Measurability of Anomalous ttH Couplings at the ILC 1.

Analysis procedure is similar with our VVH analysis.

KEK/SOKENDAI

Motivation.

- **>. The top quark is the heaviest particle in the SM.**
- **>. Its Yukawa coupling to the Higgs is also the strongest and has a important role.** = cos ^ξ) liggs is also the strongest and has a impo
costs this the continue
- >. Some deviations might appear in this ttH coupling. sho might appear in the time expling.
	- **>.** Direct measurement using ttH vertex with ILC.

 39.9 ± 0.000 case has been shown to be quite promising for this purpose λ

>. The ttH coupling with CP-mixed Higgs CP phase are parametrized;

$$
g_{\Phi tt} = -i\frac{e}{s_W} \frac{m_t}{2M_W} \big((1+a) + ib\gamma_5 \big) \qquad \phi_{CP} \equiv \arctan(b/a + 1)
$$

- >. a and b are independent.
- >. In the SM coefficients are given by a=0, b=0.
- >. A pure pseudo-scalar coupling is given by a=-1, b≠0.
- >. A mixed state of a Higgs is given by a≠-1, b≠0. ιο grvcn ω₎
∪p by a≠-1
- >. Angular observables exhibit a clear difference n n # Angular observables exhibit a clear difference
between the scalar and pseudo-scalar processes. $\frac{1}{2}$
- **>. Exploiting angular/spin correlation and polarization of a top pair is useful to extract CP information.** $\overline{2}$

Momentum/Angular Distributions.

Momentum/Angular Distributions.

Angular Distributions: ttH→4qlv+bb.

- **25 by ATLAS 25 by ATLAS 2015 The top spin info. is translated to the dists. of the decay products and it is not polluted by the effects of strong interaction.** $\frac{1}{26}$ is not uselling decay channel. The second coupling of the CP-odd $\frac{1}{2}$ 27 and 10 component of the Higgs of Subing interaction. **25 by ATLAS it is not polluted by the effects of strong interaction.**
- >. Lepton angular dists. in the decay of top is not affected by any non-SM effects in the decay vertex. 26 (h 26 (h 26) decay channel. It should be noted however that tree-level coupling of the CP-odd however that the CP-odd however that trees-level coupling of the CP-odd however that the CP-odd however the CP-odd ho 2. Lepton angular dists. In the decay of top is not allected by any non-

Pseudo-scalar: ECM vs Total σ(ee→ttH)

> Power of kinematical information for anomalous couplings.

>. Construct momentum/angular distributions.

 $@500GeV$: otth ~ 0.4(SM) *1000(L) * 20% = 80 * 1.5%(b=1) ~ 1 \rightarrow useless $@550GeV$: otth ~ 2.0(SM) *1000(L) * 20% = 400 * 4 %(b=1) ~ 16 $@600GeV$: otth ~ 4.0(SM) *1000(L) * 20% = 800 * 6 %(b=1) ~ 48 \rightarrow possible? $@$ 1 TeV: σ tth ~ 6.0(SM) *1000(L) * 20% = 1200* 15 %(b=1) ~ 180 → possible X(physsim) selection effi ratio of A

 $*$ @500GeV nominal σ (whizard) LR ~ 0.809

Rough Estimation At 500GeV.

CP-violating interaction contributes to the electric dipole moment (EDM). Any EDM constraints are not imposed here.

>. No information on momentum/angular distributions. \geq \overline{N} bin − f BSM w/accepture of BSM w/a
Distribution of BSM w/accepture of nent<mark>ı</mark>

- >. Distributions can not be available due to lack of events & low contamination of A. $\overline{1}$
- >. Intormation we can rely on is only Xsec probably.
→ Estimate sensitivity roughly **>. Information we can rely on is only Xsec probably.**
	- **>. Estimate sensitivity roughly**

!

j=1

 \sim

 for anomalous parameters in the a–b plane.

$$
\Delta \chi^2 = \left[\frac{N^{SM} \cdot \epsilon - N^{BSM} \cdot \epsilon}{\delta \sigma \cdot N^{SM} \cdot \epsilon} \right]^2
$$

 $>$. σ_{tth} (LR/RL) = 0.809. / 0.340. $>$. NSM, N^{BSM} is calculated analytically.

$$
\sigma_{tth} = (1+a)^2 \cdot \sigma_{tth}^{SM} + b^2 \cdot \sigma_{tth}^{Pseudo} + (1+a)b \cdot \sigma_{tth}^{Inter}
$$

 $>$. $\delta \sigma_{\text{tth}} \sim 10\%$ (with 500fb-1) $>$. ε \sim 20% (selection efficiency) ND(1)
Nefficiency) · Financy in Bottle

Rough Estimation At 600GeV.

