Feedback-free optical cavity with selfresonating mechanism for laser-Compton scattering sources

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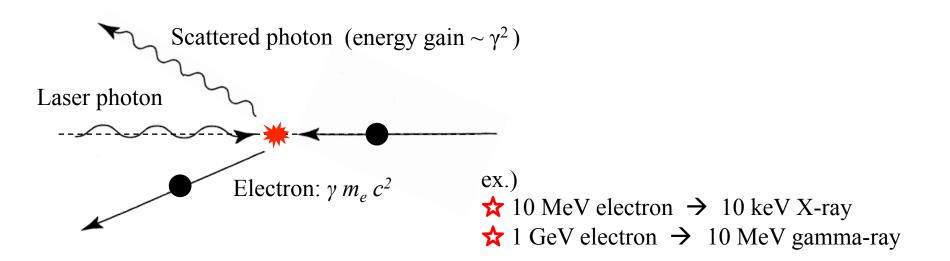
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Introduction



Laser-Compton scattering sources

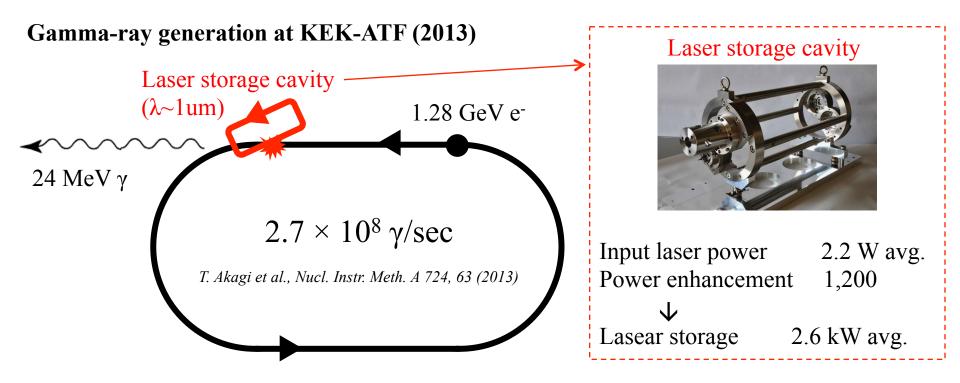
• Photon sources by the inverse-Compton scattering of laser photons from relativistic electrons

Applications

- Compact X-ray sources for a wide field of material science, industrial, medical applications etc.
- A gamma-ray driven polarized positron source for the ILC

ILC Technical Design Report: Volume 3, Part I, p.p. 137

Motivation

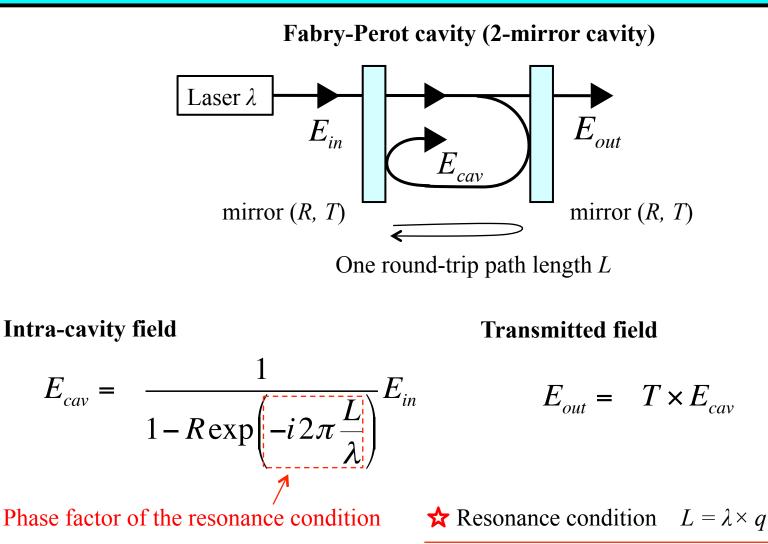


Need more gamma-rays (2 to 3 orders of magnitudes)

 \rightarrow Need more laser power storage (~ 1 MW)

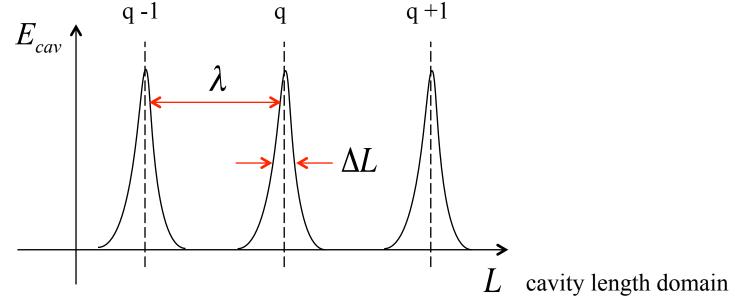
 \rightarrow Need more power enhancement of the cavity! (10k to 100k)

Optical resonant cavity



(L: round-trip path length, λ : wavelength of light, q: integer, R: reflectance, T: transmittance)

Optical resonant cavity



- Mode separation equals λ
- Resonance peaks have the "resonance width" written as ΔL (FWHM)
- "Finesse" is defined as the ratio of λ and ΔL ,

$$F = \frac{\lambda}{\Delta L}$$

• The finesse is almost same as the quality factor, Q-value

 \bigstar Finesse \propto power enhancement factor!

Issue on the development

An example of the cavity: at KEK-ATF

•	Wavelength	$\lambda = 1064 \text{ nm}$
•	Finesse	F = 4,040
•	Resonance width	$\Delta L = 260 \text{ pm}$
•	Enhancement factor	1,200

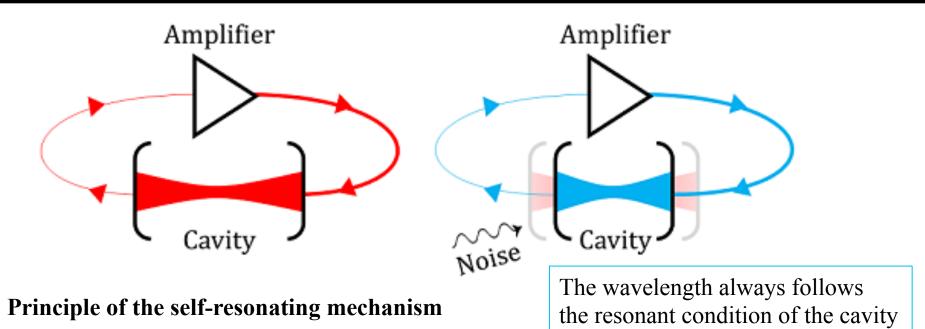
Achieved laser power storage

- Input laser power 2.2 W avg.
- laser power storage 2.6 kW avg.
- Power fluctuation 37 W
 - \rightarrow This fluctuation came from accuracy of the feedback system
 - \rightarrow It corresponds to deviation of the cavity length of 16 pm
- The enhancement increases ten times $\rightarrow \Delta L$ of 26 pm \approx feedback accuracy of 16 pm
- Precise feedback control will be a major technical difficulty for high finesse cavity





