A Spin Rotator for CLIC

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- Spin rotator design criteria
- Spin dynamics
- CLIC RTML layout and SR location
- Spin rotator lattice and options
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Spin Rotator Design Criteria

- Design Criteria (P. Emma for NLC, 1994)
 - Spin should be orientable in any direction
 - Net momentum compaction must be small such that energy fluctuations do not become longitudinal position fluctuations (less than 100 μ m bunch length @ IP for NLC)
 - It should be located such that total spin diffusion due to energy spread is small
 - System should not dilute significantly the beam transverse emittance (small energy spread)
 - System should be short, simple and robust

Spin Dynamics

Spin Precession

$$\phi_s = G \gamma_0 \alpha$$

Mean polarization:

$$< P_z> = P_0 e^{\frac{-(G\gamma_0\alpha\sigma_\delta)^2}{2}}$$

Relative depolarization:

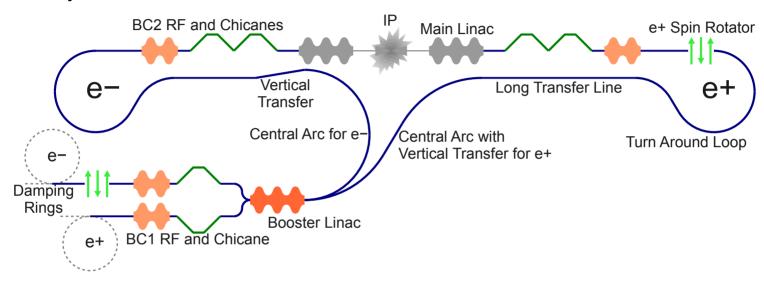
$$1 - \frac{\langle P_z \rangle}{P_0}$$

Where

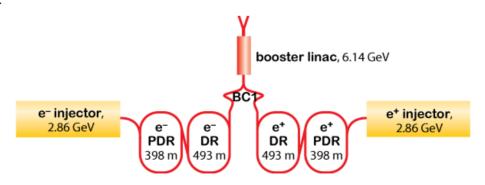
Symbol	Value	Description
G	0.00115965219	anomalous momentum of the electron
α	-	arc bending angle
γ_0	-	relativistic factor
σ_{δ}	_	energy spread

RTML Layout and Spin Rotator Location

New layout



Previous layout



Spin Precession and Depolarization in CLIC

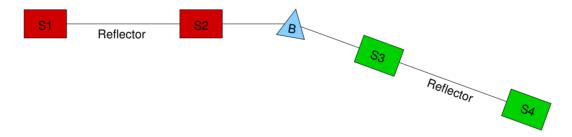
region	E_0 [GeV]	σ_{δ}	$lpha_{ m electrons}$ [rad]	$lpha_{ m positrons}$ [rad]
exit of damping rings to bc1	2.86	0.13%	0	0
exit of bc1 to booster	2.86	1.04%	0	0
exit of booster to bc2	9	0.33%	$\pi - \pi + HV\text{-doglegs} = 0$	π +HV-doglegs= π
exit of bc2 to bds	9	1.64%	0	0
exit of main linac to ip	1500	0.35%	$1\cdot 10^{-3}$	$1\cdot 10^{-3}$

region	E_0 [GeV]	σ_{δ}	$1 - \frac{\langle P_z \rangle}{P_0}$ [%]	$\phi_s = a \gamma_0 \alpha [deg]$	<i>n</i> -turns
exit of damping rings to bc1	2.86	0.13%	0	0	0
exit of bc1 to booster	2.86	1.04%	0	0	0
exit of booster to bc2 entrance	9	0.33%	0 / 2.2	$0 / 3676.4 \equiv 76.4$	0 / 10.2
exit of bc2 to bds	9	1.64%	0	0	0
exit of main linac to ip	1500	0.35%	0.007	195	0.54

