

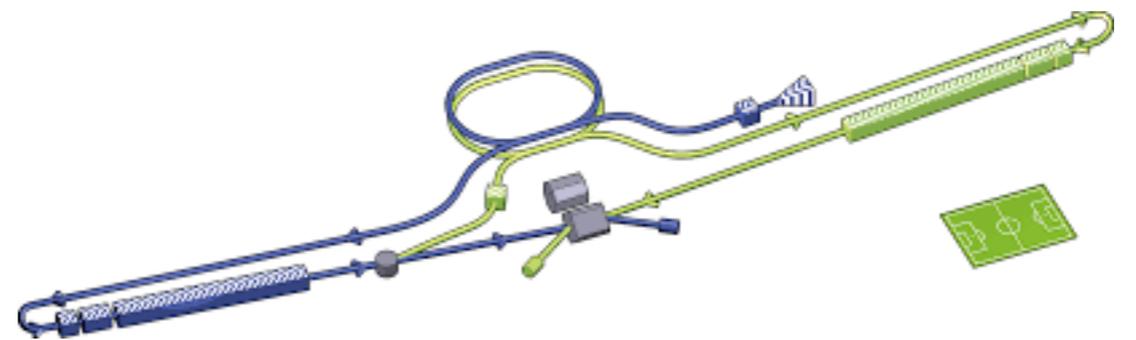
How LHC tells us there is excellent potential for ILC to discover new particles

Howard Baer
University of Oklahoma

or

Why

SUSY



While ILC250 well-motivated
to detect new physics via virtual effects
can ILC weigh in on the biggest possibility:
the discovery of SUSY?

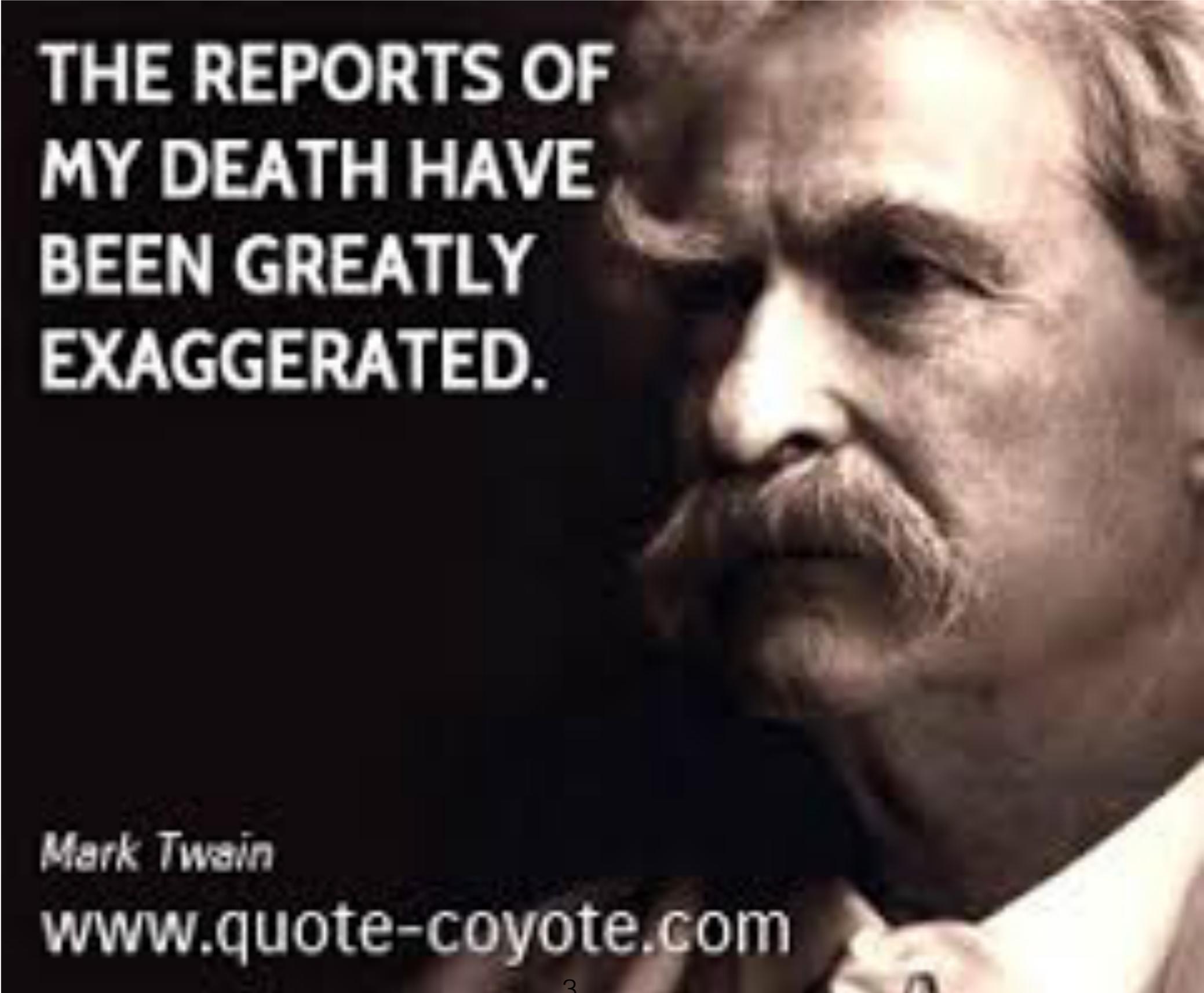
Or do recent LHC results, $m(\tilde{g}, \tilde{u}) > 2 \text{ TeV}$, $m(\tilde{t}_1) > 1 \text{ TeV}$
pre-empt any possibility?

