- the c-tagging is now necessary to identify the charm jets (bench-marking for the vertex detector)
- Two approaches are used, a cut based analysis, a multi-parameters optimization analysis IDA
- The polarization improves further the signal to background ratio

Cross-Section Precision In Production



Cross-sections [fb]
calculated up to NLO
In MC software by
Freitas et al EPJ
C21(2001)361,
EPJ C34(2004)487



The Method

$$\sigma = \frac{N - B}{\varepsilon L}$$

$$Y(M_x \sqrt{s_{th}}) = \frac{N_{th} - B_{th}}{N_{pk} - B_{pk}} = \frac{\sigma(\sqrt{s_{th}}) \varepsilon_{th} L_{th}}{\sigma(\sqrt{s_{pk}}) \varepsilon_{pk} L_{pk}}$$

s - the cross-section [fb]

N - the number of selected data events

B - number of estimated background events

s -Square of the energy in center of Mass

N_{th} , B_{th} , s_{th} at or close to production threshold

N_{pk} , B_{pk} , s_{pk} at peak value

ε_{th} and ε_{pk} - total efficiency & acceptance threshold & peak

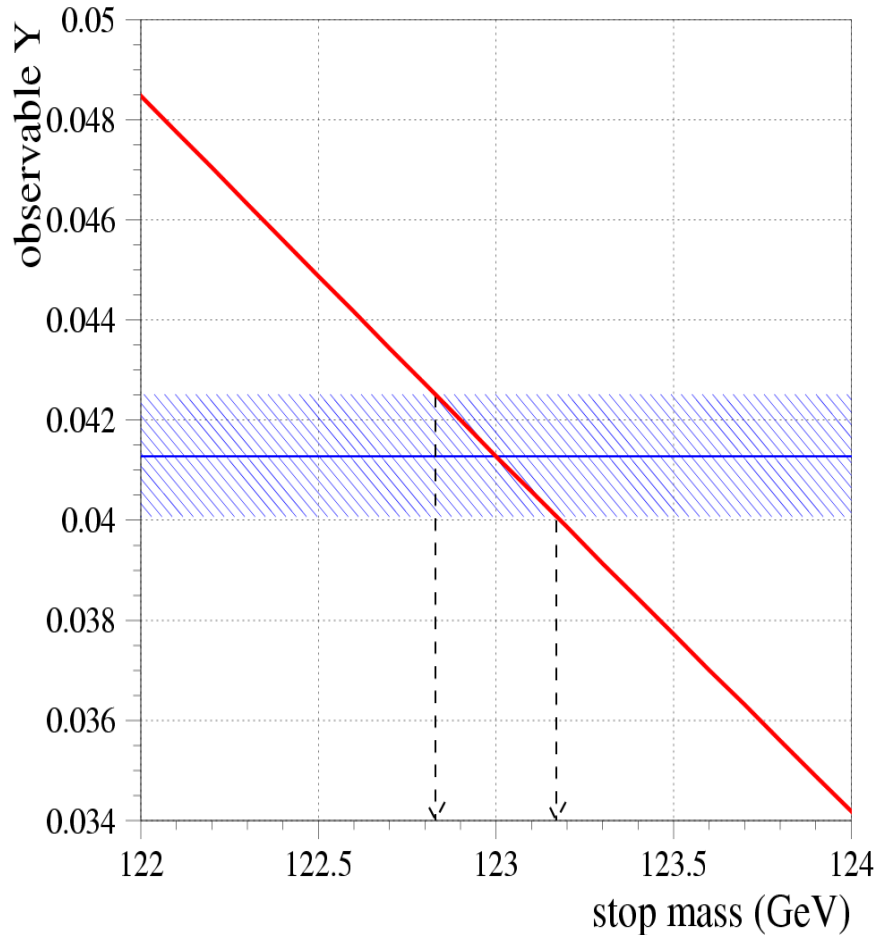
L_{th} ; and L_{pk} -Integrated luminosity

M_x : Mass to be determined with high precision.

Y - ratio of signals at threshold and peak \rightarrow Allows Reduction of systematic uncertainty as well as uncertainties from L measurement.

Remark: yield close to threshold is very sensitive to $M_x \rightarrow$ choice of N_{th} and B_{th} ..

Determination of the Stop Mass



$Y=f(M_x)$ calculated from the theoretical cross-section is drawn in Red (NLO)
Y from the data the blue line.

As an example, Assume 3% variation of Y,
The blue hashed region \rightarrow one obtains
 \rightarrow Precision $\Delta M_x \sim \pm 0.016$, the 2 vertical arrows

The Scenario depicted:

$E_{CM}=260\text{GeV}$ with $s=9.2\text{ fb}$ and $s=77\text{fb}$
at peak

Remark: Assumed luminosities

$L_{th}=50\text{fb}^{-1}$ (260 GeV), $L_{pk}=500\text{fb}^{-1}$ (500 GeV)

$$e^+ e^- \rightarrow \tilde{t}_1 \tilde{t}_1^- \rightarrow c \tilde{\chi}_0^1 \bar{c} \tilde{\chi}_0^1$$

Theoretical Motivation

- Electroweak Baryogenesis:

Sakharov Requirements:

- 1- Baryon Number Violation - (SM - Anomalous process)
- 2- C & CP violation - (SM-Quark CKM mixing)
- 3- Departure from Equilibrium - (SM-at EW phase transition)

Limitations of SM:

2) *Not Enough CP violation & 3) $\rightarrow M_{Higgs} < 40 \text{ GeV}$, LEP Bound*

$M_{Higgs} > 114.4 \text{ GeV}$

\rightarrow Supersymmetry with light scalar top, below the top mass: $m\tilde{t}_1 < m_t$

- Dark Matter

The Supersymmetric Lightest particle (LSP), in the MSSM, the neutralino X_1^0 is a candidate

However, the annihilation cross-section $s_a(X_1^0, X_1^0)$ too small

But for $m\tilde{t}_1 - m X_1^0 \sim 15\text{-}30 \text{ GeV}$, there is co-annihilation between the \tilde{t}_1 and the $X_1^0 \rightarrow s_a(X_1^0, \tilde{t}_1) + s_a(X_1^0, X_1^0)$ consistent with dark matter.

$$e^+ e^- \rightarrow \tilde{t}_1 \tilde{t}_1^- \rightarrow c \tilde{\chi}_0^1 c \tilde{\chi}_0^1$$

