

Dark Matter in the $U(1)$ -extended MSSM

Jonathan Roberts

IFT, University of Warsaw

ECFA '08, Warsaw, June 10, 2008

Based on work with Jan Kalinowski and Steve King

- 1 What's wrong with the MSSM
 - μ - the black sheep of the MSSM
- 2 Solving the μ problem
- 3 Phenomenology of the USSM
- 4 Conclusions

What's wrong with the MSSM?

The MSSM is very successful.

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass
- Gives radiative electroweak symmetry breaking

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass
- Gives radiative electroweak symmetry breaking
- Allows for unification of the gauge couplings

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass
- Gives radiative electroweak symmetry breaking
- Allows for unification of the gauge couplings
- Provides a dark matter candidate

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass
- Gives radiative electroweak symmetry breaking
- Allows for unification of the gauge couplings
- Provides a dark matter candidate
- Can account for $(g - 2)_\mu \dots$

What's wrong with the MSSM?

The MSSM is very successful.

- Stabilises the Higgs mass
- Gives radiative electroweak symmetry breaking
- Allows for unification of the gauge couplings
- Provides a dark matter candidate
- Can account for $(g - 2)_\mu \dots$

But all is not well

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$

- $\mu = m_{PI}$

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$
 - Charginos would be massless

- $\mu = m_{PI}$

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$
 - Charginos would be massless
 - Requires $\langle H_d \rangle = 0$, so we would have massless d, s, b and massless charged leptons.
- $\mu = m_{Pl}$

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$
 - Charginos would be massless
 - Requires $\langle H_d \rangle = 0$, so we would have massless d, s, b and massless charged leptons.
- $\mu = m_{Pl}$
 - Corrections to the Higgs mass are of order μ^2

The μ Problem

The MSSM has a Higgs bilinear term in the superpotential:

$$\mu \hat{H}_u \hat{H}_d$$

The μ parameter is **SUSY preserving**

The μ term only has two natural values:

- $\mu = 0$
 - Charginos would be massless
 - Requires $\langle H_d \rangle = 0$, so we would have massless d, s, b and massless charged leptons.
- $\mu = m_{Pl}$
 - Corrections to the Higgs mass are of order μ^2

We need a rationale for μ to be of the SUSY breaking scale.

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

This is stretching the case a little in the MSSM.

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus [explain the origin of the electroweak scale](#).

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus [explain the origin of the electroweak scale](#).

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

- The programme [sets](#) μ to the value where EWSB is achieved

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

- The programme **sets** μ to the value where EWSB is achieved
- Often μ needs to be **precisely tuned** to generate the EWSB scale

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

- The programme **sets** μ to the value where EWSB is achieved
- Often μ needs to be **precisely tuned** to generate the EWSB scale
- This is because it has to cancel large contributions from SUSY breaking masses in the Higgs running

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

- The programme **sets** μ to the value where EWSB is achieved
- Often μ needs to be **precisely tuned** to generate the EWSB scale
- This is because it has to cancel large contributions from SUSY breaking masses in the Higgs running
- To remove this fine-tuning we would need to link μ to the SUSY breaking scale m_{SUSY}

The μ problem take 2

One of the triumphs of the MSSM is that we can create electroweak symmetry breaking radiatively (REWSB) and thus **explain the origin of the electroweak scale**.

This is stretching the case a little in the MSSM.

In programmes such as softsusy we **require** REWSB.

- The programme **sets** μ to the value where EWSB is achieved
- Often μ needs to be **precisely tuned** to generate the EWSB scale
- This is because it has to cancel large contributions from SUSY breaking masses in the Higgs running
- To remove this fine-tuning we would need to link μ to the SUSY breaking scale m_{SUSY} , which can't be done in the MSSM.

Solving the μ problem

How do we solve the μ problem?

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$,

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- 1 Giudice-Masiero mechanism

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- 1 Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- ① Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory
- ② Generate an effective μ term **dynamically**

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- 1 Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory
- 2 Generate an effective μ term **dynamically**

Following 2: add a new superfield S with a superpotential term

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- ① Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory
- ② Generate an effective μ term **dynamically**

Following 2: add a new superfield S with a superpotential term

$$\lambda \hat{S} \hat{H}_u \hat{H}_d$$

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- ① Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory
- ② Generate an effective μ term **dynamically**

Following 2: add a new superfield S with a superpotential term

$$\lambda \hat{S} \hat{H}_u \hat{H}_d$$

Make sure the new field S gets a VeV: $\langle S \rangle = v_S / \sqrt{2}$ and thus an effective μ term: $\mu_{\text{eff}} = \lambda v_S / \sqrt{2}$

Solving the μ problem

How do we solve the μ problem?

