#### IMPROVED PREDICTION OF EVENT SHAPES FROM EFFECTIVE FIELD THEORY

Thomas Becher **Fermilab** LCWS '08, UIC, Nov. '08

# OVERVIEW

- Event shape variables
  - Definition and experimental prospects at the ILC
- Improved theoretical prediction of thrust
  - NNLO fixed order
  - N<sup>3</sup>LL resummation by RG evolution in Soft Collinear Effective Theory (SCET) TB and M.D. Schwartz, JHEP 0807:034,2008
- Phenomenological results
  - Precision determination of  $\alpha_s$  using LEP data.
  - Bound on light gluinos
  - Comparison with MC event generators

#### **EVENT-SHAPEVARIABLES**



- Parameterize geometric properties of energy and momentum flow in high energy collisions.
  - Inclusive observables: can be calculated in perturbation theory, hadronisation effects are suppressed at high energy.
- Canonical event shape is thrust T

#### THRUST T AND THRUST AXIS $\vec{n}$ .



## MEASUREMENTS OF THRUST



Based on 300'000 events. Similar precision by the other LEP experiments.

# USE OF EVENT SHAPES

#### QCD studies

- convergence of perturbation theory, validation of shower MCs, studies of hadronisation effects
- Measurement of SM parameters
  - strong coupling constant  $\alpha_s$  with  $e^+e^- \rightarrow q\bar{q}$
  - top-mass with  $e^+e^- \rightarrow t\bar{t} \rightarrow \text{Sonny Mantry's talk}$
- Discrimination against background
  - e.g. identification of energetic hadronic top-jets
- Search for new physics
  - e.g. search for light gluinos

## EVENT SHAPES AT THE ILC

- At design luminosity, the ILC produces few hundred thousand  $e^+e^- \rightarrow q\bar{q}$  events/year
  - Statistical uncertainties on extracted value of  $\alpha_s$  is below 0.5%.
  - Systematic uncertainties are expected to be ~ 1%
    Schumm '96 & Truitt '01; Burrows '01
    - "contamination" from  $e^+e^- \rightarrow (t\bar{t}, W^+W^-, ZZ)$
    - luminosity spectrum
  - Note: hadronisation effects scale as  $\sim 1/E_{c.m.}$  and are thus smaller than at LEP
    - can be further constrained by varying  $E_{c.m.}$

#### EVENT SHAPES AT NNLO



• After years of work, the NNLO calculation of  $e^+e^- \rightarrow 3$  jets has been completed.

A. Gehrmann-De Ridder, T. Gehrmann, E.W.N. Glover, G. Heinrich '07

- First time a subtraction scheme has been implemented at NNLO.
  - Real and virtual contributions are have collinear and soft divergences which cancel in the sum.
- Implemented in fixed order event generator. Can be used for NNLO evaluation of event shapes.

## $\alpha_s$ from event shapes at Lep I



 $\alpha_s(M_Z^2) = 0.1240 \pm 0.0008 \,(\text{stat}) \pm 0.0010 \,(\text{exp}) \pm 0.0011 \,(\text{had}) \pm 0.0029 \,(\text{theo})$ 

# RESUMMATION

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All-order formalism for resummation of thrust distribution

N<sup>3</sup>LL resummation

Comparison with fixed order

#### LOGARITHMICALLY ENHANCED CONTRIBUTIONS

• The LO thrust distribution has the form

$$\frac{1}{\sigma_0} \frac{d\sigma}{d\tau} = \frac{2\alpha_s}{3\pi} \left[ -\frac{3}{\tau} + 6 + 9\tau + \frac{(6\tau^2 - 6\tau + 4)}{(1 - \tau)\tau} \ln \frac{1 - 2\tau}{\tau} \right]$$
$$= \frac{2\alpha_s}{3\pi} \left[ \frac{-4\ln\tau - 3}{\tau} + d_{regular}(\tau) \right]$$
singular terms  
• Integral over the end-point is
$$R(\tau) = \int_0^\tau d\tau' \frac{1}{\sigma_0} \frac{d\sigma}{d\tau'} = \frac{2\alpha_s}{3\pi} \left[ -2\ln^2\tau - 3\ln\tau + \dots \right]$$
Sudakov double logarithm

#### SINGULAR TERMS DOMINATE



- Singular terms are predicted (and later resummed to all orders) using Soft-Collinear Effective Theory.
- Regular terms (difference of blue and red) are added back after resummation.

