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# Particle-In-Cell & Molecular Dynamics MODELLING OF VACUUM ARCS

## VACUUM ARCS

Vacuum discharges occur in a wide range of modern technology, either utilised in a controlled way, as in electrical discharge machining, or as an undesirable phenomenon from fusion reactors over satellite systems to future linear collider components. Within the CLIC project we now coordinate experiments and theory to explain vacuum arcs.



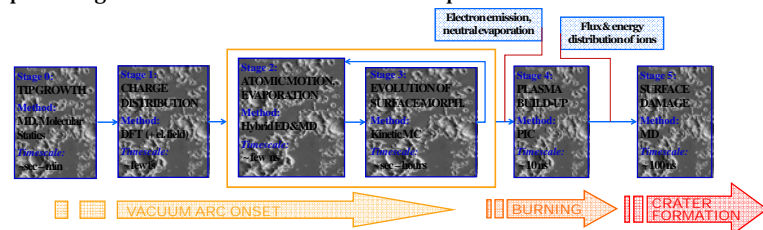
Vacuum arcs decrease the efficiency of CLIC accelerating structures (left).

The DC spark setup at CERN (right) serves to explore the properties of vacuum arcs in well-defined conditions.



## MULTISCALE MODEL

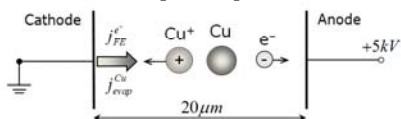
Our group at the Helsinki Institute of Physics is using a multiscale approach to model vacuum arcs. We simulate separately the triggering of the breakdown ("onset" phase), the evolution of the plasma ("burning" phase), and the resulting surface damage ("crater formation" phase), providing a feedback between the different phases.



The three phases of a vacuum arc: Onset, burning and crater formation.

## A PLASMA FORMS

We used a 1d3v electrostatic Particle-in-Cell (PIC) code with Monte Carlo collision scheme (MCC) to simulate the early stage of plasma build-up in a DC discharge between two infinite electrodes. Three species have been taken into account:  $e^-$ , Cu, and  $\text{Cu}^+$ . Simulations assume the same conditions (voltage, gap, pulse energy) as there are in DC spark experiments at CERN.

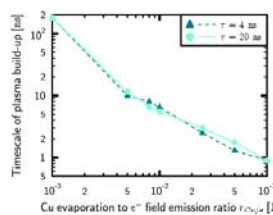


PIC simulated system. Starting from electron field emission and neutral evaporation, and taking into account impact ionisation, a plasma builds up in-between the electrodes.

As an initial condition, a field emitter with a given field enhancement factor is assumed to be present at the anode. The electron field emission (FE) current follows the Fowler-Nordheim equation.

$$j_{FE}(E) = a_{FN} \frac{(e \cdot \beta \cdot E)^2}{\phi \cdot t(y)^2} \exp\left(-b_{FN} \frac{\phi^{3/2} \cdot v(y)}{e \cdot \beta \cdot E}\right), \quad \text{where } y \propto \frac{\beta \cdot E}{\phi^2}$$

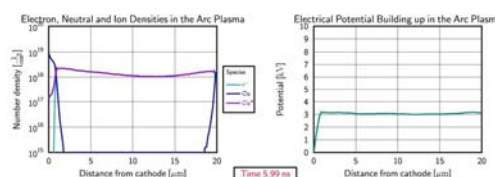
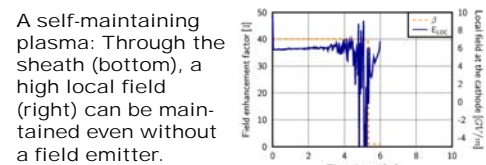
A field assisted thermal evaporation of neutrals from the field emitter following the FE current in a constant ratio  $r_{\text{Cu}/e}$  is assumed too.



How the Cu evaporation to  $e^-$  FE ratio influences the timescale of plasma build-up.

## BUILD-UP CRITERIA

Two criteria for breakdown can be formulated with the aid of our model: (i) a high enough *initial local field* – acting at a field emitter tip – to ensure growing FE current despite of space charge effects and (ii) reaching a *critical neutral density* – determined by the ionisation cross section and system length – to induce an ionisation avalanche.

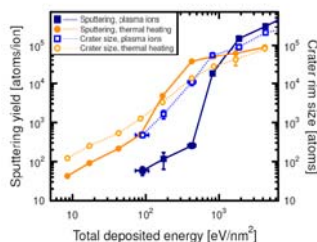


[H. Timko et al., *A One-Dimensional Particle-in-Cell Model of Plasma Build-up in Vacuum Arcs*, *Contrib. Plasma Phys.* (accepted 2010)]

## SURFACE DAMAGE

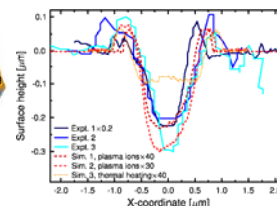
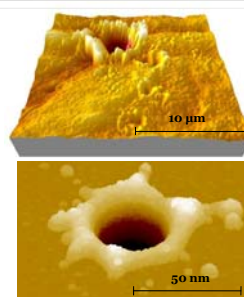
Molecular Dynamics (MD) simulations have been carried out to determine the surface damage caused by arcs. A perfect (100) Cu surface has been bombarded with  $\text{Cu}^+$ . The flux and the energy distribution of incident ions have been previously calculated with PIC, under conditions typical for the DC spark setup.

Simulations using different doses of incident ions showed that above a given threshold, the whole volume into which energy is deposited, gets melted and an enhanced sputtering starts, dominated by cluster emission.



Sputtering yield and crater rim size as a function of deposited energy. The same energy has been deposited in two different ways: arc plasma (blue) and thermal deposition (orange).

[H. Timko et al., *Mechanism of surface modification in the plasma-surface interaction in electrical arcs*, *Phys. Rev. B* **81**, 184109 (2010)]



Comparison of experiment (upper left) and simulation (lower left). Although simulations are limited to smaller scales, the crater depth to width ratio (right) stays constant over several orders of magnitude, in experiments as well as in simulations.

