ILC Higgs Coupling Precision and Systematic Errors

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Outline

- Review of ILC Higgs Coupling Precisions
- Experimental and Theoretical Systematic Errors
- Limits on BSM decays and the Ultimate Higgs Coupling Precision

ILC Measurement of
$$\sigma(e^+e^- \rightarrow ZH)$$
 $\sqrt{s} = 250 \text{ GeV}$

Higgs Recoil Measurement of Higgs Mass and Higgstrahlung Cross Section



e⁺ Z^{*} Z^{*} Z^{*} H

ILC: $\Delta M_{H} = .025 \text{ GeV}, \ \Delta \sigma_{HZ} / \sigma_{HZ} = 1.4\% \text{ for L} = 500 \text{ fb}^{-1}$ $\Delta M_{H} = .013 \text{ GeV}, \ \Delta \sigma_{HZ} / \sigma_{HZ} = 0.7\% \text{ for L} = 2000 \text{ fb}^{-1}$ $\sigma_{HZ} \sim g_{HZZ}^{2}$ $\Rightarrow \Delta g_{HZZ} / g_{HZZ} = 0.7\% (0.35\%) \text{ for L} = 500 (2000) \text{ fb}^{-1}$ [from leptonic recoil alone]

ILC $\sigma \times BR$ measurements using $e^+e^- \rightarrow ZH$ $\sqrt{s} = 250 \text{ GeV}$



$$e^+e^- \rightarrow ZH$$
, $\nu\nu$ H $\sqrt{s} = 350$ GeV



All of the $\sigma \cdot BR$ Higgstrahlung studies that were done at $\sqrt{s} = 250$ GeV can also be done at $\sqrt{s} = 350$ GeV. Precisions for $\sigma \cdot BR$ are comparable, as is the precision for $\sigma(ZH)$ once $Z \rightarrow q \bar{q}$ decays are included.

WW fusion production of the Higgs at $\sqrt{s} = 350$ GeV provides a much better measurement of g_{HWW} compared to $\sqrt{s} = 250$ GeV. This gives a much improved estimate of the total Higgs width Γ_H which in turn significantly improves the coupling errors obtained from $\sigma \cdot BR$ measurements made at $\sqrt{s} = 250$ GeV. *WW* fusion also provides additional $\sigma \cdot BR$ measurements.

The recoil Higgs mass measurement is significantly worse at $\sqrt{s} = 350$ GeV with respect to $\sqrt{s} = 250$ GeV. However, there is hope that direct calorimeter Higgs mass measurements using $e^+e^- \rightarrow vvH$ will recover the precision (two ongoing studies were presented at this conference)

 $e^+e^- \rightarrow ZH$, $\nu\nu H \sqrt{s} = 350 \text{ GeV}$



Felix Müller | ILD analysis meeting | 07.10.2015 | 8





The g_{HWW} coupling can also be measured well at $\sqrt{s} = 500$ GeV through WW fusion production of the Higgs. Also the measurement of $\sigma(e^+e^- \rightarrow vvH) \times BR(H \rightarrow X)$ can be made for many Higgs decay modes $H \rightarrow X$.

Through $e^+e^- \rightarrow ttH$ the top Yukawa coupling can be measured to $\Delta y_t / y_t = 18\%$ with 500 fb⁻¹ at $\sqrt{s} = 500$ GeV. With same luminosity at $\sqrt{s} = 550$ GeV the precision is $\Delta y_t / y_t = 7.2\%$.

The ZHH channel is open at $\sqrt{s} = 500$ GeV. The Higgs self coupling can be measured to 27% with 4 ab⁻¹ assuming the true value is the SM value.