A scan in the super-symmetry parameter space

(*hep-ph/0403224v2-2004*) C. Balazs, M. Carena, C. Wagner)

Baryogenesis \rightarrow ($m_{\tilde{t}_1} < m_{top}$ && $m_{\tilde{t}_1} > 120$ GeV) ; Higgs involved in the symmetry breaking mechanism $m_{Higgs} = 114.4$ GeV

\rightarrow Our points $m_{\tilde{t}_1} = 122.5$ GeV; $m_{\tilde{\chi}_0^1} = 107.2$ GeV ; $\Delta m = 15.3$ GeV

Events Final State :

- Stop Hadronization \rightarrow the final state jets smeared :
due to Radiation + Fragmentation
- Soft Multi-jets in the final state
- Stop Hadronization \rightarrow the final state jets smeared :
due to gluon radiation + fragmentation
- At ECM=260 GeV mostly 2 jets, carry the charm.
- At ECM=500 GeV 2jets \rightarrow 2,3,4 jets (more energy available in the CM)
 \rightarrow the Charm tagging (*T. Kuhl*) a necessary tool
to identify the charm jets (Vertex bench-marking)
- Analysis uses N-tuple tool incorporating jet finding algorithm (*T. Kuhl*)

Simulation Characteristics

- Signal and Background generated with: Pythia (6.129)
Simdet (4-0-3)– Circe(1.0)
 - Hadronisation and fragmentation of the \tilde{t} and the fragmentation of the c quark from the Lund string fragmentation Pythia uses Peterson fragmentation (*Peterson et al PR D27:105*)
 - The \tilde{t} fragmentation is simulated using Torbjorn 's code
[//http://www.thep.lu.se/torbjorn/pythia/main73.f](http://www.thep.lu.se/torbjorn/pythia/main73.f)
- The \tilde{t}_1 quark is **set stable** until **after fragmentation** where it is Allowed to **decay again** as described in (*Kraan, EPJ C37:91*)
 - *The stop fragmentation parameter is set relative to the bottom fragmentation Parameter, $\tilde{\alpha}=e_b*m_b^2/m_{\tilde{t}}^2$; $e_b=-0.0050 \pm 0.0015$ following (*OPAL,EPJ C6:225*)*
 - *Newer:*
 - $e_b=-0.0031 \pm 0.0006$ -*ALEPH-phys.lett B152,30(2001)*
 - $e_b=-0.0041 \pm 0.0004$ (*OPAL*)-*Eu.Phys.J,C29,463(2003)*
 - *Variation in e_b was too big by a factor 4, it will be corrected for in the systematics*
- Signal and Background are generated in each channel for the given luminosity in conjunction to the cross-sections

The cross-sections

| Process | s[pb] at ECM=260GeV | | | s[pb] at ECM=500GeV | | |
|-----------------------------|---------------------|-----------|-----------|---------------------|-----------|-----------|
| P(e-)/ P(e+) | 0/0 | -80%/+60% | +80%/-60% | 0/0 | -80%/+60% | +80%/-60% |
| $\tilde{t}_1 \tilde{t}_1^*$ | 0.032 | 0.017 | 0.077 | 0.118 | 0.072 | 0.276 |
| W W | 16.9 | 48.6 | 1.77 | 8.6 | 24.5 | 0.77 |
| Z Z | 1.12 | 2.28 | 0.99 | 0.49 | 1.02 | 0.44 |
| Wenu | 1.73 | 3.04 | 0.50 | 6.14 | 10.6 | 1.82 |
| eeZ | 5.1 | 6.0 | 4.3 | 7.5 | 8.5 | 6.2 |
| qq, qq \neq tt | 49.5 | 92.7 | 53.1 | 13.1 | 25.4 | 14.9 |
| tt | 0.0 | 0.0 | 0.0 | 0.55 | 1.13 | 0.50 |
| 2 γ ($p_t > 5$ GeV) | 786 | | | 936 | | |

Table 1

A. Freitas et al EPJ C21(2001)361, EPJ C34(2004)487 and GRACE and COMPHEP -Next to leading order, assuming a stop mixing angle (0.01)

Pre-Selection Cuts

- A short list of the sequential cuts applied as a pre-selection first, allowed larger samples to be produced
- The pre-selection cuts are the same at the 500 and 260 GeV unless listed in parenthesis for 500 GeV

Pre-selection: 260GeV ;(500 GeV)

- $4 < \text{Number of Charged tracks} < 50$
- $P_t > 5 \text{ GeV}$
- $\cos\theta_{\text{Thrust}} < 0.8$
- $|P_l / P_{\text{tot}}| < 0.9$
- $E_{\text{vis}} < 0.40 \text{ ECM}; (E_{\text{vis}} < 0.76 \text{ ECM})$
- $M(\text{inv}) < 200 \text{ GeV}$

The cuts were refined further at Selection as shown next

Selection Cuts at $E_{CM}=260, 500$ GeV

| Variable | ECM 260 GeV | ECM 500 GeV |
|-------------------------------------|---|--|
| Number of jets | $N_{\text{jets}} = 2$ | $N_{\text{jets}} = 2 \ \& \ E_n < 25 \text{ GeV}$ $n=3,4$ |
| Number of charged tracks | $5 = N_{\text{tracks}} = 25$ | $5 = N_{\text{tracks}} = 20$ |
| Transverse Momentum p_t | $15 < p_t < 45 \text{ GeV}$ | $22 < p_t < 50 \text{ GeV}$ |
| Thrust T | $0.77 < T < 0.97$ | $0.55 < T < 0.90$ |
| Longitudinal Momentum | $ p_L / p_{\text{tot}} < 0.85$ | $ p_L / p_{\text{tot}} < 0.85$ |
| Visible Energy E_{vis} | $0.1 < E_{\text{vis}} / \text{ECM} < 0.3$ | $0.1 < E_{\text{vis}} / \text{ECM} < 0.3$ |
| Acoplanarity F_{acop} | $ \cos(\text{acop}) < 0.9$ | $ \cos(\text{acop}) < 0.9$ |
| Invariant mass of jet pair m_{jj} | $m_{jj}^2 < 5500 \text{ GeV}^2$ or $m_{jj}^2 > 8000 \text{ GeV}^2$ | $m_{jj}^2 < 5500 \text{ GeV}^2$ or $m_{jj}^2 > 10000 \text{ GeV}^2$ |
| Charm tagging likelihood P_c | $P_c > 0.6$ | $P_c > 0.6$ |
| <i>Signal Efficiency</i> | <i>0.340</i> | <i>0.212</i> |

Table 2

In order to optimize the cancellation of the systematics we aim to have a selection as similar as possible at the two energies. (cancellation in $Y=(N_{\text{th}}-B_{\text{th}})/(N_{\text{pk}}-B_{\text{pk}})$)
The two-photons background did require a 5GeV p_t pre-selection cut.

