

PHYSICS AT LC - II

TOP & COLLIDER OPTIONS

SAURABH D. RINDANI

P.R.L. AHMEDABAD

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INTERACTION MEETING ON LINEAR COLLIDER,
NEW DELHI

OUTLINE

I. TOP PHYSICS

- MEASUREMENTS AT LC
- ISSUES ON TOP POLARIZATION
- ISSUES ON BEAM POLARIZATION
- SOME RECENT THEORETICAL WORK

II. $\gamma\gamma$ OPTION

- HIGGS PHYSICS
- OTHER ELECTRO-WEAK PROPERTIES
- EXTRA DIMENSIONS

III. $e\gamma, e^-e^-$ OPTIONS

- ANOMALOUS COUPLINGS
- EXTRA DIMENSIONS.

I. TOP PHYSICS

INTEREST IN TOP QUARKS:

- HEAVIEST QUARK ($m_t \sim 175 \text{ GeV}$)
- MASS CLOSE TO EWSB SCALE
- LARGE YUKAWA COUPLING IN SM ($Y_t \sim 1$)
- DECAYS FASTER THAN HADRONIZATION TIME SCALE

$$\Gamma_t \sim \text{GeV}$$

$$T \sim 3.6 \times 10^{-24} \text{ s.} < T_H \sim 10^{-23} \text{ s.}$$

TOP QUARK AT e^+e^- COLLIDER: ($\sqrt{s} \gtrsim 350 \text{ GeV}$)

- PRODUCTION THROUGH $e^+e^- \rightarrow t\bar{t}$
- CROSS SECTION $\sim 500 \text{ fb}$ ($\sqrt{s} = 500 \text{ GeV}$)
- ABOUT $10^5 t\bar{t}$ / YEAR
($\int L dt \sim 200 \text{ fb}^{-1}$)

- MEASUREMENTS THROUGH

SEMILEPTONIC CHANNELS (B.R. $2 \times 4/27$)

$$e^+e^- \rightarrow t\bar{t} \rightarrow (b\ell^+\bar{\nu}_\ell)(\bar{b}q\bar{q}') \rightarrow (b\bar{q}q')(\bar{b}\ell^-\nu_\ell)$$

PURELY LEPTONIC CHANNELS (B.R. $4/81$)

$$e^+e^- \rightarrow t\bar{t} \rightarrow (b\ell^+\bar{\nu}_\ell)(\bar{b}\ell^-\nu_\ell)$$

MEASUREMENTS NEAR THRESHOLD :

- TOP MASS DETERMINATION :

SHAPE OF TOTAL PRODUCTION C.S.
IN THE THRESHOLD REGION

- MORE DETAILED FIT TO

$\sigma_{t\bar{t}}$, p_t DISTRIBUTION , F-B ASYM

GIVES m_t (1 σ) , α_s (M_Z) , Γ_t , $g_{t\bar{t}H}$

- F-B ASYMM. MEASURES INTERFERENCE BETWEEN V & A CONTRIBUTIONS:
S-WAVE & P-WAVE ^{RESONANCE} CONTRIBUTIONS
- POLARIZATION OF TOP QUARK CAN BE MEASURED TO STUDY ITS PROPERTIES

MEASUREMENTS IN THE CONTINUUM

PRECISION MEASUREMENTS OF THE PROPERTIES OF TOP (MASS, WIDTH, DECAY BRANCHING RATIOS,...) AND ITS COUPLINGS CAN BE MADE.

→ TEST OF STANDARD MODEL.

ELECTROWEAK COUPLINGS

COUPLINGS OF THE TOP QUARK TO E.W.

GAUGE BOSONS (γ, Z, W) MAY BE WRITTEN AS:

$$\mathcal{L}_{\gamma, Z}^{\text{eff.}} = \sum_{i=\gamma, Z} e \left[\bar{t} \gamma_{\mu} (F_V^i + F_A^i \gamma_5) t G^{\mu\lambda} + \frac{i}{2m_t} \bar{t} \sigma_{\mu\nu} (c_M^i + \gamma_5 c_d^i) t F_i^{\mu\nu} \right]$$

$F_{V,A}^i$: FORM FACTORS (V,A)

c_M^i : MAGNETIC DIPOLE COUPLING

c_d^i : ELECTRIC DIPOLE COUPLING (P, T VIOLATING)

$$\mathcal{L}_{tbW}^{\text{eff}} = -\frac{g}{\sqrt{2}} V_{tb} \bar{t} \left[\gamma^{\mu} (f_{tL} P_L + f_{tR} P_R) W_{\mu}^i - \frac{i}{m_W} \sigma^{\mu\nu} (f_{2L} P_L + f_{2R} P_R) W_{\mu\nu}^i \right] b$$

SIMILARLY FOR $\bar{t} b \bar{W}$

SM: $f_{tL} = 1$, $\bar{f}_{tL} = 1$, REST ZERO.

FORWARD-BACKWARD ASYMMETRY CAN BE USED
TO DETERMINE/CONSTRAIN $F_{V,A}^{\gamma,Z}$

FOR OTHER FORM FACTORS, MORE DETAILED
INFORMATION IS NEEDED: TOP POLARIZATION
AND/OR DECAY DISTRIBUTIONS (ENERGY, ANGULAR)
(PARTICULARLY FOR STUDYING CP VIOLATION)
FULL SPIN DENSITY MATRIX INFORMATION
 \equiv DECAY DISTRIBUTION INFORMATION
(IFF DECAY IS GOVERNED BY SM).

SOME ISSUES RELATED TO TOP POLARIZATION

1. TOP POLARIZATION CAN BE MEASURED
ONLY THROUGH DECAY DISTRIBUTIONS
(J. KÜHN 1984; CZARNECKI, JEZABEK
& KÜHN 1994)
(NEEDS ACCURATE KNOWLEDGE OF DECAY)
2. PREDICTIONS FOR POLARIZATION DEPEND
ON CHOICE OF SPIN BASIS
(S. PARKE & Y. SHADMI 1996; G. MAHLON & PARKE 1997)

ALTERNATIVE APPROACH :

MAKE PREDICTIONS FOR DECAY LEPTON
ANGULAR DISTRIBUTIONS

ADVANTAGE

- INDEPENDENT OF ANOMALOUS $t\bar{b}W$ COUPLINGS
IN THE LINEAR APPROX. OF ANOMALOUS COUPLINGS,
(GRZADKOWSKI & HIOKI, SDR)
(FIRST SHOWN TO HOLD FOR $m_b \rightarrow 0$. THEN
FOUND TO BE TRUE FOR $m_b \neq 0$.)
- THIS PROPERTY DOES NOT DEPEND ON PRODUCTION
PROCESS FOR $t\bar{e}$
(FOR $\gamma\gamma \rightarrow t\bar{e}$: OHKUMA 2002 ;
GODBOLE, SDR, SINGH 2003.
FOR ANY PROCESS : GRZADKOWSKI & HIOKI)
- DOES NOT NEED PRECISE RECONSTRUCTION
OF TOP ENERGY-MOMENTUM.
- TOP SPIN BASIS NOT RELEVANT
- INCIDENTALLY, ALSO INDEPENDENT OF
QCD CORRECTIONS TO $t\bar{b}W$ VERTEX

HOW CAN ONE STUDY DECAY PROPERTIES/
ANOMALOUS $t \rightarrow W$ COUPLINGS?

ANSWER: USE b QUARK DISTRIBUTION
(DIFFICULT IN PRACTICE)

OR
USE l^\pm ENERGY DISTRIBUTION
(WITH/WITHOUT ANGULAR DIST.)