New idea: self-resonating mechanism



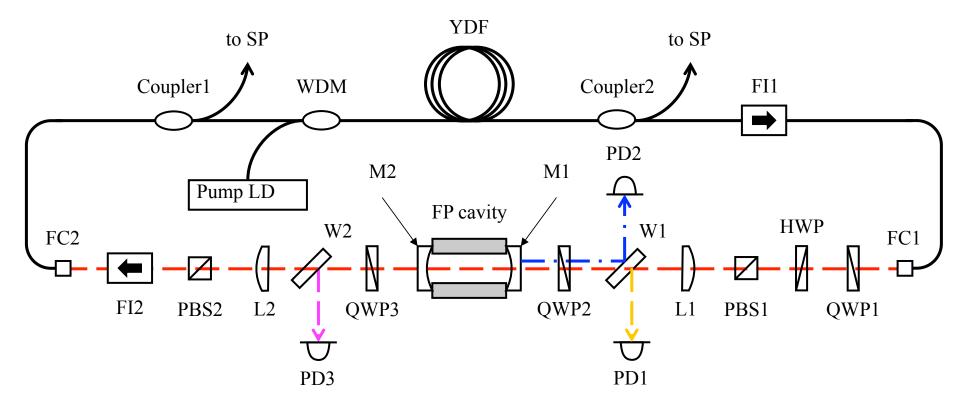
- 1. An optical amplifier emits seed (noise) light
- 2. A part of seed light, which is satisfied the resonance condition, transmits the cavity
- 3. That light doesn't feel the optical loss \rightarrow reaches the laser oscillation

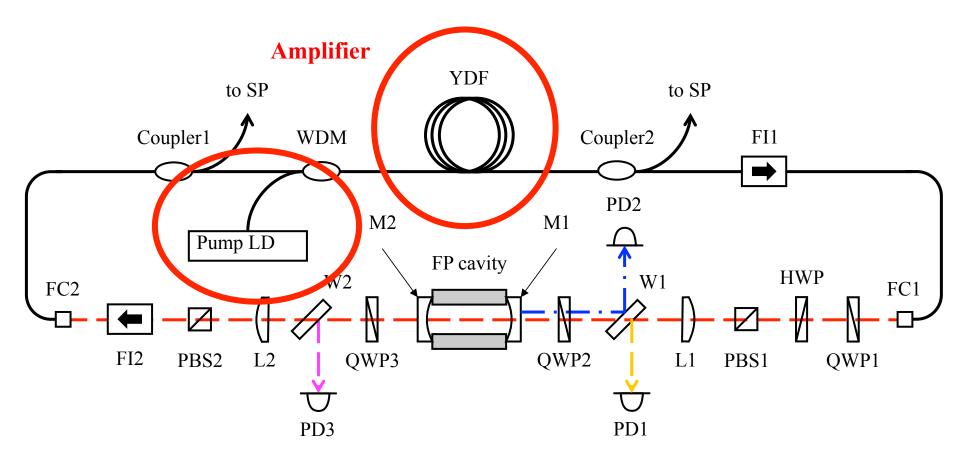
★ Laser oscillation condition: (gain - losses) > 1

 \bigstar Laser oscillation \Leftrightarrow laser storage

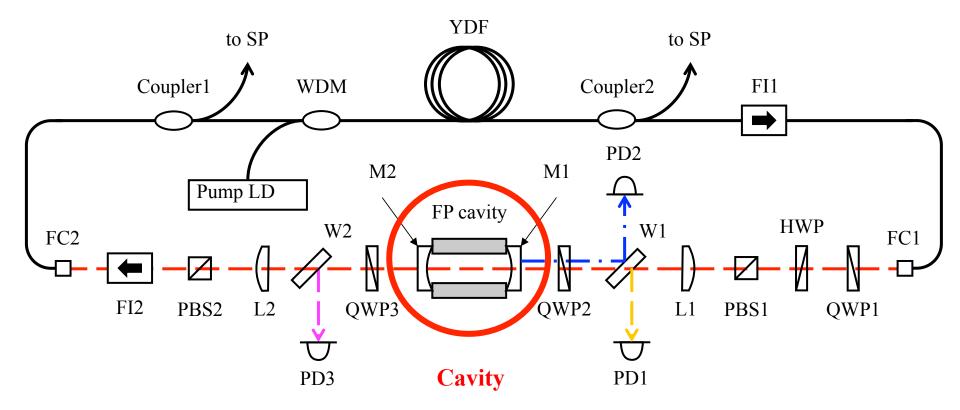
The laser oscillation continues without the feedback system to maintain the cavity length!