Events Generated and After Sequential cuts

| | L=50fb ⁻¹ at ECM=260GeV | | | L= 500fb ⁻¹ at ECM=500GeV | | |
|-----------------------------|------------------------------------|------------|------------|--------------------------------------|--------------|-------------|
| P (e-)/ P(e+) | Generated | 0/0 | +80%/-60% | Generated | 0/0 | +80%/-60% |
| $\tilde{t}_1 \tilde{t}_1^*$ | 50000 | 543 | 1309 | 50000 | 12514 | 29270 |
| WW | 180000 | 38 | 4 | 210000 | 91 | 8 |
| ZZ | 30000 | 8 | 7 | 30000 | 90 | 81 |
| Wenu | 210000 | 208 | 60 | 210000 | 18540 | 5495 |
| eeZ | 210000 | 2 | 2 | 210000 | <18 | <15 |
| qq, q≠t | 350000 | 42 | 45 | 350000 | 37 | 43 |
| tt | - | 0 | 0 | 180000 | 18 | 17 |
| 2-Photons | 1.6 10 ⁶ | 53 | 53 | 8.5x10 ⁶ | 31 | 31 |
| <u>Total backgrd</u> | - | <u>351</u> | <u>171</u> | - | <u>18807</u> | <u>5781</u> |
| <u>S/B</u> | | 1.5 | 7.6 | | 0.7 | 5.2 |

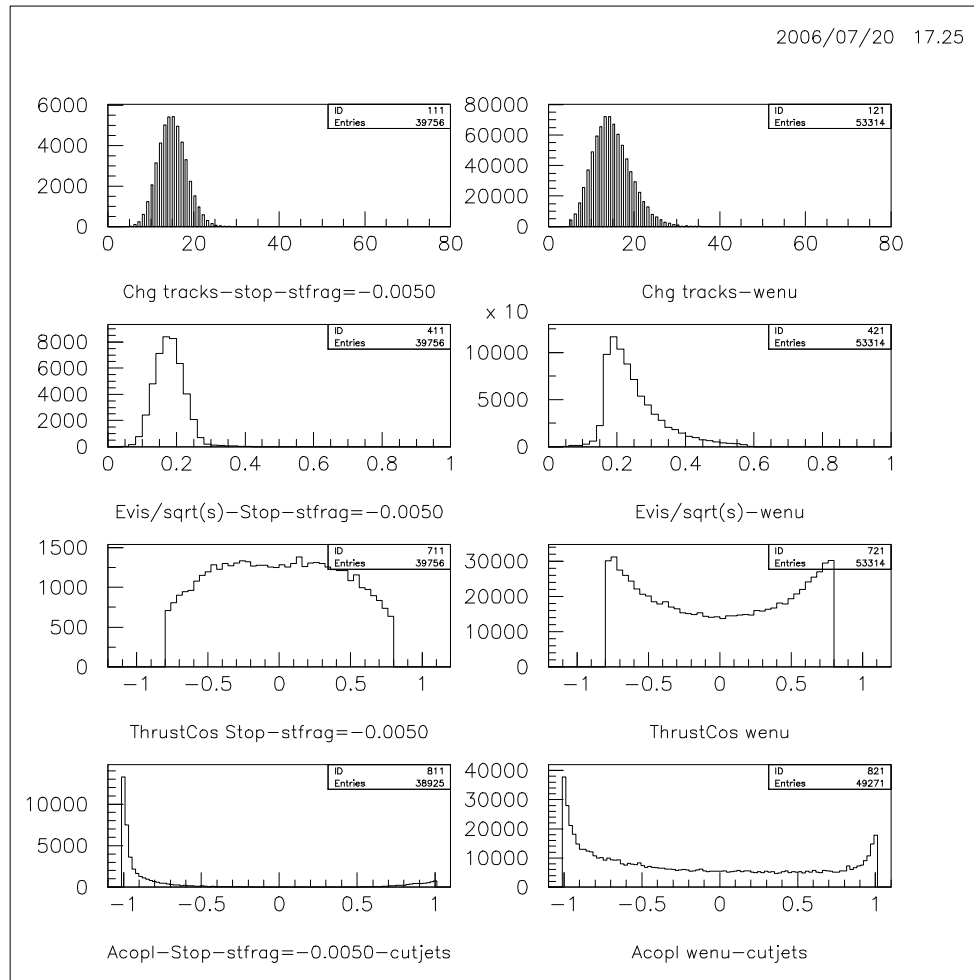
0/0 polarization beam → Unambiguous discovery

+80%/-60% polarization → Precision Measurement

Remark: \tilde{t}_1 fragmentation → the separation from the Wenu more difficult (we had 5044 Wenu at 0/0 polarization)