(GRZADKOWSKI & HIOKI, SDR)

e^-/e^+ BEAM POLARIZATION

e^- LONGITUDINAL POLARIZATION OF
80% EXPECTED TO BE AVAILABLE

e^+ LONG. POL. MAY BE LIMITED TO 60%

IN PRINCIPLE, LONG. POLARIZATION CAN BE
CONVERTED TO $\overset{N}{\wedge}$ TRANSVERSE.

- LONGITUDINAL POLARIZATION CAN HELP BY
EITHER SUPPRESSING COUPLINGS OR
ENHANCING COUPLINGS

→ HELPS TO SUPPRESS BACKGROUND

→ CAN HELP TO ENHANCE SENSITIVITY

- FLIPPING OF SIGN OF POLARIZATION CAN
GIVE ADDITIONAL INDEPENDENT DATA.

SOME ISSUES RELATED TO TRANSVERSE POLARIZATION

TRANSVERSE e^- AND/OR e^+ POL. CAN HELP
TO ENHANCE SENSITIVITY TO CERTAIN COUPLINGS

RECENT EXAMPLES:

T. RIZZO : IN THE CONTEXT OF TEV-SCALE
(2002) GRAVITY & EXCHANGE OF MASSIVE
SPIN-2 GRAVITONS

J. FLEISCHER : ANOMALOUS COUPLINGS IN
K. KOLODZIEJ $e^+e^- \rightarrow W^+W^-$
F. JEGERLEHNER
(1994)

2 DIFFERENT ROLES PLAYED BY TRANSVERSE POL.

1. IN $e^+e^- \rightarrow t\bar{t}$ CP VIOLATION (OR T VIOLATION)

NEEDS EITHER POLARIZED BEAMS OR
MEASUREMENT OF TOP POLARIZATION,

IF TRANSVERSE POLARIZATION IS PRESENT
TRIPLE PRODUCTS (CP & T - ODD) CAN BE

SEARCHED FOR (WITHOUT t DECAY DISTRIBUTIONS)

$$\text{e.g.}, \quad \vec{p}_{e^-} \times (\vec{S}_{e^-} - \vec{S}_{e^+}) \cdot \vec{p}_t$$

(B. ANANTHANARAYAN & SDR, 2003)

2. ONE CAN USE AZIMUTHAL DISTRIBUTIONS TO
LOOK FOR VIOLATION OF CHIRAL INVARIANCE.

FOR $e^+e^- \rightarrow f\bar{f}$,

$$\begin{aligned}\frac{d\sigma}{d\Omega} \propto & \frac{1}{4} \left(|T_{+-}|^2 + |T_{++}|^2 + |T_{--}|^2 + |T_{-+}|^2 \right) \\ & + 2P\bar{P} \operatorname{Re} \left(e^{-2i\phi} [T_{+-}^* T_{-+}] + [T_{++}^* T_{--}] \right) \\ & + 2P \operatorname{Re} \left(e^{-i\phi} [T_{+-}^* T_{--} + T_{++}^* T_{-+}] \right) \\ & + 2\bar{P} \operatorname{Re} \left(e^{-i\phi} [T_{+-}^* T_{++} + T_{--}^* T_{-+}] \right)\end{aligned}$$

(HIKASA 1986)

CHIRAL INVARIANCE \Rightarrow $T_{++} = 0$ (EQUAL
($m_e = 0$) $T_{--} = 0$ e^- & e^+
HELICITIES)

\rightarrow NO $\cos\phi$, $\sin\phi$ TERMS

ASYMMETRY LIKE $\phi \rightarrow -\phi + \pi$ ($\cos\phi \rightarrow -\cos\phi$)
OR $\phi \rightarrow 2\pi - \phi$ ($\sin\phi \rightarrow -\sin\phi$)

CAN SEARCH/MEASURE T_{++} OR T_{--} AMPLITUDES

THESE DO NOT INTERFERE WITH ONLY
LONGITUDINAL POLARIZATION

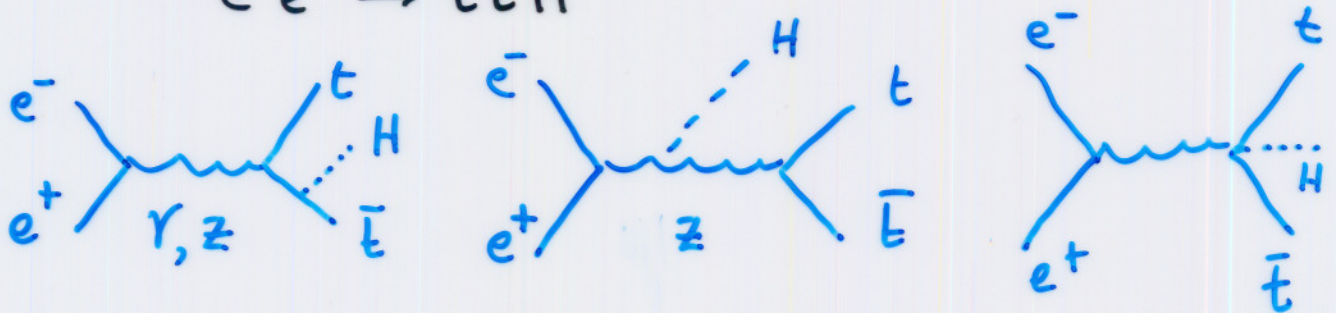
THIS IS WHERE TRANSVERSE POL. CAN HELP

OTHER TOP COUPLINGS

- $\bar{t}tH$ Yukawa coupling.

CAN BE MEASURED IN THE PROCESS

$$e^+e^- \rightarrow t\bar{t}H$$

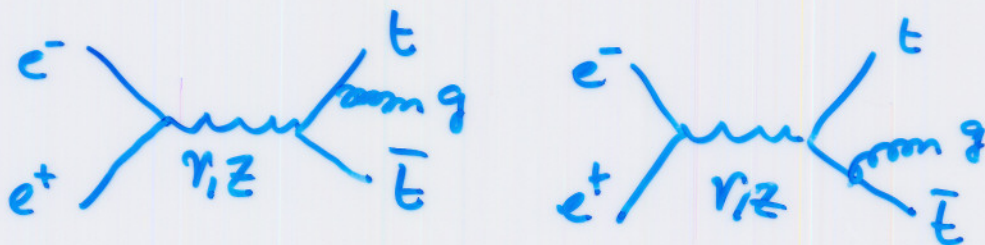


DEPENDING ON THE MASS OF THE HIGGS THIS
PROCESS WOULD HAVE VARYING SENSITIVITY

DETAILED STUDY OF BACKGROUND & LOOP CORRECTIONS
NEEDED

- GLUONIC COUPLINGS OF TOP

AN e^+e^- COLLIDER IS A GOOD PLACE TO MEASURE
THE $t\bar{t}g$ COUPLING THROUGH $e^+e^- \rightarrow t\bar{t}g$
SINCE ONLY t & \bar{t} CAN EMIT A GLUON (NOT e^+e^-
OR γ, Z)



BEYOND SM : SOME RECENT WORK

CP VIOLATION IN TOP PRODUCTION & DECAY :