Y. Honda, et al., Proc. 7th Annual Meeting of PASJ, 1102 (2010)

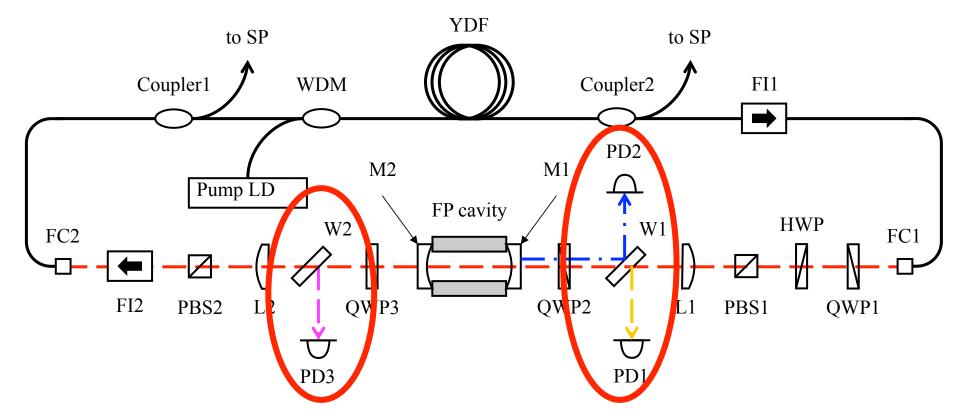




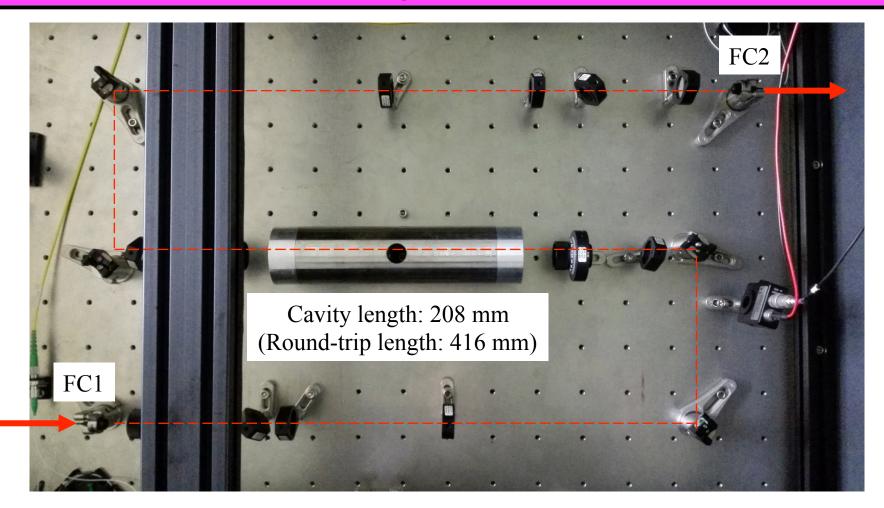
• The optical amplifier consists of a laser medium, the Ytterbium-doped fiber (YDF), and a pumping light sources, the laser-diode (Pump LD)



• The optical cavity consists of two concave mirrors with R > 99.999%

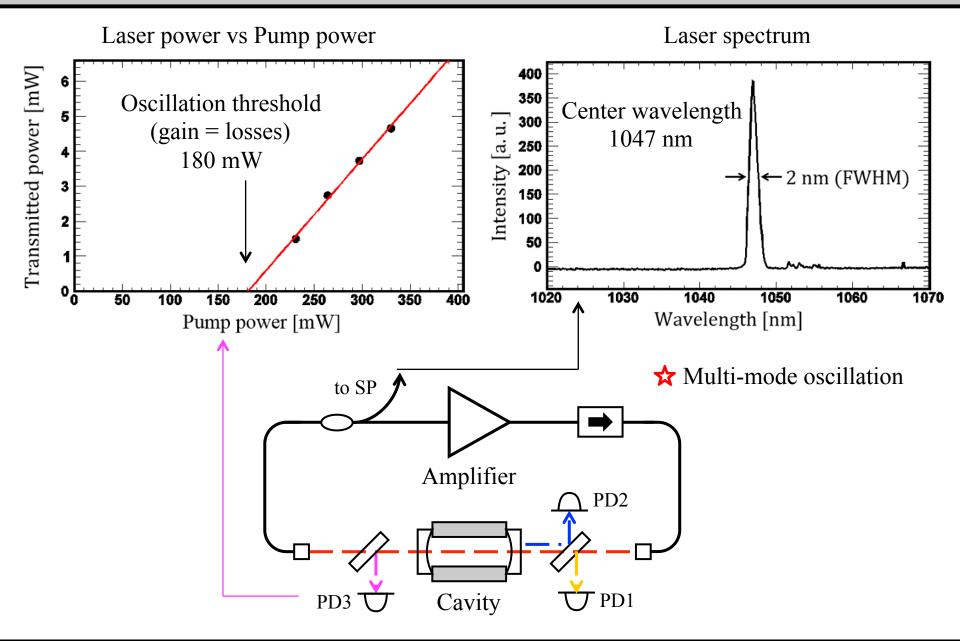


• The intra-loop laser power around the cavity is monitored by wedge plates (W1, 2) and photo-diode detectors (PD1, 2, 3)

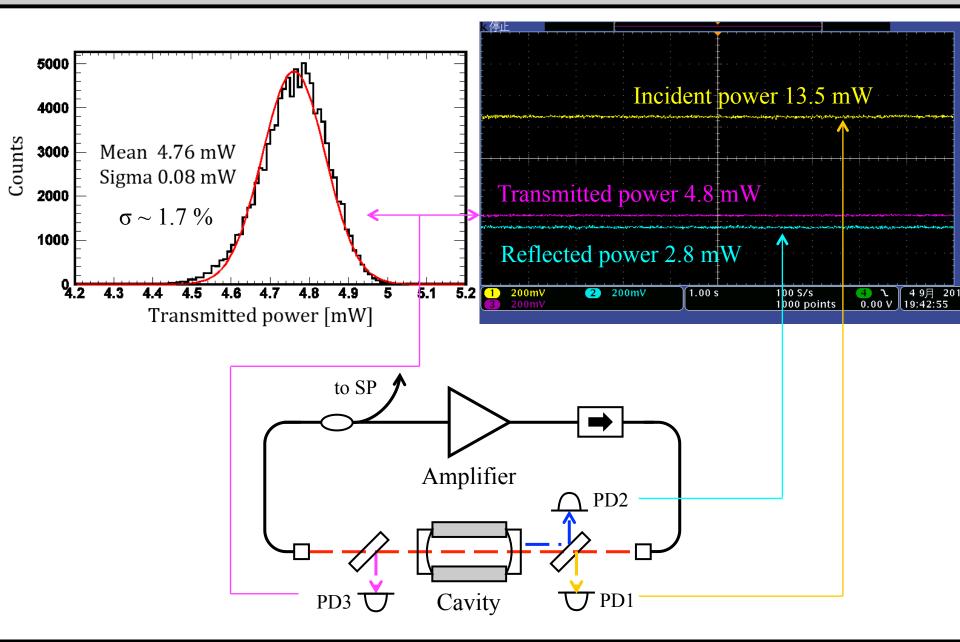


- Mirrors are mounted on a super-invar alloy tube by aluminum holders
- All component are in the air with the room temperature