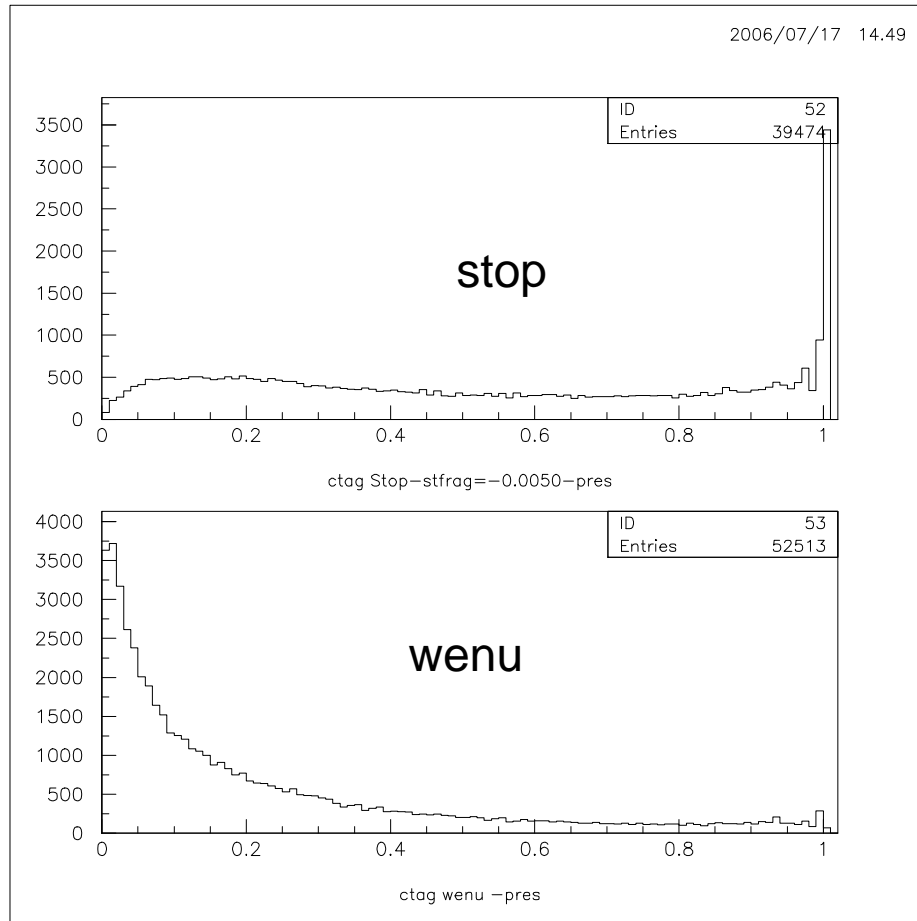
C. Milsténe

Stop/wenu- Variables Distributions



Left column: Stop
Right column: wenu (main Bg)

Charm-tagging

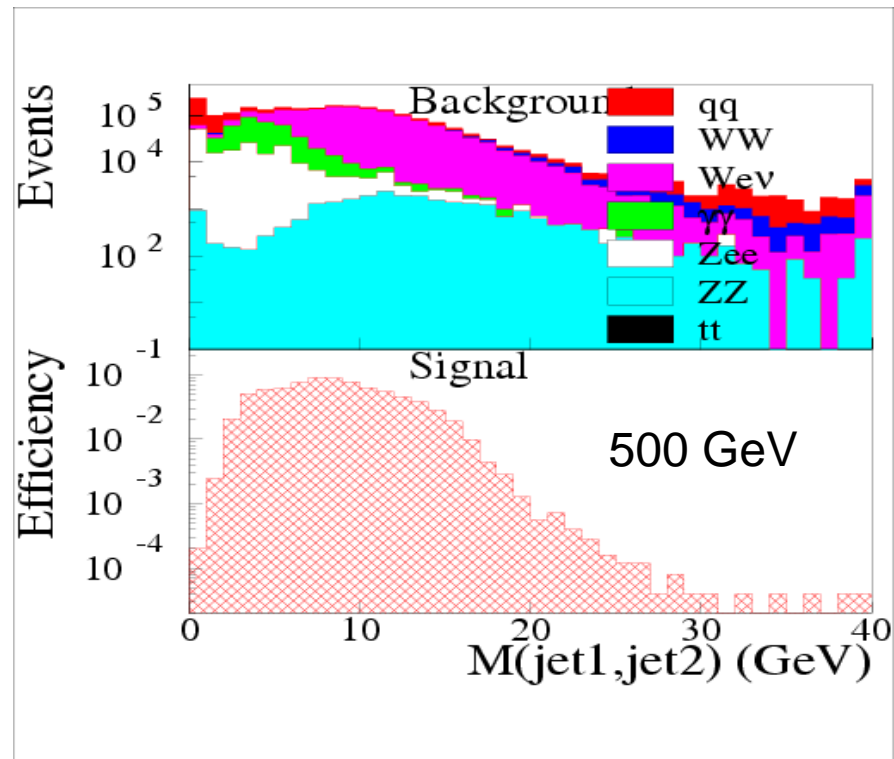
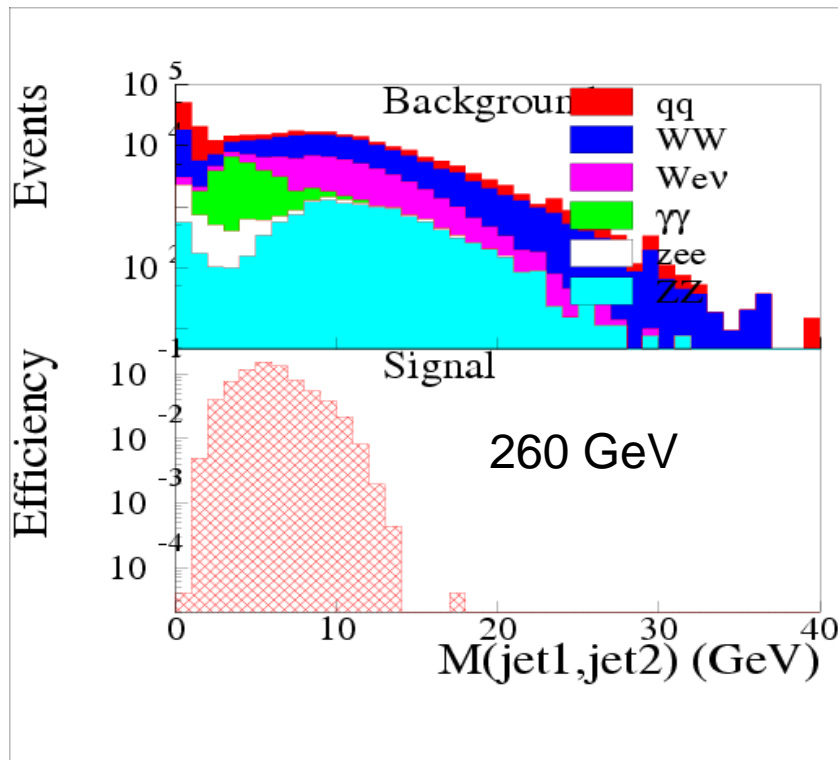


- The charm tagging provides
A good cut between signal
And wenu background
- It has been used here as a
tool to find the charm jets in
The multi-jet event