- SEVERAL PROPOSALS HAVE EXISTED
- INVOLVE MEASUREMENT OF TOP POLARIZATION OR DECAY PRODUCT CORRELATIONS/DISTRIBUTIONS (KANE, LADINSKY & YUAN; SCHMIDT & PESKIN; ... BERNREUTHER ET AL. ;)
- RECENT WORK INVOLVING DECAY DISTRIBUTIONS :
 - CHARGED LEPTON ANGULAR DISTRIBUTION: (P. POULOUSE & SDR, B. GRZADKOWSKI & Z. HIOKI, SDR)
 - MEASURE OF ELECTRIC & WEAK DIPOLE COUPLINGS
 - INDEPENDENT OF CP VIOLATION IN DECAY
 - INCLUDING EFFECT OF LONGITUDINAL POL.
 - INCLUDING $O(\alpha_s)$ QCD CORRECTIONS
 - CP-VIOLATING ANGULAR ASYMMETRY LEADS TO BEST LIMITS OF $(1-2) \times 10^{-18} \text{ e cm}$
(90% CL, $\sqrt{s} = 500 \text{ GeV}$, $\int L dt = 500 \text{ fb}^{-1}$)
 - CLOSE TO OPTIMISTIC VALUES FROM POPULAR EXTENSIONS OF SM
 - CHARGED LEPTON ENERGY DISTRIBUTIONS & b-QUARK ENERGY & ANGULAR DISTRIBUTIONS: (B. GRZADKOWSKI, Z. HIOKI, SDR)
 - CONTRIBUTIONS FROM BOTH DIPOLE COUPLINGS & CP-VIOLATION IN DECAY.

SINGLE TOP PRODUCTION

→ SINGLE TOP PRODUCTION IS POSSIBLE THROUGH

$$e^+e^- \rightarrow e^- \bar{\nu}_e t \bar{b}$$

(S. AMBROSANO & B. MELE 1994;
N. DOKHOLIAN, G.V. JIKIA 1994)

A LARGE NUMBER OF DIAGRAMS CONTRIBUTE

SOME CORRESPOND TO $e^+e^- \rightarrow t \bar{t}$ FOLLOWED BY
DECAY, AND THESE HAVE TO BE SUBTRACTED

RIGHT-HANDED POLARIZATION ($e_R^- e_R^+$) ~~TO~~

HELPS TO REMOVE THE $t \bar{t}$ BACKGROUND

$$\sigma(e_R^- e_R^+) \sim \text{FEW fb AT } \sqrt{s} > 500 \text{ GeV}$$

AND CAN BE MEASURED

(E. BOOS ET AL. 2001)

→ ANOMALOUS $t b W$ COUPLINGS CAN BE CONSTRAINED

$$-0.05 \leq f_{2R} \leq 0.05$$

$$\text{FOR } \sqrt{s} = 500 \text{ GeV, } L = 500 \text{ fb}^{-1}$$

→ FLAVOUR-CHANGING COUPLINGS TO Z & γ
CAN PRODUCE SINGLE TOP:

$$e^+e^- \rightarrow t \bar{q}$$

(V. OBRAZTSOV ET AL. 1998

T. HAN & J. HEWETT, 1999

S. BAR-SHALOM & J. WUDKA, 1999)

CAN BE USED TO BOUND $Z t \bar{q}$, $\gamma t \bar{q}$ COUPLINGS
BEAM POLARIZATION CAN HELP REDUCE BACKGROUND.

J. AGUILAR-SAAVEDRA, 2001.

II. PHOTON COLLIDER : $\gamma\gamma$ OPTION

HIGH-ENERGY PHOTONS PRODUCED BY
COMPTON BACKSCATTERING OF INTENSE
LASER BEAMS OFF HIGH-ENERGY ELECTRONS
ENERGY SPECTRUM OF PHOTONS DETERMINED
BY

$$x = \frac{4 E_b \omega_0}{m_e^2} = 15.3 \left(\frac{E_b}{\text{TeV}} \right) \left(\frac{\omega_0}{\text{eV}} \right)$$

E_b : ELECTRON ENERGY

ω_0 : LASER ENERGY

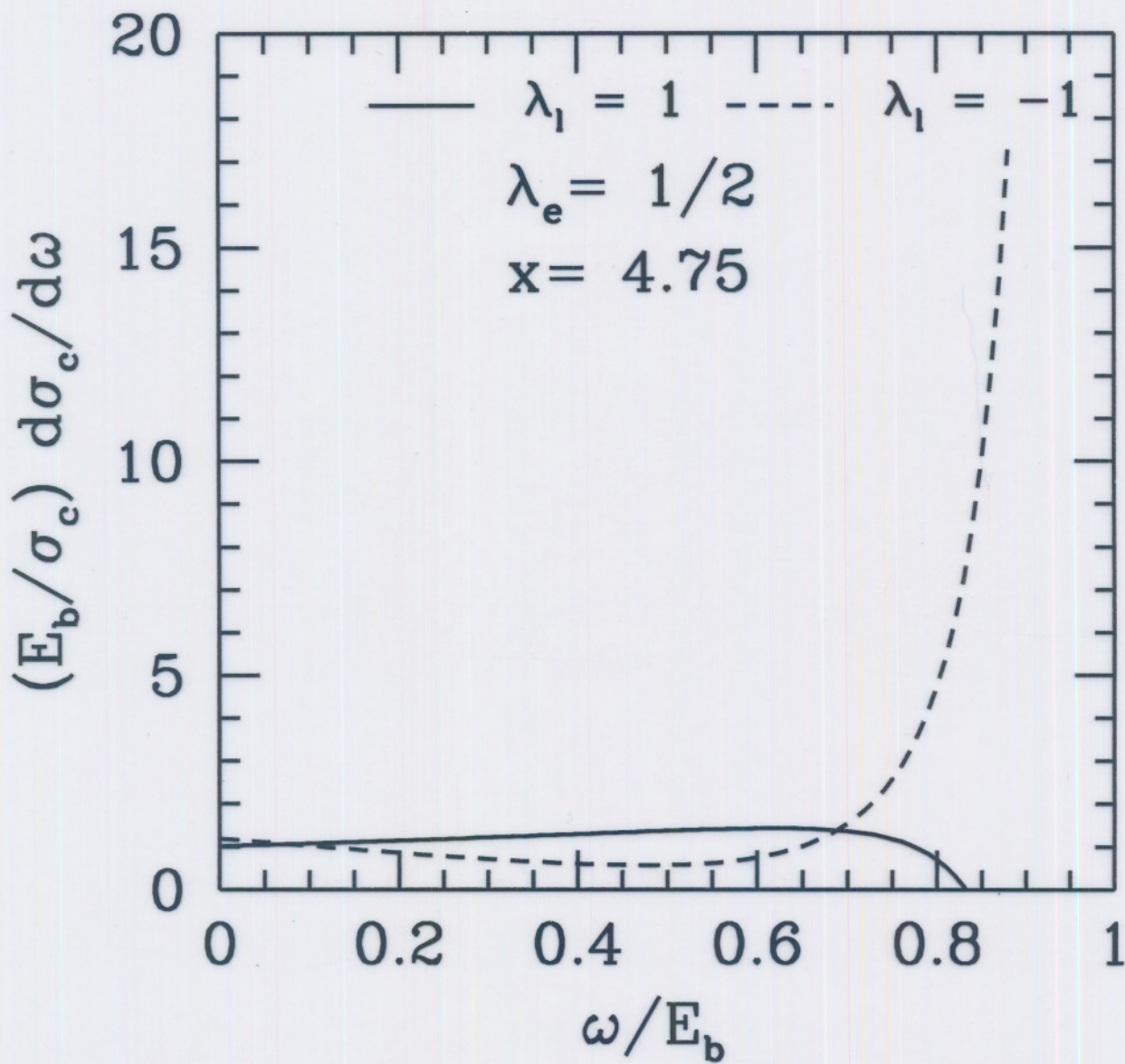
AND POLARIZATION OF ELECTRONS & LASER γ 'S

