

Top pair and single top cross sections at Tevatron and LHC energies

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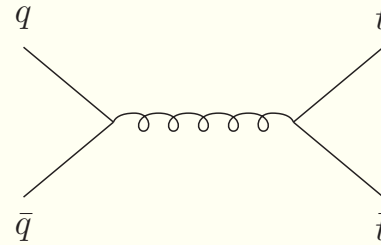
- $t\bar{t}$ and single top production channels
- Higher-order two-loop corrections
- $t\bar{t}$ cross section at Tevatron and LHC
- Top quark p_T and Y distribution at Tevatron and LHC
- t -channel production at Tevatron and LHC
- s -channel production at Tevatron and LHC
- Associated production of a top with a W^- or H^-

Partonic processes at LO

Top-antitop pair production

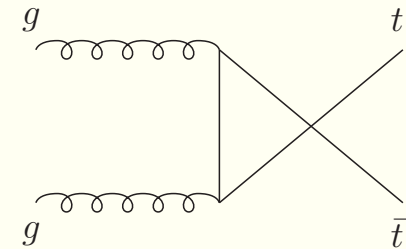
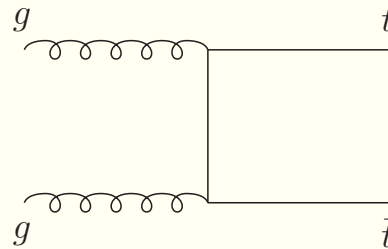
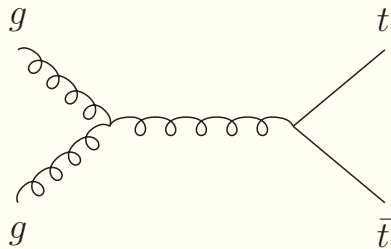
- $q\bar{q} \rightarrow t\bar{t}$

dominant at Tevatron



- $gg \rightarrow t\bar{t}$

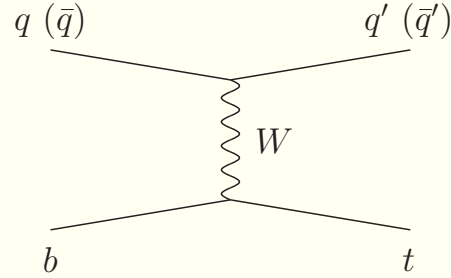
dominant at LHC



Single top quark production

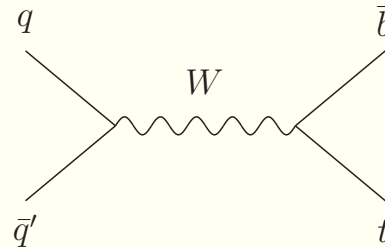
- **t channel:** $qb \rightarrow q't$ and $\bar{q}b \rightarrow \bar{q}'t$

dominant at Tevatron and LHC



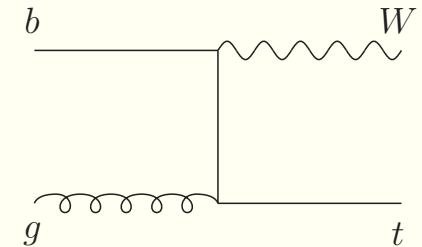
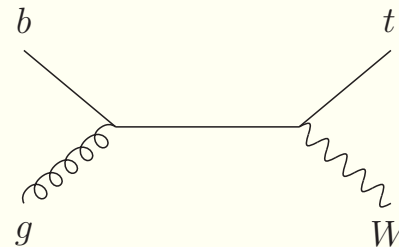
- **s channel:** $q\bar{q}' \rightarrow \bar{b}t$

small at Tevatron and LHC



- **associated tW production:** $bg \rightarrow tW^-$

very small at Tevatron, significant at LHC



- **Related process:** $bg \rightarrow tH^-$

Higher-order corrections

QCD corrections significant for top pair and single top quark production

NLO corrections fully known

Soft-gluon corrections from incomplete cancellations of infrared divergences between virtual diagrams and real diagrams with soft (low-energy) gluons

Soft corrections $\left[\frac{\ln^k(s_4/m^2)}{s_4} \right]_+$ with $k \leq 2n - 1$ and s_4 distance from threshold

Soft-gluon corrections are dominant near threshold

Resum (exponentiate) these soft corrections

At NLL accuracy requires one-loop calculations in the eikonal approximation

New results at NNLL—two-loop calculations completed

Approximate NNLO cross section from expansion of resummed cross section

Essential ingredient: two-loop soft anomalous dimension

Allows NNLL resummation

Resummed cross section

Resummation follows from factorization properties of the cross section

- performed in moment space

Use RGE to evolve function associated with soft-gluon emission

H: hard-scattering function

S: soft-gluon function

$$\hat{\sigma}^{res}(N) = \exp \left[\sum_i E_i(N) \right] \exp \left[\sum_j E'_j(N) \right] \text{tr} \{ H(\alpha_s) \} \\ \times \exp \left[\int_{\sqrt{s}}^{\sqrt{s}/\tilde{N}} \frac{d\mu}{\mu} \Gamma_S^+(\alpha_s(\mu)) \right] S \left(\alpha_s \left(\frac{\sqrt{s}}{\tilde{N}} \right) \right) \exp \left[\int_{\sqrt{s}}^{\sqrt{s}/\tilde{N}} \frac{d\mu}{\mu} \Gamma_S(\alpha_s(\mu)) \right]$$

where

Γ_S is the soft anomalous dimension - a matrix in color space

and a function of kinematical invariants s, t, u

Calculate Γ_S in eikonal approximation

Calculation is at differential cross section level

kinematics refer to partonic threshold (not just absolute threshold)

Eikonal approximation

Feynman rules for soft gluon emission simplify

$$\bar{u}(p) (-ig_s T_F^c) \gamma^\mu \frac{i(\not{p} + \not{k} + m)}{(p+k)^2 - m^2 + i\epsilon} \rightarrow \bar{u}(p) g_s T_F^c \gamma^\mu \frac{\not{p} + m}{2p \cdot k + i\epsilon} = \bar{u}(p) g_s T_F^c \frac{v^\mu}{v \cdot k + i\epsilon}$$

with $p \propto v$, T_F^c generators of SU(3)

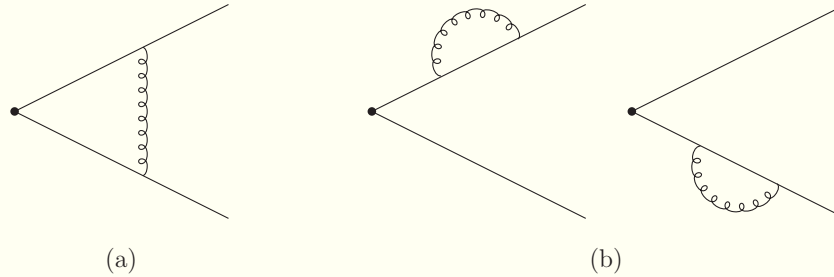
Perform calculation in momentum space and Feynman gauge

Complete two-loop results for

- soft (cusp) anomalous dimension for $e^+ e^- \rightarrow t \bar{t}$
- $t \bar{t}$ hadroproduction
- t -channel single top production
- s -channel single top production
- $bg \rightarrow tW^-$ and $bg \rightarrow tH^-$