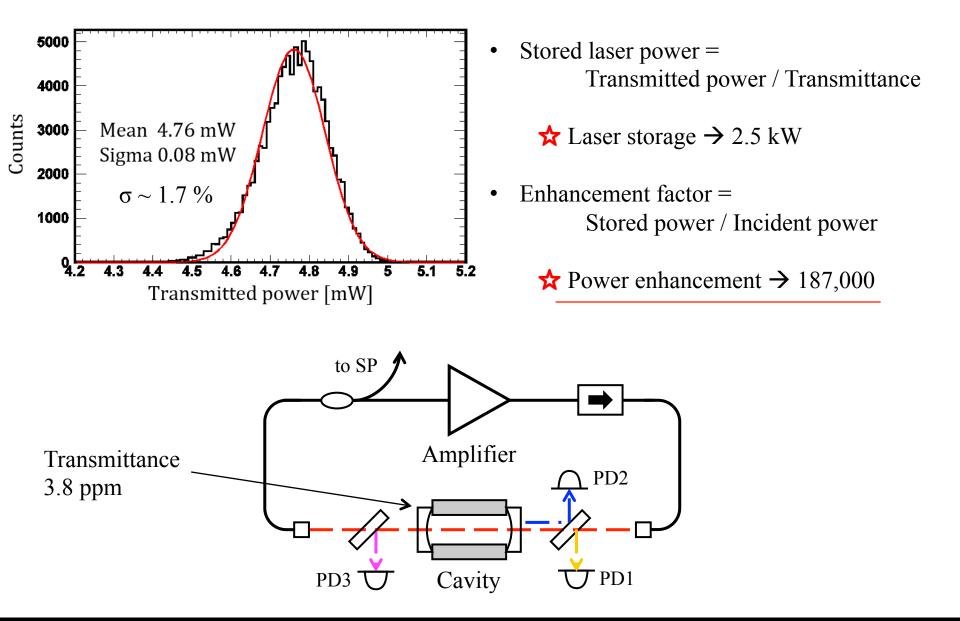
Laser oscillation and laser power storage



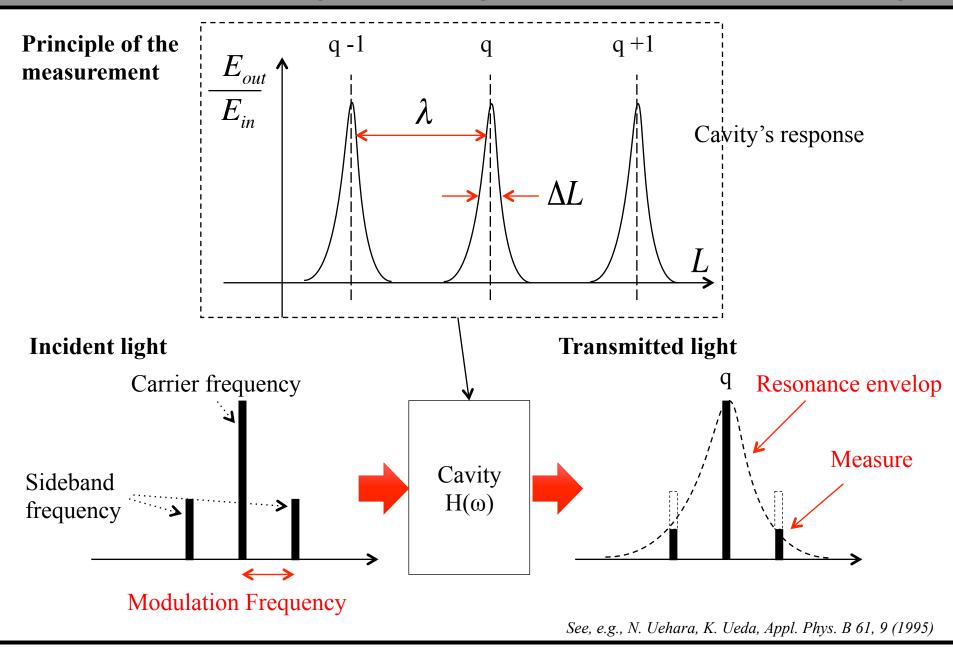
Laser oscillation and laser power storage



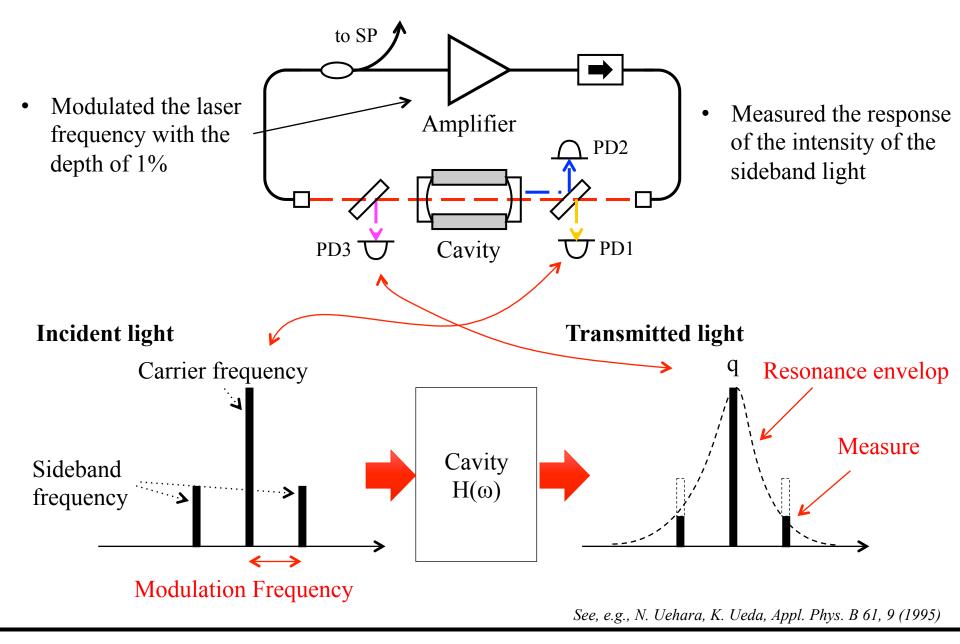
Laser oscillation and laser power storage