First set $\mu = 0$, then take one of two options:

- ① Giudice-Masiero mechanism
 - Generate a μ term from specific hidden sector theory
- ② Generate an effective μ term **dynamically**

Following 2: add a new superfield S with a superpotential term

$$\lambda \hat{S} \hat{H}_u \hat{H}_d$$

Make sure the new field S gets a VeV: $\langle S \rangle = v_S / \sqrt{2}$ and thus an effective μ term: $\mu_{eff} = \lambda v_S / \sqrt{2}$

$\langle S \rangle$ is set by the running of soft SUSY breaking parameters

Do we know what we've started?

Can we just add fields without any consequences?

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

- NMSSM: **N**ext to **M**inimal **SSM**:
 - Add a cubic term $\kappa \hat{S}^3$ to break the symmetry
 - Left with a Z_3 symmetry which can create dangerous domain walls

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

- **NMSSM: Next to Minimal SSM:**

- Add a cubic term $\kappa \hat{S}^3$ to break the symmetry
- Left with a Z_3 symmetry which can create dangerous domain walls
- [Djouadi et al, 08; Belanger, Hugonie, Pukhov 07; Cerdeno et al '07; Belanger, Boudjema, Hugonie, Pukhov, Semenov, 05]

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

- **NMSSM: Next to Minimal SSM:**

- Add a cubic term $\kappa \hat{S}^3$ to break the symmetry
- Left with a Z_3 symmetry which can create dangerous domain walls
- [Djouadi et al, 08; Belanger, Hugonie, Pukhov 07; Cerdeno et al '07; Belanger, Boudjema, Hugonie, Pukhov, Semenov, 05]

- **USSM: $U(1)'$ -extended MSSM:**

- Gauge the $U(1)'$ to avoid axions and domain walls
- Includes a Z' with a mass related to v_S , and therefore to μ .

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

- NMSSM: **N**ext to **M**inimal **SSM**:
 - Add a cubic term $\kappa \hat{S}^3$ to break the symmetry
 - Left with a Z_3 symmetry which can create dangerous domain walls
 - [Djouadi et al, 08; Belanger, Hugonie, Pukhov 07; Cerdeno et al '07; Belanger, Boudjema, Hugonie, Pukhov, Semenov, 05]
- USSM: **U**(1)'-extended MSSM:
 - Gauge the $U(1)'$ to avoid axions and domain walls
 - Includes a Z' with a mass related to v_S , and therefore to μ .
 - [Barger, Langacker, Shaughnessy 07; de Carlos, Espinosa 97]

Do we know what we've started?

Can we just add fields without any consequences?

No - the extra field gives an extra $U(1)$ which would give a massless axion when S gets a VeV.

- NMSSM: **N**ext to **M**inimal **SSM**:
 - Add a cubic term $\kappa \hat{S}^3$ to break the symmetry
 - Left with a Z_3 symmetry which can create dangerous domain walls
 - [Djouadi et al, 08; Belanger, Hugonie, Pukhov 07; Cerdeno et al '07; Belanger, Boudjema, Hugonie, Pukhov, Semenov, 05]
- USSM: **U**(1)'-extended MSSM:
 - Gauge the $U(1)'$ to avoid axions and domain walls
 - Includes a Z' with a mass related to v_S , and therefore to μ .
 - [Barger, Langacker, Shaughnessy 07; de Carlos, Espinosa 97]

Here we gauge the $U(1)'$

The USSM

The USSM

The USSM has the features:

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S (H_u H_d) + h.c.)$$

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S (H_u H_d) + h.c.)$$

- 3 New particle content:

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S(H_u H_d) + h.c.)$$

- 3 New particle content:
 - Z' (that mixes a *little* with the Z)

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S (H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S (H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)
- Two new neutralinos - a **bino'** and a **singlino**

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S (H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)
- Two new neutralinos - a **bino'** and a **singlino**

What doesn't work?

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S(H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)
- Two new neutralinos - a **bino'** and a **singlino**

What doesn't work?

- Need exotic fermions to cancel anomalies

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S(H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)
- Two new neutralinos - a **bino'** and a **singlino**

What doesn't work?

