

Status $H \rightarrow \gamma\gamma$ and $H \rightarrow \mu\mu$

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Update on $H \rightarrow \mu\mu$ with improved analysis methods

- Evaluated $\sigma(M(\mu^+, \mu^-))$ event by event.
 - Distribution similar for signal and main background.
 - Receive suggestion by Fuji san to compare this variable with the expected Higgs mass resolution.
- Plan to simulate aa_4f samples (fast simulation on DBD analysis).
 - This contribution was not negligible on previous analysis.
 - I want to check it with realistic samples.
 - Old simulated samples already available but i found one bug (produced with same seed: duplicated).

$H \rightarrow \gamma\gamma$

- Writing note documentation analysis until now.
 - LCWS13 proceedings as subset of this note.

BACKUP

Weight on $\sigma(M(\mu^+, \mu^-))$

Mikael Suggestion

- Dimuon mass given by muon momentum measurement.
 - Variance on $M(\mu^+, \mu^-)$ \rightarrow error propagation from momentum uncertainty.
- Well measured signal events have small $M(\mu^+, \mu^-)$ errors.
- Badly measured background events lying around signal window would have larger $M(\mu^+, \mu^-)$ errors.

\rightarrow Variable $1/\sigma(M(\mu^+, \mu^-))$ could improve S/B.

$(M(\mu^+, \mu^-))$ depends simetrically on both muons)

$$\sigma^2(M(\mu^+, \mu^-)) = \frac{1}{M(\mu^+, \mu^-)^2} [P_1^T \Sigma_2 P_1 + P_2^T \Sigma_1 P_2]$$

(*) P_i column vector $x^i = (E, px, py, pz)$, P_i^T , row vector $x_i = g_{ij}x^j = (E, -px, -py, -pz)$

(g_{ij} metric in the Minkowski space: [+1 -1 -1 -1]), Σ_i covariance matrix of muon i.

There was a typo in this expression in previous meeting (Feb 7th 2014). I updated new slides fixing that typo.

Weight on $\sigma(\mathbf{M}(\mu^+, \mu^-))$

$$\sigma^2(\mathbf{M}(\mu^+, \mu^-)) = \frac{1}{\mathbf{M}(\mu^+, \mu^-)^2} [\mathbf{P}_1^T \Sigma_2 \mathbf{P}_1 + \mathbf{P}_2^T \Sigma_1 \mathbf{P}_2]$$

(*) P_i column vector $x^i = (E, px, py, pz)$, P_i^T , row vector $x_i = g_{ij}x^j = (E, -px, -py, -pz)$
(g_{ij} metric in the Minkowski space: [+1 -1 -1 -1]), Σ_i covariance matrix of muon i .

- I assumed $\text{Var}(m_i) = 0$ (m_i , muon mass).
- Even without mass constraint, at 1 TeV, i would expect $\lim \frac{\text{Var}(m)}{\text{Var}(p)} \approx 0$

ReconstructedParticle.getCovMatrix() returns null

- This covariance matrix is not filled (why not?).
 - Need to perform additional algebra.
 - Covariance matrix from associated track is available.
 - In helix parameterization basis: $d_0, z_0, \Omega, \phi, \tan\lambda$
 - If \mathfrak{J} is the Jacobian matrix from helix parameters to (px, py, pz, E)
 - $\Sigma'_i = \mathfrak{J} \Sigma_i \mathfrak{J}^T$, covariance matrix in momenta space.

Jacobian helix parameters to momenta space

- After some derivative exercises ...

$$\mathfrak{J} = \begin{bmatrix} \frac{\partial P_x}{\partial \tan \lambda} & \frac{\partial P_y}{\partial \tan \lambda} & \frac{\partial P_z}{\partial \tan \lambda} & \frac{\partial E}{\partial \tan \lambda} \\ \frac{\partial P_x}{\partial \Omega} & \frac{\partial P_y}{\partial \Omega} & \frac{\partial P_z}{\partial \Omega} & \frac{\partial E}{\partial \Omega} \\ \frac{\partial P_x}{\partial d_0} & \frac{\partial P_y}{\partial d_0} & \frac{\partial P_z}{\partial d_0} & \frac{\partial E}{\partial d_0} \\ \frac{\partial P_x}{\partial z_0} & \frac{\partial P_y}{\partial z_0} & \frac{\partial P_z}{\partial z_0} & \frac{\partial E}{\partial z_0} \\ \frac{\partial P_x}{\partial \phi} & \frac{\partial P_y}{\partial \phi} & \frac{\partial P_z}{\partial \phi} & \frac{\partial E}{\partial \phi} \end{bmatrix} = \frac{-1}{\Omega} \begin{bmatrix} 0 & 0 & -\Omega P_T & -\frac{P_z^2 \Omega}{E \tan \lambda} \\ P_x & P_y & P_z & \frac{P^2}{E} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ P_y \Omega & -P_x \Omega & 0 & 0 \end{bmatrix}$$

→ $\Sigma'_i = \mathfrak{J} \Sigma_i \mathfrak{J}^T$, covariance matrix in momenta space.

- Could be useful storing this covariance matrix for charged particles.
- so that, we could retrieve it by calling `ReconstructedParticle.getCovMatrix()`.
(Currently returning null matrix)