Measurement of the top quark mass and couplings at Linear Colliders

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International Linear Collider (ILC)





ILD and SiD detectors are optimised for: Particle Flow Algorithm (PFLOW)

Electron positron collisions

Superconducting acceleration technology

Ecms tuneable between 200 GeV and 500GeV, posible upgrade to 1TeV

Integrated $\mathcal{L} = 500 \text{ fb}^{-1}$ (2 years of running)

Beams are polarised: P(e⁻) ≈ ± 80% , P(e⁺) ≈ ± 30%





The energy of charged hadrons will be measured by the tracking detectors

The energy of photons will be measured by the electromagnetic calorimeter

The hadronic calorimeter is then used only to measure the energy of neutral hadrons



Compact Linear Collider (CLIC)





Electron-positron collider in the **multi-TeV** energy range

About 48 km site length

The c.o.m. energy: $\sqrt{s} = 3 \text{ TeV}$ (default design) 500 GeV - 1.5 TeV

Luminosity: *L* = 2x10³⁴ cm⁻²s⁻¹

Intense R&D in the CLIC collaboration to fully develop two-beam acceleration at high gradients





A **CLIC_ILD** and **CLIC_SiD** detector concepts have been developed from the ILD and SiD detector concepts for ILC

Modifications motivated by the more challenging experimental conditions at CLIC and by the higher collision energy



Motivation

arXiv:1205.6497



Top quark electroweak couplings

- Learn about BSM physics from the deviations observed on Higgs and top EW couplings.
- LHC cannot achieve enough accuracy in the measurement of the coupling deviations -> ILC accuracies are needed to access to fully significant deviations arXiv:1403.2893



Precision in the measurement of the top quark mass



- At linear colliders there are two techniques to determine the mass of the top quark
 - 1. Direct reconstruction of top from its decay products (above threshold)

Experimentally well-defined but the generated mass is **not well-defined theoretically** and non-perturbative corrections could be substantial

2. A scan of the top pair production threshold

Marcel Vos, Topics in top physics 18 High degree of precision using a theoretically well-defined top mass (1*S* mass, can be transformed into other mass schemes). Precise top mass measurement with well-controlled theory uncertainties

See details about the **1S mass scheme**: arXiv:hep-ph/9904468v2



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The $t\bar{t}$ threshold



- Top mass input **174 GeV** in the **1S** mass scheme and $\alpha_s = 0.118$
- NNLO calculations provided by the code *TOPPIK*
- Corrections for ISR and luminosity
 spectrum
- These corrections result in a smearing of the cross section peak at threshold

• The **smearing** is due to the **statistical efficiency**, reduced by the luminosity spectra. Not affected by systematics.



Event generation, detector simulation and reconstruction

 $e^+e^- \rightarrow t\bar{t}$ production at threshold CLIC@352 GeV

$m_t = 174.0 \text{ GeV}$ $\Gamma_t = 1.37 \text{ GeV}$

1. Generated events (signal + background)

Pythia: $e^+e^- \rightarrow t\bar{t}$, WW, ZZ

WHIZARD: $e^+e^- \rightarrow q\bar{q}, q\bar{q}e^+e^-, q\bar{q}e_{\nu}$

2. Simulation of the detector

Full simulation with high level of realism

3. Reconstruction

Standard algorithms Kinematic fitting: Grouping W-bosons and b-jets into top quarks

> Katja Seidel, Frank Simon, Michal Tesar, Stephane Poss Eur. Phys. J. C73 (2013) 2530



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- The cross-section depends on the top mass, so measuring the cross-section the top mass can be extracted (Also the α_s , Yukawa coupling, top width...)
- Inclusion of higher-order QCD contributions are needed for a correct description of the cross section
- Determination of event selection efficiency and background contamination
- Threshold scan with 10 energy points spaced by 1 GeV from 344 GeV to 353 GeV with an integrated luminosity of 10 fb⁻¹



Threshold scans at ILC and CLIC



The cross section for ILC rises faster due to the luminosity peak is narrower

But it does not result in a significant difference of the precision of the top quark mass measurement



Measurement of the top mass and α_s at CLIC and ILC

- Statistical uncertainty of top mass around 30 MeV (CLIC ~ 20% larger than ILC due to different luminosity spectrum)
- In addition: Experimental and theoretical systematics, and uncertainties from the conversion to the MS mass scheme. Total uncertainty below 100 MeV within reach.



100 MeV



Top quark electroweak couplings at the ILC

- The process $e^+e^- \rightarrow t\bar{t}$ involves only $t\bar{t}Z_o$ and $t\bar{t}\gamma$ primary vertices
- A way to describe the current at the $t\bar{t}X$ vertex:
- See details in: arxiv.org/abs/hep-ph/0601112

$$\left[\Gamma^{ttX}_{\mu}(k^2, q, \overline{q}) = ie\left\{\gamma_{\mu}\left(\widetilde{F}^X_{1V}(k^2) + \gamma_5 \widetilde{F}^X_{1A}(k^2)\right) + \frac{(q - \overline{q})_{\mu}}{2m_t}\left(\widetilde{F}^X_{2V}(k^2) + \gamma_5 \widetilde{F}^X_{2A}(k^2)\right)\right\}$$

where:

- V = Vector coupling
- **A** = Axial coupling
- $\mathbf{X} = Z, \gamma \qquad \qquad F_{1V}^{\gamma} \quad F_{1A}^{\gamma} \quad F_{2V}^{\gamma}$
 - F_{1V}^Z F_{1A}^Z F_{2V}^Z