- Need exotic fermions to cancel anomalies
 - ...but these appear in the full-blown E_6 SSM

The USSM

The USSM has the features:

- 1 Extra superpotential term:

$$W_{USSM} = W_{MSSM}(\mu = 0) + \lambda \hat{S} \hat{H}_d \hat{H}_u$$

- 2 Extra soft terms:

$$\mathcal{L}_{soft} \in m_S^2 |S|^2 + (\lambda A_\lambda S(H_u H_d) + h.c.)$$

- 3 New particle content:

- Z' (that mixes a *little* with the Z)
- Extra Higgs boson (usually h_3 , dominantly singlet with mass $m_S \approx m_{Z'}$)
- Two new neutralinos - a **bino'** and a **singlino**

What doesn't work?

- Need exotic fermions to cancel anomalies
 - ...but these appear in the full-blown E_6 SSM
- Small amount of tuning remaining because $v_S \approx 10 \times v$

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

The important features are:

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

The important features are:

- New sector is almost decoupled from MSSM neutralinos

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

The important features are:

- New sector is almost decoupled from MSSM neutralinos
- No singlino mass term

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

The important features are:

- New sector is almost decoupled from MSSM neutralinos
- No singlino mass term
 - See-saw structure for singlino/bino' mass

Dark matter: New Neutralinos

The two new neutralinos mean we have a 6×6 neutralino mass matrix:

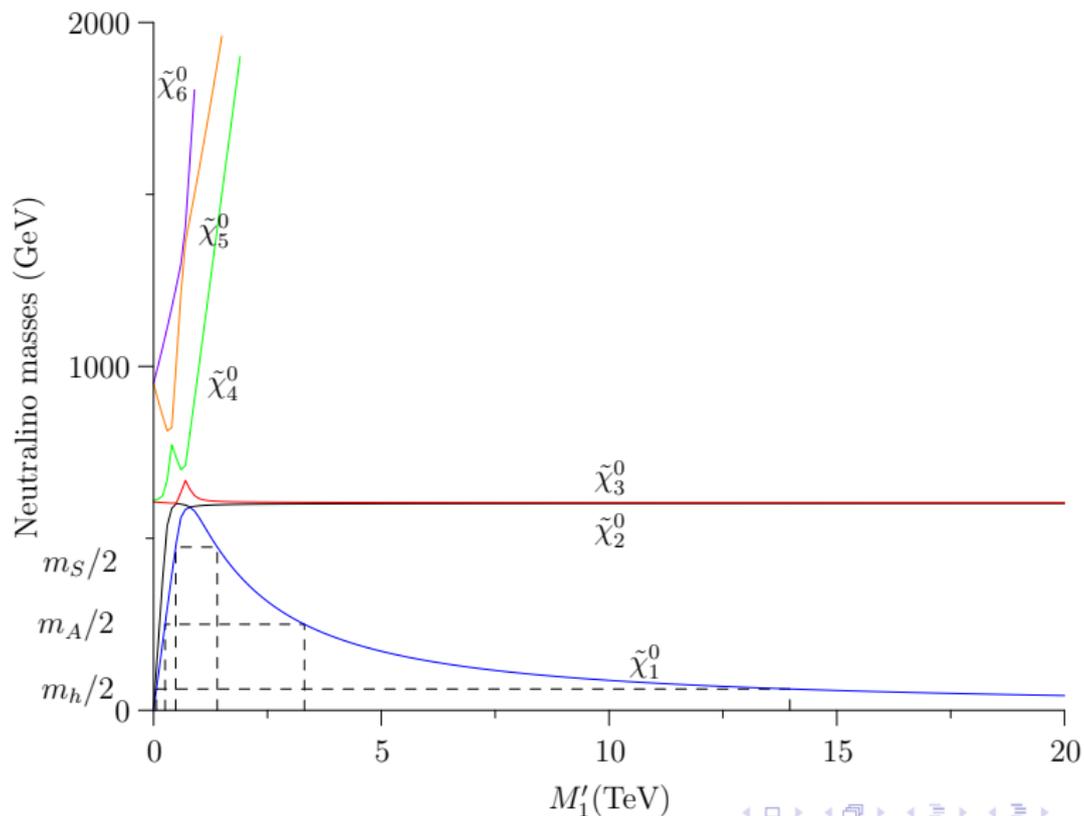
$$M_{\tilde{\chi}^0} = \left(\begin{array}{cccc|cc} M_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W & 0 & M_K \\ 0 & M_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W & 0 & 0 \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\mu & -\mu_\lambda s_\beta & Q'_1 g'_1 v c_\beta \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\mu & 0 & -\mu_\lambda c_\beta & Q'_2 g'_1 v s_\beta \\ \hline 0 & 0 & -\mu_\lambda s_\beta & -\mu_\lambda c_\beta & 0 & Q'_5 g'_1 v s \\ M_K & 0 & Q'_1 g'_1 v c_\beta & Q'_2 g'_1 v s_\beta & Q'_5 g'_1 v s & M'_1 \end{array} \right)$$