Summary of ILC Higgs Measurement Precisions

From "500 GeV ILC Operating Scenarios" arXiv:1506.07830

$\int \mathscr{L} dt$ at \sqrt{s}	$250 {\rm fb}^{-1}$ at	250 GeV	$330{\rm fb}^{-1}$ at 3	350 GeV	500fb^{-1} at 500GeV				
$P(e^-, e^+)$	(-80%,+30%)								
production	Zh	$v\bar{v}h$	Zh	$v\bar{v}h$	Zh	$v\bar{v}h$	tīh		
$\Delta\sigma/\sigma$	[39] 2.0%	-	[10,40] 1.6%	-	3.0	-	-		
BR(invis.) [41]	< 0.9%	-	< 1.2%	-	< 2.4%	-	-		
decay	$\Delta(\sigma \cdot BR)/(\sigma \cdot BR)$								
$h \rightarrow b\bar{b}$	1.2%	10.5%	1.3%	1.3%	1.8%	0.7%	28%		
$h \rightarrow c\bar{c}$	8.3%	-	9.9%	13%	13%	6.2%	-		
$h \rightarrow gg$	7.0%	-	7.3%	8.6%	11%	4.1%	-		
$h \rightarrow WW^*$	6.4%	-	6.8%	5.0%	9.2%	2.4%	-		
$h ightarrow au^+ au^-$	[42] 3.2%	-	[43] 3.5%	19%	5.4%	9.0%	-		
$h \rightarrow ZZ^*$	19%	-	22%	17%	25%	8.2%	-		
$h ightarrow \gamma \gamma$	34%	-	34%	[44] 39%	34%	[44] 19%	-		
$h \rightarrow \mu^+ \mu^-$ [45]	72%	-	76%	140%	88%	72%	-		

H-20: Preferred 20 year Running Scenario

	\sqrt{s}	∫Ldt	Lpeak	Ramp				Т	$T_{\rm tot}$	Comment
	[GeV]	$[fb^{-1}]$	$[fb^{-1}/a]$	1	2	3	4	[a]	[a]	
Physics run	500	500	288	0.1	0.3	0.6	1.0	3.7	3.7	TDR nominal at 5 Hz
Physics run	350	200	160	1.0	1.0	1.0	1.0	1.3	5.0	TDR nominal at 5 Hz
Physics run	250	500	240	0.25	0.75	1.0	1.0	3.1	8.1	operation at 10 Hz
Shutdown								1.5	9.6	Luminosity upgrade
Physics run	500	3500	576	0.1	0.5	1.0	1.0	7.4	17.0	TDR lumi-up at 5 Hz
Physics run	250	1500	480	1.0	1.0	1.0	1.0	3.2	20.2	lumi-up operation at 10 Hz

Table 7: Scenario H-20: Sequence of energy stages and their real-time conditions.

H-20

	first phase	lumi upgrade	total
250 GeV	500 fb⁻1	1500 fb-1	2 ab⁻1
350 GeV	200 fb-1		0 .2 a b⁻¹
500 GeV	500 fb⁻1	3500 fb-1	4 ab⁻¹
time	8.1 yrs	10.6 yrs	20.2 yrs*

ILC Higgs Coupling Precision vs Time



ILC Higgs Coupling Precisions

		H20 @ 8yrs	H20 @ 20yrs		
Topic	Parameter	Initial Phase	Full Data Set	units	ref.
Higgs	m_h	25	15	MeV	[51]
	g(hZZ)	0.58	0.31	%	[8]
	g(hWW)	0.81	0.42	%	[8]
	$g(hbar{b})$	1.5	0.7	%	[8]
	g(hgg)	2.3	1.0	%	[8]
	$g(h\gamma\gamma)$	7.8	3.4	%	[8]
		1.2	1.0	%, w. LHC results	[52]
	g(h au au)	1.9	0.9	%	[8]
	$g(hc\bar{c})$	2.7	1.2	%	[8]
	$g(ht\bar{t})$	18	6.3	%, direct	[8]
		20	20	%, $t\bar{t}$ threshold	[53]
	$g(h\mu\mu)$	20	9.2	%	[8]
	g(hhh)	77	27	%	[8]
	Γ_{tot}	3.8	1.8	%	[8]
	Γ_{invis}	0.54	0.29	%, 95% conf. limit	[8]

- [8] D. M. Asner et al., "ILC Higgs White Paper," arXiv:1310.0763 [hep-ph].
- [51] H. Li, arXiv:1007.2999 [hep-ex].
- [52] M. E. Peskin, in the Proceedings of the APS DPF Community Summer Study (Snowmass 2013), arXiv:1312.4974 [hep-ph].
- [53] T. Horiguchi, A. Ishikawa, T. Suehara, K. Fujii, Y. Sumino, Y. Kiyo and H. Yamamoto, arXiv:1310.0563 [hep-ex].