>. My interest 600GeV. But no official samples.

>. What I can do is to assume reasonably and estimate it roughly.

>. At 600 GeV pseudo-scalar component could be relatively large.

- >. Angular dist. will be observed because remaining events will be much.
	- $>$. σ_{tth} (LR/RL) = 3.84./1.51 (Physsim), (which are 5 times better than that of 500.)
- $>$. $\delta \sigma_{\text{tth}} \sim 10\%$ (@500GeVwith 500fb-1) / **Sqrt(5)**.
- ϵ \sim 20% (selection efficiency for N).

 $N_{\text{remaining}} = \sigma_{\text{tth (h-} > bb)}$ 2.30 (eL0.8pR0.3) * 1000(L) * 0.2 (ε) \sim 450 (180+180+90)

- >. Angular dist. is divided into 8 bins. (3*3 for 2d)
- >. Acceptance for angular info. is 20% for each bin.
- >. Bkgs on each bin are asuumed to be 10 time more. **το pe 10**
(a = 1) Notime more.
- >. NErr on each bin is estimated by $\sqrt{(S+B)}$.

$$
\chi^2 = \sum_{i=1}^n \sum_{j=1}^n \left[\frac{N^{SM}(x_{ij}) \cdot f_{ij} - N^{BSM}(x_{ij}; a, b) \cdot f_{ij}}{\delta N^{SM}(x_{ij})} \right]^2 + \left[\frac{N^{SM} \cdot \epsilon - N^{BSM} \cdot \epsilon}{\delta \sigma \cdot N^{SM} \cdot \epsilon} \right]^2
$$

Rough Estimation At 600GeV.

 $\mathbf \Omega$

 -2

-3

 $-0.3 -0.2 -0.1$

20

 14

 12

 10

 $\overline{\mathbf{8}}$

l6

4

 $\overline{2}$

>. Only 1 angular info. is used (Ph).

>. 2d info. is used. (Ph,cosθ). (Ph,ΔΦ).

 $\overline{\mathbf{0}}$

 0.1

ArcTan(1.2/0.95) ~ 0.90 CP-phase Φcp: ~ 0.29π

 $ArcTan(0.5/0.985) \sim 0.47$ CP-phase Φcp: ~ 0.15π

Move on 1TeV.

Ecm = 1TeV.

>. Sigs.

- \Rightarrow . ttH→4qlv+bb / 6q+bb / 2q2l2v+bb
- $>$. σ_{tth} (LR/RL) = 5.89697. / 2.65115.

$$
\geq \text{Beam (Pe}^-\text{,Pe}^+) = (-0.8, +0.3).
$$

 $> L=1000fb^{-1}$

>. Bkgs.

>. Interfering BGs (same final state: ttbb).

- $-EW : ttZ \rightarrow ttbb$
- $-$ QCD: ttg $-$ ttbb (g $-$ bb: dominant)

>. Non-interfering BGs (but, huge cross sections).

- ttbar

(Hard gluon emission from bottom quarks.)

(Fraction of mis-reconstruction and/or failure of b-tag lead to significant BGs)

@ 1 TeV information http://www-jlc.kek.jp/~miyamoto/CDS/mc-dbd.log/generated/1000-B1b_ws/tth/

Ecm = 1TeV.

>. Reconstruction chain.

```
 <processor name="myRootFileProcessor"/>
<processor name="myTauJetFinder"/> (suehara's tau finding)<br><processor name="myIsoLepExtractor"/> (old method with measured E)
 <processor name="myIsoLepExtractor"/> (old method with measured E)
<processor name="MyFastJetProcessor"/>
 <processor name="MyUndoJetProcessor"/>
```

```
 <processor name="MyVertexFinder"/>
      <processor name="MyJetClusteringAndFlavorTag_6Jets"/> or
      <processor name="MyJetClusteringAndFlavorTag_8Jets"/> or
      <processor name="MyJetClusteringAndFlavorTag_4Jets"/>
 <processor name="MyThrustReconstruction"/>
 <processor name="MySphericityProcessor"/>
 <processor name="myTTHAnalyzer"/>
```
>. #IsoLeptons.