- ⇒ From the point of view of the spin dynamics, ideal location for the spin rotators would probably be: before bc1 for the electrons, before bc2 for the positrons
- \Rightarrow Notice that, in case of a symmetric RTML where both spin rotators are placed before bc1 and assuming that the beam experiences a total bending angle $\alpha = \pi/2_{\mathrm{booster} \to \mathrm{bc2}}$ for each line, the total depolarization per beam is 0.56% per line. (with a precession of 5.1 n-turns)

Spin Rotator Lattice

- Spin Rotation is achieved by two solenoids with a bending magnet in between
- Each solenoid is split in two parts separated by a reflector $\begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix}$ to correct for couplings \Rightarrow there are four solenoids in total
- The central bending section must rotate the spin by 90 degrees
- This configuration allows arbitrary spin orientation



Description

- Reflector beamline : four FODO cells with 90 degrees phase advance in X and 45 degrees phase advance in Y
- Bend section: mini arc composed by three FODO cells with 90 degrees phase advance in X and Y (can be shortened)

Solenoid Strength

ullet Each of the four solenoids must be capable of providing a maximum of ± 45 degrees spin rotation

$$\psi_{\rm spin} = \pi/4$$
, with $\psi_{\rm beam} = \psi_{\rm spin}/2$

- Solenoid strength

$$k = \frac{\psi_{\text{spin}}}{2L} = \frac{B_z}{2(B_0 \rho)}$$

Assuming 2.6 meters long solenoids (like ILC)

$$k = \frac{\pi/4}{2} \frac{1}{(L = 2.6 \text{ m})} = 0.15104 \text{ m}^{-1}$$

⇒ The maximum longitudinal field is:

$$B_{z,max} = 2 \cdot k \cdot (B_0 \rho) = 2 \cdot k \cdot \frac{E_0}{ec} = 2 \cdot 0.15104 \text{ m}^{-1} \cdot \frac{E_0}{ec}$$

required magnetic field at 2.86 or 9 GeV is:

$$B_{z,max} @ 2.86 \text{ GeV} = \mathbf{2.9} \text{ T}$$

 $B_{z,max} @ 9 \text{ GeV} = \mathbf{9.1} \text{ T}$

Bending Arc

• The bending section should rotate the spin by 90 degrees

$$\phi_s = a \gamma_0 \alpha = \frac{\pi}{2}$$

$$\alpha @ 2.86 \text{ GeV} = \frac{\pi/2}{a (\gamma_0 = 2.86e3/0.511)} = 0.24202 \text{ rad} = 13.867 \text{ degrees}$$

$$\alpha @ 9 \text{ GeV} = \frac{\pi/2}{a (\gamma_0 = 9e3/0.511)} = 0.076908 \text{ rad} = 4.4065 \text{ degrees}$$

Magnetic strength:

$$B\rho @ 2.86 \text{ GeV} = \frac{pc}{ec} = \frac{2.86 \text{ GV}}{c} = \frac{2.86 \text{ GV}}{2.997925 \cdot 10^8 \text{ m/s}} = 9.5 \text{ T m}$$

$$B\rho @ 9 \text{ GeV} = \frac{pc}{ec} = \frac{9 \text{ GV}}{c} = \frac{9 \text{ GV}}{2.997925 \cdot 10^8 \text{ m/s}} = 30 \text{ T m}$$

Bending Magnets and Longitudinal Motion

• Assuming to be using 6, 1 meter long magnets, this corresponds to a bending radius

$$\rho$$
 @ 2.86 GeV = $\frac{L}{\alpha} = \frac{6 \cdot 1 \text{ m}}{0.24202 \text{ rad}} = 24.792 \text{ m}$

$$\rho$$
 @ 9 GeV = $\frac{L}{\alpha} = \frac{6 \cdot 1 \text{ m}}{0.076908 \text{ rad}} = 78.015 \text{ m}$

