Higgs Cubic Coupling and Electroweak Phase Transition

Maxim Perelstein, Cornell LCWS-08 Workshop, Chicago, November 17 2008

Andrew Noble, MP, arXiv0711:3018 [hep-ph], in print in PRD

ElectroWeak Phase Transition

• In our world, EW gauge symmetry is broken: $SU(2) imes U(1)_Y o U(1)_{
m em}$

- At high temperature, symmetry is restored (in most models)
- Early universe: electroweak phase transition at $T \sim 100 \text{ GeV}$
- How much can we learn about the dynamics of this transition? First-order ("boiling") or second-order ("quasiadiabatic") transition?
- Has implications for electroweak baryogenesis (1st order required to satisfy Sakharov's out-of-eq. condition!)

EWPT and Collider Data

- Direct relics from the transition in the early universe unlikely to survive (possibly gravitational waves?)
 [see e.g. Grojean and Servant, hep-ph/0607107]
- However, finite-T physics is described by the same Lagrangian as the T=0 physics we will study at colliders
- Only weak-scale states are relevant for the EW phase transition ("decoupling")
- Determine the TeV Lagrangian at the LHC, ILC
 learn the order of the transition, critical temperature, etc.
- Note: standard FRW cosmology at the phase transition time is assumed
- What measurements will be necessary to address this?
- We'll try to approach this question in a fairly modelindependent way

Finite-T Effective Potential

- Assume weakly coupled physics at the TeV scale (otherwise this analysis would require lattice simulations!)
- Assume single physical higgs *h* participates in the transition (easy to generalize away from this assumption)
- One-loop effective potential has the form

 $V(h;T) = V_{t}(h) + V_{T=0}(h) + V_{T}(h)$

• T=0 (Coleman-Weinberg) part:

$$V_{T=0}(h) = \sum_{i} \frac{g_i(-1)^{F_i}}{64\pi^2} M_i^4 \left(\log\frac{M_i^2}{\mu^2} + C_i\right)$$

• Finite-T part:

$$V_T(h) = \sum_i \frac{g_i(-1)^{F_i} T}{2\pi^2} \int dk k^2 \log[1 - (-1)^{F_i} \exp(-\beta \sqrt{k^2 + M_i^2})]$$

• In a renormalizable theory $M_i^2 = M_{i,0}^2 + a_i h^2$

Effective Potential: Ring Terms

• Beyond one-loop, include "ring" contributions



• Can be summed up to yield: [Carrington, PRD 1992]

$$V_r(h,T) = \sum_b \frac{T}{12\pi} \left[M_b^3 - (M_b^2 + \Pi_b(0))^{3/2} \right]$$

- Only bosons contribute (due to IR divergence)
- Important for the first-order EWPT since at high T, $V_r \propto T |h|^3$ which can produce the desired "dip"
- Ring terms typically controlled by the same parameters as the one-loop effective potential

First-Order Phase Transition



 $T > T_c$



$$T = T_c$$



critical temperature: $V_{\text{eff}}(0) = V_{\text{eff}}(v(T_c))$ nucleation temperature:





strong first-order transition:

$$\xi = \frac{v(T_c)}{T_c} \ge 1$$

Effective Potential from Colliders?

- So, to reconstruct finite-T potential, we need to know the following:
 - Higgs zero-temperature tree-level potential: vev, mass
 - Full spectrum of states (SM and BSM) with significant couplings to the Higgs and masses up to ~few 100 GeV
 - Their fermion numbers and state multiplicities
 - Their masses and couplings to the Higgs:

 $M_i^2 = M_{i,0}^2 + a_i h^2$

- This is definitely difficult, and may be impossible: e.g. $V = V_{\rm SM}(H) + \frac{1}{2}M_0^2S^2 + \zeta |H|^2S^2 \text{ with } m_S > \frac{m_h}{2}$
- [still, may be possible in specific models future work!]