>. I did not optimize carefully this time, there is still room for improvement of results.

Event Selection.

- **>. Strong information to distinguish anomalous couplings is angular dists. → No angular cut.**
- **>. Momentum is also useful info. to distinguish it**
	- **→ Correlated observables with it should not be used.**

>. What we can use is event topology.

- >. Cut region is determined by a scan so as to maximize signal significance.
	- >. Significance (ttH \rightarrow 4qlv+bb) : S/ $\sqrt{(S + B)} = 4.84$ effi: 19.84
	- >. Significance (ttH→6q +bb) : S/ $\sqrt{(S + B)} = 5.79$ effi: 22.21
	- >. Significance (ttH \rightarrow 2q2l2v+bb) : S/ $\sqrt{(S + B)} = 2.71$ effi: 17.59

The selection efficiencies for arXiv:1409.7157v4 T. Price^{a, 1}, P. Roloff^{b, 2}, J. Strube^{c, 3, 5}, T. Tanabe^{d, 4}

signal events are 33.1% (6jets) and 56.0% (8jets)

Correlation between Phiggs vs observables.

Overall Acceptance: f.

>. Need to get acceptance.

$$
\chi^2 = \sum_{i=1}^n \left[\frac{N^{SM}(x_i) \cdot f_i - N^{BSM}(x_i; a, b) \cdot f_i}{\delta N^{SM}(x_i)} \right]^2
$$

1d

$$
N^{Reco}(x_j^{Reco}) = \sum_i f(x_j^{Reco}, x_i^{Gene}) \cdot N^{Gene}(x_i^{Gene})
$$

=
$$
\sum_i f_{ji} \cdot N_i^{Gene}
$$

=
$$
\sum_i \bar{f}_{ji} \cdot \eta_i \cdot N_i^{Gene}
$$

$$
\eta_i \equiv \frac{N_i^{Accept}}{N_i^{Gene}}
$$
 (Event acceptance.)

$$
\bar{f}_{ji} \equiv \frac{N_j^{Accept}}{N_i^{Accept}}
$$
 (Detector response function.)

2d

$$
N^{Reco}(x_{j\beta}^{Reco}) = \sum_{i} \sum_{\alpha} \bar{f}_{j\beta i\alpha} \cdot \eta_{i\alpha} \cdot N_{i\alpha}^{Gene}
$$

\n
$$
\eta_{i\alpha} \equiv \frac{N_{i\alpha}^{Accept}}{N_{i\alpha}^{Gene}}
$$
 (Event acceptance.)
\n
$$
\bar{f}_{j\beta i\alpha} \equiv \frac{N_{j\beta i\alpha}^{Accept}}{N_{i\alpha}^{Accept}}
$$
 (Detector response function.)

Error of #N on Each Bin: *δN***.**

Dists. on Bkgs, Sig, Event/Overall Acceptance.

Sensitivity to Anomalous Parameters.

- **>. Contour plots on sensitivity for anomalous parameters in the a–b plane.** $\frac{1}{2}$ in the a $\frac{1}{2}$ plane.
- >. Process is **ttH→4qlv+bb** and assumed L is 1000fb-1. or anomarous parameters in a
d assumed L is 1000fb-1
	- >. Use only momentum/angular distributions.

Sensitivity to Anomalous Parameters. ty to Anomalous ty to Anomalous

- **2. Contour plots on sensitivity for anomalous parameters in the a–b plane.**
- >. Process is **ttH→4qlv+bb** and assumed L is 1000fb-1.
- >. Use only momentum/angular distributions. **+ Cross section effect.**

>. Excluded Region on Φcp $\phi_{CP} \equiv \arctan(b/a + 1)$

Excluded Region on Φcp

- **>. Contour plots on sensitivity for anomalous parameters in the a–b plane.**
	- >. Each process on **ttH** and assumed L is 1000fb-1.
	- >. **Use the only higgs momentum distribution. + Cross section effect.**

Excluded Region on Φcp: Combined

- **>. Contour plots on sensitivity for anomalous parameters in the a–b plane.**
	- >. Each process on **ttH** and assumed L is 1000fb-1.
	- >. **Use the only higgs momentum distribution. + Cross section effect.**

>. Φcp > 0.29π is excluded using the only one distribution.