where $\mu = \lambda v_s / \sqrt{2}$, $\mu_\lambda = \lambda v / \sqrt{2}$

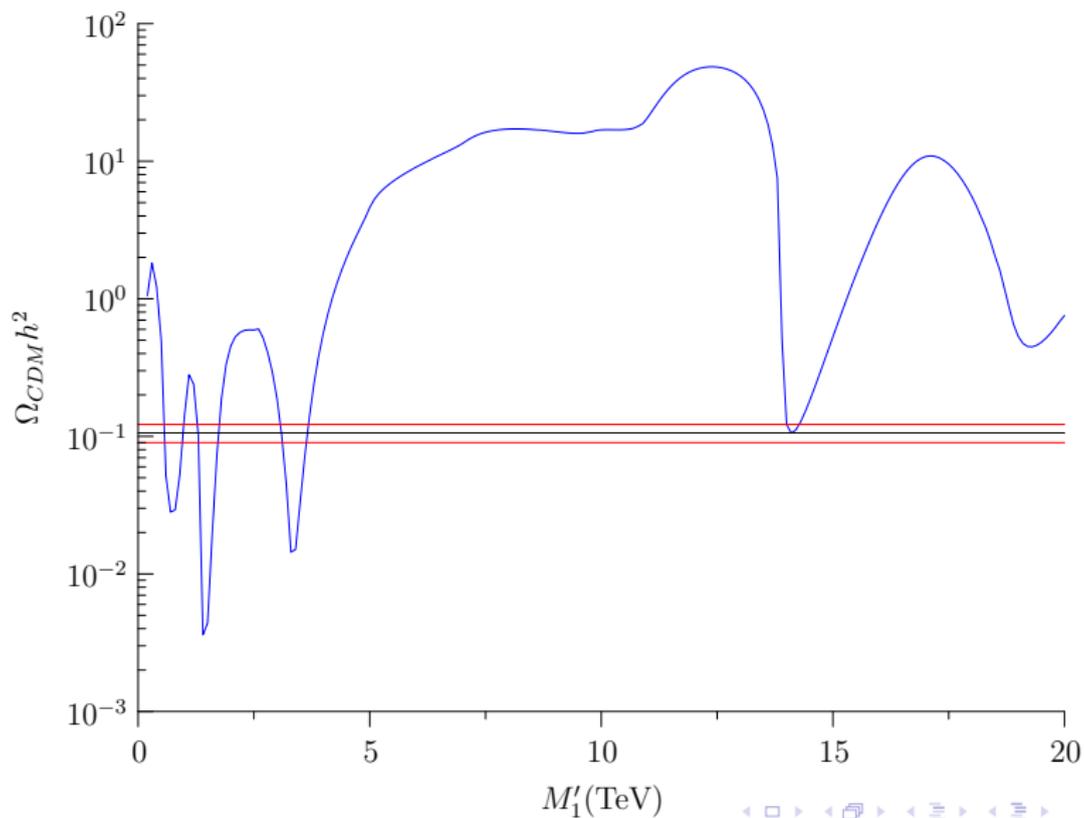
The important features are:

- New sector is almost decoupled from MSSM neutralinos
- No singlino mass term
 - See-saw structure for singlino/bino' mass
 - Never have a dominantly bino' LSP