Non CP violating top quark couplings





 $e^+e^- \rightarrow t\bar{t}$ semi-leptonic channel at ILC@500 GeV

1. Generated events (signal)

WHIZARD: Generate 6 fermions final state $e^+e^- \rightarrow q\bar{q}b\bar{b}l_{\nu}$ **Pythia**: Parton shower and hadronisation

Beams are polarised

(2 samples)

1. $e_L^- e_R^+$ P(e⁻) \approx -80%, P(e⁺) \approx +30% 2. $e_R^- e_L^+$ P(e⁻) \approx +80%, P(e⁺) \approx -30%

2. Simulation of the detector

Full realistic simulation of the ILD detector concept

M.S. Amjad, M. Boronat, et al http://www-flc.desy.de/lcnotes/ LC-REP-2013-007

+ WbWb and beam backgrounds added



Three-dimensional image of a 500-GeV $t\overline{t}$ event simulated in the ILD detector



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-1 -0.8 -0.6 -0.4 -0.2 0 0.2 0.4 0.6 0.8 1 cos(θ)

on

3. Reconstruction

Standard algorithms for event selection

Signal reconstruction: **combination of** *b* quark jet and *W* boson that minimises the following equation

$$d^{2} = \left(\frac{m_{cand.} - m_{t}}{\sigma_{m_{t}}}\right)^{2} + \left(\frac{E_{cand.} - E_{beam}}{\sigma_{E_{cand.}}}\right)^{2} + \left(\frac{p_{b}^{*} - 68}{\sigma_{p_{b}^{*}}}\right)^{2} + \left(\frac{\cos\theta_{bW} - 0.23}{\sigma_{\cos\theta_{bW}}}\right)^{2}$$

Efficiency of selection

51.9% for P,P' = -1,+1 (Left-handed electrons) **55.0%** for P,P' = +1, -1 (Right-handed electrons)





Observables

Standard model values

Total cross section (σ)	Observables	e⁻∟e⁺ _R	e⁻ _R e⁺ _L
The Forward-Backward Asymmetry (A _{FB} top)	σ(fb)	1564	724
	A _{FB}	0.38	0.47
The slope of the distribution of the helicity angle (λ_t)	$F_{\rm R} = (1 + \lambda_t)/2$	0.25	0.76

But actually there are 6 independent observables = 3 observables x 2 polarisations

So we can obtain the following CP conserving 6 couplings of the top to Z and γ

$$\begin{aligned} \sigma(+) & A_{FB}(+) & \lambda_{hel}(+) & (+ = e_{R}^{-}) \\ \sigma(-) & A_{FB}(-) & \lambda_{hel}(-) & (- = e_{L}^{-}) \end{aligned} \Rightarrow \begin{cases} F_{1V}^{\gamma} & F_{1A}^{\gamma} & F_{2V}^{\gamma} \\ F_{1V}^{Z} & F_{1A}^{Z} & F_{2V}^{Z} \end{cases} \end{cases}$$

* $F_{1A}\gamma = 0$ because of the gauge invariance





[1] <u>arXiv:hep-ph/0601112</u>

500/fb at 500 GeV yields 1-2 orders of magnitude better sensitivity than the LHC (300/fb at 14 TeV)

Coupling	SM value	LHC $[1]$	e^+e^-	[ILC DBD]		
		$\mathcal{L} = 300 \text{ fb}^-$	$\mathcal{L} =$	500 fb^{-1}		
			\mathcal{P},\mathcal{P}' :	$=\pm 0.8, \mp 0.3$		
$\Delta \widetilde{F}_{1V}^{\gamma}$	0.66	$^{+0.043}_{-0.041}$		$+0.002 \\ -0.002$		
$\Delta \widetilde{F}^Z_{1V}$	0.23	$^{+0.240}_{-0.620}$		$+0.003 \\ -0.003$		
$\Delta \widetilde{F}^Z_{1A}$	-0.59	$^{+0.052}_{-0.060}$		$+0.005 \\ -0.005$		
$\Delta \widetilde{F}_{2V}^{\gamma}$	0.015	$^{+0.038}_{-0.035}$		$+0.003 \\ -0.003$		
$\Delta \widetilde{F}^Z_{2V}$	0.018	$^{+0.270}_{-0.190}$		$+0.006 \\ -0.006$		
	LHC studies (Snowmass 2005)			Present study denoted as ILC DBD		



Summary

Top mass at threshold

Top electroweak couplings

Statistical uncertainty of top mass around 30 MeV

(CLIC ~ 20% larger than ILC due to different luminosity spectrum).

Total uncertainty below 100 MeV in

reach, expected to be dominated by theory systematics.





Conclusions

- In a threshold scan, the top mass can be determined in a theoretically well defined way, using 1S mass scheme
- These studies confirm the expectation that a linear e⁺e⁻ collider will be capable of measuring the mass of the top quark with 30 MeV error
- Polarisation allows to double the number of observables
- It is a powerful tool for analysis because it also allows full separation between axial and vectorial couplings and between tt̄Z and tt̄γ vertices

 In LC with polarised beams we can measure with accuracies one or two orders of magnitude better than LHC



THANK YOUR FOR YOUR ATTENTION



BACKUP SLIDES



CLIC@CERN

Slide by Steinar Stapnes, CERN



Ties Bennke, 27.5.2014

ILC - ILD



ILC@Japan



tt decay modes

 $e^+e^- \rightarrow t\overline{t}$ gives three different final states:

Fully leptonic (10.3%) 2*jets* + 2 *leptons* + 2 *neutrinos* Semi-leptonic (43.5%) 4 jets + lepton + neutrino

Fully hadronic (46.2%) 6 jets at final state