[So far, the bulk of LHC searches take place within
either unnatural or simplified models]



Recent episode Big Bang Theory:
Physics is dead because
LHC didn't discover SUSY:
is this true? **NO!**

Mark Twain, 1835-1910 (or SUSY)



**THE REPORTS OF
MY DEATH HAVE
BEEN GREATLY
EXAGGERATED.**

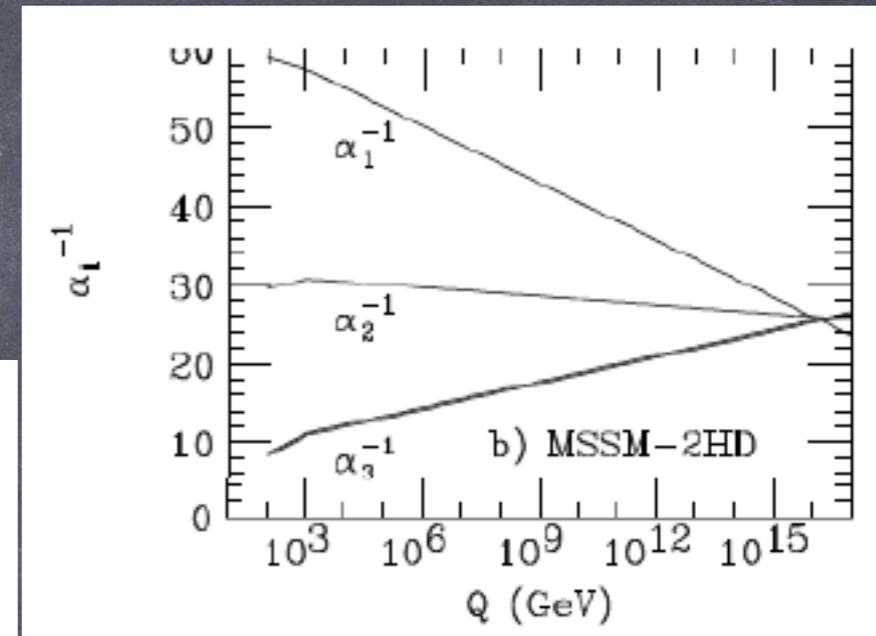
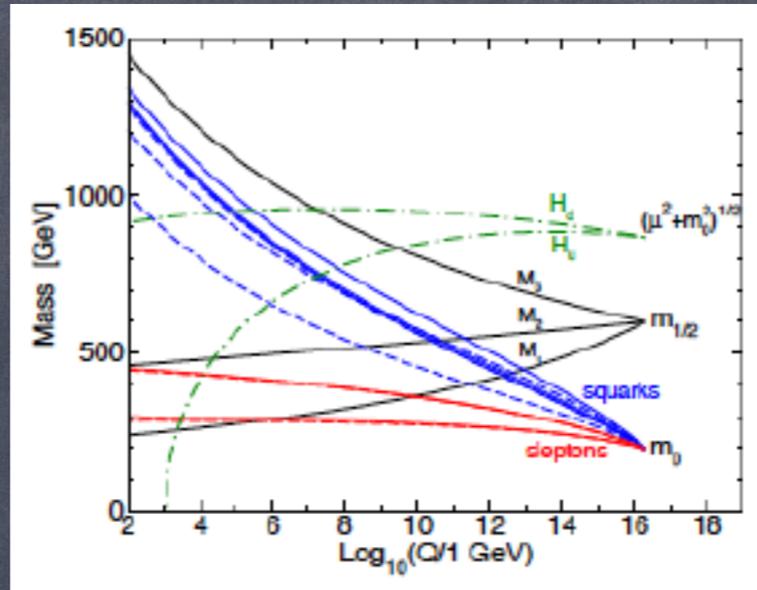
Mark Twain

