

Hidden sector particle production at the ILC

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1. Introduction

New Physics (NP) models

New Phys. sector
(Hidden sector)

Standard Model
sector

messenger

Completely
or
Partly } SM gauge singlet

Integrating out
or
Be just effective cutoff

Effective Lagrangian: $\mathcal{L} = \frac{c_n}{\Lambda^n} \mathcal{O}_{NP} \mathcal{O}_{SM}$

Candidate of New Particles: scalar, fermion, vector and tensor

We consider \rightarrow scalar

$$\mathcal{L} \sim \frac{c}{\Lambda^{d_{SM}-3}} X \mathcal{O}_{SM} \quad \left\{ \begin{array}{l} X: \text{SM singlet scalar} \\ \Lambda \sim 1\text{TeV} \end{array} \right.$$

Above higher dim. couplings become strong!
 \rightarrow It is Possible to product the hidden scalar at the ILC

But How we can get $\Lambda \sim 1\text{TeV}$?

Ex.) Can be realized by brane world scenario

Analysis of the $X \rightarrow$ **Can apply to many models**

Ex.) Singlet in NMSSM, Radion, Polonyi,...

These NP models have some impact on Higgs pheno. by the X .

2. Possible operators and our studies

Derivative couplings

With gauge boson:

Like Higgs →

$$\frac{X}{\Lambda} G_{\mu\nu}^a G^{a\mu\nu}$$

$$\frac{X}{\Lambda} F_{\mu\nu} F^{\mu\nu}$$

$$\frac{X}{\Lambda} Z_{\mu\nu} F^{\mu\nu}$$

$$\mathcal{L}_{\text{SM}} \sim \frac{\alpha_s}{4\pi v} h G_{\mu\nu}^a G^{a\mu\nu}$$

$$\mathcal{L}_{\text{SM}} \sim \frac{\alpha}{4\pi v} h F_{\mu\nu} F^{\mu\nu}$$

$$\mathcal{L}_{\text{SM}} \sim \frac{eg}{16\pi^2 v} h Z_{\mu\nu} F^{\mu\nu}$$

Derivative

Not like Higgs →

$$\frac{X}{\Lambda} W_{\mu\nu}^+ W^{-\mu\nu}$$

$$\frac{X}{\Lambda} Z_{\mu\nu} Z^{\mu\nu}$$

Non derivative

$$\mathcal{L}_{\text{SM}} \sim 2 \frac{m_W^2}{v} h W_{\mu}^+ W^{-\mu}$$

$$\mathcal{L}_{\text{SM}} \sim \frac{m_Z^2}{v} h Z_{\mu} Z^{\mu}$$

With Higgs boson: $m X H^\dagger H$ $\frac{X}{\Lambda} (H^\dagger H)^2$: Mixing with Higgs boson

With SM fermion: $\frac{X}{\Lambda} \bar{\Psi} \gamma^\mu \partial_\mu \Psi \rightarrow \frac{X}{\Lambda} m_f \bar{\Psi} \Psi$

→ Not important for light fermions

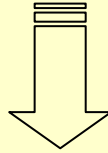
Our studies

Concentrate the X-gauge-gauge couplings

- 1) X-gluon-gluon, X-photon-photon, X-photon-Z
- 2) X-W-W, X-Z-Z

Important point on this study

The X behaves like Higgs in its production process



Interesting issue

How to distinguish X from Higgs at ILC ?