Soft (cusp) anomalous dimension

One-loop eikonal diagrams



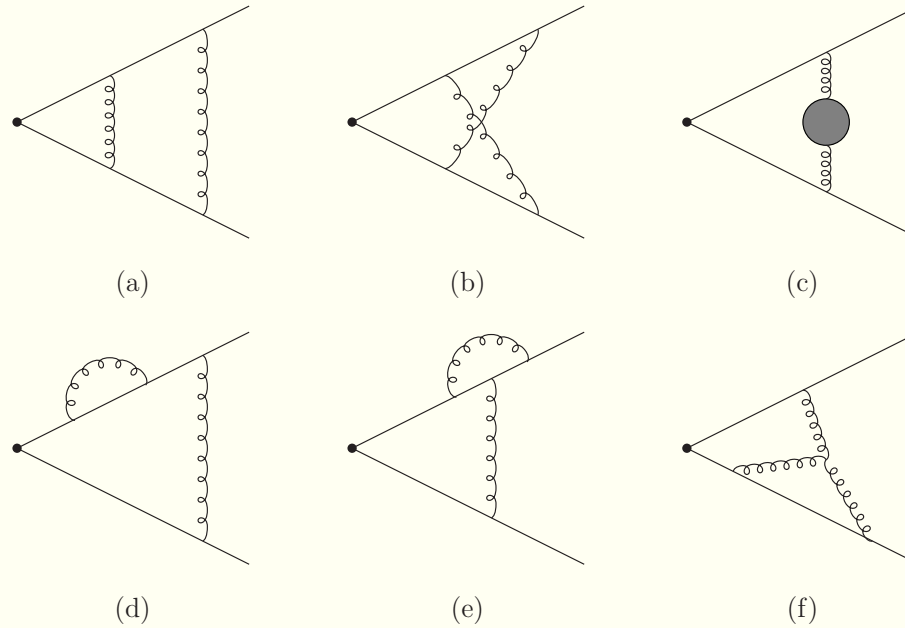
$$\Gamma_S = \frac{\alpha_s}{\pi} \Gamma_S^{(1)} + \frac{\alpha_s^2}{\pi^2} \Gamma_S^{(2)} + \dots$$

The one-loop soft anomalous dimension, $\Gamma_S^{(1)}$, can be read off the coefficient of the ultraviolet (UV) pole of the one-loop diagrams

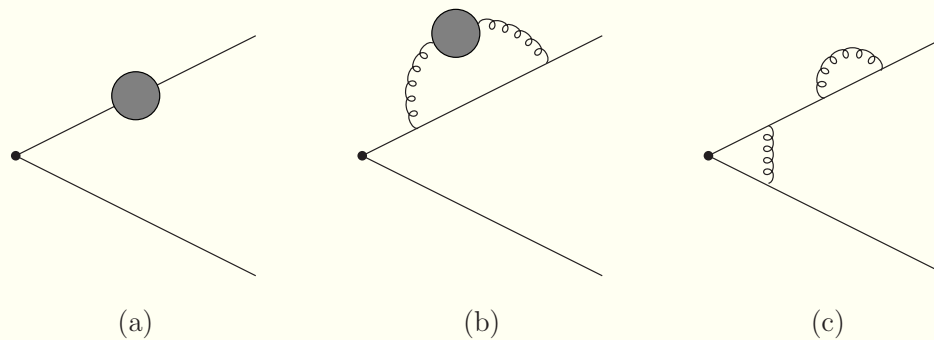
$$\Gamma_S^{(1)} = C_F \left[-\frac{(1+\beta^2)}{2\beta} \ln \left(\frac{1-\beta}{1+\beta} \right) - 1 \right] \quad \text{with} \quad \beta = \sqrt{1 - \frac{4m^2}{s}}$$

Two-loop eikonal diagrams

Vertex correction graphs



Heavy-quark self-energy graphs



Include counterterms for all graphs and multiply with corresponding color factors

Determine two-loop soft anomalous dimension from UV poles of the sum of the graphs

$$\Gamma_S^{(2)} = \frac{K}{2} \Gamma_S^{(1)} + C_F C_A M_\beta = \frac{K}{2} \Gamma_S^{(1)} + C_F C_A \left\{ \frac{1}{2} + \frac{\zeta_2}{2} + \frac{1}{2} \ln^2 \left(\frac{1-\beta}{1+\beta} \right) \right. \\ \left. - \frac{(1+\beta^2)^2}{8\beta^2} \left[\zeta_3 + \zeta_2 \ln \left(\frac{1-\beta}{1+\beta} \right) + \frac{1}{3} \ln^3 \left(\frac{1-\beta}{1+\beta} \right) + \ln \left(\frac{1-\beta}{1+\beta} \right) \text{Li}_2 \left(\frac{(1-\beta)^2}{(1+\beta)^2} \right) - \text{Li}_3 \left(\frac{(1-\beta)^2}{(1+\beta)^2} \right) \right] \right. \\ \left. - \frac{(1+\beta^2)}{4\beta} \left[\zeta_2 - \zeta_2 \ln \left(\frac{1-\beta}{1+\beta} \right) + \ln^2 \left(\frac{1-\beta}{1+\beta} \right) - \frac{1}{3} \ln^3 \left(\frac{1-\beta}{1+\beta} \right) + 2 \ln \left(\frac{1-\beta}{1+\beta} \right) \ln \left(\frac{(1+\beta)^2}{4\beta} \right) \right. \right. \\ \left. \left. - \text{Li}_2 \left(\frac{(1-\beta)^2}{(1+\beta)^2} \right) \right] \right\}$$

where $K = C_A(67/18 - \zeta_2) - 5n_f/9$

N. Kidonakis, Phys. Rev. Lett. 102, 232003 (2009), arXiv:0903.2561 [hep-ph]

$\Gamma_S^{(2)}$ vanishes at $\beta = 0$, the threshold limit, and diverges at $\beta = 1$, the massless limit

If one quark is massless and one is massive

$$\Gamma_S^{(2)} = \frac{K}{2} \Gamma_S^{(1)} + C_F C_A \frac{(1-\zeta_3)}{4}$$

QCD processes: Color structure gets more complicated with more than two colored partons in the process - Cusp anomalous dimension an essential component of other calculations

Top-antitop production in hadron colliders

The soft anomalous dimension matrix for $q\bar{q} \rightarrow t\bar{t}$ is

$$\Gamma_{S q\bar{q}} = \begin{bmatrix} \Gamma_{q\bar{q} 11} & \Gamma_{q\bar{q} 12} \\ \Gamma_{q\bar{q} 21} & \Gamma_{q\bar{q} 22} \end{bmatrix}$$

At one loop

$$\begin{aligned} \Gamma_{q\bar{q} 11}^{(1)} &= -C_F [L_\beta + 1] & \Gamma_{q\bar{q} 21}^{(1)} &= 2 \ln \left(\frac{u_1}{t_1} \right) & \Gamma_{q\bar{q} 12}^{(1)} &= \frac{C_F}{C_A} \ln \left(\frac{u_1}{t_1} \right) \\ \Gamma_{q\bar{q} 22}^{(1)} &= C_F \left[4 \ln \left(\frac{u_1}{t_1} \right) - L_\beta - 1 \right] + \frac{C_A}{2} \left[-3 \ln \left(\frac{u_1}{t_1} \right) + \ln \left(\frac{t_1 u_1}{s m^2} \right) + L_\beta \right] \end{aligned}$$

where $L_\beta = \frac{1+\beta^2}{2\beta} \ln \left(\frac{1-\beta}{1+\beta} \right)$ with $\beta = \sqrt{1 - 4m^2/s}$