www.quote-coyote.com

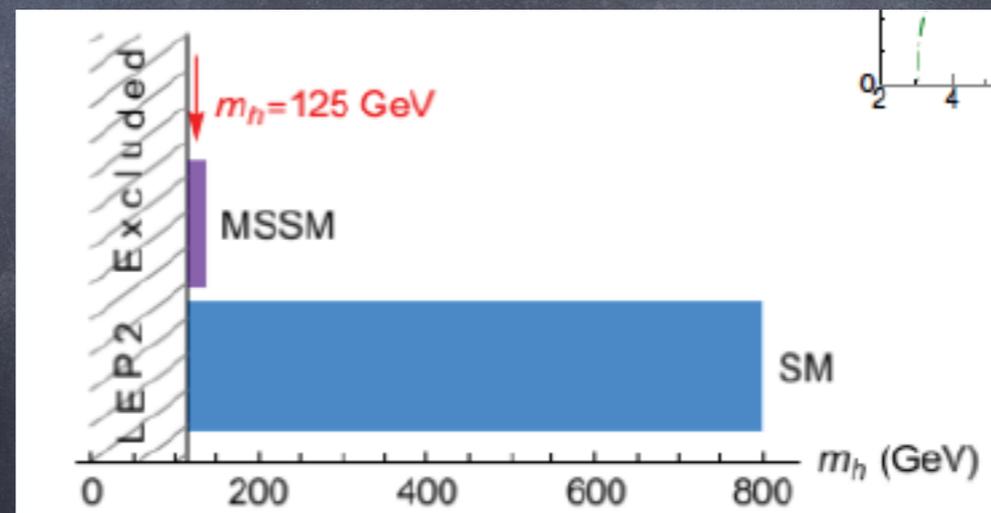
Three MSSM success stories:

Unification of gauge couplings

$m(t) \sim 150-200$ GeV
required for EWSB



$m(h)$ just right

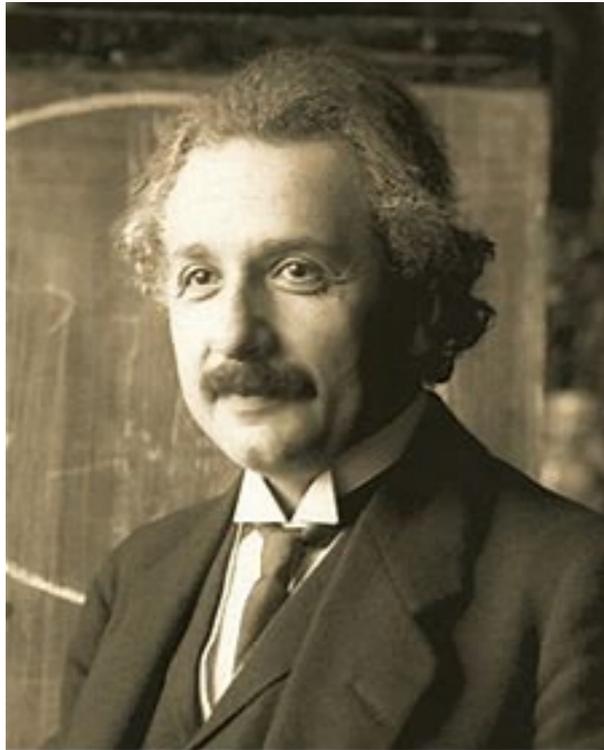


But: is there now a Little Hierarchy problem, i.e.
is SUSY now unnatural in light of strong LHC sparticle mass limits?

SUSY motivation: simplicity and naturalness



“The appearance of fine-tuning in a scientific theory is like a cry of distress from nature complaining that something needs to be better explained”
S. Weinberg



“Everything should be made as simple as possible, but not simpler”
A. Einstein

Needed: a natural theory which contains SM: the MSSM

“...settling the ultimate fate of naturalness is perhaps the most profound theoretical question of our time”



Arkani-Hamed et al.,
arXiv:1511.06495

“Given the magnitude of the stakes involved,
it is vital to get a clear verdict