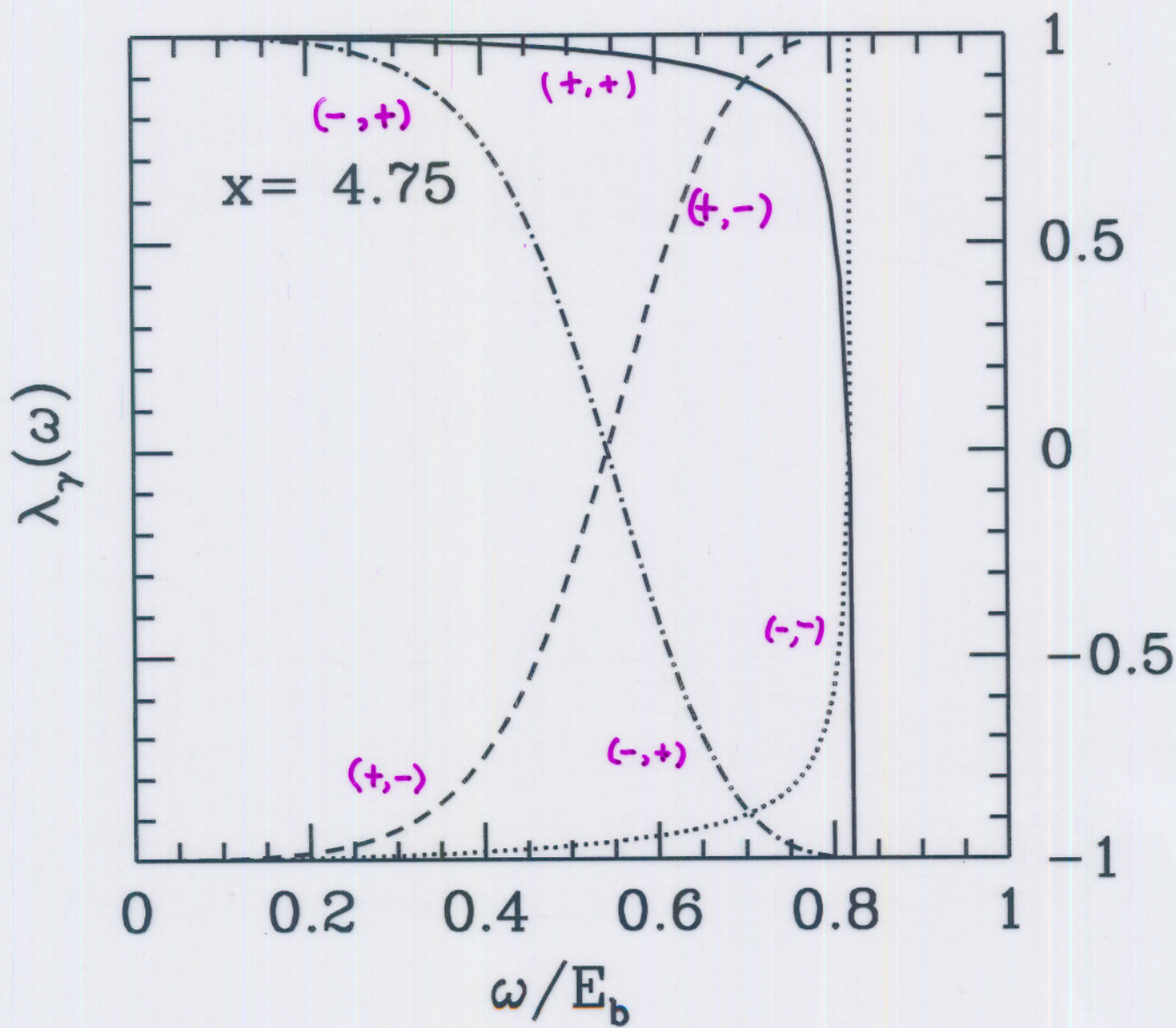
LARGE FRACTION OF THE ELECTRON
ENERGY & POLARIZATION TRANSFERRED

TO THE HIGH-ENERGY PHOTON. $\left(\frac{\omega_{\max}}{E_b} = \frac{x}{x+1} \right)$

(FOR OPPOSITE HELICITIES)

FOR $\lambda_e \lambda_\gamma = -\frac{1}{2}$, BOTH INTENSITY AND
POLARIZATION PEAK AT LARGE ENERGY





APPLICATIONS: TWO-PHOTON ACCESSIBLE MORE DIRECTLY

HIGGS PHYSICS

e^+e^-H COUPLING SUPPRESSED BY m_e/E

$\gamma\gamma H$ COUPLING POSSIBLE AT ONE LOOP

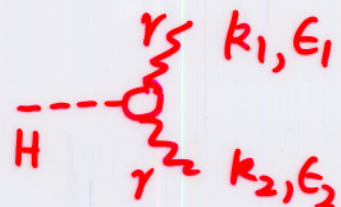


- HIGGS CAN BE PRODUCED AS RESONANCE.
- ACCURATE DETERMINATION OF $H\gamma\gamma$ COUPLING
(AND HENCE $\Gamma(H \rightarrow \gamma\gamma)$ POSSIBLE)
(IMPORTANT FOR DETERMINATION OF PARAMETERS)
- CP PROPERTIES OF HIGGS CAN BE DETERMINED
USING POLARIZED PHOTONS
(GRZADKOWSKI & GUNION, 1992)

EFFECTIVE COUPLINGS:

SCALAR: $(R_1 \cdot R_2 E_1 \cdot E_2 - E_1 \cdot R E_2 \cdot R_1)$

PSEUDOSCALAR: $\epsilon_{\mu\nu\alpha\beta} E_1^\mu E_2^\nu R_1^\alpha R_2^\beta$



BEYOND SM : (MULTI-HIGGS, MSSM)

- MORE THAN ONE HIGGS
- CP VIOLATION

CP PROPERTIES OF HIGGS CAN BE DETERMINED ALSO FROM FINAL STATE POLARIZATION IN

$$\gamma\gamma \rightarrow t\bar{t}$$

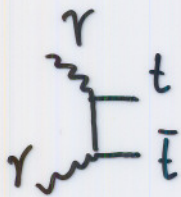
(E. ASAKAWA ET AL., 2000, 2003)

~~S. CHOI & K. HAGIWARA, 1995; M.S. BAEK ET AL. 1997~~

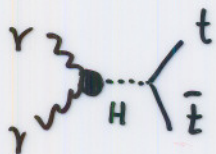
t, \bar{t} POLARIZATION INFORMATION CAN BE BYPASSED USING CHARGED LEPTON ANGULAR ASYMMETRIES

(R. GODBOLE, SDR, R.K. SINGH, 2003)

MODEL INDEPENDENT ANALYSIS :



$$V_{t\bar{t}H} = -ie \frac{m_t}{m_W} (S_t + i\gamma^5 P_t)$$



$$V_{\gamma\gamma H} = -\frac{i\alpha\sqrt{s}}{4\pi} \left[S_\gamma \left(\epsilon_1 \cdot \epsilon_2 - \frac{1}{k_1 \cdot k_2} \epsilon_1 \cdot k_2 \epsilon_2 \cdot k_1 \right) - P_\gamma \frac{1}{k_1 \cdot k_2} \epsilon_{\mu\nu\alpha\beta} \epsilon_1^\mu \epsilon_2^\nu k_1^\alpha k_2^\beta \right]$$