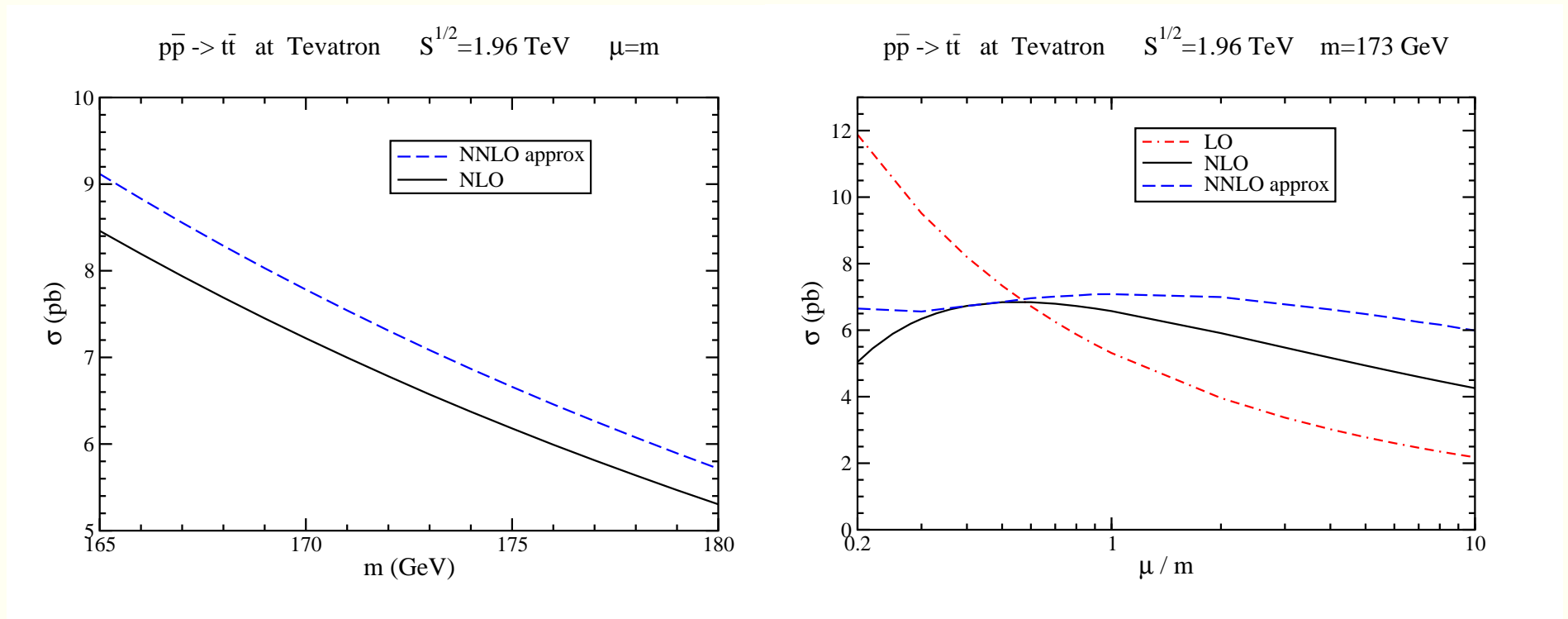
Write the two-loop cusp anomalous dimension as $\Gamma_S^{(2)} = \frac{K}{2} \Gamma_S^{(1)} + C_F C_A M_\beta$. Then at two loops

$$\begin{aligned} \Gamma_{q\bar{q} 11}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q} 11}^{(1)} + C_F C_A M_\beta & \Gamma_{q\bar{q} 22}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q} 22}^{(1)} + C_A \left(C_F - \frac{C_A}{2} \right) M_\beta \\ \Gamma_{q\bar{q} 21}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q} 21}^{(1)} + C_A N_\beta \ln \left(\frac{u_1}{t_1} \right) & \Gamma_{q\bar{q} 12}^{(2)} &= \frac{K}{2} \Gamma_{q\bar{q} 12}^{(1)} - \frac{C_F}{2} N_\beta \ln \left(\frac{u_1}{t_1} \right) \end{aligned}$$

with N_β a subset of terms of M_β

Similar results for $gg \rightarrow t\bar{t}$ channel N. Kidonakis, Phys. Rev. D 82, 114030 (2010), arXiv:1009.4935 [hep-ph]

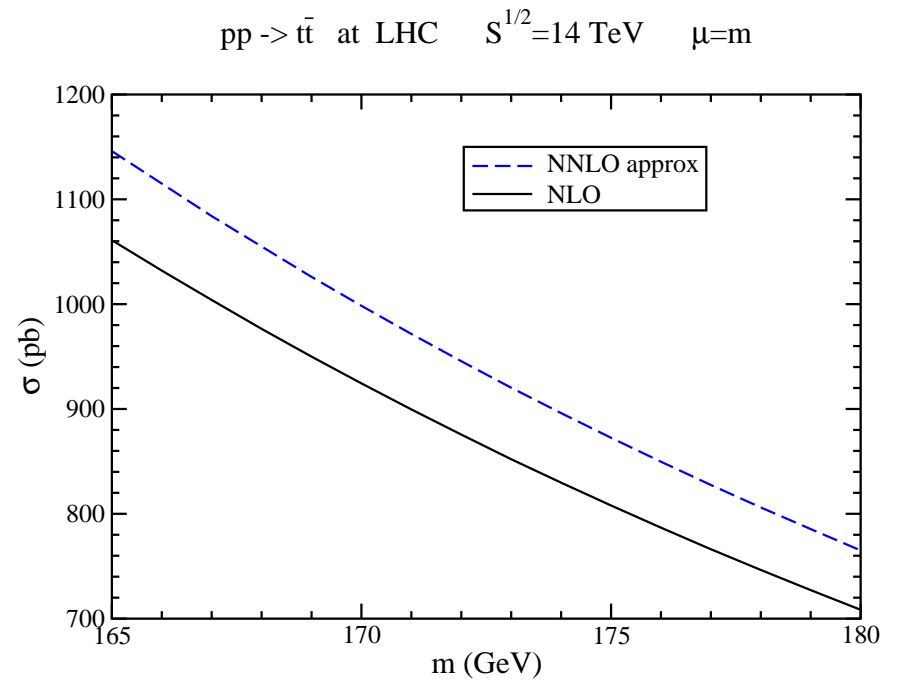
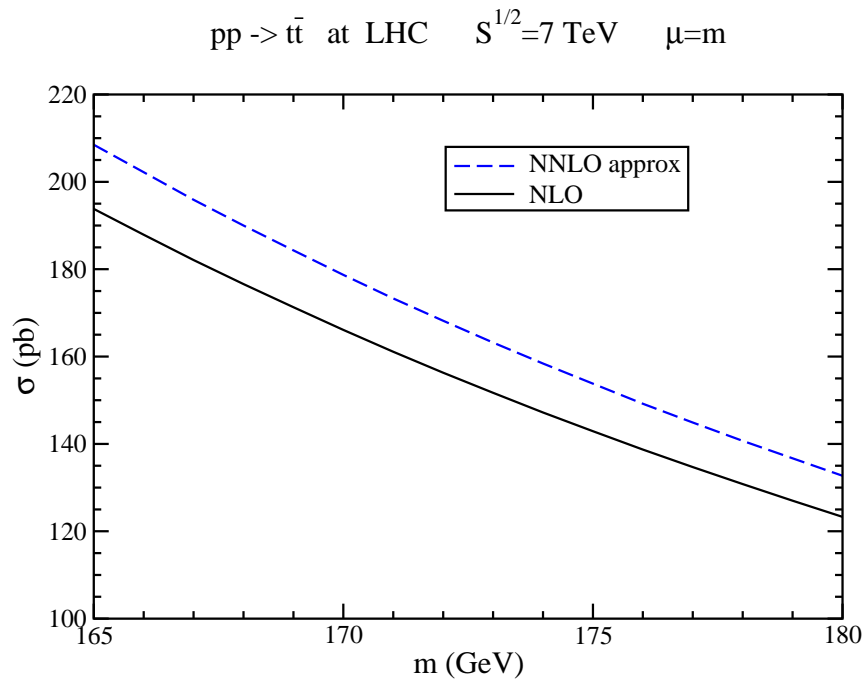
$t\bar{t}$ cross section at the Tevatron