- >. If MELA is applied, sensitivity will much better because all dist. is calculated based on mom.
- >. 600GeV might has similar sensitivity.

ATLAS Studies on Anomalous ttH. to be twice the signal in the invariant mass window specified above. the formalism of anomalous *H* boson couplings is discussed. Monte Carlo simulation with the JHU generator is Λ Setudios on Anomalous ttHe matrix elements technique and the MELA framework and the MELA framework are discussed in Section IV. Analogue and the MELA framework are discussed in Section IV. Analogue are discussed in Sec studies on Anoniques tun with data-driven methods suggest larger background in the *tt*tle the *t*tle the *t*tle the *t* **to be twice the signal in the signal in the signal in the international in the international in** *^A*(*Hf* ¯*f*) = *m^f v* ¯*^f* (*^f* +i˜*^f* 5) *^f ,* (1) In ATI AS Studies on Anomalous ffH $S_{\rm eff}$ prediction. The new physics effects on the new physics effects on the parameterised by a minimal by a minima to **A T H A Q decays.** These matrix element to the optimal analysis of the production of t and decay processes. Such techniques have been proposed to enhance signal over background in application to *ttH*¯ production [40, 63, 64], and we employ them to probe anomalous *Hf* ¯*f* couplings for the first time. We define the

A. Study of the *ttH*¯ process

to determine with the jet reconstruction techniques and therefore the *CP* discriminants are not used in this analysis.

techniques for *HV V* coupling measurements [60–62] to *Hf* ¯*f* couplings in *ttH*¯ , *b*¯*bH*, and *tqH* production2, as well as

4 Conclusion

II. PARAMETERIZATION OF HIGGS BOSON COUPLINGS BOSON COUPLINGS BOSON COUPLINGS BOSON COUPLINGS BOSON COUPLINGS

where $\mathcal{S}_{\mathcal{S}}$ takes the value $\mathcal{S}_{\mathcal{S}}$

to additional radiation and loop corrections.

$$
\mathcal{L} \supset -\frac{y_t}{\sqrt{2}} \overline{t} (\cos\xi + i \gamma^5 \sin\xi) t h, \qquad \qquad \mathop{\stackrel{\scriptscriptstyle{0.0}}{}}\raisebox{-2.5pt}{$\scriptstyle{+} \atop \scriptstyle{|\chi - \text{best fit}|}}}
$$

Figure 2. The Higgs data constraints on the anomalous couplings C_S and C_P at the LHC and the **Exercise** \mathcal{L}_{eq} **Conclusion**
 Exercise \mathcal{L}_{eq} **Conclusion** Conclusion
to 68% and 95% C.L. respectively. The shadowed region represents the expected measurement uncertainty at HL-LHC. **discriming the up and down flavors of the up and down flavors of the decay chain is chain is known.** The latter is discriminate in In the *H* in the *H* indicate the indicate the signal and background. The signal and background. Figure 11 shows the signal and background. The signal and background. The signal and background. The signal and μ is seen ΔP distribution in the *H* ΔP . ΔP for order theory. The shadowed region represents the expected measurement uncertainty at HL-LHC. to 68% and 95% C.L. respectively. The shadowed region represents the expected measurement
uncertainty at HL-LHC. value of the Section value of the Higgs field. It is useful to the Section value of the shadowed representing at HL-LHC.

The shadowed representing at HL-LHC.