Higgs Physics Systematic Errors

Given that the statistical errors of many of the Higgs cross section and σ •BR reach the several per-mil level for the full H20 program, systematic errors must typically be 0.1% or less.

The following systematic errors have been considered:

- Flavor Tagging
- Luminosity
- Polarization
- Model Independence of ZH Recoil Measurements
- Theory Error

Higgs Physics Systematic Errors

Luminosity, Polarization, & Flavor Tagging Systematic Errors Assumed in 2013 Snowmass Higgs White Paper:

	Baseline	LumUp
luminosity	0.1%	0.05%
polarization	0.1%	0.05%
b-tag efficiency	0.3%	0.15%

$$\frac{\Delta L}{L} = \frac{2\Delta\theta}{\theta_{min}}, \quad \theta_{min} = 46 \text{ mrad} \quad \Delta\theta < 0.02 \text{ mrad} \quad \Delta R(\text{sensors}) < 30 \,\mu\text{m}$$

polarization obtained from polarimeters upstream and downstream of IP + physics processes such as $e^+e^- \rightarrow W^+W^-$

b-tag efficiency errors obtained from a quick studying using $e^+e^- \rightarrow ZZ \rightarrow l^+l^-b\overline{b}$ as a control sample; could be improved with additional control sample processes

Higgs Physics Systematic Errors

Model Independence of ZH Recoil Measurements

In order to use the hadronic ZH recoil measurement in our Higgs analyses we have to quantify the penalty associated with the fact that $\sigma(ZH \rightarrow q \,\overline{q} + X)$ is "almost model independent". By how much must we blow up $\Delta\sigma(ZH \rightarrow q \,\overline{q} + X)$ to account for the fact that the efficiencies differ by as much as 7%?



★ Combining visible + invisible analysis: wanted M.I.

i.e. efficiency independent of Higgs decay mode

Decay mode	$arepsilon_{\mathscr{L}>0.65}^{\mathrm{vis}}$	$arepsilon_{\mathscr{L}>0.60}^{\mathrm{vis}}$	$\varepsilon^{\rm vis} + \varepsilon^{\rm invis}$	_	
$H \rightarrow invis.$	<0.1%	22.0%	22.0 %		
${ m H} ightarrow { m q} \overline{ m q}/{ m gg}$	22.2%	<0.1 %	22.2 %		
$\rm H \rightarrow WW^*$	21.6%	0.1~%	21.7 %		
${ m H} ightarrow { m ZZ}^*$	20.2%	1.0%	21.2%		very similar
$H\to\tau^+\tau^-$	24.7 %	0.3 %	24.9 %		efficiencies
${ m H} ightarrow \gamma\gamma$	25.8%	<0.1%	25.8 %		
$H \to Z \gamma$	18.5 %	0.3 %	18.8 %		
$H \rightarrow WW^* \rightarrow q\overline{q}q\overline{q}$	21.3%	<0.1 %	21.3 %	٦	
$H \to WW^* \to q\overline{q} l\nu$	21.9%	<0.1 %	21.9 %		
${ m H} ightarrow { m W}{ m W}^* ightarrow { m q}\overline{ m q} au{ m v}$	22.1 %	<0.1 %	22.1 %	L	Look at wide
$H \rightarrow WW^* \rightarrow l \nu l \nu$	24.8 %	0.1 %	25.0 %		range of WW
$H \to WW^* \to l\nu\tau\nu$	20.5 %	0.8~%	22.1 %		tonologico
$H \to WW^* \to \tau \nu \tau \nu$	16.4 %	2.5 %	18.9 %		topologies

It is not sufficient to vary the SM Higgs branching ratios to estimate this systematic error. The problem is the BSM decays; they cannot be accounted for in this way.

To handle the BSM decays we have used an approach where we use all of our $\sigma \cdot BR$ measurements for SM Higgs decays to obtain an estimate of the average signal efficiency for $\sigma(ZH \rightarrow q \,\overline{q} + X)$. It is then straightforward to propagate the $\sigma \cdot BR_i$ errors, as well as the systematic errors on the individual decay mode efficiencies for the $\sigma(ZH \rightarrow q \,\overline{q} + X)$ selection, to the error on $\sigma(ZH \rightarrow q \,\overline{q} + X)$.

