

Implementation of Covariance Matrix on ReconstructedParticle

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Motivation

- `ReconstructedParticle.getCovMatrix` is not implemented in current ILCSoft release (return 0, $\forall p \in \text{PandoraPFOs}$).
- This method provide covariance matrix of the reco. particle 4 vector {px,py,pz,E}.
- I was suggested to apply this cov. matrix to obtain dimuon mass error event-by-event basis.

- I have written Marlin processor adding new LCCollection to the event.
- This collection is a copy of PandoraPFOs but with filled cov. matrix.
- Code is available here: <http://www-jlc.kek.jp/jlc/en/node/209>

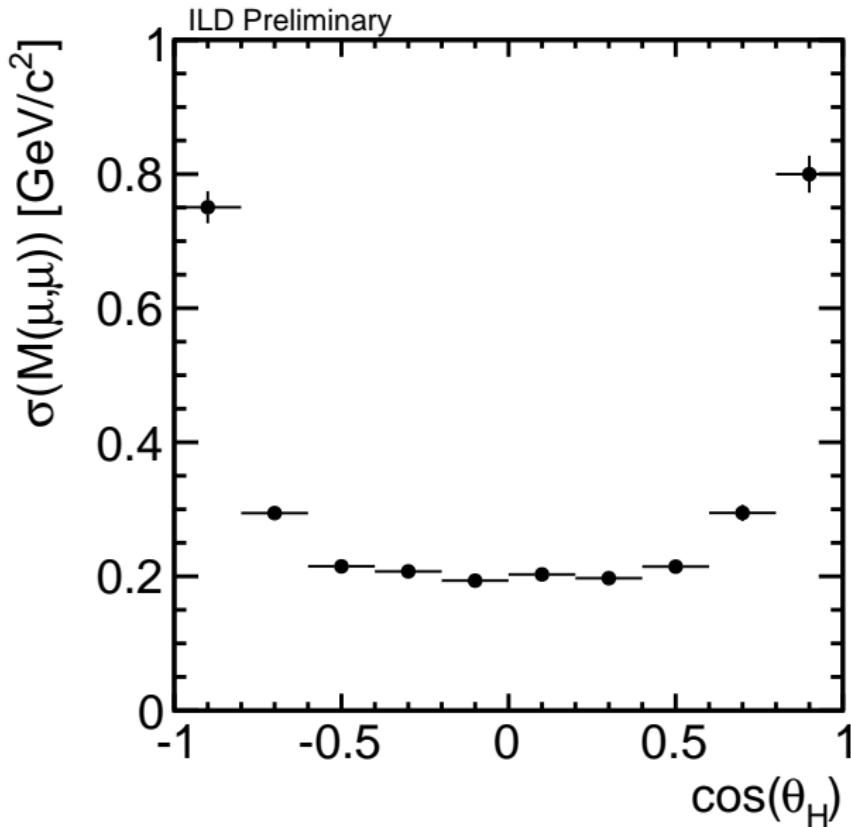
Calculation

$$\Sigma_i = \Sigma_i(\tan\lambda, \phi, \Omega), \Sigma'_i = \Sigma'_i(px, py, pz, E)$$