S_γ, P_γ : COMPLEX FORM FACTORS

ANOMALOUS $t\bar{b}W$ INTERACTIONS DONT CONTRIBUTE

$$6 \text{ PARAMETERS : } \begin{cases} S_t, P_t \\ \text{Re } S_\gamma & \text{Im } S_\gamma \\ \text{Re } P_\gamma & \text{Im } P_\gamma \end{cases}$$

CHOOSE $2\lambda_e\lambda_\gamma = -1$ (HARD PHOTON SPECTRO

& $\lambda_{e^-} = \lambda_{e^+}$ (HIGGS COUPLE TO
PHOTONS OF EQUAL
HELICITY ; SM : OPPOS

NOW $\lambda_{e^-} = +\frac{1}{2}$ OR $-\frac{1}{2}$

CORRESPONDING TO l^+ OR l^- IN FINAL
STATE, WE HAVE 4 COMBINATIONS

$\sigma(+,+)$ $\sigma(+,-)$ $\sigma(-,+)$ $\sigma(--)$

CAN CONSTRUCT 6 ASYMMETRIES:

2 CP-VIOLATING ~~CP~~

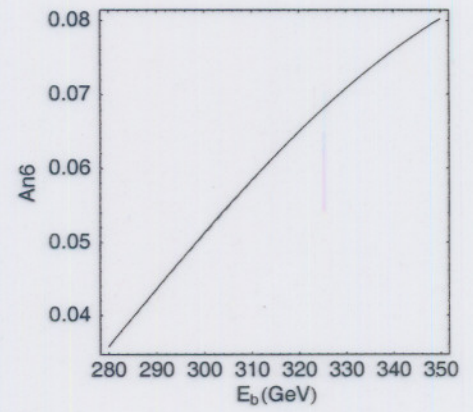
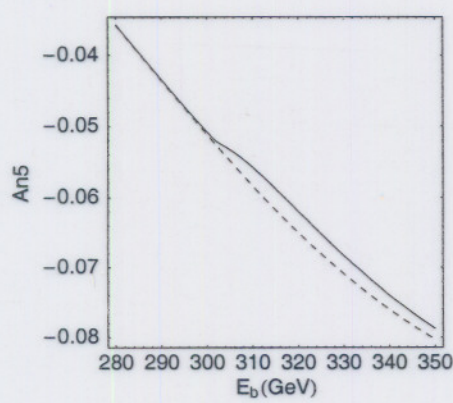
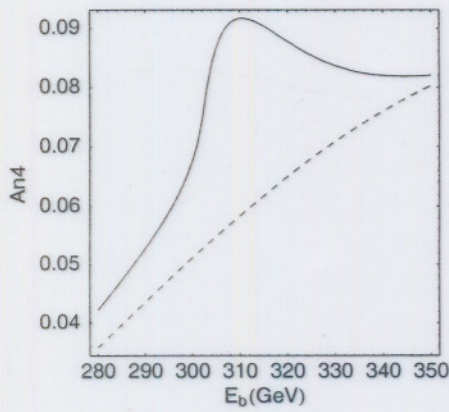
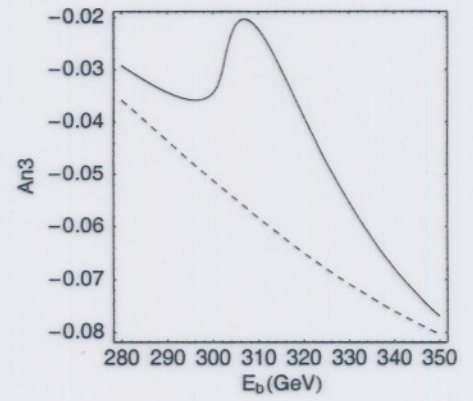
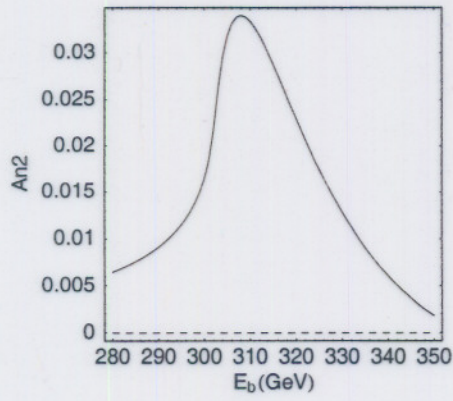
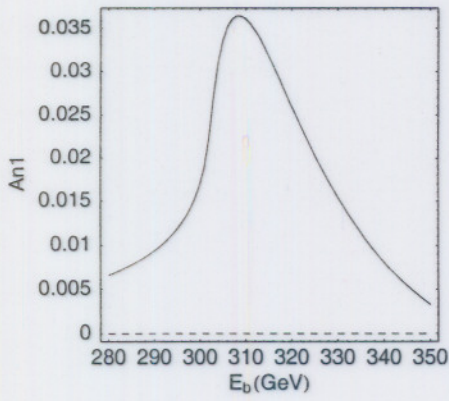
2 CHARGE ASYMMETRIES