EWPT and Higgs Cubic Coupling

- Idea: look for simple observables that are correlated with the order of the EWPT in a reasonably model-independent framework
- Proposal: use Higgs boson cubic self-coupling λ_3
- Heuristic explanation:
 - In the SM transition is 2nd order for $m_h > 114 \text{ GeV}$
 - New physics must change the shape of V(h) at T_c
 - This changes the shape of V(h) at T=0 \Rightarrow different $\lambda_3(v, m_h)$
- Models with 1st order phase transition exhibit large (typically 20-100%) deviations of λ_3 from its SM value
- Evidence: analysis of a series of toy models designed to mimic the known mechanisms for getting a first-order PT

Toy Model I:"Quantum" EWPT

- Single Higgs doublet, SM couplings to SM states, add a real scalar field *S*
- Scalar potential: $V = V_{SM}(H) + \frac{1}{2}M_0^2S^2 + \zeta |H|^2S^2$
- Assume positive $M_0^2, \zeta \implies \langle S \rangle = 0$
- Compute effective Higgs potential
- At high T, $V_{\text{eff}}(h;T) = (\mu^2 + DT^2)h^2 + ET|h|^3 + \lambda h^4 + \dots$
- Look for minima: $\partial V_{\rm eff}/\partial h = 0$
- If h = 0 and $h \neq 0$ minima coexist, 1 st order transition
- Scan $m_h, M_0, \zeta \implies$ find points with first-order EWPT
- Physical Higgs boson cubic self-coupling:

$$\lambda_3 = \frac{d^3 V_{\text{eff}}(v; T=0)}{dh^3}$$

Quantum EWPT: Results



Exp. prospects: 23% for a 120-GeV Higgs at a 500-GeV ILC, 1 ab-1 [Castanier, Gay, Lutz, Orloff, hep-ex/0101028] 20-30% for 160-180 GeV Higgs at SLHC 8-25% for 150-200 GeV Higgs at 200 TeV VLHC [Baur, Rainwater, Plehn, hep-ph/0206024, 0211224]

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Quantum EWPT: Extensions

- Same calculations apply in a model with identical N real (or N/2 complex) scalars simple scaling argument: $\xi \rightarrow \xi N^{1/4}$
- One-loop analysis is independent of the scalar's gauge charges - could be e.g. stops of the MSSM (in the decoupling limit - one Higgs), weak triplets, etc.
- Same picture in a model with 2 independent (non-identical) scalars (N ind. scalars is a reasonable conjecture)
- If scalar replaced with a fermion, no points with first-order EWPT found, due to the different structure of the fermion contribution to $V_{\rm eff}$ (no ring terms \rightarrow no $|h|^3$)
- A more interesting case: add a scalar-fermion pair ("supermultiplet") with same coupling to the Higgs, different masses

Quantum EWPT with BF Pair



Accidental Cancellation between B and F contributions at T=0 can result in near-SM value of λ_3 - counterexample to our claim! [But SUSY is broken by strong coupling to the SM via h!] $\Delta M \ge 100 \text{ GeV} \Rightarrow \frac{\delta \lambda_3}{\lambda_3} \ge 7\%$

TM 2: "Non-renormalizable" EWPT

• An alternative way to get 1st-order EWPT: add a nonrenormalizable operator to the SM Higgs potential

 $V = \mu^{2} |H|^{2} + \lambda |H|^{4} + \frac{1}{\Lambda^{2}} |H|^{6}$

[Grojean et al, 2004]

- Reasonable EFT if $v \ll \Lambda \implies |\lambda| \ll 1$
- First-order transition can occur for $\mu^2 > 0, \lambda < 0$



TM 3: Higgs-Singlet Mixing

• As in TMI, add I real scalar, but with a more general potential:

 $V(H,S) = \mu^2 |H|^2 + \lambda |H|^4 + \frac{a_1}{2} |H|^2 S + \frac{a_2}{2} |H|^2 S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$

- Generically, both H and S get vevs at zero temperature
- EWPT involves both H and S changing, order parameter is a linear combination of H and S
- Effective potential for order parameter contains tree-level cubic terms from possible strongly first-order EWPT
- Zero-T spectrum: two "higgses" (mixed H and S)
- Only H enters Yukawa couplings pon-SM Yukawas!
- Cubic self-coupling of the "H-like" higgs = λ_3

Higgs-Singlet Mixing: Results



Partial scan of the 6-dim parameter space Both suppression and enhancement of λ_3 is possible

Small correction to λ_3 seems only possible if there's an accidental cancellation of two large contributions



- Higgs boson cubic self-coupling is correlated with the order of EWPT
- Stronger I-st order phase transition \triangleleft larger deviation in λ_3
- Typical deviations large enough to be seen at the ILC or the SLHC/VLHC
- Correlation seen in 3 toy models, illustrating different mechanisms for getting a first-order EWPT
- All examples (known to us) violating this conclusion involve accidental cancellations of two large corrections to λ_3
- Observing SM value would strongly disfavor first-order phase transition (and hence EW baryogenesis)
- Caution: Models with 2nd order EWPT can still produce large deviations, though!