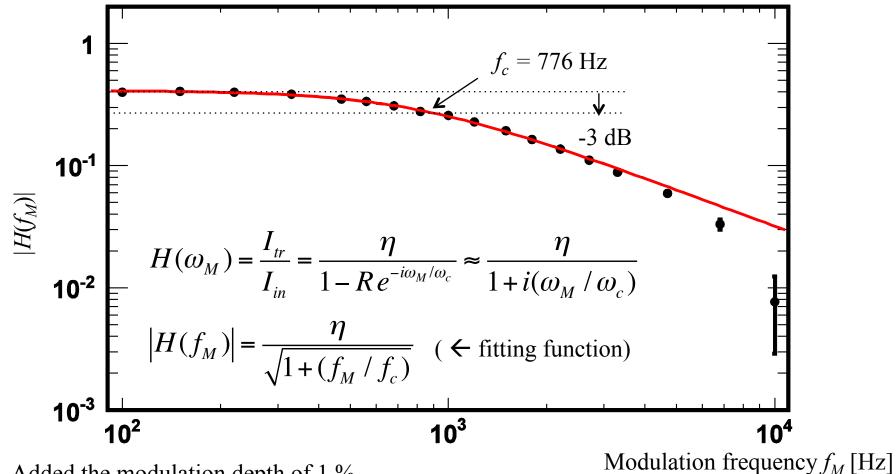
Finesse measurement by a frequency response function of the cavity



Finesse measurement by a frequency response function of the cavity



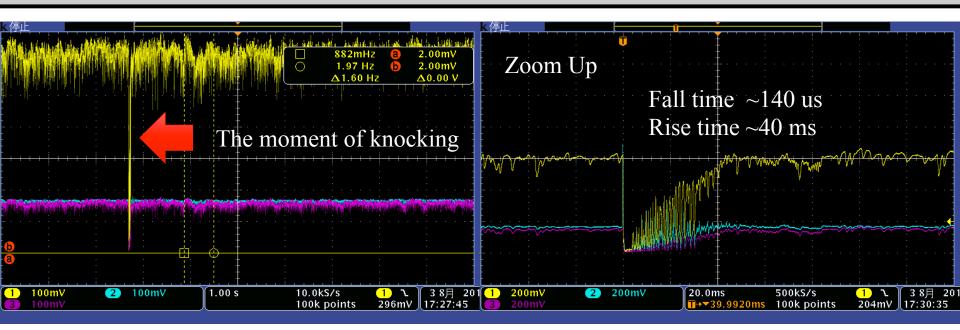
Measured frequency response function of the sideband light



- Added the modulation depth of 1 %
- Measured response function has a low-pass filter shape
- The resonance width was 1.55 kHz which was twice of the cut-off frequency, $f_c = 776$ Hz

★ Resonance width = 2.3 pm ★ Finesse = 465,000

Stability of the oscillation



Watched the stability of the resonance condition

- If we added some mechanical noise to the optical table, the laser power and spectrum were not change; the laser oscillation continued
- The oscillation stopped when the cavity was knocked, however, it recovered soon
 - > The fall time of ~140 us \rightarrow decay time constant of the cavity
 - > The rise time of ~40 ms \rightarrow maybe, dumping time of the super-invar alloy tube

★ Very robust system!

Summary

- We are working on a development of the optical resonant cavity for laser-Compton scattering sources
- Developed a new laser storage system which doesn't need a feedback system to maintain the resonance condition
- Succeeded the stable and continuous laser storage in the high finesse cavity with the finesse of 465,000 and the resonance width of 2.3 pm
- Achieved the power enhancement factor of 187,000
- R/D for the high peak-power storage (mode-lock pulse oscillation) \rightarrow on going!

backups

Introduction

The point of issue in the development

- Increase the brightness of X/gamma-rays
- Need 2 to 3 orders

Requirement

 \bigstar Electron beam \rightarrow Develop a high current and low emittance accelerator

 \bigstar Laser light \rightarrow High leaser power and small spot size at the IP

- Use an optical resonant cavity
- Store laser light in the cavity
- ➤ Enhance the laser power
- ➢ Forces the laser light about 10 micro meters
- We are developing the optical cavity
- The target storage power is about 1 MW in average
 e.g.) 100 W incident laser 10,000 enhancement cavity = 1 MW storage

Optical Resonant Cavity

An example of the cavity: at KEK-ATF

•	Wavelength	$\lambda = 1064 \text{ nm}$
•	Input laser power	2.2 W avg.
•	Round-trip path length	<i>L</i> = 1.68 m
•	Finesse	F = 4,040
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- Linewidth of the cavity $\Delta L = 260 \text{ pm}$
- Enhancement factor 1,200

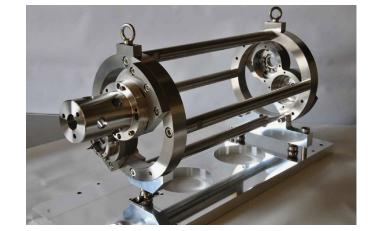
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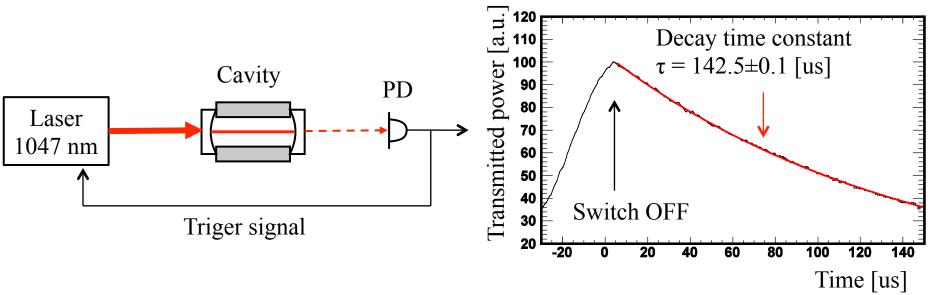
☆ The enhancement increases ten times \rightarrow 26 pm linewidth ~ 16 pm feedback accuracy

• Development of a high accuracy feedback system has technical difficulty...