on naturalness from experiment”

This should be matched by theoretical scrutiny
of what we mean by naturalness

naturalness:

all independent contributions to some observable O
should be comparable to or less than O

otherwise:

if one contribution is $\gg O$, then some other
must be fine-tuned to large opposite-sign to maintain O
at its measured value: this is unnatural
(i.e. highly implausible; likely wrong)

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \Sigma_u^u - \mu^2$$

MSSM+naturalness=> light higgsinos with mass $\sim 100-200$ GeV

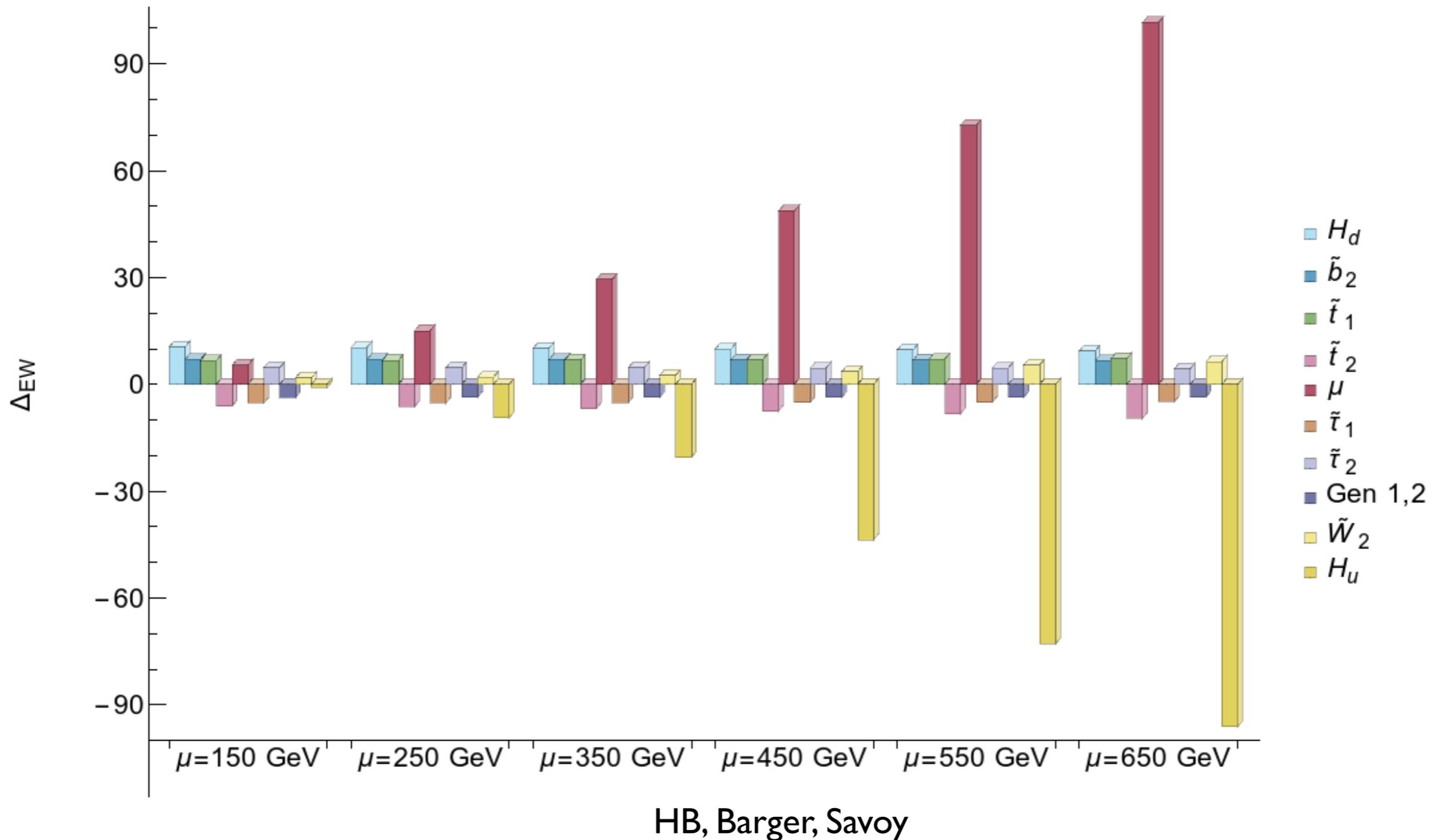
- $m_{H_u}^2$ is driven small negative (radiatively driven naturalness)
- Higgs/higgsino mass $\mu \sim 100 - 200$ GeV (the smaller the better)
- radiative corrections to m_Z^2 relation minimized for highly mixed TeV-scale top squarks