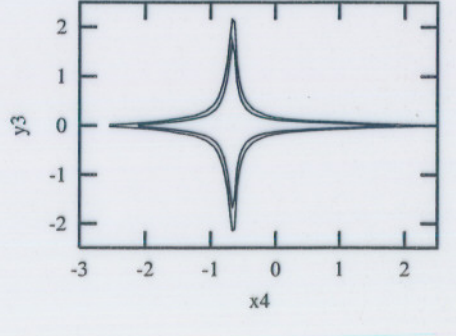
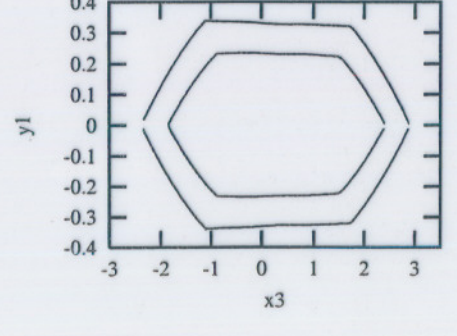
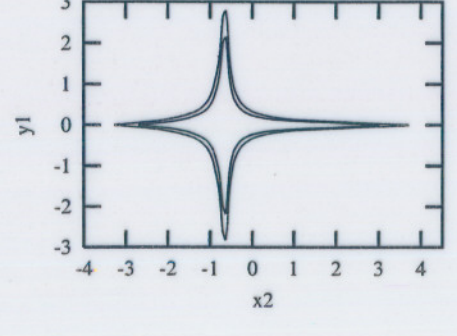
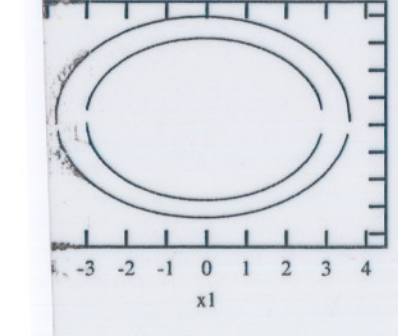
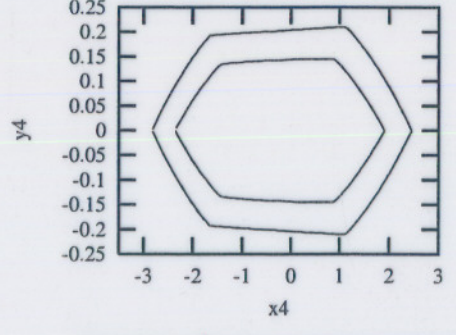
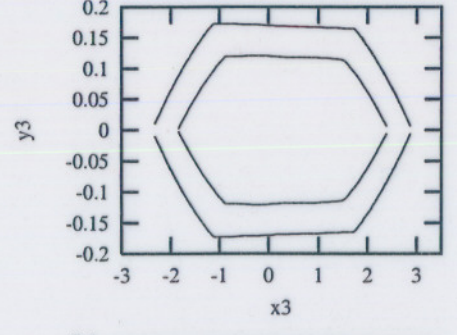
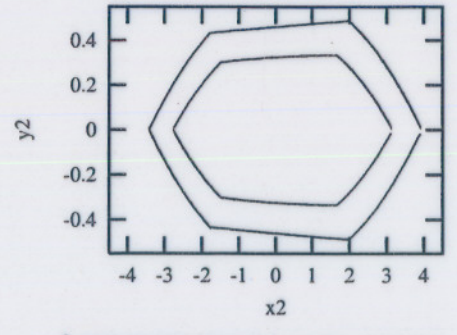
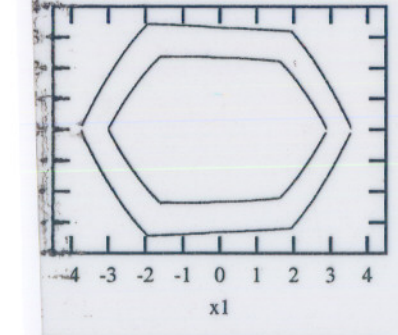
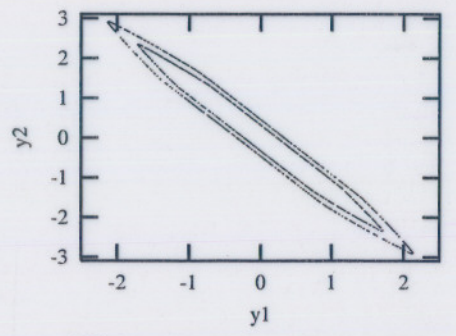
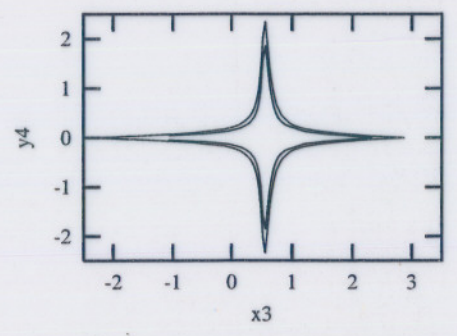
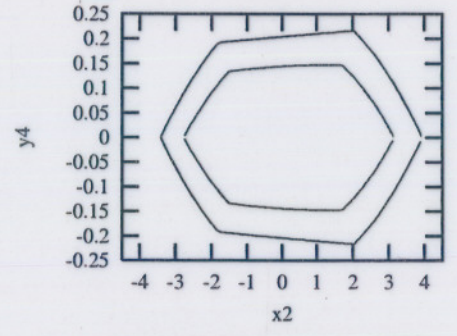
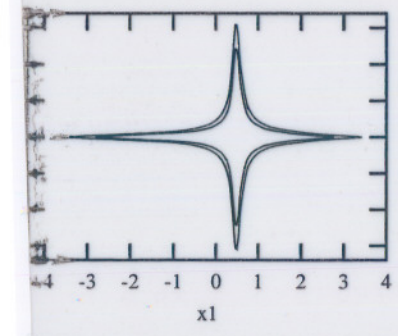
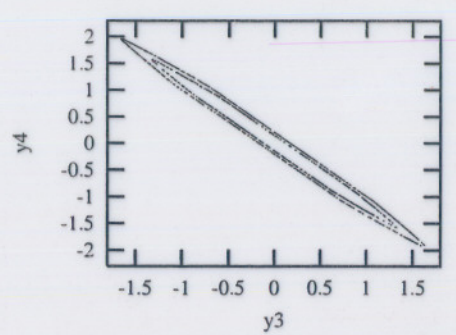
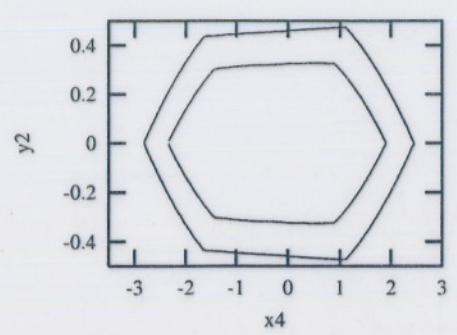
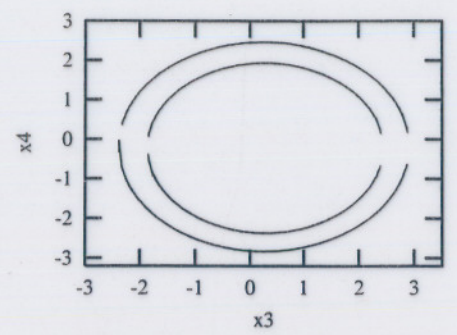
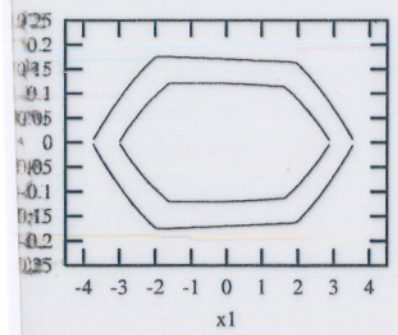
2 POLARIZATION ASYMMETRIES

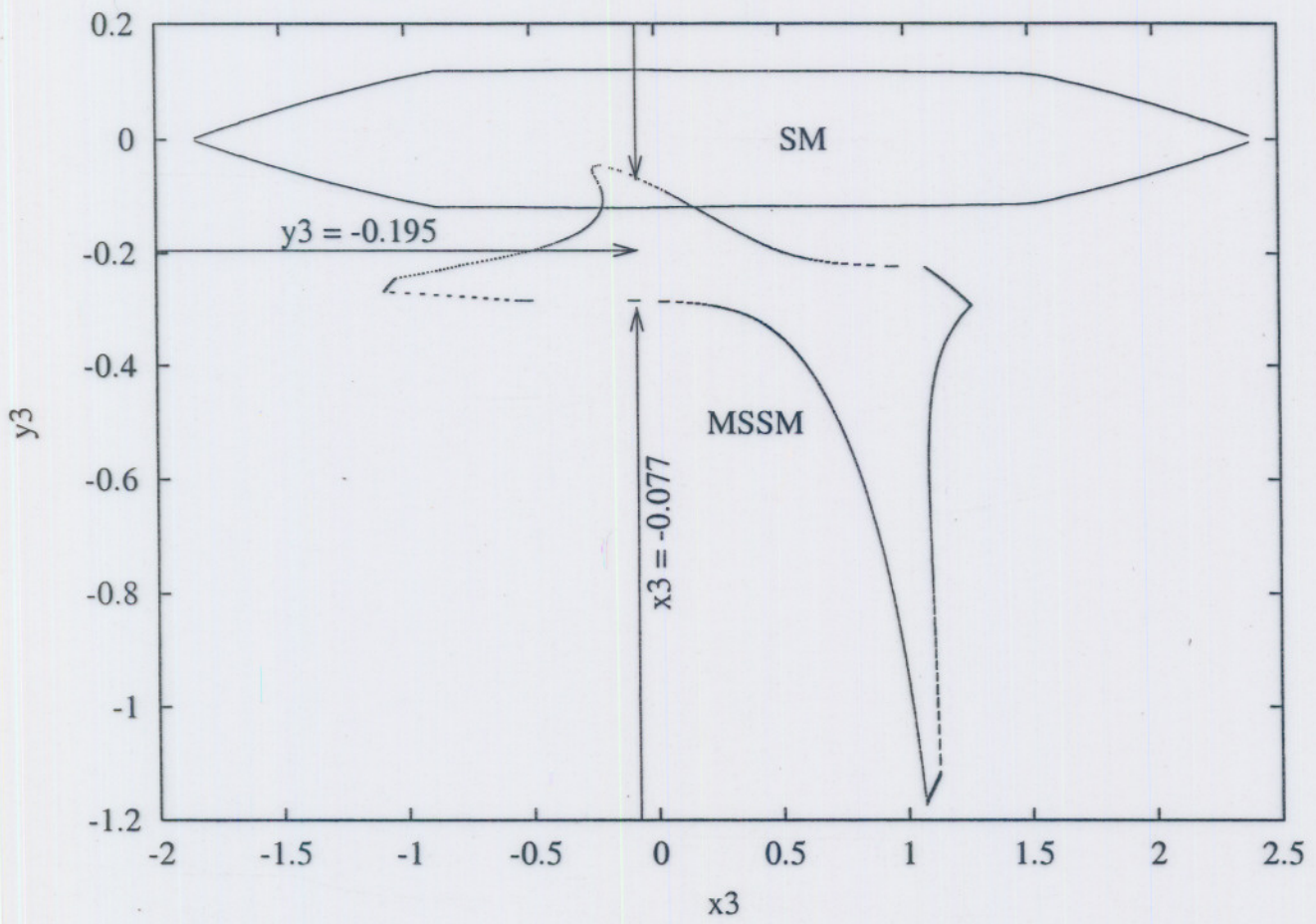
CAN FIND REGIONS IN PARAMETER SPACE
WHERE THE THEORY CAN BE DISTINGUISHED
FROM S.M. AT 95% C.L.