Ring cavity at KEK-ATF T. Akagi et al., Nucl. Instr. Meth. A 724, 63 (2013)

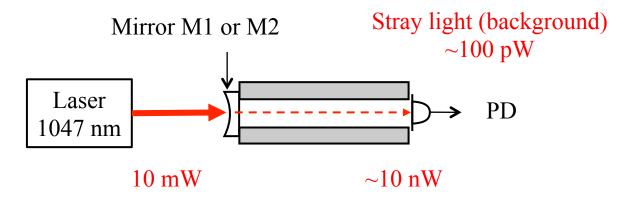


Measured the cavity finesse

- Use an external laser source with 1047 nm wavelength
- Switch off the laser light when the transmitted power increases (the cavity reaches the resonance condition)
- Intra-cavity light goes out with the exponential decay
- Finesse: $F = 2\pi c\tau / L$ (L: cavity length, τ : decay time constant)

★ Measured finesse: F = 650,000 which corresponds to R = 99.99952 %

☆ Cavity linewidth: $\Delta L = 1.6 \text{ pm}$



Measured the transmittance of mirrors

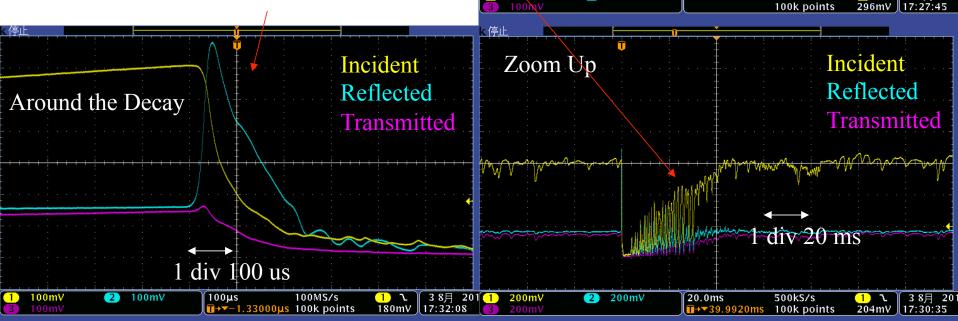
- Use an external laser source with 1047 nm wavelength
- The transmittance would be ~1 ppm (10⁻⁶) \rightarrow be careful the SNR in the measurement
 - \bigstar M1 4.2 ppm (input-side mirror)
 - \bigstar M2 3.8 ppm (output-side mirror)

Oscillation with High Finesse

Nocking the cavity, the oscillation is broken in a moment.

The oscillation seems stable when increasing the power at reach some threshold power.

Reflected power increases since the cavity is in the off resonance.



100mV

2 100mV

882mHz

1.97 Hz

∆1.60 Hz

Incident

Reflected

Pedestal Line

Transmitted

2.00m\

A0.00 V

 \bigcirc

10.0kS/s

1 div 1

1.00 s

Oscillation with High Finesse

Relaxation Oscillations

In the stable condition, the fluctuation of the laser power didn't have spikes with the ringing like as the case of the low finesse setup.

The power spectrum at the free running condition has a LPF-like shape.

At the Incident, LD Current = 500 mAPower [A.U.] 10⁻² 10^{-3} 10⁻⁴ 10^{2} 10 10 10 Fourier Frequency [Hz] At the Reflected, LD Current = 500 mAPower [A.U.] 10⁻² 10⁻³ 10⁻⁴ 10^{3} 10^{2} 10 Fourier Frequency [Hz]

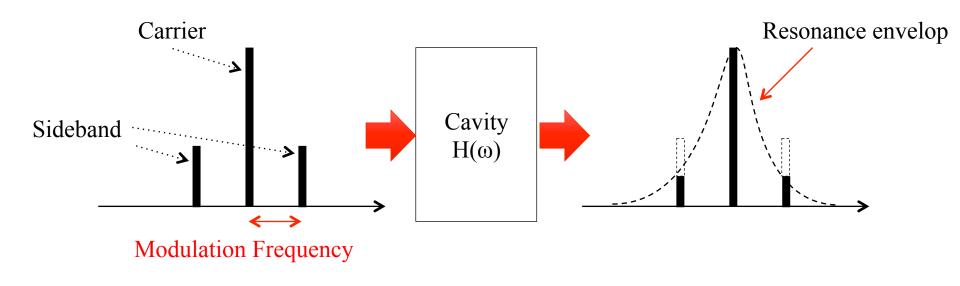
Finesse measurement

Measure a frequency response function of the cavity

• In a conventional cavity system (laser oscillator + resonant cavity), the frequency response of the cavity is measured by using amplitude modulation

See, e.g., N. Uehara, K. Ueda, Appl. Phys. B 61, 9 (1995)

- Add the sideband frequency to the carrier frequency of the incident laser light
- Measure a frequency response of the intensity of the sideband resonance
- The modulation depth must be small because large modulation depth would break the saturated laser oscillation state

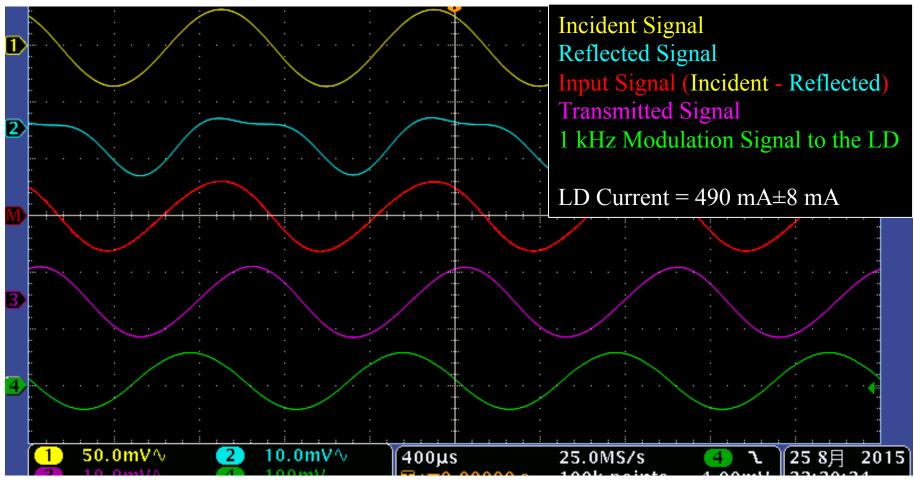


Pump rate: r Photon density: q $B(\omega)$ A(w) $\blacktriangleright \mathbf{x}(\omega)$ $u(\omega)$ + $H(\omega)$ If the frequency modulation by LD current modulation can be expressed as linear equations, the ratio between ►y(ω) response of the transmitted power and response of the incident power equals $\frac{x}{u} = BG_1 = \frac{BA}{1 - AH}$ the transfer function of a cavity, $H(\omega)$. For the photon density *q*, a linear $\frac{y}{u} = BG_2 = \frac{BAH}{1 - AH}$ equation is shown as below, (for the 4-level system) U 10 $\frac{G_2}{G_1} = \frac{x}{v} = H$

$$\frac{d^2 \delta q}{dt^2} + r\gamma_2 \frac{d \delta q}{dt} + (r-1)\gamma_c \gamma_2 \delta q = 0$$