Let

- $\Psi \equiv \sigma(ZH \to q\,\overline{q} + X)$
- Ω = Number of signal + background events in $\sigma(ZH \rightarrow q \,\overline{q} + X)$ analysis
- B = Predicted number of background events in $\sigma(ZH \rightarrow q \,\overline{q} + X)$ analysis
- Ξ = Average efficiency for signal events to pass $\sigma(ZH \rightarrow q \,\overline{q} + X)$ analysis L = luminosity

$$\Psi = \frac{\Omega - B}{L\Xi} = \frac{1}{\Xi} \sum_{i} \psi_i \xi_i = \sum_{i} \psi_i \quad \text{where}$$

 $\psi_i = \sigma(ZH) \cdot BR_i$

 $\xi_i = e$ fficiency for events from Higgs decay i to pass $\sigma(ZH \rightarrow q \,\overline{q} + X)$ analysis

$$\Xi = \frac{\sum_{i} \psi_i \xi_i}{\sum_{i} \psi_i}$$

$$\psi_i = \frac{\omega_i - \beta_i}{L\eta_i}$$

 ω_i = Number of signal + background events in $\sigma(ZH) \cdot BR_i$ analysis β_i = Predicted number of background events in $\sigma(ZH) \cdot BR_i$ analysis η_i = efficiency for Higgs decay i to pass $\sigma \cdot BR_i$ analysis

- K_i = number of signal + background events common to had Z recoil and $\sigma \cdot BR_i$ analyses
- E = number of signal + background events unique to had Z recoil analysis ε_i = number of signal + background events events unique to $\sigma \cdot BR_i$ analysis

$$\Omega = E + \sum_{i} K_{i} \qquad S \equiv \Omega - B \qquad T \equiv \frac{\sqrt{S + B}}{S}$$

$$\omega_{i} = K_{i} + \varepsilon_{i} \qquad S_{i} \equiv \omega_{i} - \beta_{i} \qquad \tau_{i} \equiv \frac{\sqrt{S_{i} + \beta_{i}}}{S_{i}}$$

$$\lambda_{i} \equiv \frac{K_{i}}{\omega_{i}} \qquad N \equiv L \sigma_{ZH} \qquad r_{i} \equiv BR_{i} \qquad \delta_{i} \equiv \xi_{i} - \Xi$$

$$\left(\frac{\Delta \Psi}{\Psi}\right)^{2} = T^{2} \left\{ 1 + \frac{N^{2}}{\Omega} \sum_{i} r_{i}^{2} \left[\tau_{i}^{2} \left(\delta_{i}^{2} - 2\lambda_{i}\eta_{i}\delta_{i} \right) + \Delta \xi_{i}^{2} \right] \right\} \qquad \text{This is our result for the error on}$$

$$\sigma(ZH \to q \overline{q} + X)$$

$$\left(\frac{\Delta\sigma(ZH\to q\,\overline{q}+X)}{\sigma(ZH\to q\,\overline{q}+X)}\right)^2 = \mathrm{T}^2 \left\{ 1 + \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right] \right\} \quad \text{i.e. sys err} = \frac{1}{2} \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right]$$

Assume $\sqrt{s} = 350$ GeV and L=500 fb⁻¹

$$\mathsf{N} = L\,\sigma_{ZH} = 45383 \quad r_i = BR_i = (1 - BR_{BSM})BR_i(SM) \quad \tau_i(SM) = \frac{\Delta\sigma \bullet \mathsf{BR}_i(SM)}{\sigma \bullet \mathsf{BR}_i(SM)} = \frac{\sqrt{s_i + \beta_i}}{s_i}$$

Assume
$$T = \frac{\sqrt{S+B}}{S} = 0.014$$
 $\Omega = S+B = 17738$ and $\xi_i(SM)$ given in the table four pages back.