- STRONG LIMITS ON CP-VIOLATING COMBINATIONS OF PARAMETERS (VARIED 2 AT A TIME)
- DETERMINATION OF ALL PARAMETERS NEEDS LINEAR POLARIZATION

ASYMMETRIES IN MSSM FOR A CHOICE OF PARAMETERS.



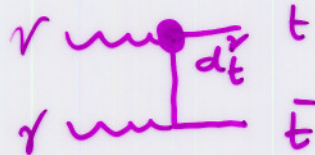




OTHER ELECTROWEAK PROPERTIES :

$\gamma\gamma$ CAN HELP MEASURE SEPARATELY

ELECTRIC DIPOLE MOMENTS d^{γ} (INDEPENDENT OF d^Z)



USE OF TOP POLARIZATION TO STUDY CP-VIOLATION

DIPOLE COUPLINGS OF TOP:

S.Y. CHOI & K. HAGIWARA (1995)

M.S. BAEK, S.Y. CHOI & C.S. KIM (1997)

USE OF DECAY-LEPTON ASYMMETRIES

TO MEASURE TOP DIPOLE COUPLING:

P. POULOSE & SDR (1998)

INCLUDING ANOMALOUS $\gamma\gamma Z$ COUPLING WITH d^{γ}

P. POULOSE & SDR (1999)

INCLUDING $\gamma\gamma H$ COUPLING AS WELL AS d_t^{γ} :

B. GRZADKOWSKI, Z. HIOKI, K. OHKUMA, J. WUDKA (2003)

OTHER THAN TOP: CHOI, HAGIWARA & BAEK (1996) ($\gamma\gamma \rightarrow W^+W^-$)

J.L. HEWETT & F.J. PETRIELLO (2001) ($\gamma\gamma \rightarrow \gamma\gamma, \gamma Z, Z\gamma$)

G. GOUNARIS & F.M. RENARD (1996).....

EXTRA DIMENSIONS :

- GRAVITONS HAVE TREE-LEVEL COUPLING TO TWO PHOTONS



- CAN CONTRIBUTE TO REAL GRAVITON EMISSION AS WELL AS VIRTUAL GRAVITON EXCHANGE.

A NUMBER OF CALCULATIONS HAVE BEEN CARRIED OUT FOR VARIOUS FINAL STATES WITH GRAVITON EXCHANGE :

T. RIZZO (1999) (GAUGE BOSONS)

K. AGASHE & N.G. DESHPANDE (1999)

D. ATWOOD ET AL. (2000) (REAL GRAVITONS)

H. DAVOUDIASHI (1999)

K.Y. LEE ET AL. (1999) (W^+W^-) ($t\bar{t}$)

* P. MATHEWS, P. POULOSE, K. SRIDHAR (1999) ($t\bar{t}$)

D.K. GHOSH, P. MATHEWS, P. POULOSE, K. SRIDHAR (1999)

M. DONCHESKI & R.W. ROBINETT (2000) \rightarrow (2 JETS)

S.R. CHOUDHURY, A.S. CORNELL, G.C. JOSHI (2002) ($\gamma\gamma \rightarrow ZZ$)

* FOR EXAMPLE, FROM TOTAL CROSS SECTIONS

& RAPIDITY DISTRIBUTIONS,

LIMIT ON M_S OF 1.6 TeV ($\sqrt{s} = 500$ GeV) ($\int_{|t|}^{10} \frac{d\sigma}{dt} = 10$ fb)

IS POSSIBLE.

\rightarrow IMPROVES TO 1.95 TeV OR 2.5 TeV WITH POLARIZATION.

CALCULATIONS WITH REALISTIC SPECTRUM

IDEAL COMPTON SCATTERED SPECTRUM IS VALID

- IF
- 1) LASER POWER IS SMALL
 - 2) PRIMARY ELECTRON BEAMS SUFFICIENTLY WIDE

1) \Rightarrow NON-LINEAR EFFECTS SMALL

2) \Rightarrow EFFECTS RELATED TO PHOTON SCATTERING ANGLE CAN BE NEGLECTED.

SIMULATION IN A REALISTIC CASE DONE BY

V. TELNOV, AND ALSO BY P. CHEN ET AL.
("CAIN")

A. ZARNECKI HAS PROVIDED A PARAMETRIZATION WHICH AGREES WITH TELNOV SIMULATION AT THE HIGH ENERGY END.

ZARNECKI'S PARAMETRIZATION ^{"COMP AZ"} HELPS TO INTERPOLATE FOR VARIOUS BEAM ENERGIES

WORK IS IN PROGRESS TO REVISE LIMITS ON VARIOUS NEW PHYSICS IN $\gamma\gamma \rightarrow e\bar{e}$ USING

CompAZ. (R.M. GODBOLE, P. MATHEWS, P. POULOSE, SDR, R.K. SINGH)