$$\sigma_{t\bar{t}}^{\text{NNLOapprox}}(m_t = 173 \text{ GeV}, 1.96 \text{ TeV}) = 7.08^{+0.00+0.36}_{-0.24-0.27} \text{ pb}$$

NNLO approx: 7.8% enhancement over NLO

$t\bar{t}$ cross section at the LHC

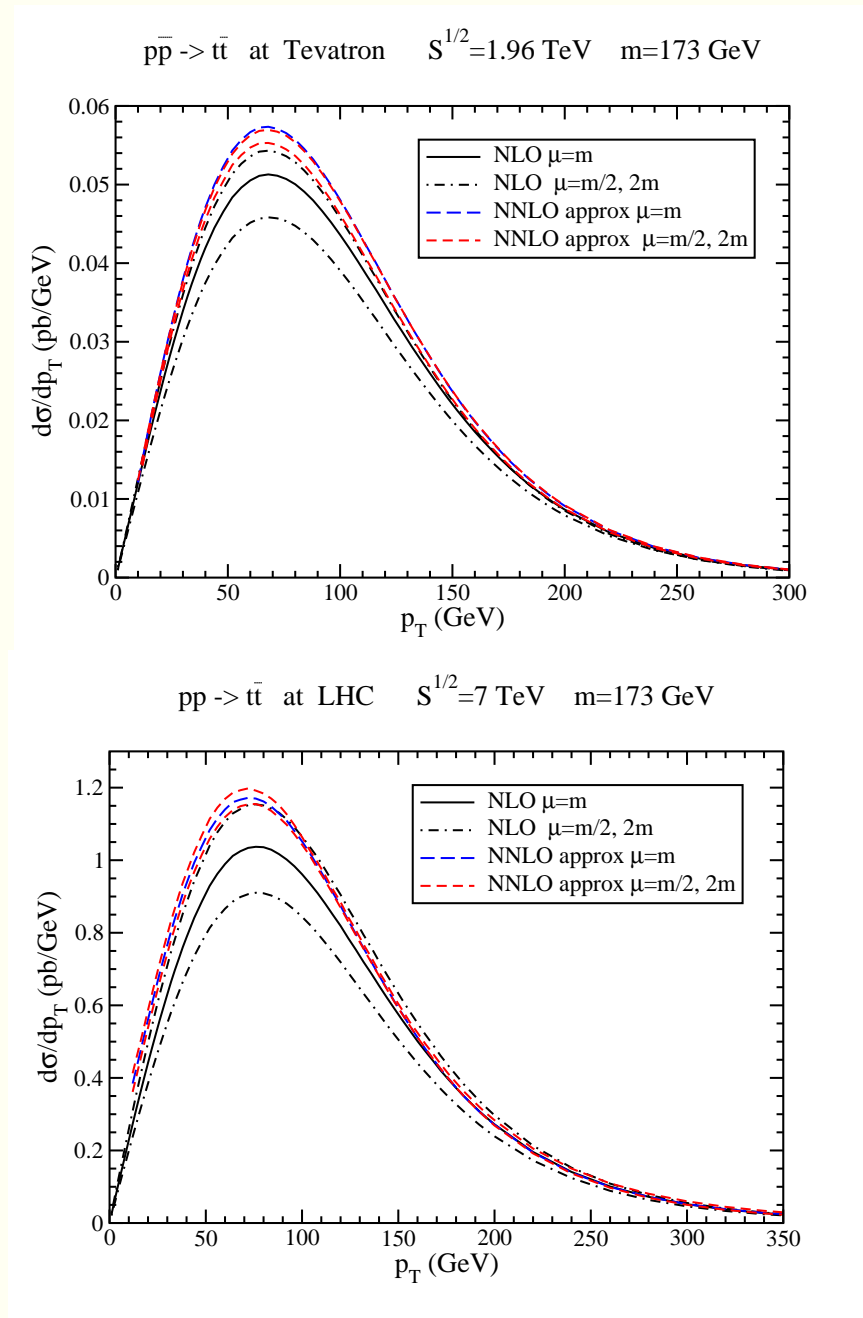


$$\sigma_{t\bar{t}}^{\text{NNLOapprox}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 163_{-5}^{+7+9} \text{ pb}$$

$$\sigma_{t\bar{t}}^{\text{NNLOapprox}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 920_{-39}^{+50+33} \text{ pb}$$

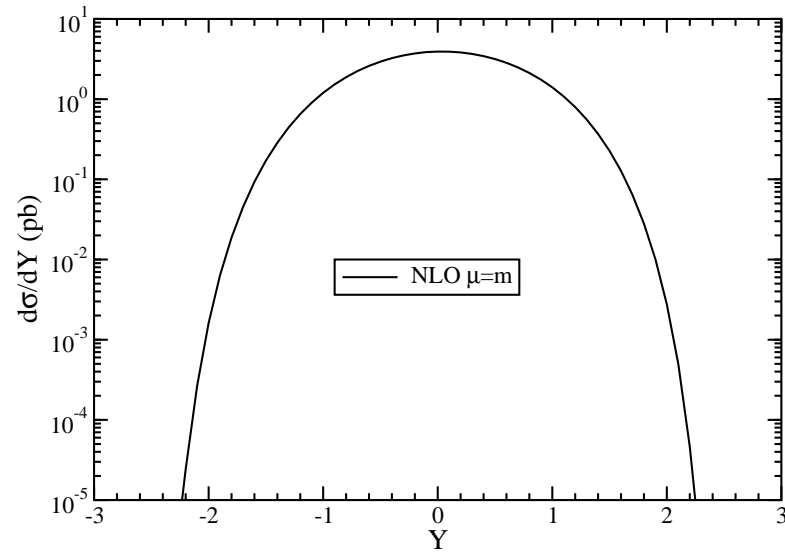
NNLO approx: enhancement over NLO is 7.6% at 7 TeV; 8.0% at 14 TeV