In this paper, we have obtained constraints on the \mathcal{CP} -violating top-Higgs couplings using excluded at 95% C.L.. We expected TLEP to improve this exclusion region to $|\xi| > 0.07\pi$. the current Higgs data and found that values of \mathcal{CP} -violating phase $|\xi| > 0.6\pi$ are already $\overline{}$ particular is performed with PYTHIA8. Similar to the study presented in Section V, the NLO \sim 1. the current Higgs data and found that values of \mathcal{CP} -violating phase $|\xi| > 0.6\pi$ are already $\frac{f}{f}$ = 0.077.
 *Letter with the formal in the <i>f*_C $\frac{f}{f}$ and *H* $\frac{f}{f}$ and *L*(*H_f* $\frac{f}{f}$) = $\frac{f}{f}$ In this paper, we have obtained constraints on the *CP*-violating top-Higgs couplings using
the connect Higgs detection of the top lease of CP sidetime above $|c| > 0$ for any close decay μ excluded at 95% C.L.. We expected TLEP to improve this exclusion region to $|\xi| > 0.07\pi$. In this paper, we have obtained constraints on the \mathcal{CP} -violating top-Higgs couplings using the current Higgs data and found that values of \mathcal{CP} -violating phase $|\mathcal{E}| > 0.6\pi$ are already excluded at 95% C.L.. We expected TLEP to improve this exclusion region to $|\xi| > 0.07\pi$. In this paper, we have obtained constraints on the CP -violating top-Higgs couplings using as CS $_{\odot}$ e

$\mathsf{arXiv}\cdot\mathsf{1606.03107v2.}$ (MFLA) which allows a connection to be made between the HL-- and α f and α and α In the *H* ! channel, we use the invariant mass *m* to separate the signal and background. Figure 11 shows the **D0** distribution in the *H* \rightarrow 0.0 and background distributions are shown. *f<i>ME* $>$. arXiv:1606.03107v2 (MELA) $\psi_{CP} > 0.1611$ We therefore conclude that discrimination power of the MELA approach is guaranteed even when \mathcal{A} \mathbb{O}_C violations contribute to the electric dipole moment (EDM). However, the bounds of \mathbb{O}_C

0.99 l. Φcp > ~0.18π(my cas on the probability functions for \mathbf{P} parameterized with temperated with temperated with temperated with generated with temperated with temperated with generated with generated with generated with generated with gener Φ and Φ \sim 0.18 π (my case) comparable to the current precision with the *HV V* measurements, but provides a fundamentally di↵erent approach $p_{\text{max}} = 0.48$

 \overline{r} $\frac{1}{2}$ uniquely defined \mathbf{v} where the f_{CD} parameter is conveniently bounded between 0 and 1 is uniquely defined and can be interpreted as the $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$ versus $\frac{1}{2}$ $\frac{$ meter defined for the HVV cripplings 13. 5. 62. V coupling conveniency bounded between 0 and 1, is uniquely defined, and can be interpreted as the convenients of the unitary security of the unitary from the unitary state of the unitary of the unitary of the unitary of the observable effects. It is a convenient counterpart of the f_{a3} parameter defined for the $H\overrightarrow{V}V$ couplings [3, 5, 62]. While cross section fraction corresponding to the pseudoscalar coupling, and therefore is directly related to experimentally where the f_{CP} parameter is conveniently bounded between 0 and 1, is uniquely defined, and can be interpreted as the on the coupling C^P depend on the assumption of Higgs couplings to other light fermions value. In the SM, the couplings³ have the values *^f* = 1 and ˜*^f* = 0. Any deviation from these values indicates the where the $f_{\rm CP}$ parameter is conveniently bounded between 0 and 1, is unique

Summary.

- **>. Based on classical method, anom couplings on ttH was studied.**
- **>. Any angular distribution can not be constructed @500GeV because of …**
- **>. Angular/Momentum dists. @1TeV are very useful.**
	- **and Φcp > ~0.18π is excluded if we use 2d/1d dimensional dists for all modes.**
- **>. However #N are not many, 2d/1d are limits.**
- **>. MELA(multi dimensional analysis) is the best way for CP of ttH.**
- **>. 600GeV might have similar power with 1TeV.**

Plan.

- **>. IsoLep finding should be improved.**
- **>. Go back to VVH analysis. We decided a plan of analysis for MELA.**
- **>. Come back later.**

>. Significance: $S/\sqrt{(S + B)} = 4.84$

The selection efficiencies (purities) for signal events are 33.1% (27.7%) and 56.0% (25.2%) for the six- and eight- jets analyses in ILD, respectively,

ttH→4qlv+bb ¹⁰⁰

>. Significance: $S/\sqrt{(S + B)} = 5.79$

ttH→2q2l2v+bb ¹⁰²

Nsig integral(restricted range) : 35.0 # Nbck integral(restricted range) : 132.4 $\frac{1}{2.71}$