with same masses

We have two methods to distinguish them

1) Big enhancement of the branching ratios

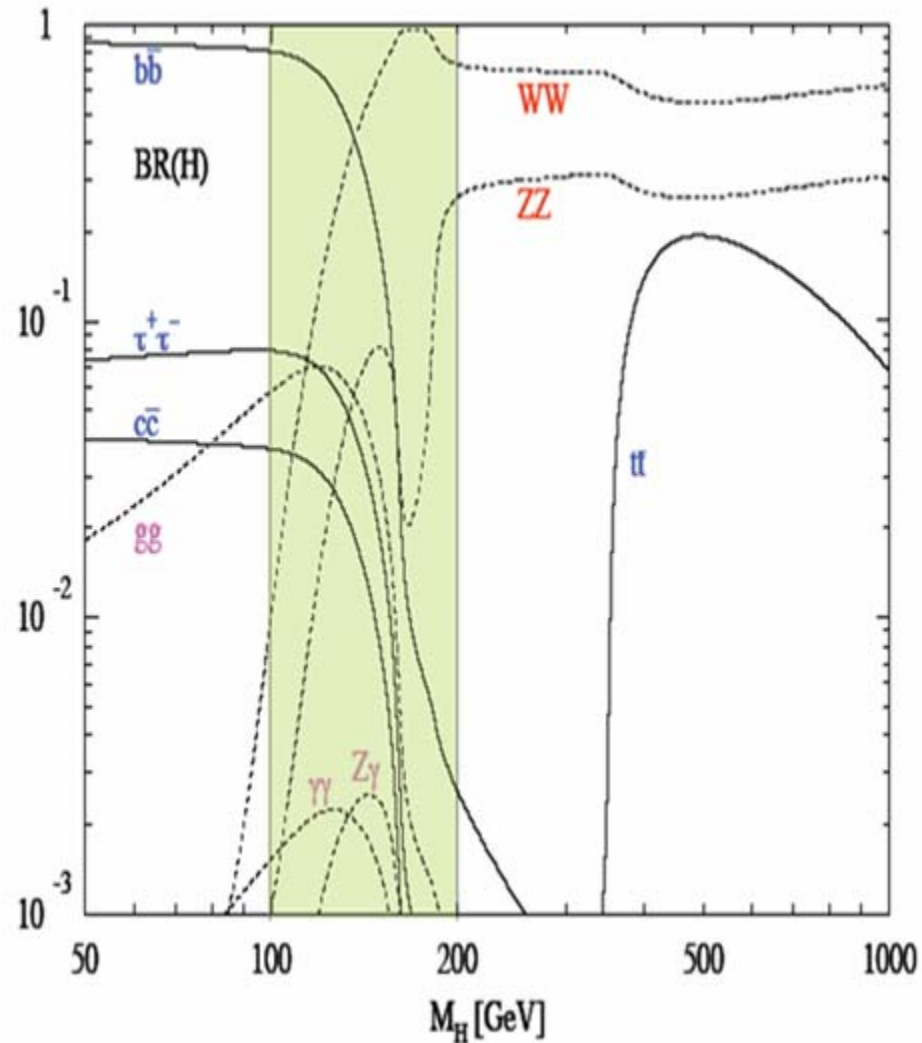
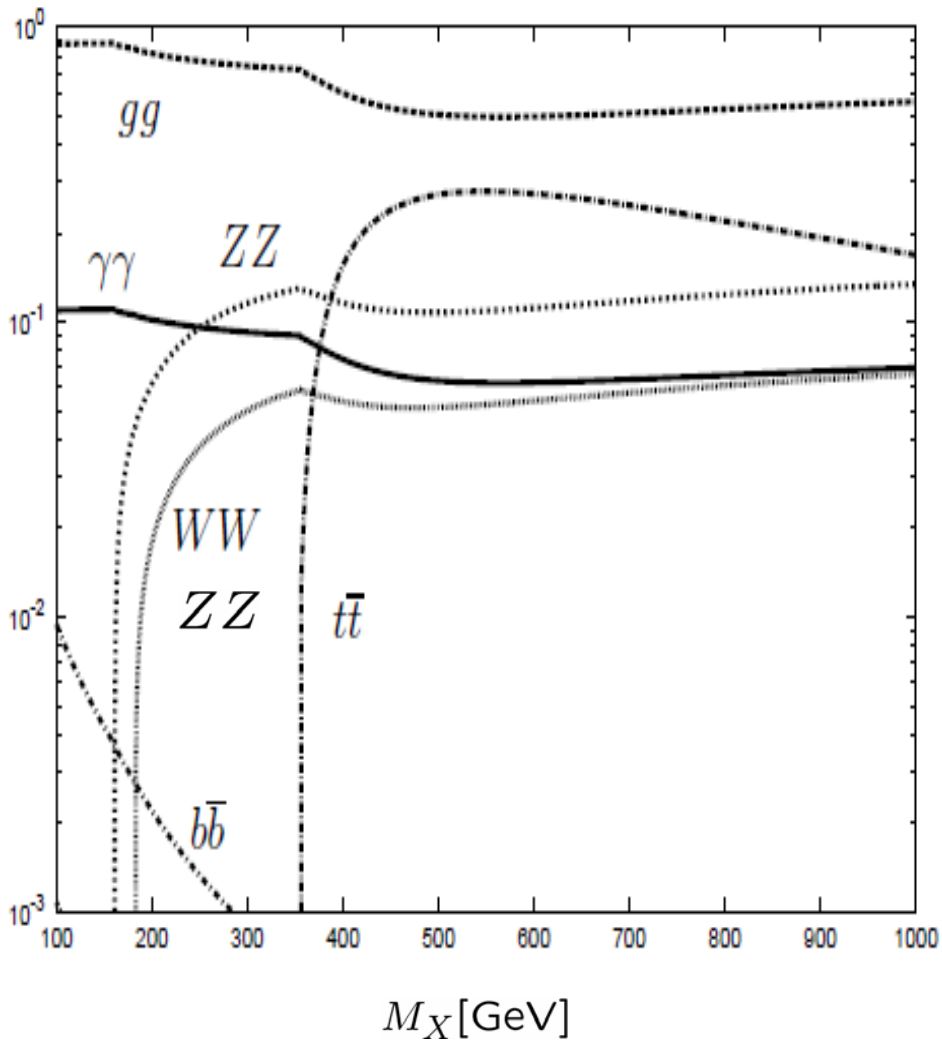
In the SM:

$$\left. \begin{aligned} \mathcal{L}_{SM} &\sim \frac{\alpha_s}{4\pi v} h G_{\mu\nu}^a G^{a\mu\nu} \\ \mathcal{L}_{SM} &\sim \frac{\alpha}{4\pi v} h F_{\mu\nu} F^{\mu\nu} \end{aligned} \right\} \begin{array}{l} \text{No tree level couplings} \\ \text{Loop suppressed} \end{array}$$

For X case

$$\left. \begin{aligned} \frac{X}{\Lambda} G_{\mu\nu}^a G^{a\mu\nu} \\ \frac{X}{\Lambda} F_{\mu\nu} F^{\mu\nu} \end{aligned} \right\} \begin{array}{l} \text{Higher dimensional OPs., but tree level} \\ \\ \frac{\alpha_s}{4\pi v} \sim \frac{1}{10 \text{ TeV}} \quad \text{If } \Lambda \sim 10 \text{ TeV} \end{array}$$

Branching ratio is very much different



Result for universal order one couplings of higher dim. ops.
(In this situation, $Z\gamma$ mode disappear)

2) Difference of coupling manner

In the SM

Higgs mechanism

$$\mathcal{L}_{SM} \sim 2 \frac{m_W^2}{v} h W_\mu^+ W^{-\mu}$$

$$\mathcal{L}_{SM} \sim \frac{m_Z^2}{v} h Z_\mu Z^\mu$$

Mostly couples with longitudinal mode

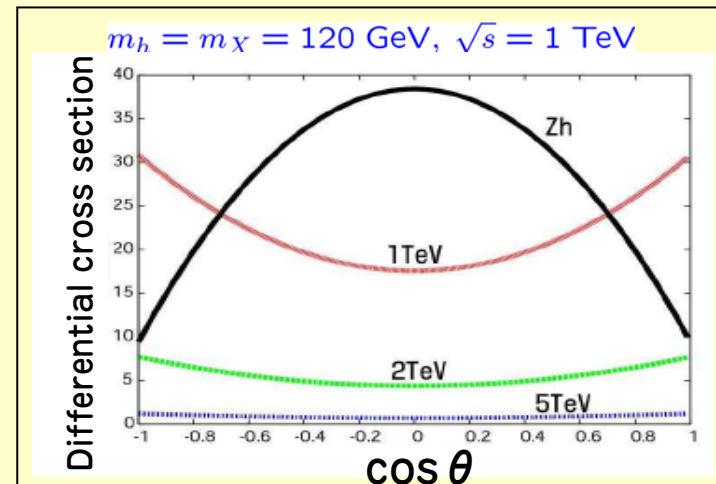
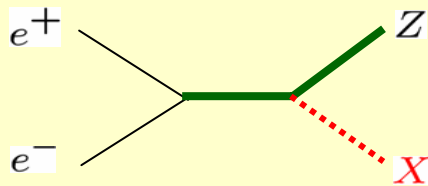
For X case

$$\frac{X}{\Lambda} W_{\mu\nu}^+ W^{-\mu\nu}$$

$$\frac{X}{\Lambda} Z_{\mu\nu} Z^{\mu\nu}$$

Mostly couples to transverse mode

X or Higgs production at ILC



$\Lambda \sim 1 \text{ TeV} \rightarrow X$ strongly couples to gluons and photons

\rightarrow Tevatron bound is severe

Higgs search through two photon decay mode

Merrana & Wells, PRD 63 015006 (2000)

$$R_{gg} = \frac{\sigma_X \times Br(X \rightarrow \gamma\gamma)}{\sigma_H \times Br(H \rightarrow \gamma\gamma)}$$