12 0

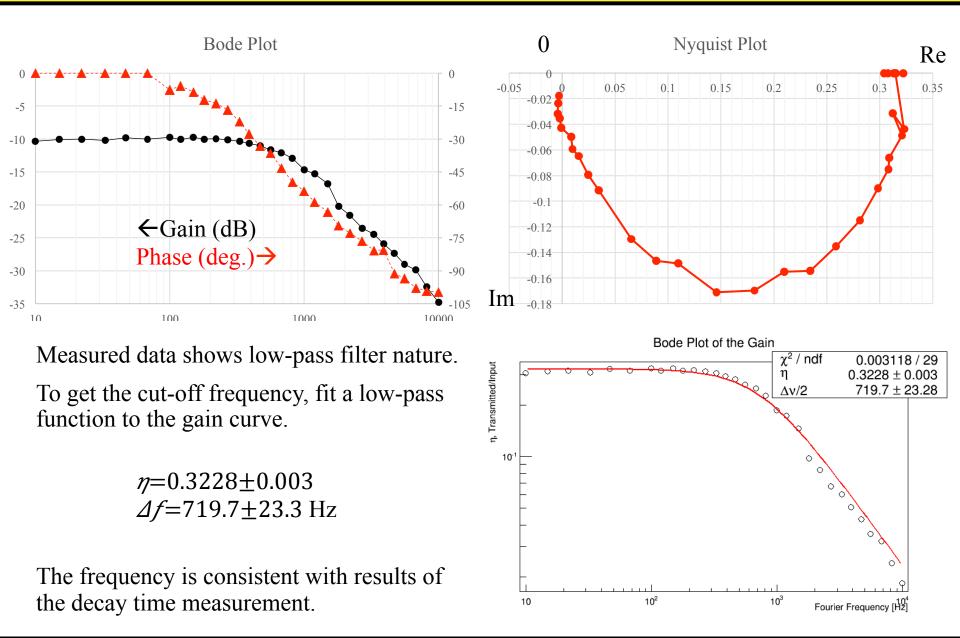
High Finesse Measurement



Measured the ratio of the amplitude and the phase difference between the input signal and transmitted signal.

Note that the ration of the amplitude between the input signal and the modulation signal is also changes because they are photon density and the population inversion.

High Finesse Measurement

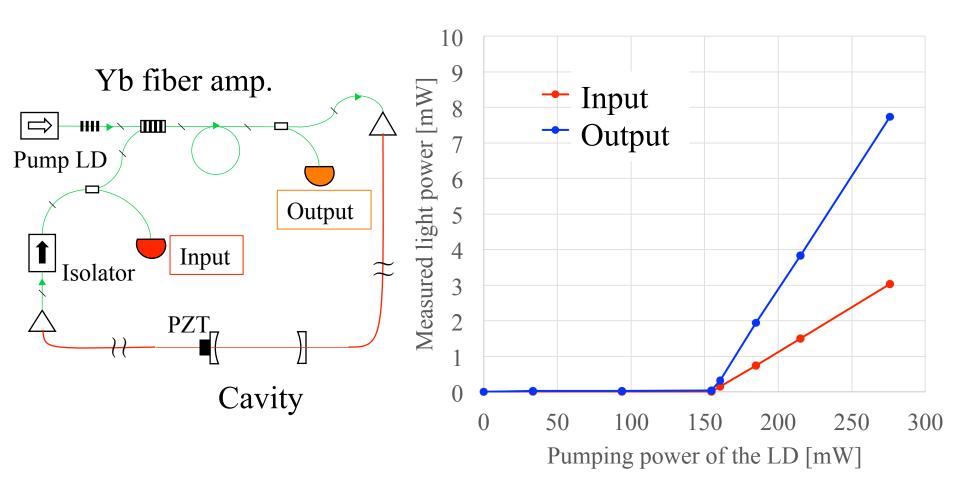


Finesse Measurement Using Intensity Modulation

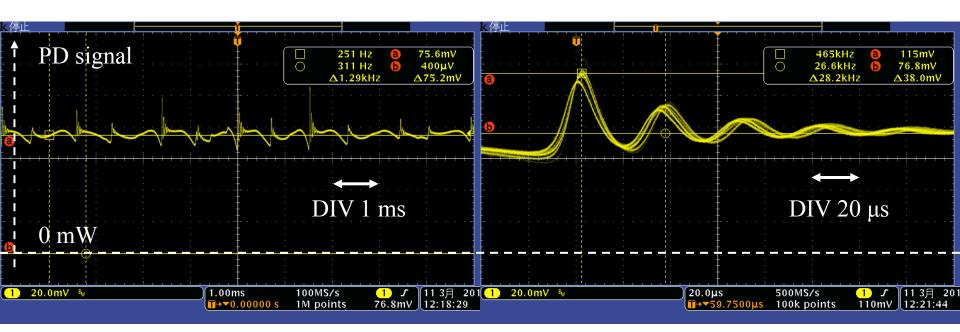
From the results of the frequency response function measurement

Cavity linewidth 1.55 kHz
Finesse 465,000
Incident / Transmitted ratio $\eta = 0.412$ Enhancement factor 190,000

*Considering the errors of detectors, which had an error of 0.5 % in its linearity and an error of 5 % in its absolute value, this results is consistent with the above results



- Threshold $\sim 150 \text{ mW}$ (pumping), Gain = 4.8 dB
- Lasing wavelength of 1034 nm with fluctuation $\pm 1~\text{nm}$



- Lasing power was kept up at 8 mW
- \checkmark Resonance condition was maintained without any cavity-control
- There were spikes in time scale of acoustic noise
- All spikes had a ringing with a time interval of ${\sim}30~{\rm kHz}$
- \checkmark The origin of time scale of this ringing?
 - \succ Decay time of the cavity: ~ 15 MHz
 - ≻ Fluorescence lifetime of a Yb fiber: ~ 1 ms