Top quark p_T distribution at Tevatron and LHC

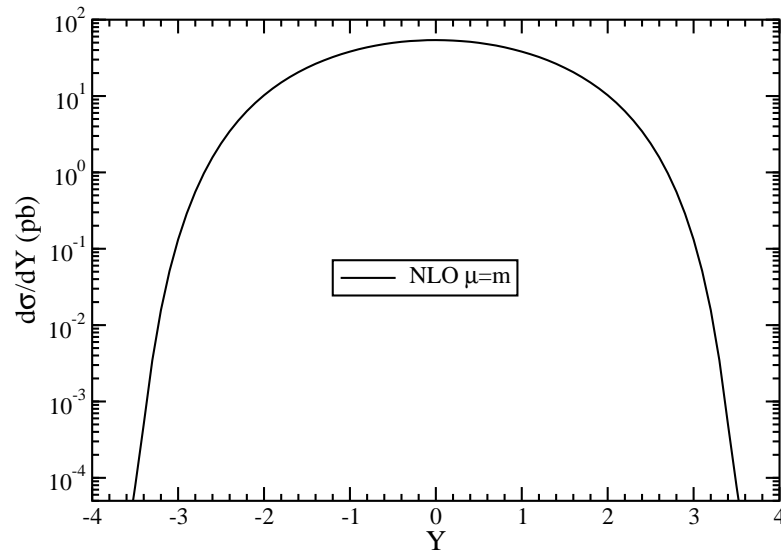


Top quark rapidity distribution at Tevatron and LHC

Top quark rapidity at Tevatron $S^{1/2}=1.96$ TeV $m=173$ GeV



Top quark rapidity at LHC $S^{1/2}=7$ TeV $m=173$ GeV



Single top quark production - t channel

Dominant single top production channel at both Tevatron and LHC energies

Soft anomalous dimension for t -channel single top production

One loop

$$\Gamma_{S11}^{(1)} = C_F \left[\ln \left(\frac{-t}{s} \right) + \ln \left(\frac{m_t^2 - t}{m_t \sqrt{s}} \right) - \frac{1}{2} \right]$$

$$\Gamma_{S21}^{(1)} = \ln \left(\frac{u(u - m_t^2)}{s(s - m_t^2)} \right)$$

$$\Gamma_{S12}^{(1)} = \frac{C_F}{2N_c} \Gamma_{S21}^{(1)}$$

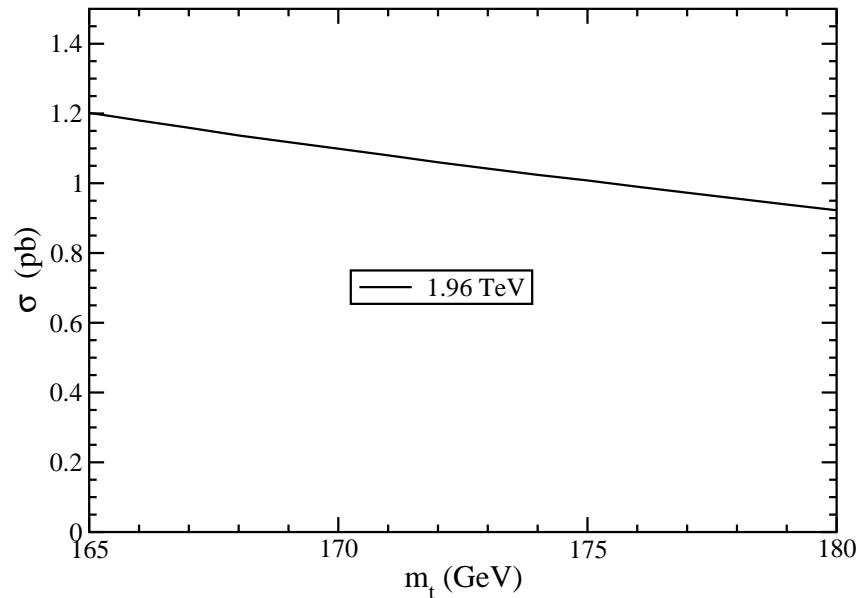
Two loops

$$\Gamma_{S11}^{(2)} = \frac{K}{2} \Gamma_{S11}^{(1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

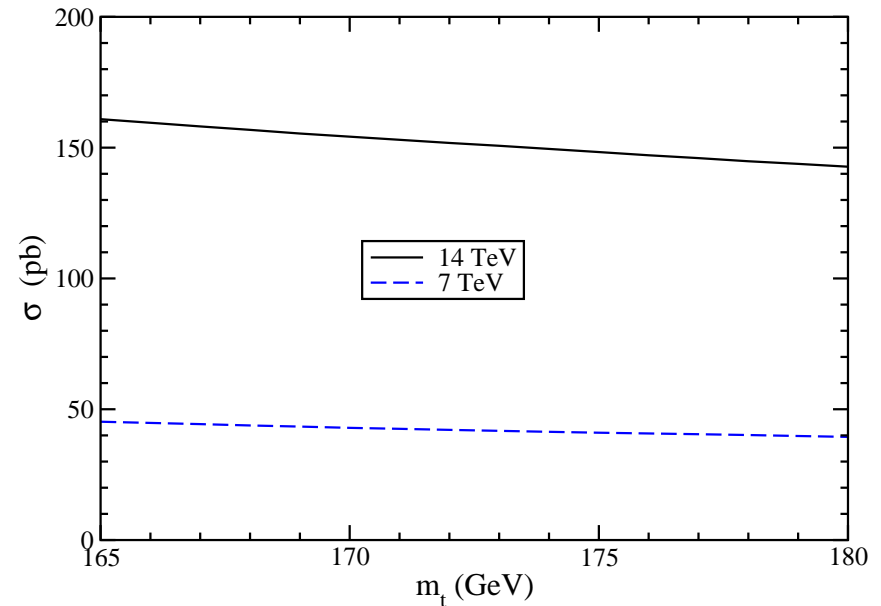
N. Kidonakis, arXiv:1103.2792 [hep-ph]