We assume that the vis+invis efficiency values in the table four pages back cover all possible BSM decays since they cover all SM decays from completely invisible to fully hadronic multi-jet decays. Assuming no knowledge of the properties of the BSM decays we can then set

$$\xi_{BSM} = 0.5 * [\xi_{vis+invis}(max) + \xi_{vis+invis}(min)] = 0.5 * [0.258 + 0.188] = 0.22$$

$$\Delta \xi_{BSM} = 0.5 * [\xi_{vis+invis}(max) - \xi_{vis+invis}(min)] = .035$$

$$\left(\frac{\Delta\sigma(ZH\to q\,\overline{q}+X)}{\sigma(ZH\to q\,\overline{q}+X)}\right)^2 = \mathrm{T}^2 \left\{ 1 + \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right] \right\} \quad \text{i.e. sys err} = \frac{1}{2} \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right]$$

We next obtain the error $\frac{\Delta \sigma \bullet BR_{BSM}}{\sigma \bullet BR_{RSM}}$ from Michael Peskin's Higgs coupling fit program. We do not use the $\sum_{i} BR_{i} = 1$ constraint, and to begin with we only use the leptonic recoil σ_{ZH} measurement. This provides a model independent measurement of g_{BSM} . For $\sqrt{s} = 350$ GeV, L=500 fb⁻¹ Michael's program gives $\frac{\Delta g_{BSM}}{g_{BSM}} = 0.032$ which we multiply by two to get $\frac{\Delta \sigma \bullet BR_{BSM}}{\sigma \bullet BR_{BSM}} = 0.064$. We take this error to mean that $0 < BR(H \rightarrow BSM) < 2 \times 0.064$, and set the measured $BR(H \rightarrow BSM) = 0.064$. This gives a model independent hadronic recoil cross section error of $\frac{\Delta\sigma(ZH \to q\,\overline{q} + X)}{\sigma(ZH \to q\,\overline{q} + X)} = 0.014 * 1.27 = 0.018.$

We then add this new model indepdendent hadronic recoil $\sigma_{\rm ZH}$ measurement as input to Michael's program to obtain a new error $\frac{\Delta \sigma \bullet BR_{BSM}}{\sigma \bullet BR_{RSM}} = 0.041$. Setting $BR(H \to BSM) = 0.041$ we then obtain a new model independent hadronic recoil σ_{ZH} error of $\frac{\Delta\sigma(ZH \rightarrow q\,q + X)}{\sigma(ZH \rightarrow q\,\overline{q} + X)} = 0.014 * 1.12 = 0.016$.

Iterating again we arrive at $BR(H \rightarrow BSM) = 0.039$ and $\frac{\Delta\sigma(ZH \rightarrow q\,q+X)}{\sigma(ZH \rightarrow q\,\overline{q}+X)} = 0.014 \times 1.11 = 0.016$. Further interations don't give any improvement.

$$\left(\frac{\Delta\sigma(ZH\to q\,\overline{q}+X)}{\sigma(ZH\to q\,\overline{q}+X)}\right)^2 = \mathrm{T}^2 \left\{ 1 + \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right] \right\} \quad \text{i.e. sys err} = \frac{1}{2} \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta\xi_i^2\right] \right\}$$

We have shown that $\frac{1}{2} \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta \xi_i^2 \right] = 0.11$ for $\sqrt{s} = 350$ GeV, L=500 fb⁻¹.

How does this scale with luminosity?

$$\frac{N^2}{\Omega} \propto L \quad \tau_i^2 \propto L^{-1} \quad r_i^2 \text{ is independent of lumi except } r_{BSM}^2 = \tau_{BSM}^2 \propto L^{-1} \text{ .}$$

If we assume $\Delta \xi_i = 0$ except $\Delta \xi_{BSM} = 0.035$ then
$$\frac{1}{2} \frac{N^2}{\Omega} \sum_i r_i^2 \left[\tau_i^2 \delta_i^2 + \Delta \xi_i^2 \right] = 0.11 \text{ independent of the luminosity at } \sqrt{s} = 350 \text{ GeV.}$$

Caveats for hadronic recoil systematic error calculation:

(1) This systematic error analysis was only done at $\sqrt{s} = 350$ GeV; it has not yet been done for $\sqrt{s} = 250$ & 500 GeV

(2) These results assume that the true $r_{BSM} = BR(H \rightarrow BSM)$ is small. As the true r_{BSM} grows we need to keep the product $r_{BSM}\Delta\xi_{BSM}$ constant to maintain the same systematic error, where ξ_{BSM} is the efficiency for BSM Higgs decays to pass the hadronic recoil analysis. For example

true r_{BSM} required $\Delta \xi_{BSM}$

.05	0.027
.10	0.014
.15	0.0091
.20	0.0068

These $\Delta \xi_{BSM}$ requirements may seem stringent for the larger values of true r_{BSM} . However as r_{BSM} grows we will have more *BSM* decays to analyze and the required improvement in Monte Carlo modelling of the *BSM* decays should follow.