Tevatron Run II

$$\mathcal{L} \sim 1 \text{ fb}^{-1} \quad R_{gg} < 50$$

Small couplings with gauge bosons

\rightarrow Become consistent with Tevatron bound

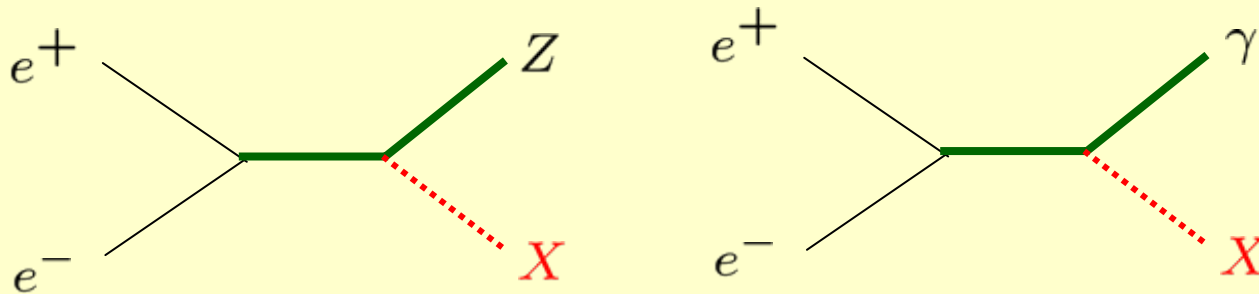
However

Small X-Z-Z coupling \rightarrow The X cannot be detected at ILC

In this time, we concentrate a gluophobic scenario
(Can be realized brane world scenario)

3. Monte-Carlo Simulations

We study Associate X production @ ILC $\sqrt{s} = 500$ GeV



$X \rightarrow \gamma\gamma$

We used fast Monte-Carlo(QuickSim)

Interesting things to see and to compare to Higgs boson

Angular distribution of Z, photon

Polarization of Z,

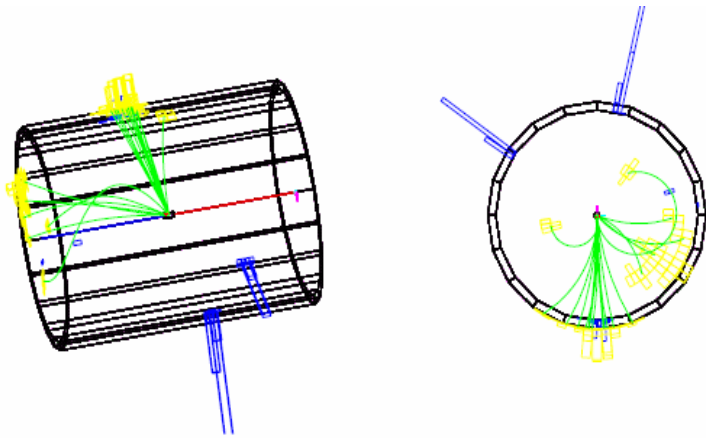
Two photons in the final state,

Less bb in the final state, etc

Our analysis is applicable to

anomalous Higgs Branching Ratio etc.

$$e^+e^- \rightarrow ZX; X \rightarrow \gamma\gamma$$



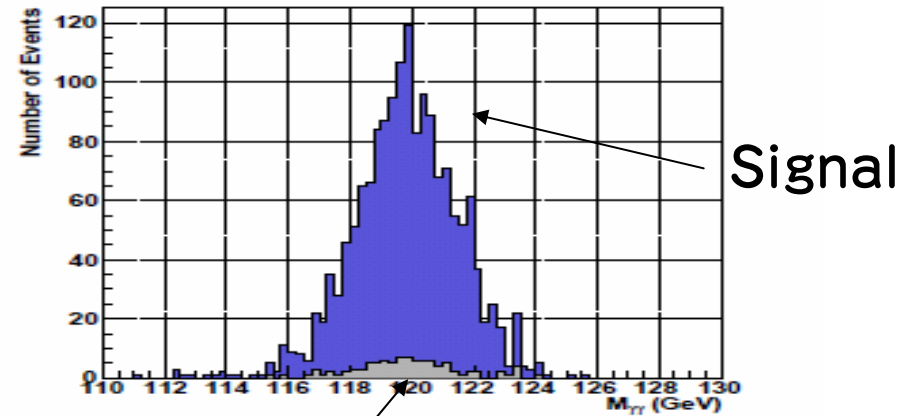
$$\sigma_h \sim 57 \text{ fb} \quad \sigma_X \sim 21 \text{ fb}$$

$$m_X = m_h = 120 \text{ GeV}$$

$$Br(X \rightarrow \gamma\gamma) = 1$$

$$Br(h \rightarrow \gamma\gamma) = 0.002$$

Invariant mass dist. from 2γ



Background from SM Higgs

This mode shows a clear peak at m_X

Can we distinguish the X or Higgs?