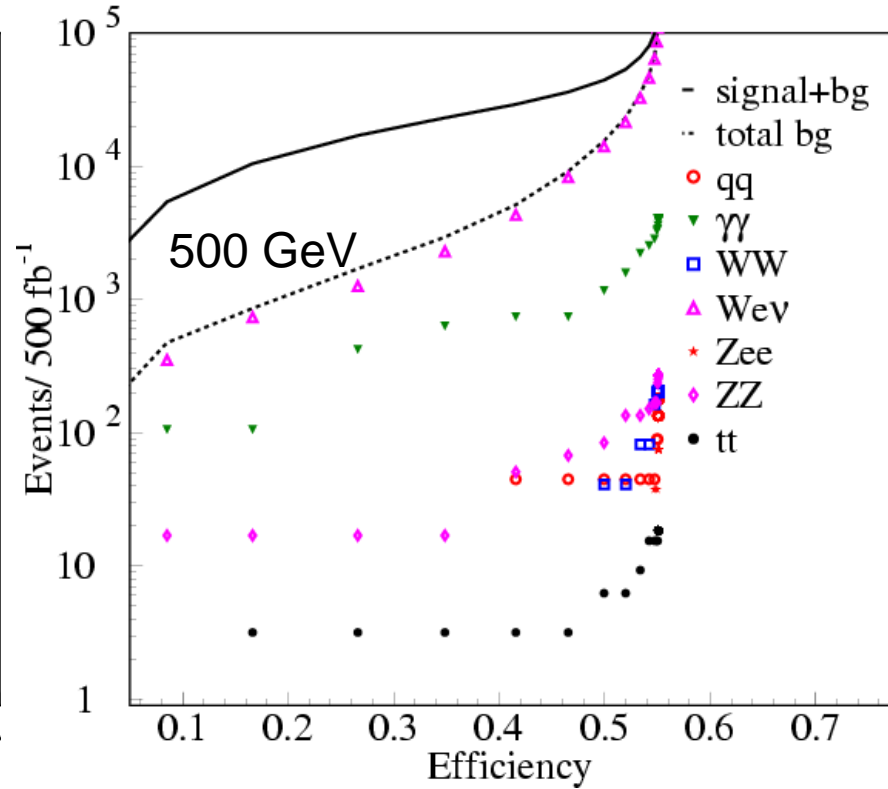
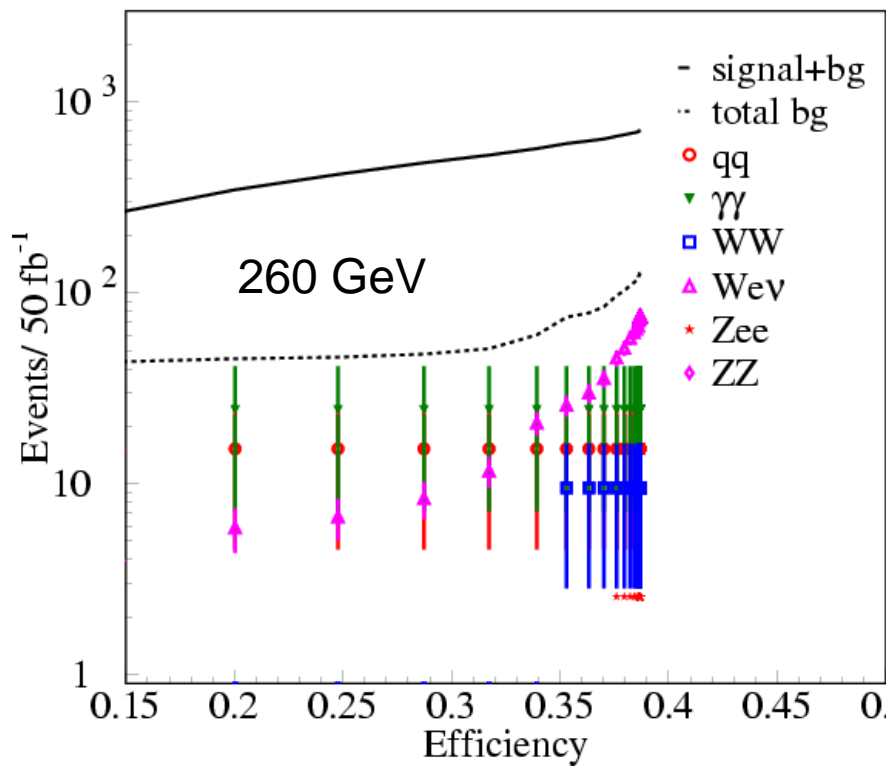
Iterative Discriminant Analysis (IDA)

- A NN approach was also used the Iterative Discriminant Analysis (IDA) . (modified Fisher Disc. Analysis)
- IDA combines the kinematic variables in parallel. The same kinematical variables we used in the cut based analysis . A non linear discriminant function followed by iterations are enhancing the separation between signal and background.
- Both the signal and background have been divided in two equally sized samples, one sample is used for training, the other as data.
- Two IDA steps have been performed, with a cut after the 1st IDA iteration keeping 99% of the signal efficiency while cutting part of the Bg.
- The performance is shown in the two next figures at 260 and 500 GeV.

Invariant Mass Di-Jets 1 Step Before Final IDA



IDA Performance



Systematic Uncertainty in Kinematics Cuts Variables

| Variables | Error on Variable | Relative shift On signal eff vs = 260 GeV | Relative shift On signal eff vs = 500 GeV | Error on Y |
|---------------------|-------------------|---|---|-------------|
| energy scale | 1% | 3.7% (3.4%) | 3.1% (1.3%) | <1% (2.1%) |
| N_{tracks} | 0.5% | Negligible | Negligible | |
| Charm tagging | 0.5% | | | |
| Luminosity | - | 0.4% | 0.2% | 0.4% |
| Charm frag | 0.011 | 0.3% (0.1%) | 0.8% (0.6%) | <1% |
| Stop frag. | 0.0015 | 2.4% (1.2%) | 1.0% (3.5%) | 2.7% (2.8%) |
| Background est. | | 0.8% | 0.1% | <1% |

All cuts are applied to hadronic and jet observables → Calibration quantities are jet energy scale & jet angle.