Single top quark production at Tevatron and LHC - t channel

Single top Tevatron t -channel NNLO approx (NNLL) $\mu=m_t$



Single top LHC t -channel NNLO approx (NNLL) $\mu=m_t$



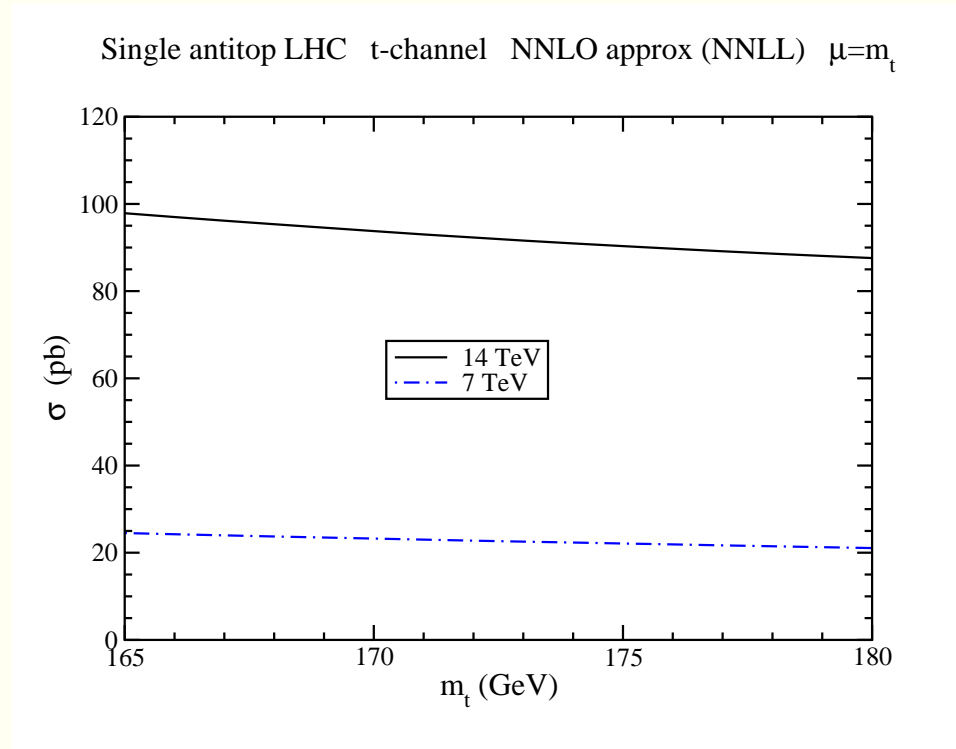
$$\sigma_{t\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 1.96 \text{ TeV}) = 1.04_{-0.02}^{+0.00} \pm 0.06 \text{ pb}$$

$$\sigma_{t\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 41.7_{-0.2}^{+1.6} \pm 0.8 \text{ pb}$$

$$\sigma_{t\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 151_{-1}^{+4} \pm 3 \text{ pb}$$

NNLO approx: 4% increase at Tevatron; 1% decrease at 7 TeV; 3% decrease at 14 TeV relative to NLO

Single antitop production at LHC - t channel



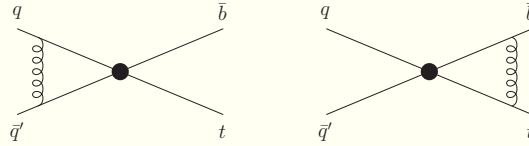
$$\sigma_{t\text{-channel}}^{\text{NNLOapprox, antitop}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 22.5 \pm 0.5_{-0.9}^{+0.7} \text{ pb}$$

$$\sigma_{t\text{-channel}}^{\text{NNLOapprox, antitop}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 92_{-1}^{+2} {}_{-3}^{+2} \text{ pb}$$

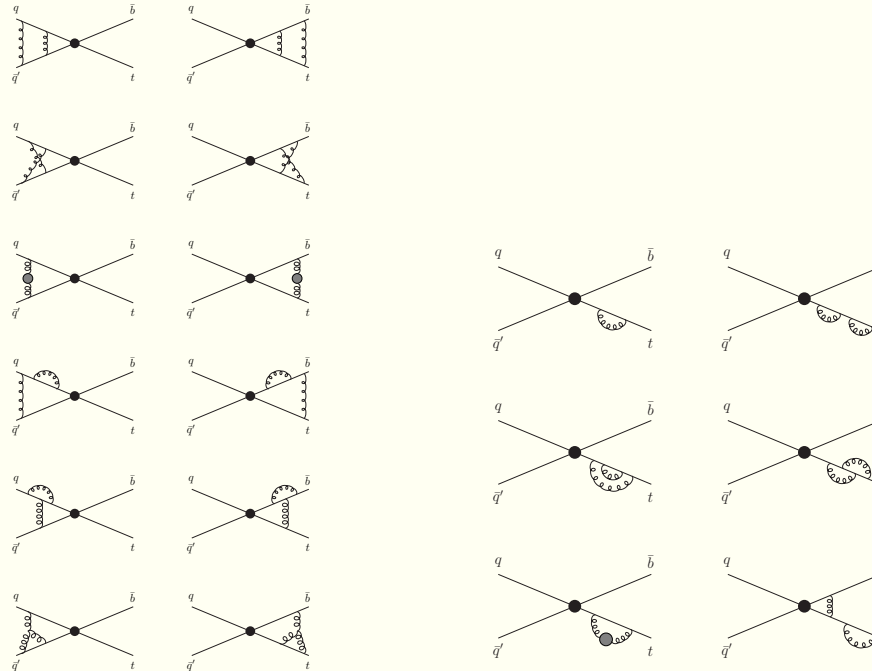
NNLO approx: 1% decrease at 7 TeV; 3% decrease at 14 TeV relative to NLO

Single top quark production - s channel

One-loop eikonal diagrams



Two-loop eikonal diagrams



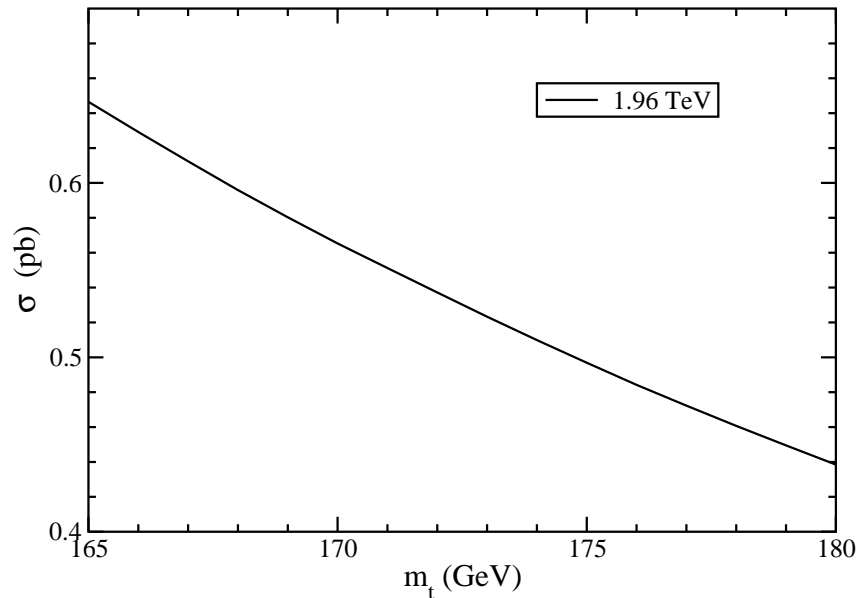
Soft anomalous dimension for s-channel single top production

$$\Gamma_{S11}^{(1)} = C_F \left[\ln \left(\frac{s - m_t^2}{m_t \sqrt{s}} \right) - \frac{1}{2} \right], \quad \Gamma_{S11}^{(2)} = \frac{K}{2} \Gamma_{S11}^{(1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