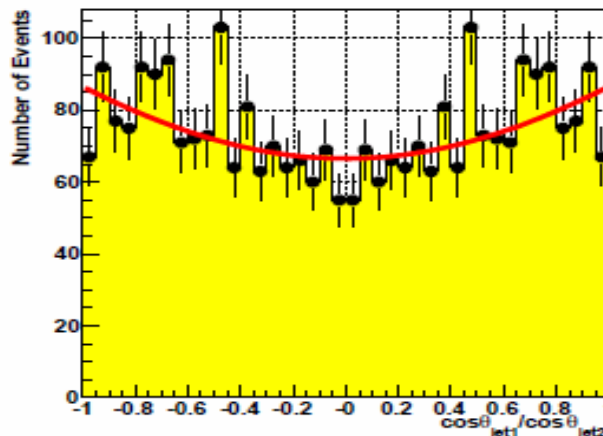
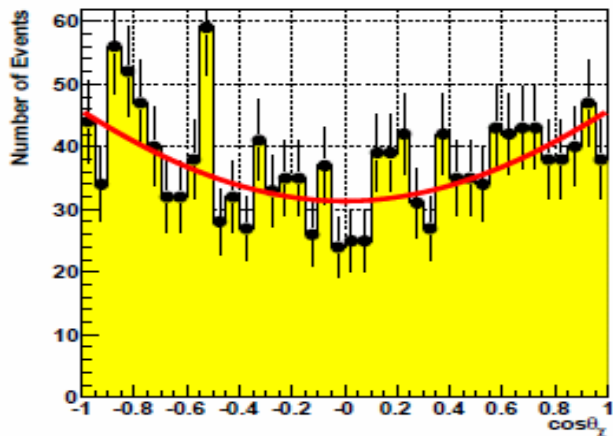
But...there is different model which give a similar result

In NMSSM, if pseudo-scalar A is almost massless,

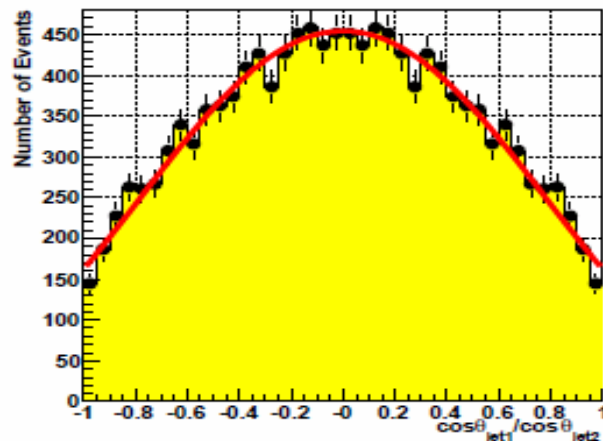
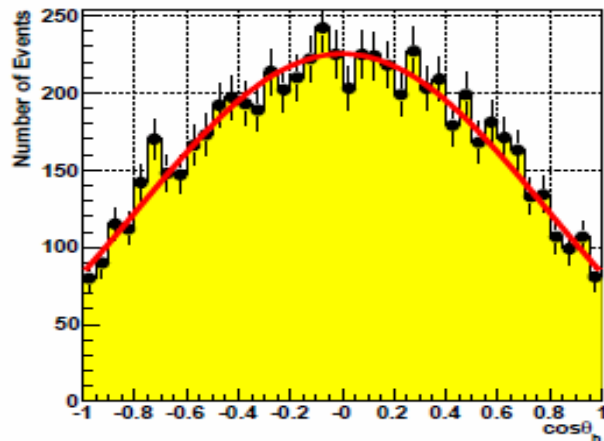
A decays to two photons: $e^+e^- \rightarrow ZH$; $H \rightarrow AA \rightarrow \gamma\gamma\gamma\gamma$

Since the pseudo-scalar is very light, two photons produced in its decay are almost collinear and will be detected as a single photon.

However angular distributions are quite different!