⇒ Magnetic field

$$B @ 2.86 \text{ GeV} = \frac{9.5 \text{ T m}}{24.792 \text{ m}} = 0.38319 \text{ T}$$

$$B @ 9 \text{ GeV} = \frac{30 \text{ T m}}{78.015 \text{ m}} = 0.38454 \text{ T}$$

 $\Rightarrow R_{56}$ for the bending section is:

$$R_{56} @ 2.86 \text{ GeV} = 60.0 \text{ mm}$$

 $R_{56} @ 9 \text{ GeV} = 6.0 \text{ mm}$

ISR-Induced Emittance Growth

The effect of incoherent synchrotron radiation (ISR) emission on the emittance growth can be estimated using

$$\Delta \gamma \epsilon = 4 \times 10^{-8} E^6 \text{ [GeV] } I_5 \text{ [m}^{-1}]$$

where

$$I_5 = \frac{4L}{|\rho|^3} \cdot \frac{\eta^2 + (\eta\alpha + \eta'\beta)^2}{\beta}$$

 \Rightarrow Case of E=2.86 GeV: using L=1 m, ρ = 24.8 m, average dispersion and its derivative η = 0.3 m and η' =0.15 rad, horizontal twiss β =22.5 m and α = \pm 3.5, and horizontal emittance $\gamma\epsilon$ = 0.68 μ m:

$$\frac{\Delta \gamma \epsilon}{\gamma \epsilon} = 0.7\%$$

 \Rightarrow Case of E=9 GeV: using L=1 m, $\rho=78.0$ m, average dispersion and its derivative $\eta=0.1$ m and $\eta'=0.05$ rad, horizontal twiss $\beta=22.5$ m and $\alpha=\pm3.5$, and horizontal emittance $\gamma\epsilon=0.68$ μ m:

$$\frac{\Delta \gamma \epsilon}{\gamma \epsilon} = 0.003\%$$

Spin Rotator and Bunch Compressor

- P. Emma, 1994: "the rotator system has very little impact on the performance of the bunch compressor"
- Longitudinal transfer matrix of the bunch compressor

$$R_{\rm BC} = \left(\begin{array}{cc} 1 + fR_{56} & R_{56} \\ f & 1 \end{array}\right)$$

- In case of full compression, ie. $1+fR_{56}=0$, adding the spin rotator changes the total transfer as follows

$$R_{\mathrm{BC}} \cdot R_{\mathrm{ROT}} = \begin{pmatrix} 1 + fR_{56} & R_{56} \\ f & 1 \end{pmatrix} \cdot \begin{pmatrix} 1 & \alpha \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 0 & R_{56} \\ f & 1 + \alpha f \end{pmatrix}$$

 \Rightarrow Bunch length after compression is unchanged by the rotator and the energy spread after compression is smaller ($f=2~\text{m}^{-1}$, $\alpha=-0.04~\text{m}$):

$$\sigma_{z,f} = \sigma_{\delta,i} R_{56}, \qquad \sigma_{\delta,f} = \sqrt{\sigma_{z,i}^2 f^2 + \sigma_{\delta,i}^2 \left(1 + \alpha f\right)}$$

- In our case, as bc1 does not fully compress,

$$R_{\text{BC}} \cdot R_{\text{ROT}} = \begin{pmatrix} 1 + fR_{56} & R_{56} + \alpha \left(1 + fR_{56}\right) \\ f & 1 + \alpha f \end{pmatrix}$$

⇒ Rotator might have an impact on the compression factor

$$\sigma_{z,f} = \sigma_{\delta,i} \left[R_{56} + \alpha \left(1 + f R_{56} \right) \right]$$
$$\sigma_{\delta,f} = \sqrt{\sigma_{z,i}^2 f^2 + \sigma_{\delta,i}^2 \left(1 + \alpha f \right)}$$

Notice that if $\alpha f < 0$ the final energy spread gets reduced

 \Rightarrow This issue can be overcome using an isochronous arc.