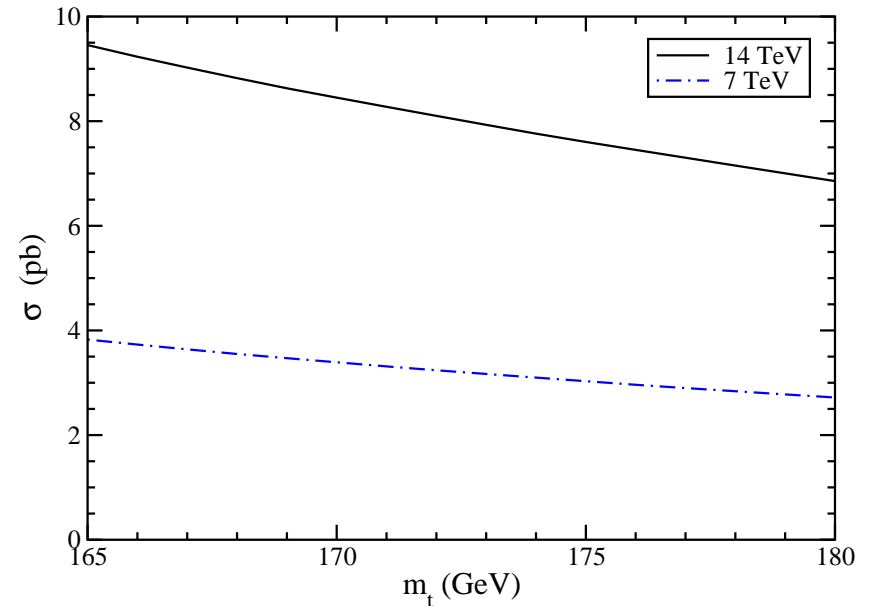
N. Kidonakis, Phys. Rev. D 81, 054028 (2010), arXiv:1001.5034 [hep-ph]

Single top quark production at Tevatron and LHC - s channel

Single top Tevatron s-channel NNLO approx (NNLL) $\mu=m_t$



Single top LHC s-channel NNLO approx (NNLL) $\mu=m_t$



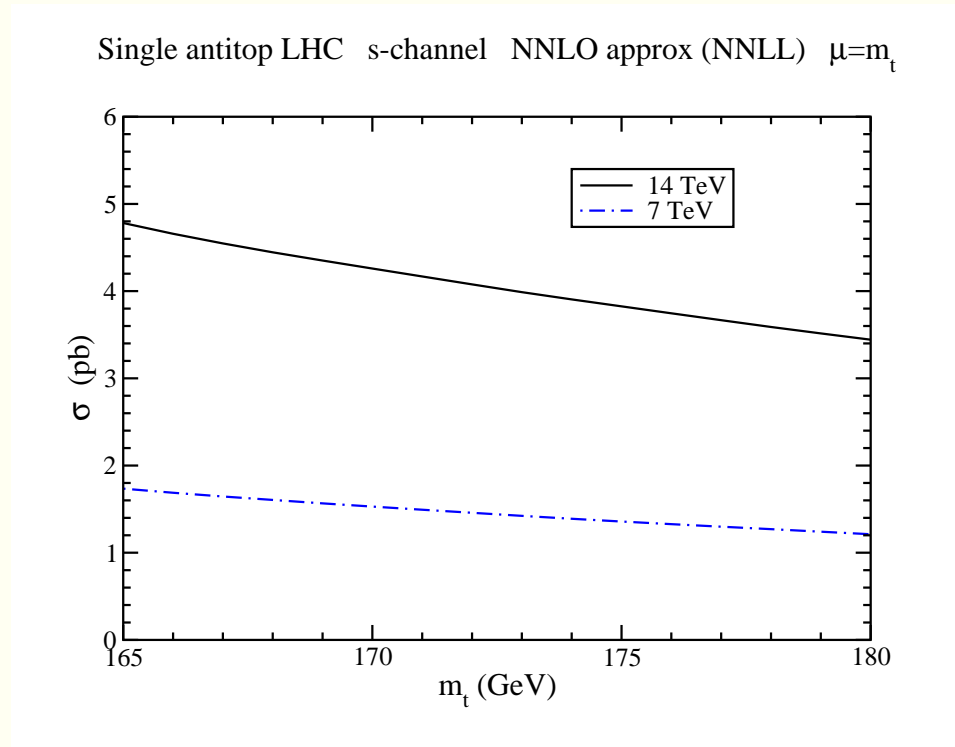
$$\sigma_{s\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 1.96 \text{ TeV}) = 0.523_{-0.005}^{+0.001+0.030} \text{ pb}$$

$$\sigma_{s\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 3.17 \pm 0.06_{-0.10}^{+0.13} \text{ pb}$$

$$\sigma_{s\text{-channel}}^{\text{NNLOapprox, top}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 7.93 \pm 0.14_{-0.28}^{+0.31} \text{ pb}$$

NNLO approx: enhancement over NLO is 15% at Tevatron; 13% at LHC

Single antitop production at LHC - s channel

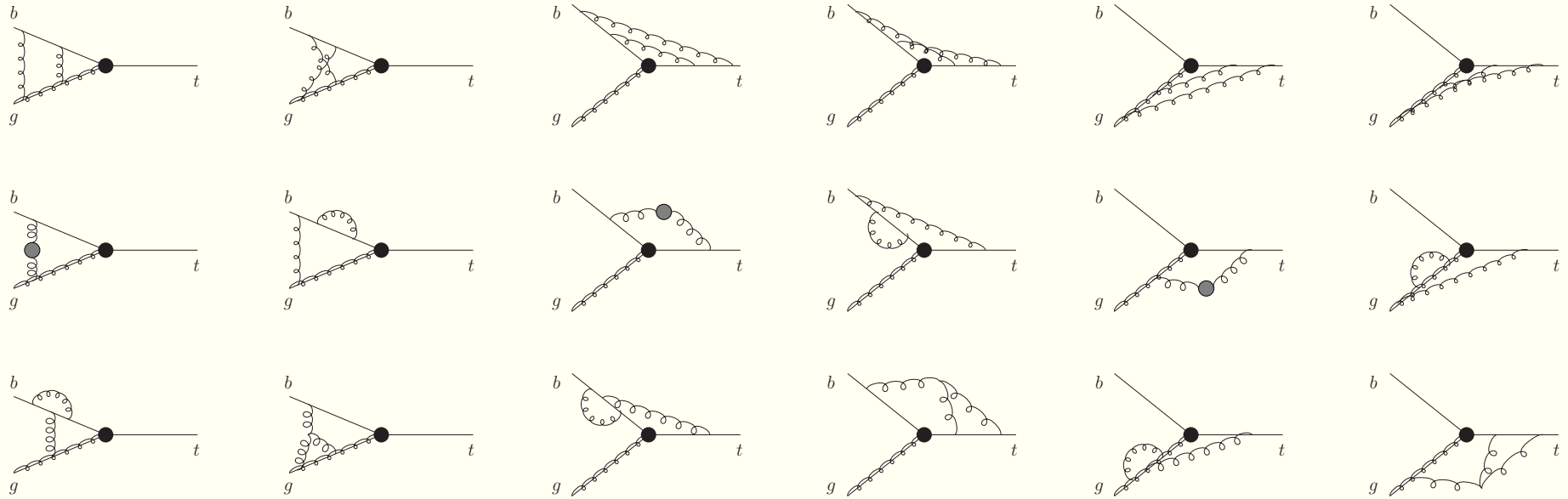


$$\sigma_{s\text{-channel}}^{\text{NNLOapprox, antitop}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 1.42 \pm 0.01_{-0.07}^{+0.06} \text{ pb}$$

$$\sigma_{s\text{-channel}}^{\text{NNLOapprox, antitop}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 3.99 \pm 0.05_{-0.21}^{+0.14} \text{ pb}$$