X production



Higgs production

production angle of scalar

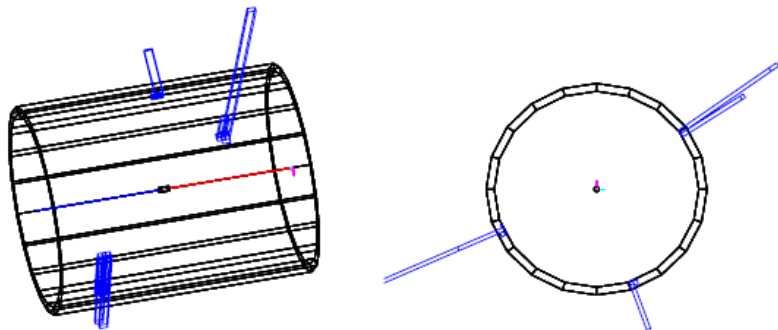
Reconstructed jets from Z

Cut	$Z\chi; \chi \rightarrow \gamma\gamma$	$Zh; h \rightarrow \gamma\gamma$	$Zh; h \rightarrow A^0 A^0$
No Cut	2187 (1.0000)	142 (1.000)	7087 (1.0000)
$N_{particles} \geq 4$	1738 (0.7947)	106 (0.747)	5692 (0.8032)
$N_{gammas} \geq 2$	1521 (0.8751)	96 (0.906)	4865 (0.8547)
Cut on $M_{\gamma\gamma}$	1499 (0.9855)	95 (0.990)	4828 (0.9924)
$N_{jets} = 2$ for Ycut = 0.004	1498 (0.9993)	95 (1.000)	4825 (0.9994)
Total Efficiency	0.6850 ± 0.0099	0.669 ± 0.040	0.6808 ± 0.0055

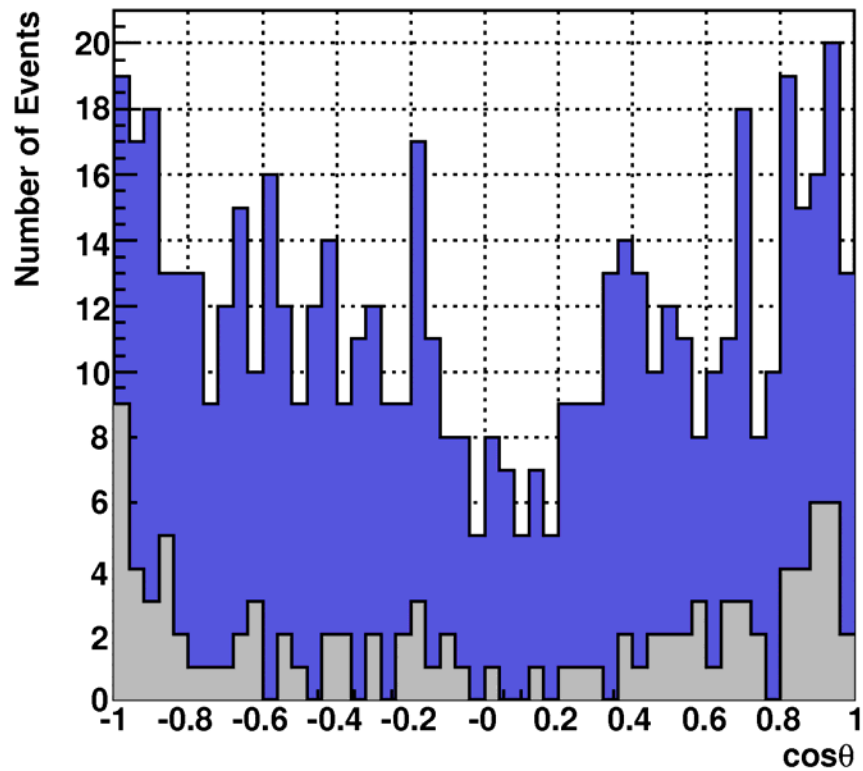
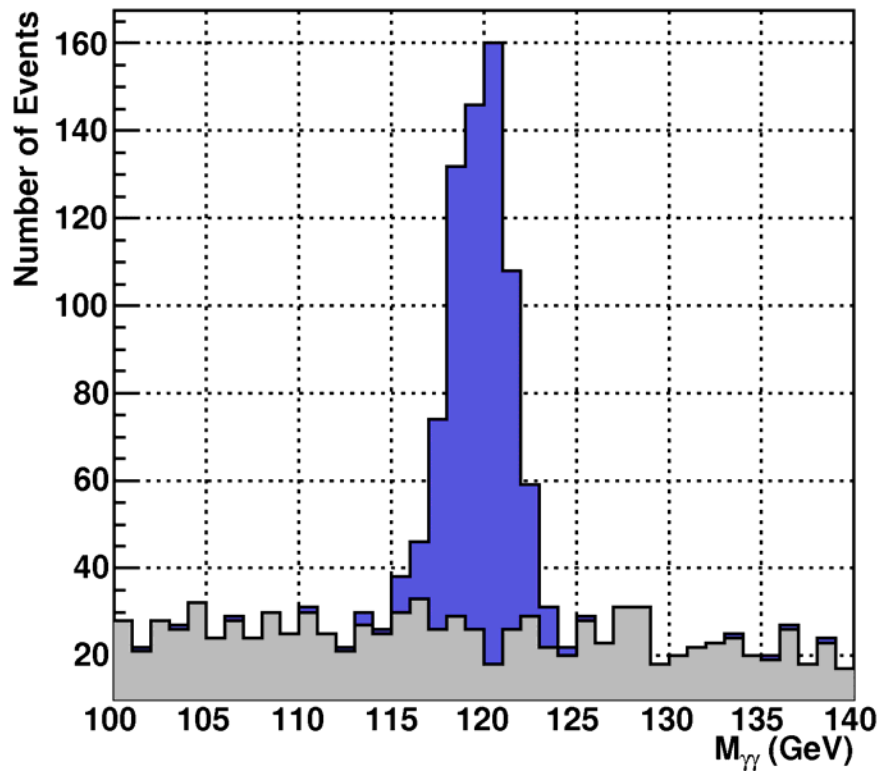
We can distinguish the X from Higgs!