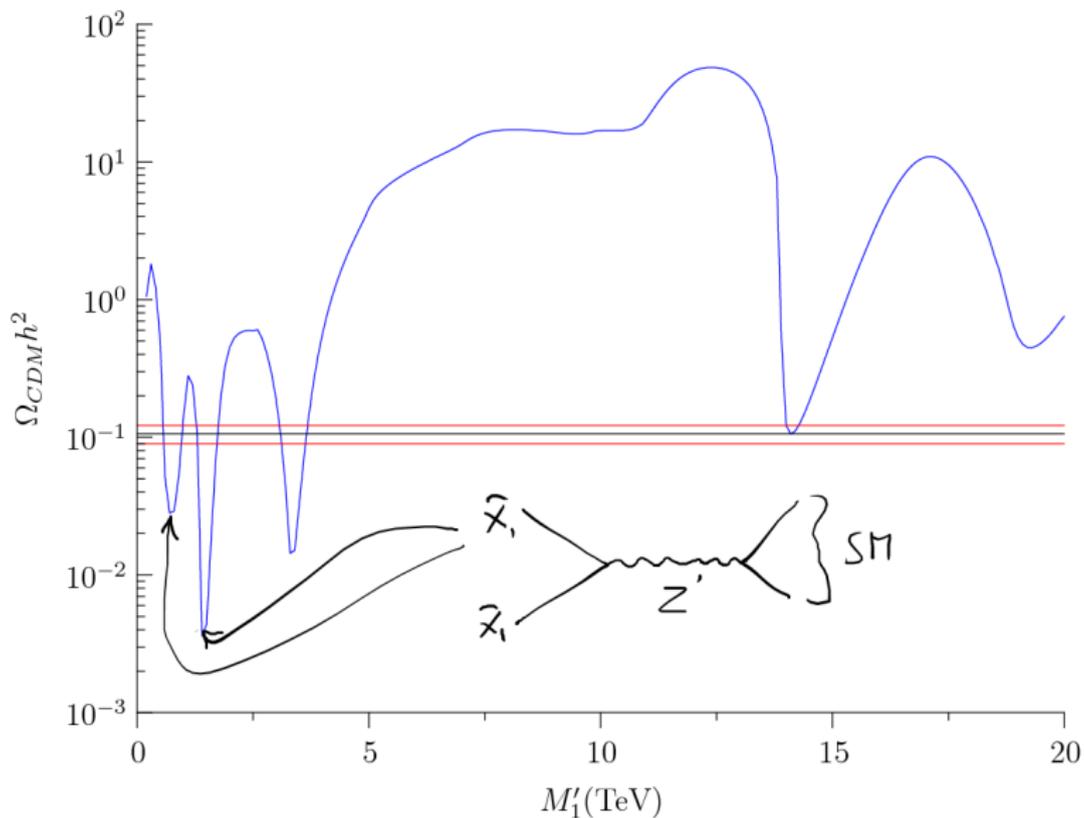
$M_1 = M'_1$: Neutralino Masses



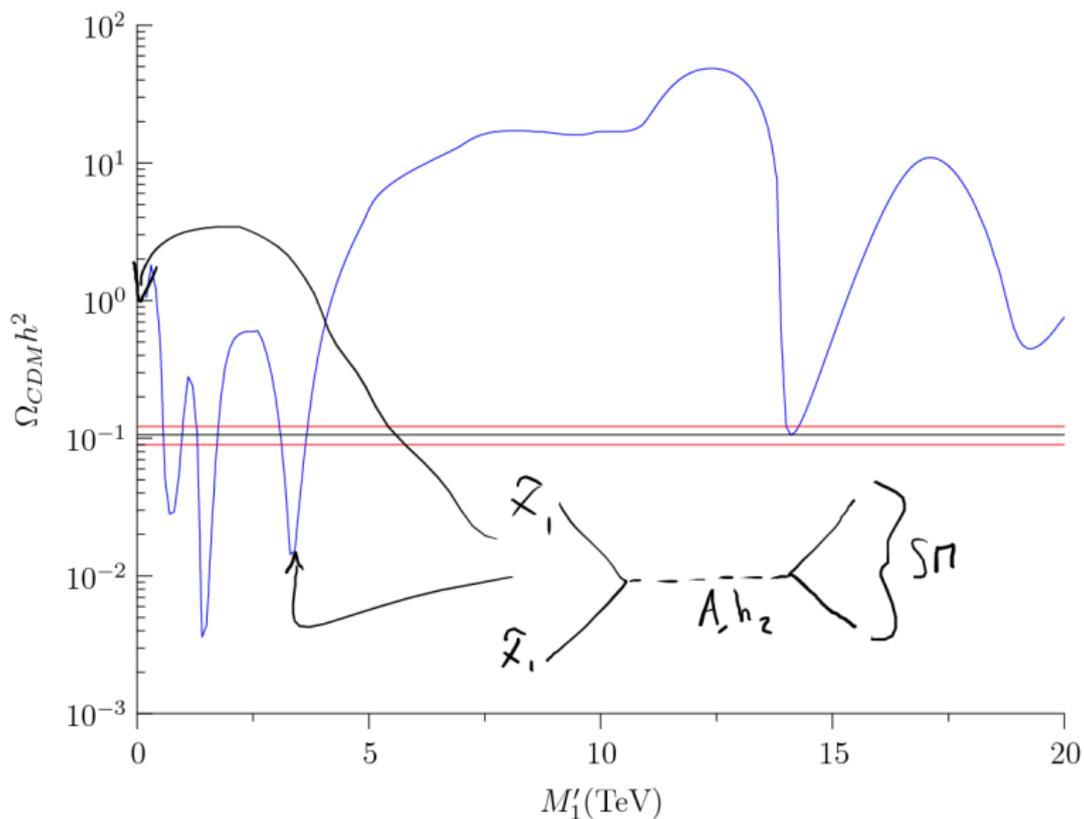
$M_1 = M'_1$: Relic Density



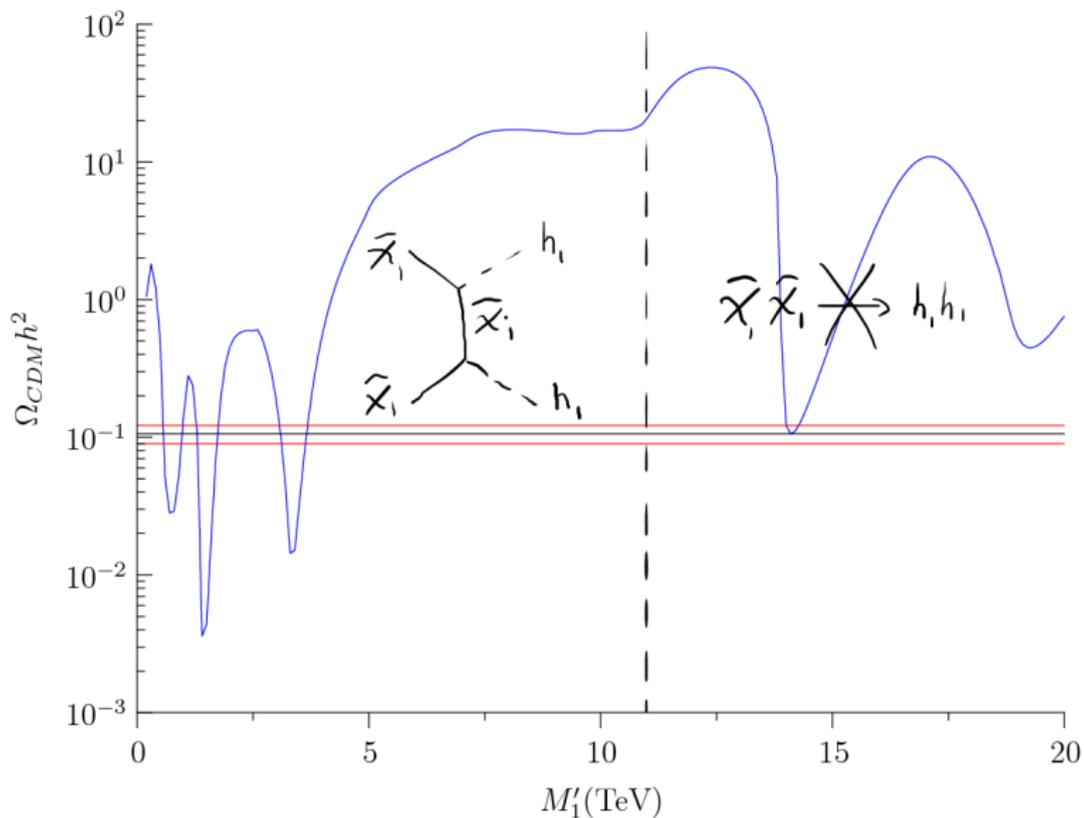
$M_1 = M'_1$: Relic Density



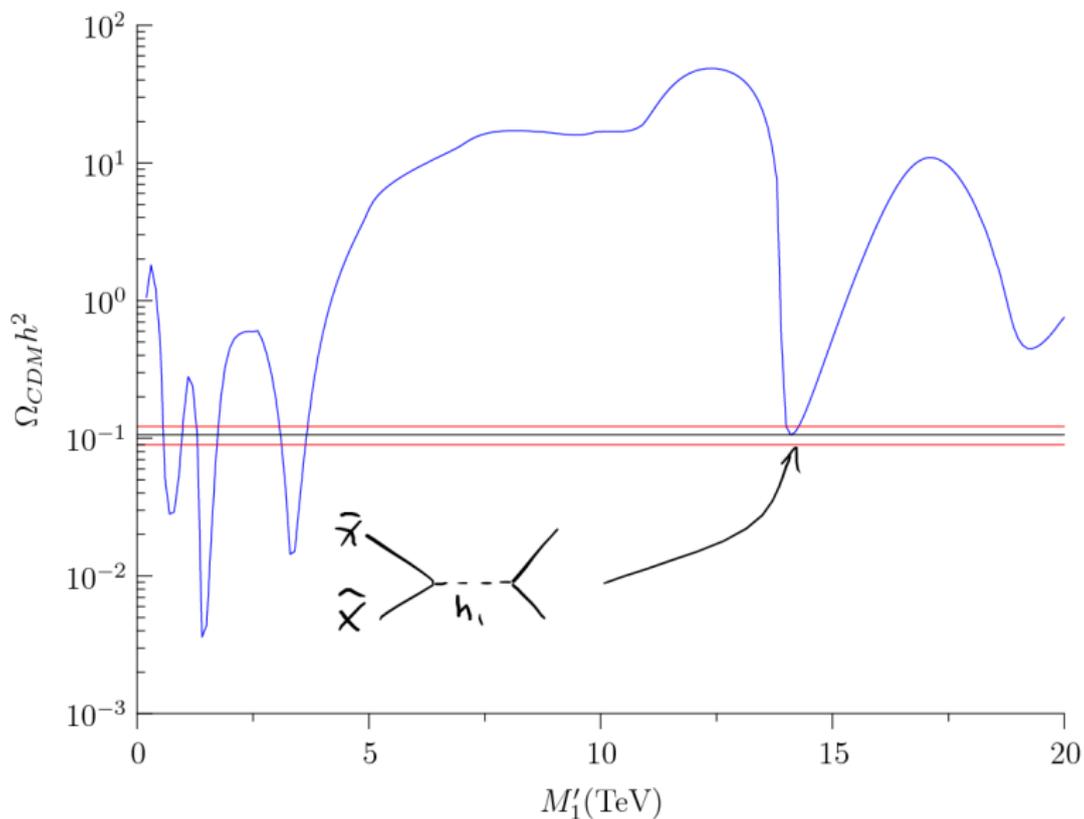
$M_1 = M'_1$: Relic Density



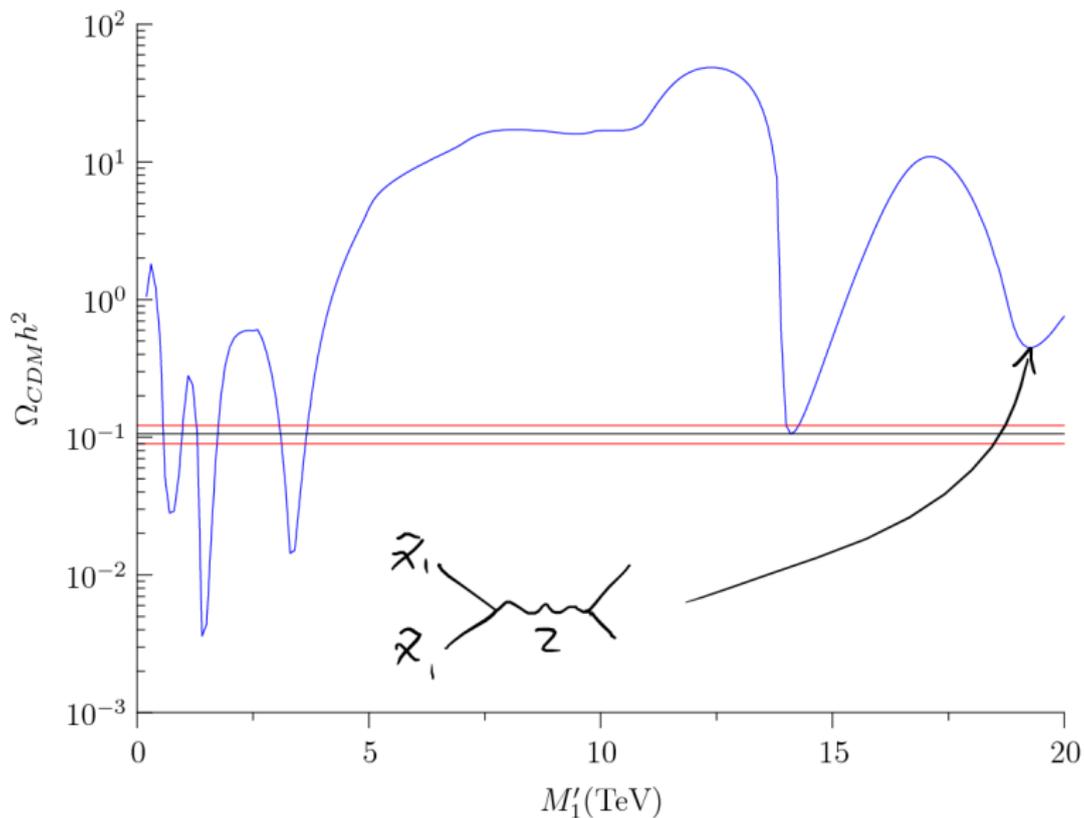
$M_1 = M'_1$: Relic Density

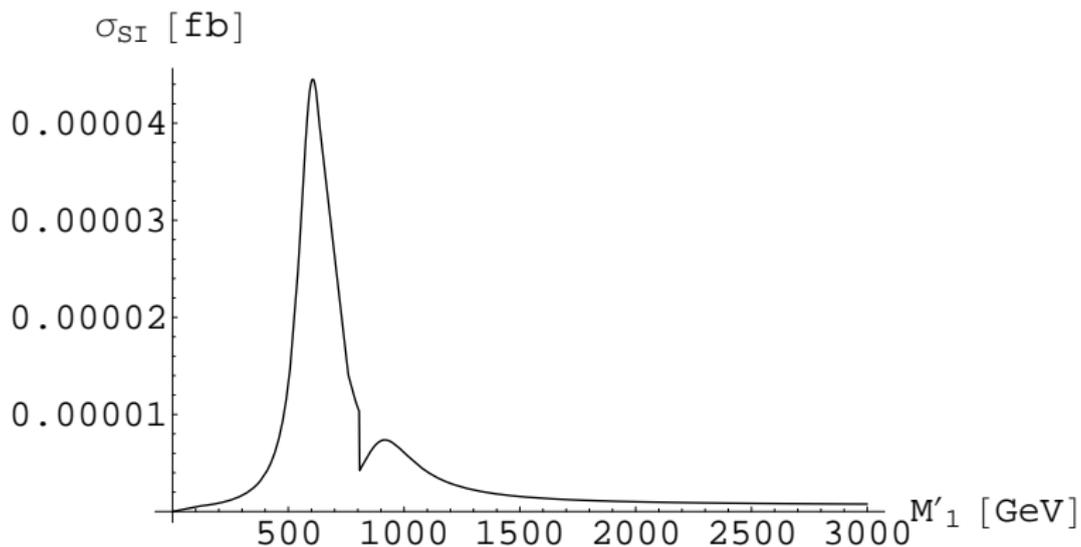


$M_1 = M'_1$: Relic Density



$M_1 = M'_1$: Relic Density



$M_1 = M'_1$: Direct Detection

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.
- These add new particle content **and** alter the dark matter phenomenology.

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.
- These add new particle content **and** alter the dark matter phenomenology.
- If they are realised in nature then we should see a spectrum of particles at the LHC.

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.
- These add new particle content **and** alter the dark matter phenomenology.
- If they are realised in nature then we should see a spectrum of particles at the LHC.
- Many models will fit this spectrum.

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.
- These add new particle content **and** alter the dark matter phenomenology.
- If they are realised in nature then we should see a spectrum of particles at the LHC.
- Many models will fit this spectrum.
- Precise data from a linear collider will allow us to narrow down the field of theories directly.

What does this have to do with a linear collider?

- A truly natural explanation of EWSB requires extensions of the MSSM.
- These add new particle content **and** alter the dark matter phenomenology.
- If they are realised in nature then we should see a spectrum of particles at the LHC.
- Many models will fit this spectrum.
- Precise data from a linear collider will allow us to narrow down the field of theories directly.
- This, along with direct detection results, will also allow us to differentiate them on the basis of their prediction of the relic density.