#### **RESUMMATION: THE TRADITIONAL WAY**

- Logarithmically enhanced contributions lead to slow convergence of perturbation theory
- The leading logarithms  $(LL) \alpha_s^n \ln^{2n} \tau$  and next-toleading log's  $(NLL) \alpha_s^n \ln^{2n-1} \tau$  can be resummed using the "coherent branching algorithm"



NLL+NNLO calculation by T. Gehrmann, G. Luisoni and H. Stenzel, arXiv:0803.0695,

#### EFFECTIVE THEORY RESUMMATION

• Using Soft Collinear Effective Theory (SCET), one can show that for  $\tau \rightarrow 0$  the rate factorizes as

$$\frac{1}{\sigma_0}\frac{d\sigma}{d\tau} = H(Q^2,\mu) \int dM_1^2 \int dM_2^2 J(M_1^2,\mu) J(M_2^2,\mu) S_T(\tau Q - \frac{M_1^2 + M_2^2}{Q},\mu)$$

Fleming, Hoang, Mantry and Stewart '07 Schwartz '07 see also: Korchemsky '98; Berger, Kucs, Sterman '03

#### • Three relevant scales:

$$\begin{array}{c|c} Q^2 & \gg & M_1^2 \sim M_2^2 \sim \tau \ Q^2 & \gg & \tau^2 Q^2 \\ \hline hard & jet & soft \end{array}$$

#### RESUMMATION

$$\frac{1}{\sigma_0}\frac{d\sigma}{d\tau} = H(Q^2,\mu) \int dM_1^2 \int dM_2^2 J(M_1^2,\mu) J(M_2^2,\mu) S_T(\tau Q - \frac{M_1^2 + M_2^2}{Q},\mu)$$

- The presence of the three separated scales leads to large perturbative logarithms.
  - Any choice of  $\mu$  will produce large logarithms in either H, J or S.
- *H* and *J* are Wilson coefficients in SCET, *S* a matrix element,
  - fulfill renormalization group equation.

#### **RESUMMATION BY RG EVOLUTION**

• Evaluate each part at its characteristic scale, evolve to common scale:



#### **RESUMMED THRUST DISTRIBUTION**

$$\frac{1}{\sigma_0} \frac{d\sigma}{d\tau} = U(\mu_h, \mu_i, \mu_s) \left(\frac{Q^2}{\mu_h^2}\right)^{-2a_{\Gamma}(\mu_h, \mu_i)} H(Q^2, \mu_h)$$
$$\times \left[\tilde{j} \left(\ln\frac{\mu_s Q}{\mu_i^2} + \partial_{\eta}, \mu_i\right)\right]^2 \tilde{s}_T \left(\partial_{\eta}, \mu_s\right) \frac{1}{\tau} \left(\frac{\tau Q}{\mu_s}\right)^{\eta} \frac{e^{-\gamma_E \eta}}{\Gamma(\eta)}$$

- $\tilde{j}$  and  $\tilde{s}_T$  are Laplace transforms of J and  $S_T$
- U is an evolution factor from solving RG eq's
- For N<sup>3</sup>LL resummation, we need:
  - 4-loop  $\Gamma_{cusp}$  (use Pade approx. for 4-loop term),
  - 3-loop γ's,
  - 2-loop H,  $\tilde{j}$  and  $\tilde{s}$ .

All ingredients known except 2-loop soft function. Obtain it numerically using EVENT2

# NNLO SINGULAR TERMS

- With 2-loop *H*, *J* and *S* and 3-loop anomalous dimension we predict all singular terms at  $\alpha_s^3$ .
- For small τ singular terms dominate full result: check of NNLO calculation of Gehrmann et al.
- In our paper arXiv:0803.0342, we found disagreement at small τ values in 2 color structures.
- In arXiv:0807.3241 Stefan Weinzierl identified a soft divergence in one of the subtraction terms used by Gehrmann et al.
  - Affects thrust at small τ. New numerical results for thrust by Weinzierl and by Gehrman et al. should soon be available.

## NNLO SINGULAR TERMS

 Nice agreement with preliminary corrected results obtained from T. Gehrmann (thanks!)



- Note: correction only affects region of very small  $\tau$ 
  - Should have negligible impact on  $\alpha_s$  extraction.

#### INDIVIDUAL COLOR STRUCTURES: SMALL $\tau$



## LEADING COLOR STRUCTURE

 Nice agreement with preliminary corrected results obtained from T. Gehrmann



# MATCHING

- Will now combine resummation and fixed order result to obtain  $\alpha_s$  from a fit to LEP data.
- Different possibilities, we use

order	fixed-order	logarithmic
	matching	accuracy
$1^{st}order$		NLL
2 <sup>nd</sup> order	LO	NNLL
3 <sup>rd</sup> order	NLO	$N^{3}LL$
4 <sup>th</sup> order	NNLO	$N^{3}LL$

• note: previous speaker G. Luisoni used NLL+NNLO

#### **RESUMMED VS. FIXED ORDER**



• For PDG value  $\alpha_s(M_Z)=0.1176$ .