\rightarrow INTERESTING EFFECTS GET DILUTED

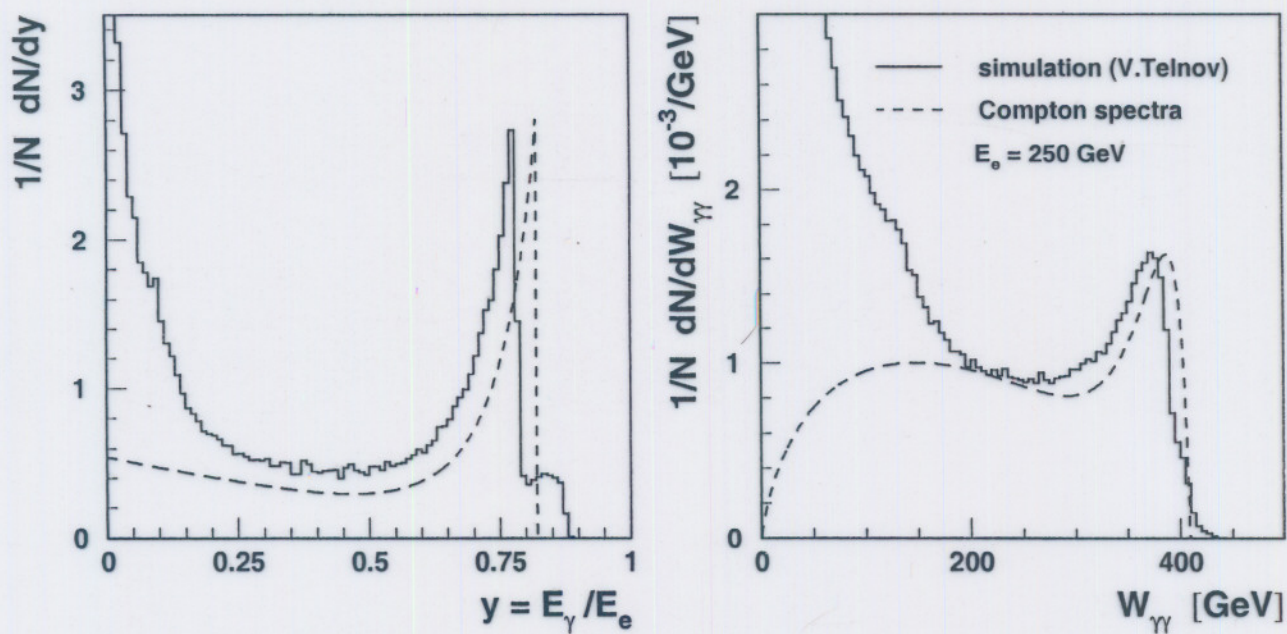


Fig. 1. Energy distribution for photons (left plot) and the $\gamma\gamma$ center-of-mass energy distribution (right plot) from full simulation of luminosity spectrum by V.Telnov [7] (solid line), compared to expectations for the simple Compton scattering (dashed line). For better comparison of shape, Compton spectra is scaled to the same height of the high energy peak.

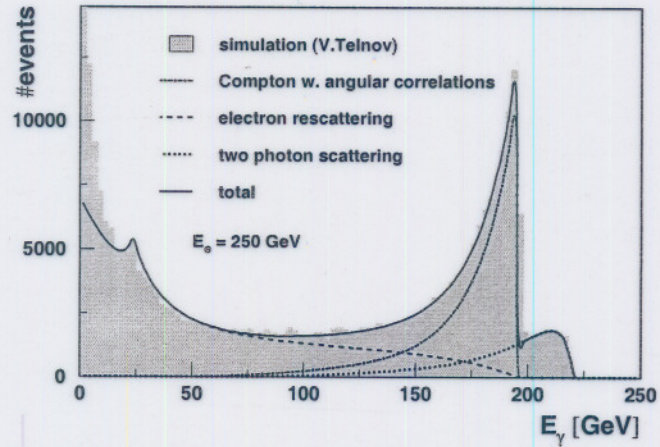


Fig. 3. Comparison of the photon energy distribution obtained from full simulation of luminosity spectrum by Telnov [7], with the fitted contributions of different processes considered in the described model, as indicated in the plot.

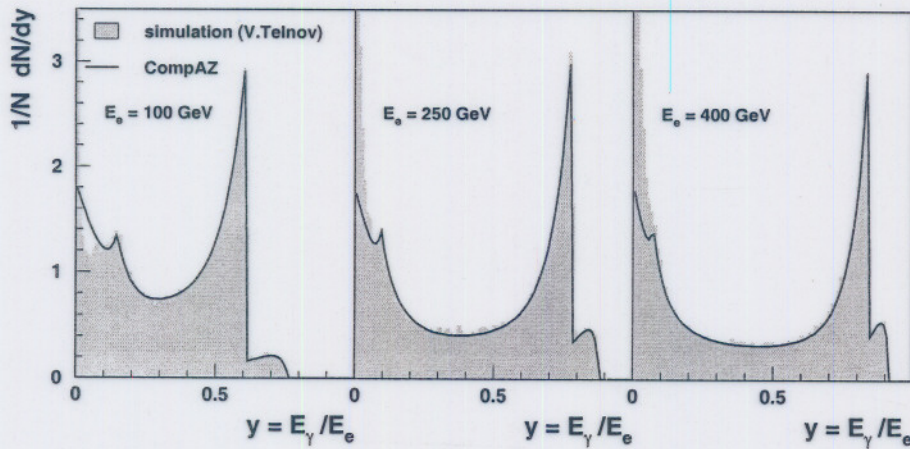


Fig. 4. Comparison of the photon energy distribution from the fitted parametrization with the distribution obtained from full simulation of luminosity spectra [7], for three electron beam energies, as indicated in the plot. Imposed cut on the energy of the second photon is 40, 150 and 260 GeV respectively.

$e\gamma$ & e^-e^- OPTIONS

$e\gamma$ OPTION MAY BE REALIZED BY USING
COMPTON BACKSCATTERING FOR ONE OF THE BEAMS.

- SOMEWHAT BETTER CONTROL ON POLARIZATION
- A DIFFERENT RANGE OF PROCESSES ACCESSIBLE (SUSY: PREVIOUS TALK)
- AN OBVIOUS PROCESS WITH $e\gamma$ OPTION IS

$$e\gamma \rightarrow e\gamma \quad (\text{COMPTON SCATTERING})$$

WHICH CAN BE USED TO CONSTRAIN NEW PHYSICS.

(H. DAVOUDIASHVILI, 1999; EXTRA DIMENSIONS)
(R. MATTHEWS, 2001; NON-COMMUTATIVE SPACE-TIME)

- $e\gamma \rightarrow eG$ (G: GRAVITON) HAS BEEN PROPOSED
FOR TEST OF WEAK SCALE GRAVITY

(D.K. GHOSH, P. POULOSE, K. SRIDHAR, 2000
D. ATWOOD, S. BAR-SHALOM, A. SONI, 2000)

STRINGENT BOUNDS MAY BE OBTAINED USING $e\gamma$
POLARIZATION

- VARIOUS ELECTROWEAK PROCESSES ARE POSSIBLE
e.g. $e^- \gamma \rightarrow \nu W$, $e^- \gamma \rightarrow e^- Z$, $e^- \gamma \rightarrow \bar{e} b \nu^*$

* G. JIKIA, 1992; E. BOOS ET AL., 1996, 1997, 2001;
J.J. CAO ET AL., 1998

ANALYSIS QUITE COMPLICATED.

AGAIN, OBVIOUS PROCESS THAT CAN BE STUDIED
WITH e^-e^- OPTION IS

$$e^-e^- \rightarrow e^-e^- \quad (\text{MØLLER SCATT.})$$

- HERE NEUTRAL PARTICLES ($\gamma, Z, Z', \text{Graviton}$)
CAN CONTRIBUTE ONLY THROUGH t, u CHANNEL
EXCHANGE.

- s CHANNEL CAN BE USEFUL FOR DOUBLY-CHARGE
PARTICLES — ~~Gauge~~ ^{HIGGS} BOSONS, "DILEPTONS", ETC.

(C.A. HEUSCH, 1995; F. CUYPERS, 1993
M. RAIDAL, 1996 ...)

- t, u CHANNEL GRAVITON EXCHANGE IS STUDIED
BY

T. RIZZO 1999; D.K. GHOSH & S. RAYCHAUDHURI
(ADD MODEL) (R-S MODEL) 2000.