Based on LEP, we assume 1% energy scale, 1 deg for jet angle

Effect on signal efficiency: Partial cancellation between 260 and 500 GeV

We assume cancellation in total luminosity in Y between 260&500GeV

In parenthesis IDA's values if different

Effect of Stop and Charm Fragmentation

Comparison of the signal generated with and without gluon radiation

→ The signal efficiency changes due to jet number cut is 2.5%

→ We assume an error of 1% for the number of jets

Charm fragmentation parameters assumed as precise as for LEP/OPAL

→ $e_c = -0.0031 \pm 0.0011$

Stop fragmentation is set relative to bottom fragmentation, $e_{\tilde{t}} = e_b (m_b/m_t)^2$

$e_{\tilde{t}} = -0.0050 \pm 0.0015$

$de_c = \pm 35\%$ → Error $dY = 0.2\%$ (from new papers of OPAL and ALEPH)

$de_{\tilde{t}} = \pm 30\%$ → Error $dY = 2.4\%$ however “

Newer analysis from ALEPH and OPAL provide

$e_b = -0.0031 \pm 0.0006$ - ALEPH-phys.lett B152,30(2001)

$e_b = -0.0041 \pm 0.0004$ (OPAL) - Eu.Phys.J, C29,463(2003)

We assume an improvement of a factor ~4 in precision as a consequence we assume that the contribution of the stop fragmentation

→ Error $dY \sim 0.7\%$

→ contribute an error $O(\text{few}\%)$

Theoretical Uncertainties

- Precise cross-section calculations are needed
- $\tilde{t}1$ production receives large corrections from QCD gluon exchange
Between the final state $\tilde{t}1$ (bigger @Threshold) \rightarrow Coulomb corr.
- NLO- QCD corrections $\sim 100\%$ @threshold down to 10% at high energies are included here
- NNLO-QCD corrections are expected of to be same order than NLO based on the results for the top quark. The missing higher order correction $\sim 7\%$ @260GeV, 2.5% @500 GeV
- It is expected that theoretical uncertainties can be brought down by a factor 2
- Here we assume an uncertainty of 3.5% @260GeV and 1% @500 GeV
- The EW corrections : NLO \sim several %, the NNLO $\sim 1\%$
- Combined $\rightarrow \sim 4\%$ @260 GeV and 1.5% @500GeV $\rightarrow dY=5.5\%$

Combined Statistical and Systematic Errors

| Error source for Y | Sequential Cuts | IDA- method |
|----------------------------|-----------------|-----------------|
| Statistical | 3.1% (0.19 GeV) | 2.7%(0.17 GeV) |
| Detector effects(syst.) | 0.6% | 2.1% |
| Charm fragmentation(syst.) | 0.5% | 0.5% |
| Stop fragmentation(syst.) | 0.7% | 0.7% |
| Background contr (syst.) | 0.8% | 0.1% |
| Sum Exp systematics | 1.3%(0.08 GeV) | 3.6%(0.14 GeV) |
| Sum of experimental errors | 3.3% (0.20GeV) | 3.5% (0.22 GeV) |
| Theory for Signal s | 5.5% | 5.5% |
| Total error dY | 6.4%(0.40 GeV) | 6.5%(0.41 GeV) |

Results

Combining the statistical and systematic errors Table 6(*)

$dY=6.4\% \rightarrow dm_{\tilde{\tau}_1} \sim 0.40 \text{ GeV}$ – a factor 3 better (*Phys. rev. D 72, 115008(2005)*)
(dominated by the theory, expected to improve for signal and background)

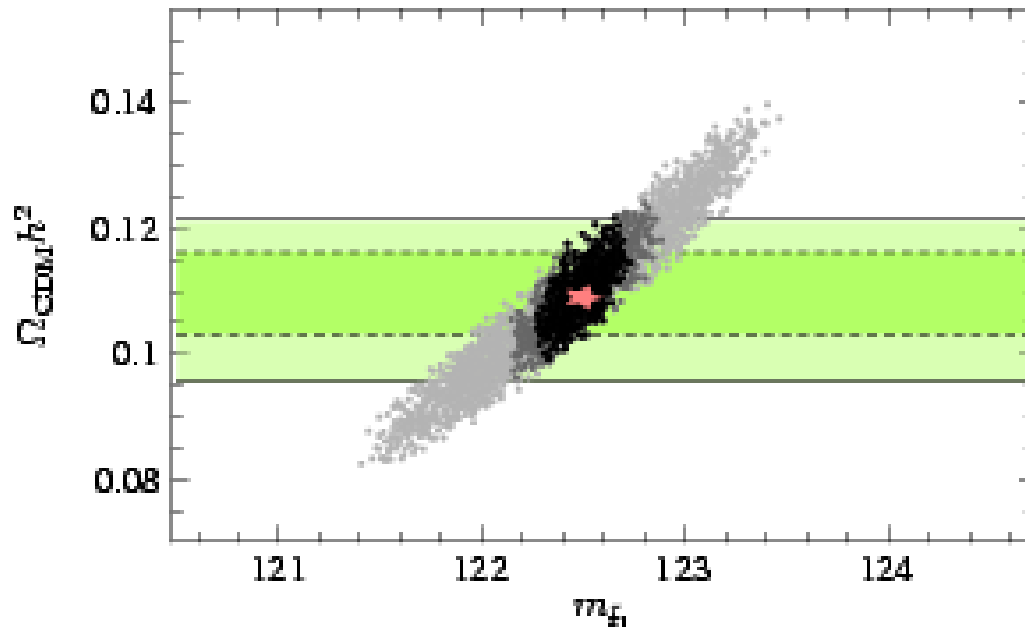
$dY=3.3\% \rightarrow dm_{\tilde{\tau}_1} \sim 0.20 \text{ GeV}$ (cut based experimental errors alone)

$dY=3.5\% \rightarrow dm_{\tilde{\tau}_1} \sim 0.22 \text{ GeV}$ (experimental errors alone & IDA)

→ Improvements in dark matter relic density due to improvement in $dm_{\tilde{\tau}_1}$
is shown in the next figure.