$$e^+e^- \rightarrow \gamma X; X \rightarrow \gamma\gamma$$

$$\sigma_X \sim 52.5 \text{ fb}$$



Cut	Signal	Background
No Cut	600 (1.0000)	100000 (1.0000)
$N_{\text{gammas}} = 3$	575 (0.9583)	3746 (0.0375)
$E_{\text{gamma}} > 1 \text{ GeV}$	575 (1.0000)	3730 (0.9959)
$ \cos(\theta_j) \leq 0.999$	575 (1.0000)	3728 (0.9992)
$ M_{\gamma\gamma} - m_X \leq 25 \text{ GeV}$	573 (0.9965)	1332 (0.3573)
$ \cos(\theta_\chi) $ and $ \cos(\theta_a) \leq 0.99$	572 (0.9983)	1269 (0.9529)
Total Efficiency	0.9533 ± 0.0086	0.0127 ± 0.0001



4. Summary

We have investigated hidden sector particle production.

- Low cutoff scale → Can produce the particle at the ILC
- Gluophobic scenario → Consistent with Tevatron bound
- Monte-Carlo sim. → Can distinguish the X from Higgs

Future work

We will investigate in a non gluophobic scenario

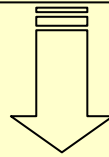


- $X \rightarrow gg$ mode
- $X \rightarrow Z \gamma$ mode

These modes are interesting and important

Buckup

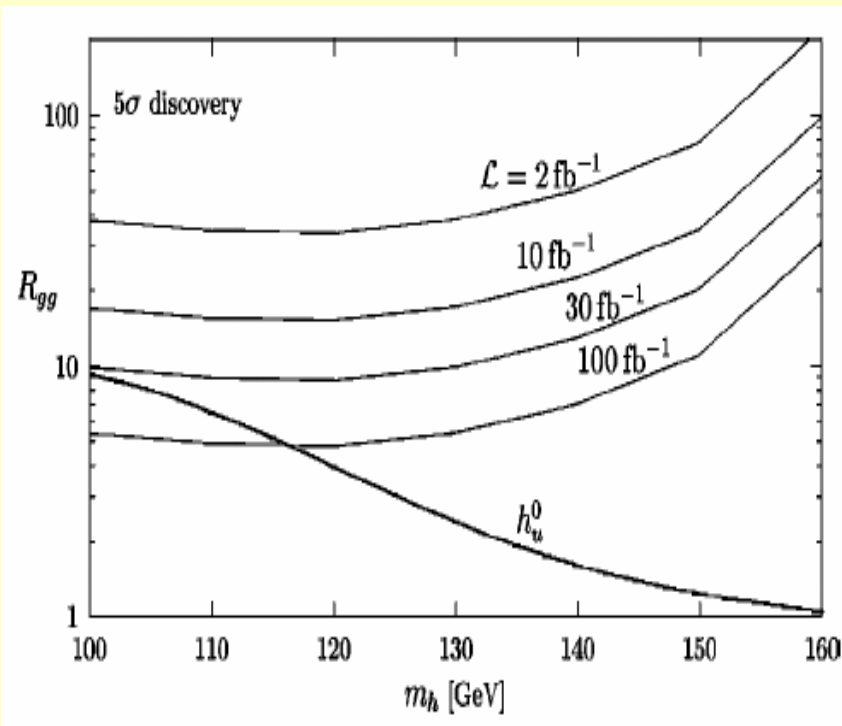
$\Lambda \sim 1 \text{ TeV} \rightarrow X$ strongly couples to gluons and photons