[$m(\text{gluino}) < \sim 4-5$ TeV; $m(\text{t1}) < 3$ TeV at little cost to naturalness]

LHC has only begun to explore natural SUSY parameter space!

light higgsinos are difficult, perhaps impossible, to see at LHC

How much is too much fine-tuning?

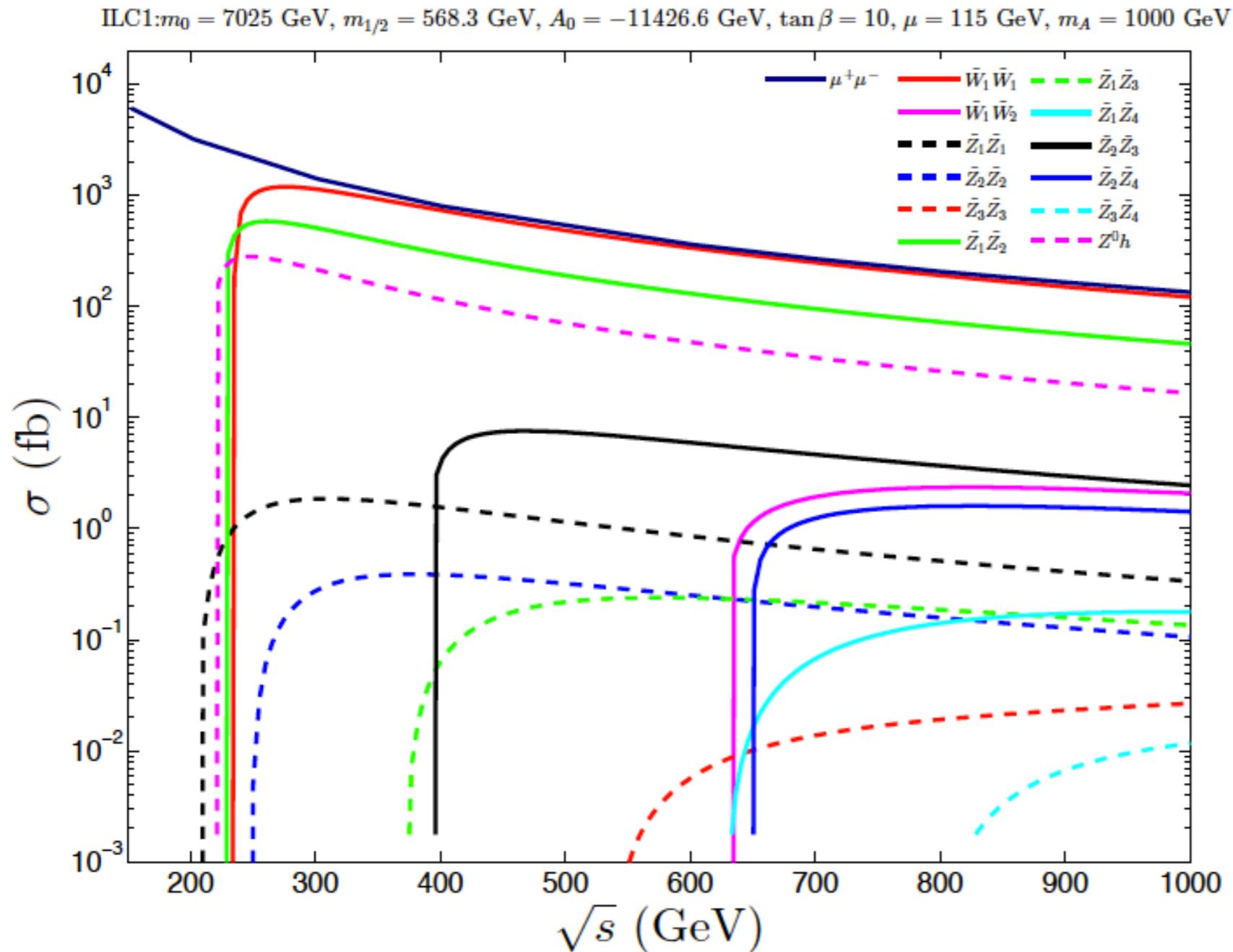


Visually, large fine-tuning has already developed by $\mu \sim 350$ or $\Delta_{EW} \sim 30$

Nature is natural $\Rightarrow \Delta_{EW} < 20 - 30$ (take 30 as conservative)

Smoking gun signature: light higgsinos at ILC:

ILC is Higgs/higgsino factory!



$$\sigma(\text{higgsino}) \gg \sigma(Zh)$$

3-15 GeV higgsino mass
gaps no problem
in clean ILC environment

HB, Barger, Mickelson, Mustafayev, Tata
arXiv:1404.7510

ILC either sees light higgsinos
or natural MSSM dead

SUSY μ problem: μ term is SUSY, not SUSY breaking:
expect $\mu \sim M_{\text{Pl}}$ but phenomenology requires $\mu \sim m(Z)$

- NMSSM: $\mu \sim m(3/2)$; beware singlets!
- Giudice–Masiero: μ forbidden by some symmetry:
generate via Higgs coupling to hidden sector
- **Kim–Nilles**: invoke SUSY version of DFSZ axion
solution to strong CP:

KN: PQ symmetry forbids μ term,
but then it is generated via PQ breaking

Little Hierarchy due to mismatch between
PQ breaking and SUSY breaking scales?

$$\mu \sim \lambda f_a^2 / M_P$$

$$m_{3/2} \sim m_{\text{hid}}^2 / M_P$$

$$f_a \ll m_{\text{hid}}$$

**Higgs mass tells us where
to look for axion!**

$$m_a \sim 6.2 \mu\text{eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$$

Little Hierarchy from radiative PQ breaking? exhibited within context of MSY/CCK model

Murayama, Suzuki, Yanagida (1992);
Gherghetta, Kane (1995)

Choi, Chun, Kim (1996)