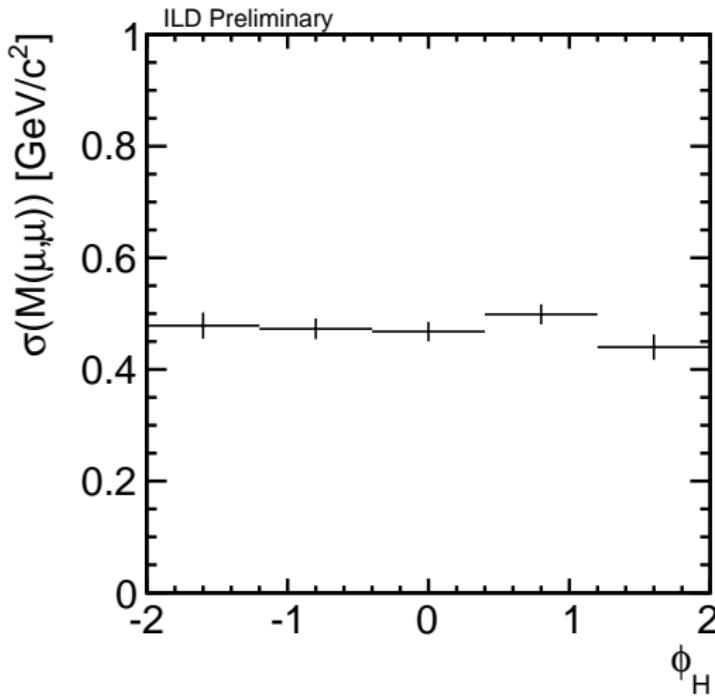
- Covariance matrix on helix parameters, Σ_i , from associated track.
- Obtain jacobian (\mathfrak{J}) and perform: $\Sigma'_i = \mathfrak{J}^T \Sigma_i \mathfrak{J}$

Dimuon invariant mass error

- Checked calculation with muons from $H \rightarrow \mu\mu$.
- The covariance matrix is used to obtain the dimuon invariant mass event-by-event.

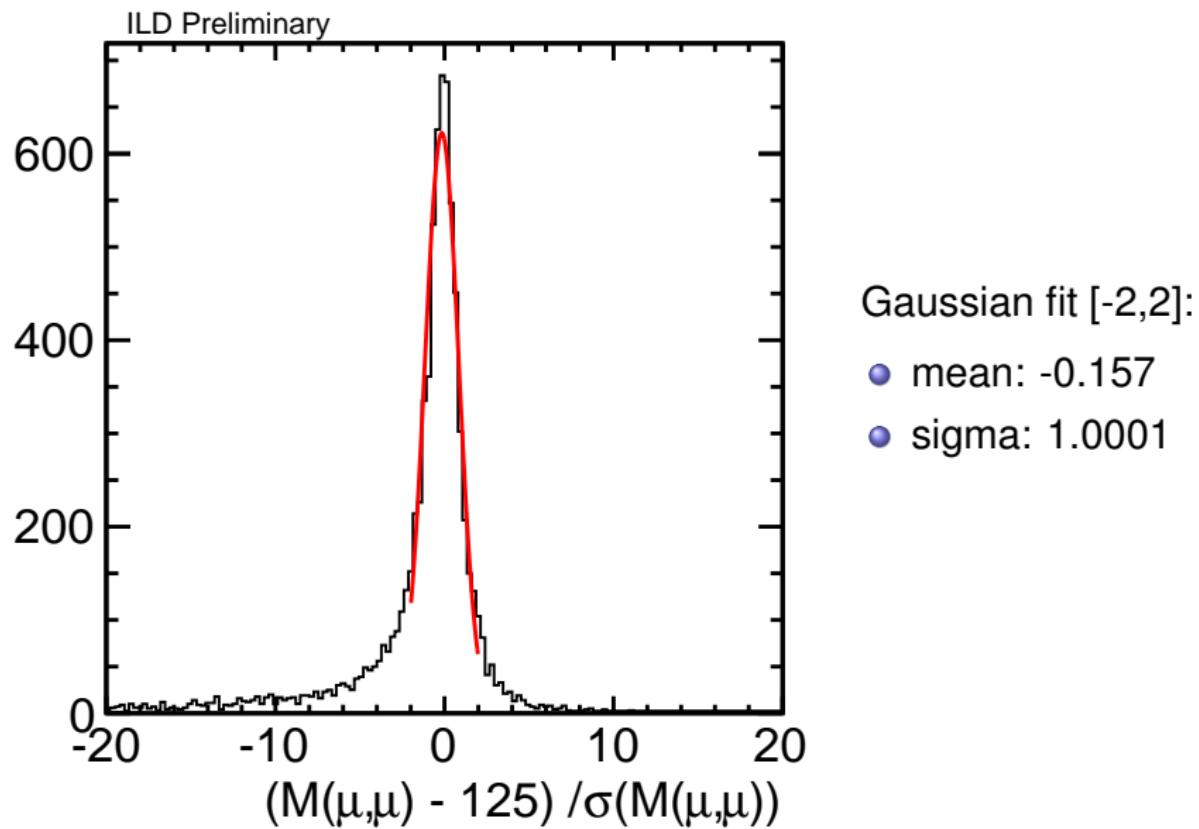


- Better precision at central region (tracks have more hits).