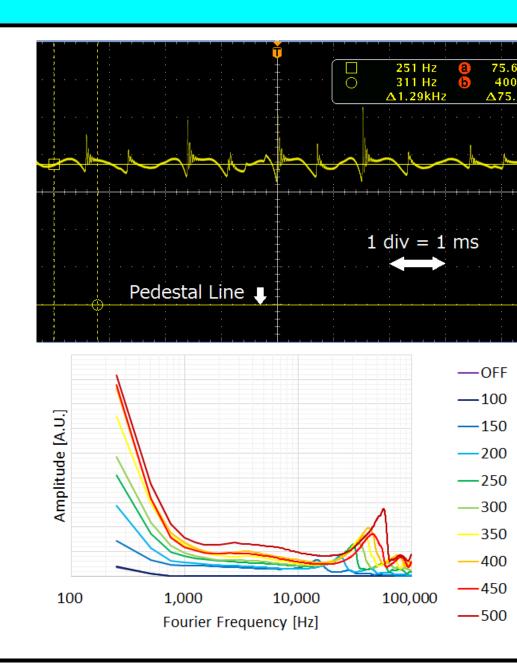
Oscillation with Low Fines

Relaxation Oscillations

Mechanical noise to the cavity

 \rightarrow a spike with the ringing occurs

The frequency of the ringing increases in proportion to increase the LD current.



Rate Equation (4-level System)

$$\frac{dN_2}{dt} = r_{pump} - \gamma_2 N_2 - KN_2 q$$
$$K = c\sigma_{emiss}^{eff} L^{eff}$$

$$\gamma_{2}$$
:自然放射 τ_{2} の逆数
 c :光速度
 σ^{eff}_{emiss} :誘導の実効反応断面積
 L^{eff} :実効光路長

励起効率

•

r

Equation of a Photon Density

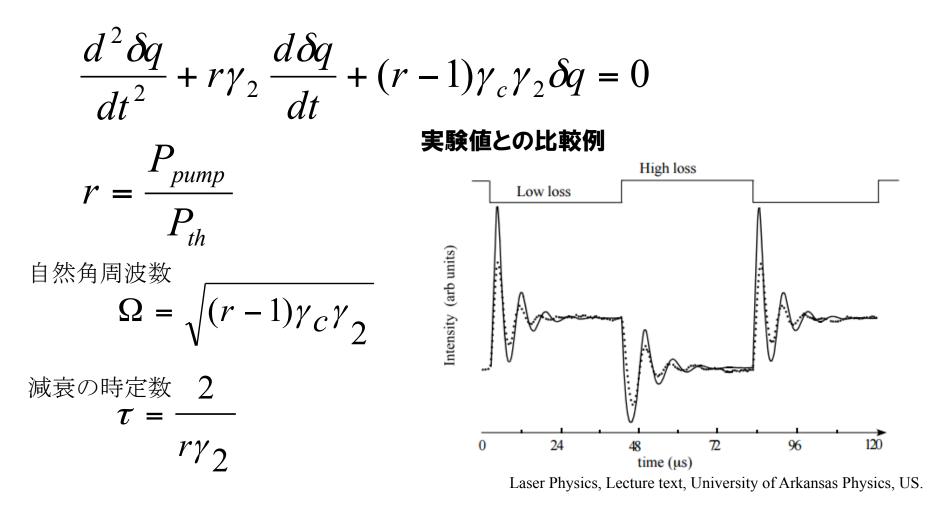
$$\frac{dq}{dt} = -\gamma_c q + KN_2 q$$

$$\tau_c : 周回光路にかかる時間$$

$$\gamma_c = \frac{\ell_c}{\tau_c}$$

Linearized Equation

定常状態からの摂動を考えて線形化する ($N_2 = N_{ss} + \delta N, q = q_{ss} + \delta q$)



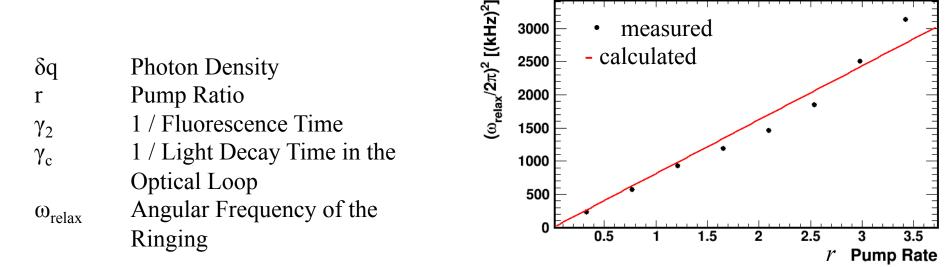
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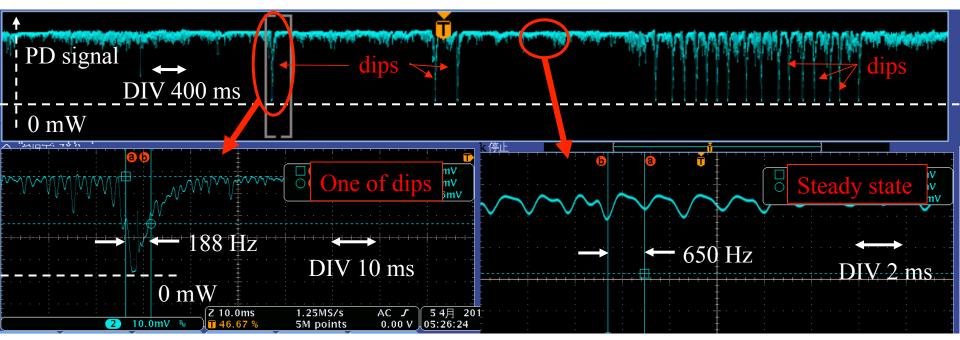
Relaxation Oscillations

Approximation of the rate equation

$$\frac{d^2 \delta q}{dt^2} + r\gamma_2 \frac{d \delta q}{dt} + (r-1)\gamma_c \gamma_2 \delta q = 0$$

$$\omega_{relax} \sim \sqrt{(r-1)\gamma_c\gamma_2}$$





- ✓ Lasing power was kept up at 4 mW
- \checkmark Laser oscillating with the high finesse cavity succeeded
- But it was <u>half</u> of the power in case of the low finesse, 8 mW
- Lasing power oscillated with about 600 Hz repetition always
- Sometime lasing power fell to zero, FWHM of a dip was 188 Hz
- \checkmark The cause of the 600 Hz oscillate and time scale of the dip?
 - \blacktriangleright Now we are looking for mechanisms of these phenomena

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Relaxation Oscillations

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