Strong constraint from Tevatron

Higgs search through two photon decay mode

Merrana & Wells, PRD 63 015006 (2000)



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Tevatron Run II

$$\mathcal{L} \sim 1 \text{ fb}^{-1} \quad R_{gg} < 50$$

For SM Higgs:

$$m_h = 120 \text{ GeV}$$

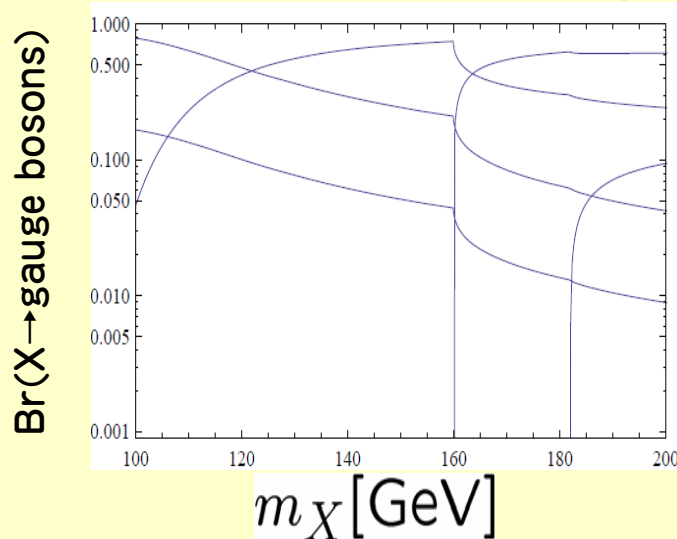
$$\sigma_H \sim 1 \text{ pb}$$

$$Br(h \rightarrow \gamma\gamma) \sim 2 \times 10^{-3}$$

Example of parameter set consistent with Tevatron bound

$$c_{gg} = 0.1, c_{\gamma\gamma} = 0.13, c_{ZZ} = 1, c_{Z\gamma} = 0.68$$

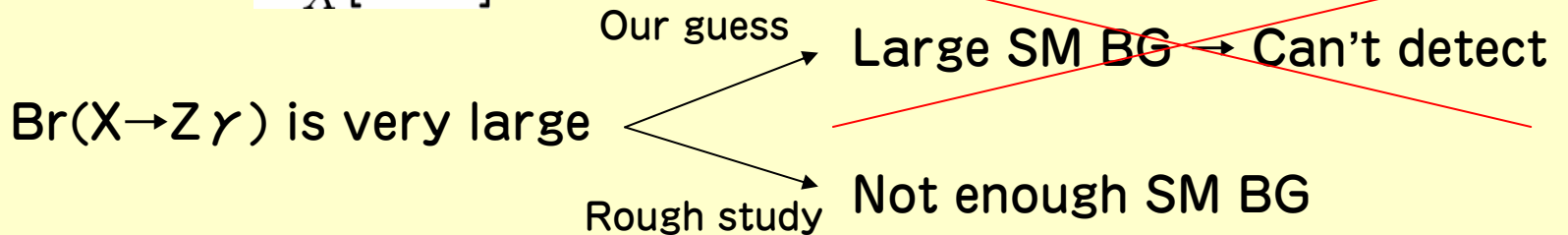
$$\Lambda = 1.56 \text{ TeV} \quad \sigma(e^+e^- \rightarrow ZX) = 21 \text{ fb}$$



For $m_X = 120 \text{ GeV}$

gg	48%
$Z\gamma$	42%
$\gamma\gamma$	10%

Be small !



If $C_{\gamma\gamma}, C_{ZZ} \sim 0.1 \rightarrow$ Disappear $Z\gamma$ mode
 But small X-Z-Z coupling } Incompatible !

We concentrate a **gluophobic scenario** on this time
 (Ex. Realized on brane world scenario)