- No dependence on azimuthal angle.

Gaussian Fit in [-2,2]



Summary / Plan

Summary

- Current ILCsoft does not provide covariance matrix on $\{px, py, pz, E\}$ for the reco. particles.
- Developed Marlin processor calculating this matrix for charged particles.

Plan

- Use cov. matrix in update of $H \rightarrow \mu\mu$ analysis.

BACK UP

Relation between variables

- Original base: $\mathfrak{A} = \{\tan \lambda, \Omega, \phi, d_0, z_0\}$
- New base: $\mathfrak{B} = \{p_x, p_y, p_z, E\}$

$$p_x = p_T \cos \phi$$

$$p_y = p_T \sin \phi$$

$$p_z = p_T \tan \lambda$$

$$p_T = \left| \frac{\kappa}{\Omega} \right|$$

$$\kappa = |a B_z| \text{ (constant)}$$

$$\begin{aligned}E^2 &= (a \frac{B_z}{\Omega \cos \lambda})^2 + m^2 \\&= \left(\frac{p_T}{\cos \lambda} \right)^2 + m^2\end{aligned}$$

Momenta does not depend on d_0, z_0

- $p_x = p_x(\tan \lambda, \Omega, \phi)$
- $p_y = p_y(\tan \lambda, \Omega, \phi)$
- $p_z = p_z(\tan \lambda, \Omega)$

Change of cov. matrix

- $\Sigma'_i = \mathfrak{J}^T \Sigma_i \mathfrak{J}$
- Σ_i cov. matrix in \mathfrak{A} .
- Σ'_i cov. matrix in \mathfrak{B} .

Jacobian helix parameters to momenta space

- After some derivative exercises ...

$$\mathfrak{J} = \begin{bmatrix} \frac{\partial P_x}{\partial \tan\lambda} & \frac{\partial P_y}{\partial \tan\lambda} & \frac{\partial P_z}{\partial \tan\lambda} & \frac{\partial E}{\partial \tan\lambda} \\ \frac{\partial P_x}{\partial \Omega} & \frac{\partial P_y}{\partial \Omega} & \frac{\partial P_z}{\partial \Omega} & \frac{\partial E}{\partial \Omega} \\ \frac{\partial P_x}{\partial d_0} & \frac{\partial P_y}{\partial d_0} & \frac{\partial P_z}{\partial d_0} & \frac{\partial E}{\partial d_0} \\ \frac{\partial P_x}{\partial z_0} & \frac{\partial P_y}{\partial z_0} & \frac{\partial P_z}{\partial z_0} & \frac{\partial E}{\partial z_0} \\ \frac{\partial P_x}{\partial \phi} & \frac{\partial P_y}{\partial \phi} & \frac{\partial P_z}{\partial \phi} & \frac{\partial E}{\partial \phi} \end{bmatrix} = \frac{-1}{\Omega} \begin{bmatrix} 0 & 0 & -\Omega P_T & -\frac{P_z^2 \Omega}{E \tan\lambda} \\ P_x & P_y & P_z & \frac{P_z^2}{E} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ P_y \Omega & -P_x \Omega & 0 & 0 \end{bmatrix}$$

→ $\Sigma'_i = \mathfrak{J}^T \Sigma_i \mathfrak{J}$, covariance matrix in momenta space.

($\Sigma'_i = \mathfrak{J} \Sigma_i \mathfrak{J}^T$ if you define jacobian as the transposed of quoted above)