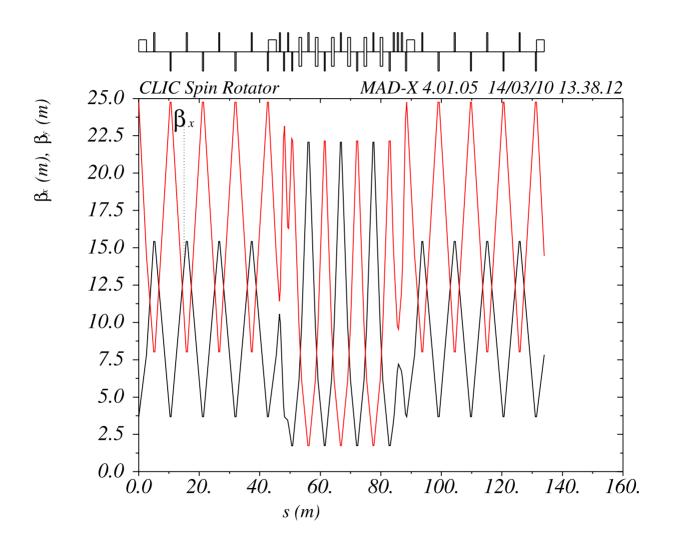
Summary Table and Conclusions

Relevant parameters with the spin rotator location, for electrons and positrons:

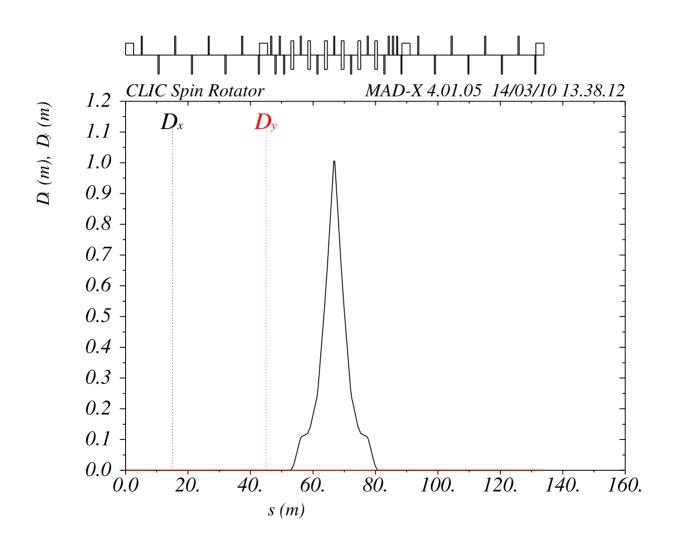
quantity	before bc1(*)	before bc2	symm.rtml	unit	remarks
beam energy	2.86	9	2.86	GeV	
bending angle	$0 (\pi)$	0	$\pi/2$	rad	
spin depolarization	0 (2.2)	0	0.56	%	bds excluded
spin precession	0 (10.2)	0	5.1	turns	11 11
solenoid field	2.9	9.1	like (*)	Т	L=2.6 m
bending angle	13.9	4.4	like (*)	deg	$L{=}1~\mathrm{m}$
bending magnet	0.38	0.38	like (*)	Т	" "
R_{56}	60.0	6.0	like (*)	mm	
$\Delta\gamma\epsilon_x$ by synrad emission	0.7	0.003	like (*)	%	negligible
total length	134.0	longer	like (*)	m	scales with the energy

- ⇒ New RTML layout: potential problem might be the large solenoid field for the positrons; positron spin rotator before bc2 would be longer; positron spin rotator before bc1: 2.2% depolarization seems to me negligible
- ⇒ Old RTML layout (symmetric): no major problems, negligible depolarization
 - Detailed beam dynamics studies have to be carried out
 - ullet Impact of R_{56} on the bunch compressor must be evaluated / use of an isochronous arc

Spin Rotator Optics



Spin Rotator Optics



Spin Rotator Optics