Other limiting factors start to interplay, e.g. the precision on the neutralino
mass $dm_{\chi_1^0} \sim 0.3 \text{ GeV}$,(hep-ph/0608255, M.Carena, A.Freitas)

Dark Matter Relic Abundance= $\Omega_{\tilde{\chi}_1^0} h^2$ ($m_{\tilde{\chi}_1^0}$)



Dark Matter relic density accounting
 The estimated experimental errors
 For stop, Chargino, neutralino and
 Higgs sector –(scan over 1s)
 versus $m_{\tilde{\chi}_1^0}$ for
 $dm_{\tilde{\chi}_1^0}=1.2$ GeV light gray dot
 Previous study
 $dm_{\tilde{\chi}_1^0}=0.40$ GeV dark gray dot
 Now this study
 $dm_{\tilde{\chi}_1^0}=0.20$ GeV black dots
 Expected this study
 with seq. cuts

$dm_{\tilde{\chi}_1^0}=0.40$ GeV $\rightarrow \Omega_{\text{CDM}} h^2 = 0.109+0.0013-0.012$ Exp. Err.+ Th. Err.
 $dm_{\tilde{\chi}_1^0}=0.20$ GeV $\rightarrow \Omega_{\text{CDM}} h^2 = 0.109+0.0010-0.0007$ Exp. Err. Seq. cuts
 $dm_{\tilde{\chi}_1^0}=0.22$ GeV $\rightarrow \Omega_{\text{CDM}} h^2 = 0.109+0.0010-0.0008$ Exp. Err. Seq. cuts
 WMAP: $\Omega_{\text{CDM}} h^2 = 0.1106+0.0056-0.0075$

Conclusion

- More realistic data were produced including hadronization/fragmentation
- The precision, however, improved by a factor three on our previous analysis with $dm_{\tilde{\tau}_1} = 0.40$ GeV
- This method could be applied to other particles e.g. to measure the Higgs mass
- The method improves the precision to the mass determination in two ways
a/ by reducing the systematics in Y- cancellation between the two energy points.
b/ by choosing the energy at threshold, Y extremely sensitive to the mass
- The polarization separates the right-handed signal $\tilde{\tau}_1$ from background.
- Due to hadronization and fragmentation the c-tagging was a necessary tool to identify the charm jets at $E_{CM} = 500$ GeV (benchmark for the vertex detector)
- IDA and the sequential cuts give almost identical results. IDA gives better statistical uncertainties but worse systematics $dm_{\tilde{\tau}_1} = 0.20$ GeV (0.22 GeV-IDA)
- Progress in the theoretical calculations is expected and partly accounted for
- With that precision we become limited by other factors.
- With this mass precision, the calculated relic density is in accordance with WMAP and SLOAN ,
 $dm_{\tilde{\tau}_1} = 0.20$ GeV $\rightarrow \Omega_{CDM} h^2 = 0.109^{+0.0010}_{-0.007}$
 WMAP: $\Omega_{CDM} h^2 = 0.1106^{+0.0056}_{-0.0075}$ C. Mistère

Backup slides

A Sample Parameter Point

- $m_{\tilde{U}_3}^2 = -99^2 \text{ GeV}^2$
- $A_t = -1100 \text{ GeV}$
- $M_1 = 112.6 \text{ GeV}$
- $M_2 = 225 \text{ GeV}$
- $|\mu| = 225 \text{ GeV}$
- $F\mu = 0.2$
- $\tan \beta = 5$

Which gives:

$m_{\tilde{t}_1} = 122.5 \text{ GeV}; m_{\tilde{t}_2} = 4333 \text{ GeV};$

$m_{\tilde{\chi}_1^0} = 107.2 \text{ GeV}; m_{\tilde{\chi}_1^+} = 162.7 \text{ GeV}; m_{\tilde{\chi}_2^0} = 170.8 \text{ GeV}$

$m_{\tilde{\chi}_3^0} = 231.0 \text{ GeV}; m_{\tilde{\chi}_2^+} = 170.8 \text{ GeV}$

$\cos\theta_{\tilde{t}} = 0.010 \sim \tilde{t} \text{ right-handed}$

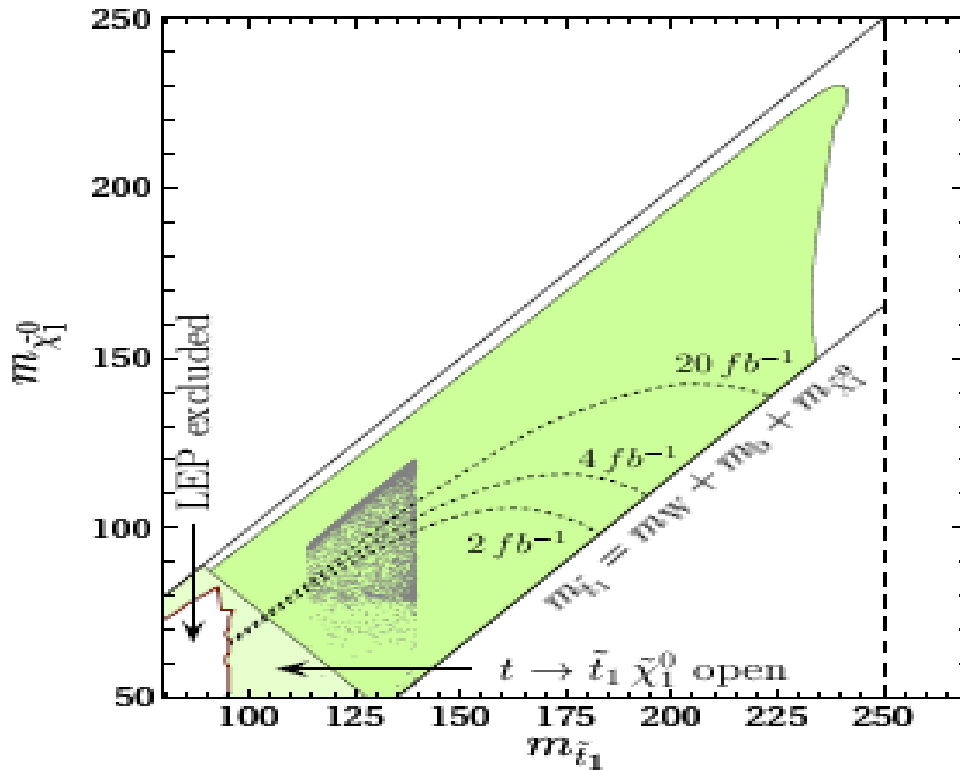
$\rightarrow \Delta m = 15.2 \text{ GeV}$

Events Generated and After IDA Selection

| | L=50fb ⁻¹ at ECM=260GeV | | L= 500fb ⁻¹ at ECM=500GeV | |
|-----------------------------|------------------------------------|-----------|--------------------------------------|-------------|
| P (e-)/ P(e+) | 0/0 | +80%/-60% | 0/0 | +80%/-60% |
| $\tilde{t}_1 \tilde{t}_1^*$ | 618 | 1489 | 24538 | 57394 |
| WW | 11 | 1 | <20 | <2 |
| ZZ | <2 | <2 | 51 | 46 |
| Wenu | 68 | 20 | 4262 | 1263 |
| eeZ | 3 | 2 | <18 | <15 |
| qq, q≠t | 16 | 17 | 45 | 52 |
| tt | 0 | 0 | 3 | 3 |
| 2-Photons | <25 | <25 | 772 | 772 |
| <i>Total background</i> | <u>125</u> | <u>67</u> | <u>5133</u> | <u>2136</u> |
| S/B | 4.9 | 22 | 4.7 | 27 |

The efficiencies improves from 34% ,21.2% cut based → 38.7% ,41.6% IDA, while the background is of the same order of magnitude.