Associated production of a top quark with a W^-

Two-loop eikonal diagrams (+ extra top-quark self-energy graphs)

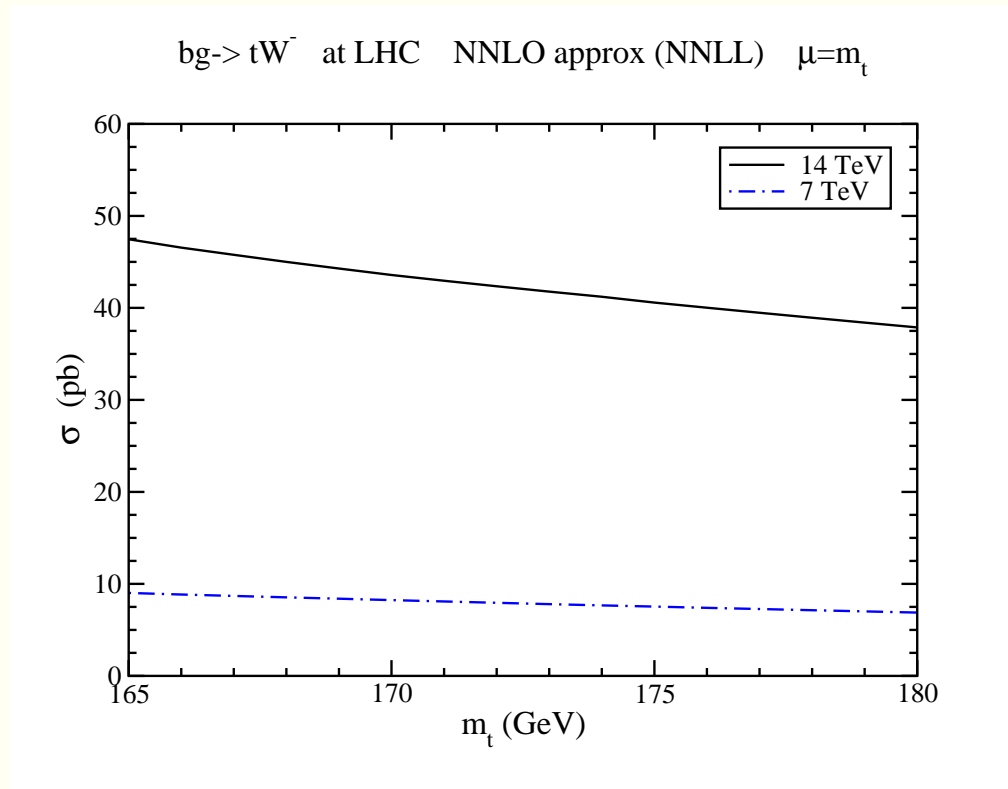


Soft anomalous dimension for $bg \rightarrow tW^-$

$$\Gamma_{S,tW^-}^{(1)} = C_F \left[\ln \left(\frac{m_t^2 - t}{m_t \sqrt{s}} \right) - \frac{1}{2} \right] + \frac{C_A}{2} \ln \left(\frac{m_t^2 - u}{m_t^2 - t} \right)$$

$$\Gamma_{S,tW^-}^{(2)} = \frac{K}{2} \Gamma_{S,tW^-}^{(1)} + C_F C_A \frac{(1 - \zeta_3)}{4}$$

Same analytical result for Γ_S for $bg \rightarrow tH^-$



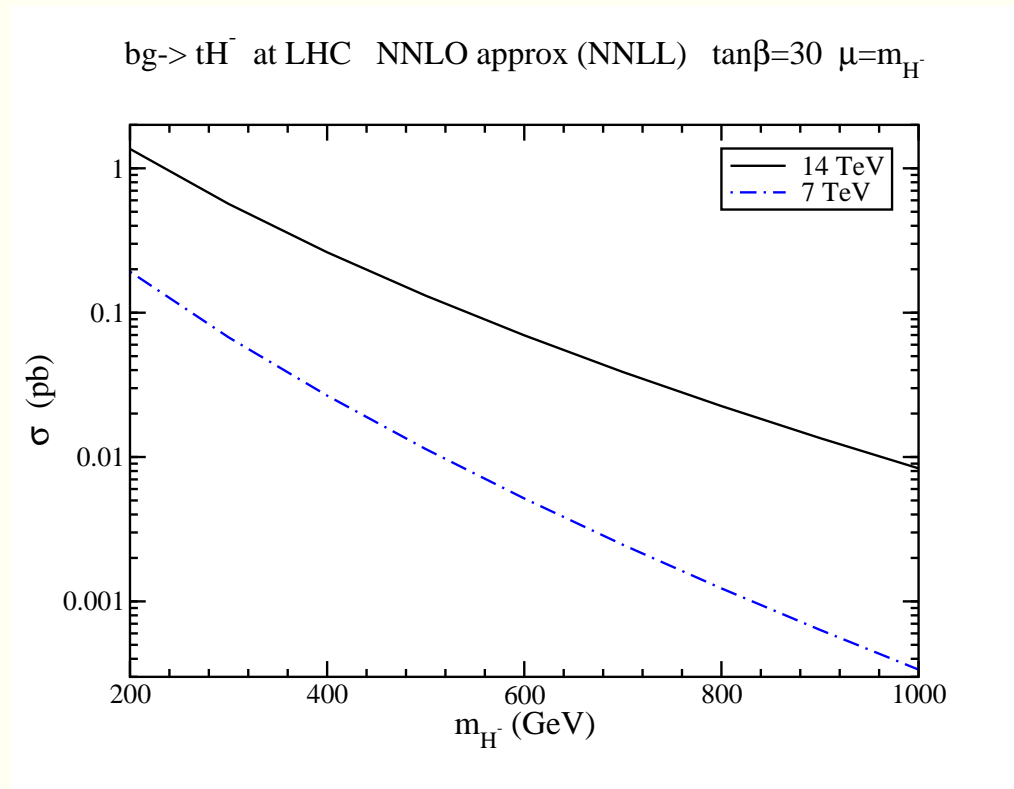
$$\sigma_{tW}^{\text{NNLOapprox}}(m_t = 173 \text{ GeV}, 7 \text{ TeV}) = 7.8 \pm 0.2_{-0.6}^{+0.5} \text{ pb}$$

$$\sigma_{tW}^{\text{NNLOapprox}}(m_t = 173 \text{ GeV}, 14 \text{ TeV}) = 41.8 \pm 1.0_{-2.4}^{+1.5} \text{ pb}$$

NNLO approx corrections increase NLO cross section by $\sim 8\%$

Cross section for $\bar{t}W$ production is identical

Associated production of a top quark with a charged Higgs



NNLO approx corrections increase NLO cross section by ~ 15 to $\sim 20\%$

Summary

- NNLL resummation for top quark pair and single top production
- $t\bar{t}$ production cross section
- top quark p_T and rapidity distributions
- t -channel and s -channel single top production cross section
- $bg \rightarrow tW^-$ and $bg \rightarrow tH^-$ at LHC
- NNLO approx corrections for top pair and single top production are significant at Tevatron and LHC