Higgs Physics Systematic Errors -- Theory Errors

ILC model independent global coupling fit using 32 σ -BR measurements Y_i and σ_{ZH} measurement Y_{33}

$$\chi^{2} = \sum_{i=1}^{i=33} \left(\frac{Y_{i} - Y_{i}^{'}}{\Delta Y_{i}}\right)^{2},$$

$$Y_i^{'} = F_i \cdot \frac{g_{HZZ}^2 g_{Hb\bar{b}}^2}{\Gamma_0}$$
, or $Y_i^{'} = F_i \cdot \frac{g_{HWW}^2 g_{Hb\bar{b}}^2}{\Gamma_0}$, or $Y_i^{'} = F_i \cdot \frac{g_{Htt}^2 g_{Hb\bar{b}}^2}{\Gamma_0}$

$$F_i = S_i G_i \quad \text{where } S_i = \left(\frac{\sigma_{ZH}}{g_Z^2}\right), \ \left(\frac{\sigma_{\nu\bar{\nu}H}}{g_W^2}\right), \text{ or } \left(\frac{\sigma_{t\bar{t}H}}{g_t^2}\right), \text{ and } G_i = \left(\frac{\Gamma_i}{g_i^2}\right).$$

The cross section calculations S_i do not involve QCD ISR. The partial width calculations G_i do not require quark masses as input.

It is felt that the total theory errors for S_i and G_i - as well as the errors on the SM calculations for σ_i and Γ_i - will be at the 0.1% level at the time of ILC running.

Higgs Physics Systematic Errors -- Theory Errors

ILC Higgs Self Coupling Measurement at $\sqrt{s} = 500$ GeV From $\sigma(e^+e^- \rightarrow ZHH)$



 g_{hZZ} fixed to value from $\sigma(ZH)$ measurement

To handle g_{hhZZ} one can

- (1) Assume g_{hhZZ} is fixed to SM value
- (2) Assume g_{hhZZ} is fixed to some value related to measured g_{hZZ}
- (3) Simultaneously measure g_{hhh} & g_{hhZZ}

Options (1) and (2) would have an additional theory systematic error associated with these assumptions

Towards the Ultimate ILC Higgs Coupling Precision or why it is important to pursue a good $BR(H \rightarrow BSM)$ measurement even if, as a result of this effort, no BSM decays are found.

Perform coupling fit with $\sum_{i} BR_{i} = 1$ including $\Delta BR(H \rightarrow BSM)$ for (the constraint $\sum_{i} BR_{i} = 1$ is model independent if $\Delta BR(H \rightarrow BSM)$ is included in the fit)

$BR(H \rightarrow BSM)$	no meas.	< 7.2%	< 3.6%	<1.8%	< 0.9%	< 0.09%
(95% CL)						
ZZ	0.31%	0.29%	0.26%	0.22%	0.20%	0.19%
WW	0.38%	0.36%	0.31%	0.25%	0.21%	0.19%
bb	0.60%	0.57%	0.52%	0.46%	0.42%	0.40%
$oldsymbol{ au}^+oldsymbol{ au}^-$	0.88%	0.86%	0.83%	0.79%	0.77%	0.76%
gg	0.92%	0.91%	0.88%	0.86%	0.85%	0.84%
сс	1.1%	1.1%	1.1%	1.1%	1.1%	1.0%
YY	3.1%	3.1%	3.1%	3.1%	3.1%	3.1%
Γ_{tot}	1.7%	1.6%	1.3%	1.0%	0.84%	0.74%

ILC Higgs Coupling Precision assuming 20 year H20 scenario

Summary

- For most of the Higgs decay modes the ILC obtains model independent errors ranging from a few per-mil to a few percent for the full H-20 program
- An error of 27% can be obtained for the Higgs self coupling assuming the full H-20 program.
- With statistical precision reaching a few per-mil, systematic errors become important. Arguments were made that both experimental and theoretical systematic errors can be held to the 0.1% level, but much work is needed to realize this.
- BSM decays, or the limits on BSM decays, play an interesting role in both the model independence of the hadronic recoil ZH cross section analysis, and in the achievement of the ultimate Higgs coupling precision at the ILC.