#### **RESUMMEDVS. FIXED ORDER**



• For PDG value  $\alpha_s(M_Z)=0.1176$ 

• This is the region relevant for  $\alpha_s$  determination

#### PHENOMENOLOGICAL APPLICATIONS

• Determination of  $\alpha_s$ 

- Scale variation, error band method
  - Fit to ALEPH and OPAL LEP data
- Bound on light gluinos Kaplan and Schwartz '08
- Comparison with event generator results at ILC energies

# THEORETICAL UNCERTAINTY

- We will assess the perturbative uncertainty in the standard way, by varying the renormalization (resp. matching) scales.
  - To the order of the calculation, the cross section is independent of these scales;
  - variation then is a measure of unknown higher order terms.
- We have four scales
  - $\mu_{hard}^2 \sim Q^2$  : scale at which *H* is evaluated
  - $\mu_{jet}^2 \sim \tau Q^2$  : scale at which *J* is evaluated
  - $\mu_{\text{soft}}^2 \sim \tau^2 Q^2$ : scale at which  $S_T$  is evaluated
  - $\mu_{match}^2$  : scale of the regular terms

# INDEPENDENT SCALE VARIATION



• Varying jet and soft scale independently by a factor 2 makes no sense at moderate  $\tau$  (leads to  $\mu_{soft} > \mu_{jet}$ , etc.), overestimates the uncertainty.

# JET AND SOFT SCALE VARIATION



Instead of independently varying the jet and soft scales, we vary as follows

• correlated:  $\mu_{jet} \rightarrow \alpha \mu_{jet}$ ,  $\mu_{soft} \rightarrow \alpha \mu_{soft}$  with  $1/2 < \alpha < 2$ 

• squeeze:  $\mu_{jet} \rightarrow \sqrt{\alpha} \, \mu_{jet}$ ,  $\mu_{soft} \rightarrow \alpha \, \mu_{soft}$  with  $1/\sqrt{2} < \alpha < \sqrt{2}$ 

#### ERROR BAND METHOD

#### Jones, Ford, Salam Stenzel & Wicke '03; adopted by ALEPH and OPAL



- Perform  $\chi^2$ -fit to the data, extract best-fit value of  $\alpha_s$ . Calculate maximum deviation from default distribution: "error band".
- To get theoretical uncertainty, calculate max. and min.  $\alpha_s$  for which theoretical distribution lies inside the error band.

#### EXPERIMENTAL UNCERTAINTY



- OPAL '05 and ALEPH '03 give results for binned thrust distributions. Do not provide correlations.
- Put only stat. err. in our  $\chi^2$ -fit. For each Q, use same fit ranges as exp. paper and use their systematic uncertainties.



 $\alpha_s(m_Z) = 0.1172 \pm 0.0010 \text{(stat)} \pm 0.0008 \text{(sys)} \pm 0.0012 \text{(had)} \pm 0.0012 \text{(pert)}$ = 0.1172 ± 0.0022.

 $[PDG: \alpha_s(m_Z) = 0.1176 \pm 0.0020]$ 

# BOUND ON LIGHT GLUINOS

#### Kaplan and Schwartz '08



- Gluinos would affect *H*, *J* and *S* functions at the twoloop level. Leading effect is  $\Delta n_f = 3$ 
  - in hard function H if  $m_{\tilde{g}} < Q$ ,
  - in jet function J if  $m_{\tilde{g}} < \sqrt{\tau}Q$ ,
  - in soft function S if  $m_{\tilde{g}} < \tau Q$ ,



#### **COMPARISON WITH PYTHIA**



• hadronic Pythia agrees perfectly with the ALEPH data

• partonic Pythia does much better than NLL

## I TEV LEPTON COLLIDER



- Partonic Pythia now looks much more NLL like.
- Will need to retune (or redesign) the shower.
  - Can tune partonic shower to our theoretical prediction.

## I TEV LEPTON COLLIDER



• Can tune partonic shower to our theoretical prediction.

# SUMMARY

- Have used effective field theory methods to resum thrust distribution to N<sup>3</sup>LL.
  - Traditional method works only up to NLL.
  - Logarithmically enhanced contributions dominate. Have evaluated all singular terms at  $\alpha_s^3$ .
    - Check of NNLO calculation of  $e^+e^- \rightarrow 3$  jets.
  - Also other event shapes can be improved beyond NLL
- Extract  $\alpha_s$  from a fit to LEP data:

 $\alpha_s(m_Z) = 0.1172 \pm 0.0010(\text{stat}) \pm 0.0008(\text{sys}) \pm 0.0012(\text{had}) \pm 0.0012(\text{pert})$ 

- Most precise determination of  $\alpha_s$  at high energies, agrees well with low energy determinations.
- Theoretical accuracy matches exp. precision at the ILC

# **EXTRA SLIDES**

## POWER CORRECTIONS

- So far, we have not included 1/Q power corrections:
  - finite b-quark mass effects  $\approx +1.5\%$  at LEP I
    - calculated perturbatively, e.g. using NLO event generator by Nason and Oleari.
    - could perform resummation for this part, using SCET, see Sonny Mantry's talk
  - hadronisation ~ -1.5% at LEP I
    - estimated using Pythia to calculate transfer matrix
    - uncertainty is estimated by comparing Pythia to Herwig and Ariadne: 2.5% at LEP I. Now the dominant uncertainty!
    - Our precise perturbative prediction can and should be used to study hadronisation effects in more detail, using also lower energy data.

#### PDG AVERAGE