Stop Discovery Reach Snowmass 2005



From Simulations:
strong green region:

$$e^+ e^- \rightarrow \tilde{t}_1 \tilde{t}_1^- \rightarrow c \tilde{\chi}_0^1 \bar{c} \tilde{\chi}_0^1$$

And Significance:

$$(S/\sqrt{S+B}) > 5$$

Background B

Signal $S = \epsilon_s L$

For e, Signal efficiency

For s, Theoretical

cross-section

dark gray region:

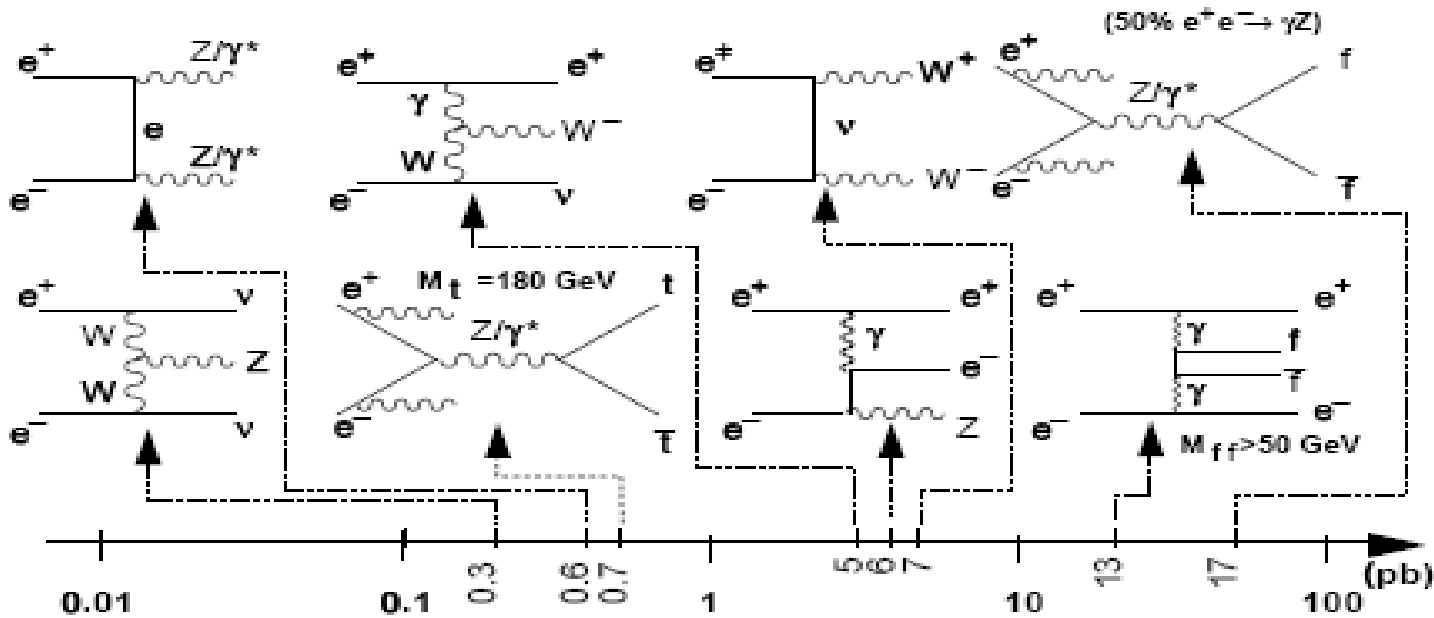
Consistent with DM

And Baryogenesis

Fig 4a-Luminosity: 500 fb⁻¹

$E_{cm} = 500 \text{ GeV}$

Background- Channels @500 GeV



Z Phys. C 76 (1997) 549- A.Bartl, H. Eberl, S. Kraml, W.Majerotto, W.Porod, A. Sopczak

C-Tagging — The Data Samples

- Neural Network (NN):

data used: 255000 stops, $M_{\text{stop}}=120-220$; $D_m=5, 10, 20$
GeV

240000 W_{ev} , the most resilient background

C-tagging-Neural Network Input

- Vertex Case 1: NN Input variables

- *Impact parameter* significance (impact parameter/error) of the 2 most significant tracks in the r-F plane (highest separation power) && their Impact parameters.

- The impact parameter significance & Impact parameters of the 2 tracks in z

- Their momenta

- The joint probability in r- F (tiny beamspot size in that plane) & z

- Vertex Case 2: NN Input variables (all of Case 1+below)

- *Decay Length* significance of the secondary vertex && Decay Length

- Momentum of all tracks associated to the secondary vertex && Multiplicity

- Pt corrected mass of secondary vertex (corrected for neutral hadrons), the pt of the decay products perpendicular to the flight direction (between primary && secondary Vertex) && joint probability in r-F and z

- Vertex Case 3: 2 secondary vertices, the tracks are assigned to the vertex closest to the primary vertex and the NN input variables are those of case 2

Systematic Uncertainty in Kinematics Cuts Variables

| Variable | Error on variable | Error on Y |
|------------------------------|-------------------|------------|
| p_t | 2% | 0.28% |
| $\cos\theta_{\text{Thrust}}$ | 1.8% | 0.18% |
| E_{vis} | 2% | 0 |
| F_{acop} | 1% | 0.08% |
| m_{jj} | 4% | 0.61% |

Table 5

- All cuts are applied to hadronic and jet observables → Calibration quantities are jet energy scale & jet angle.
- Based on LEP, we assume 2% calibration error for jets, 1 deg for jet angle
- Effect on signal efficiency: Partial cancellation between 260 and 500 GeV
- We assume cancellation in total luminosity in Y between 260&500GeV

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