Bae, HB, Serce, PRD91 (2015) 015003

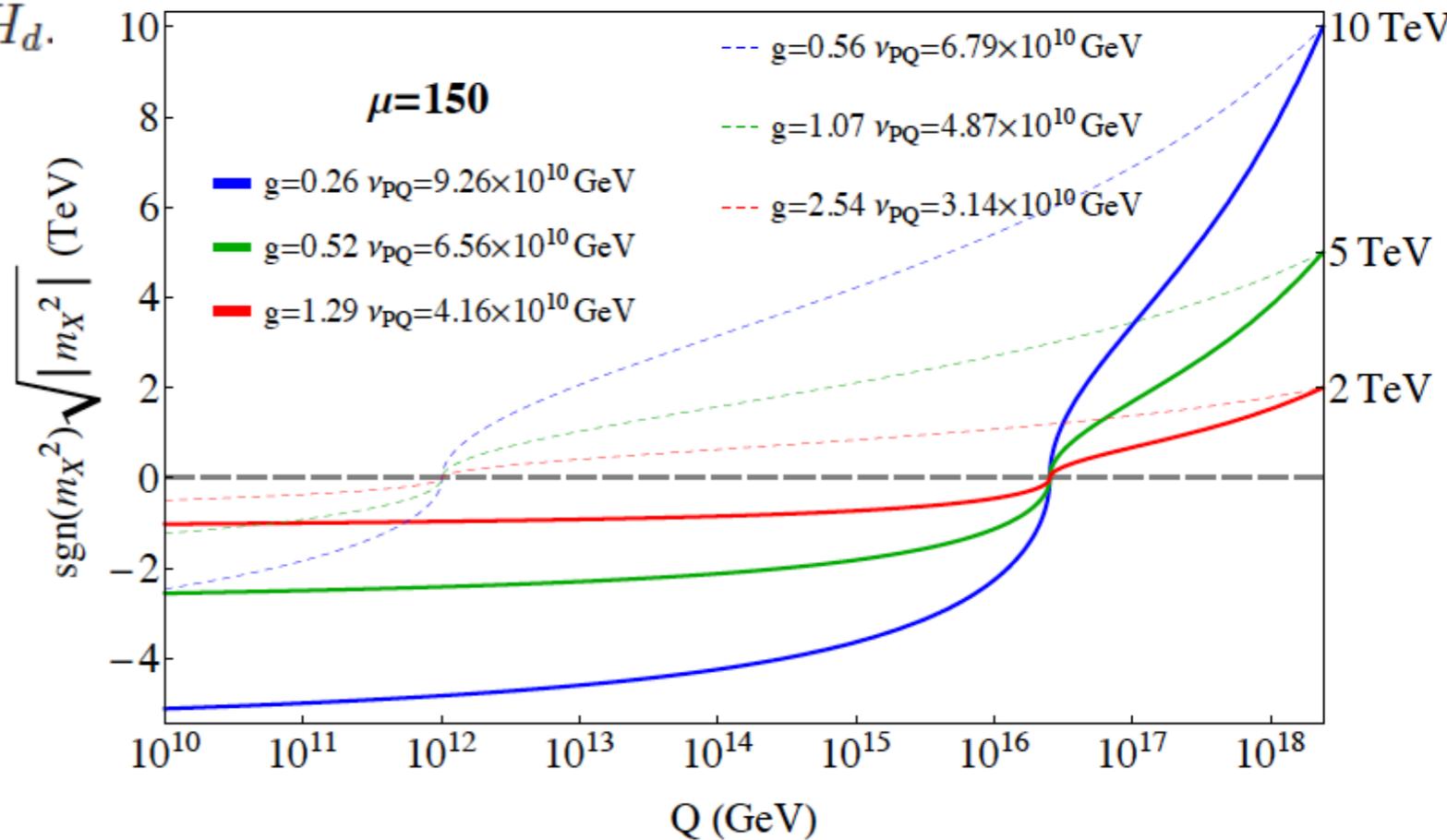
augment MSSM with PQ charges/fields:

$$\hat{f}' = \frac{1}{2} h_{ij} \hat{X} \hat{N}_i^c \hat{N}_j^c + \frac{f}{M_P} \hat{X}^3 \hat{Y} + \frac{g}{M_P} \hat{X} \hat{Y} \hat{H}_u \hat{H}_d.$$

$$M_{N_i^c} = v_X h_i |_{Q=v_X}$$

$$\mu = g \frac{v_X v_Y}{M_P}.$$

Intermediate scale SUSY
breaking generates as well
both PQ and Majorana nu scales!



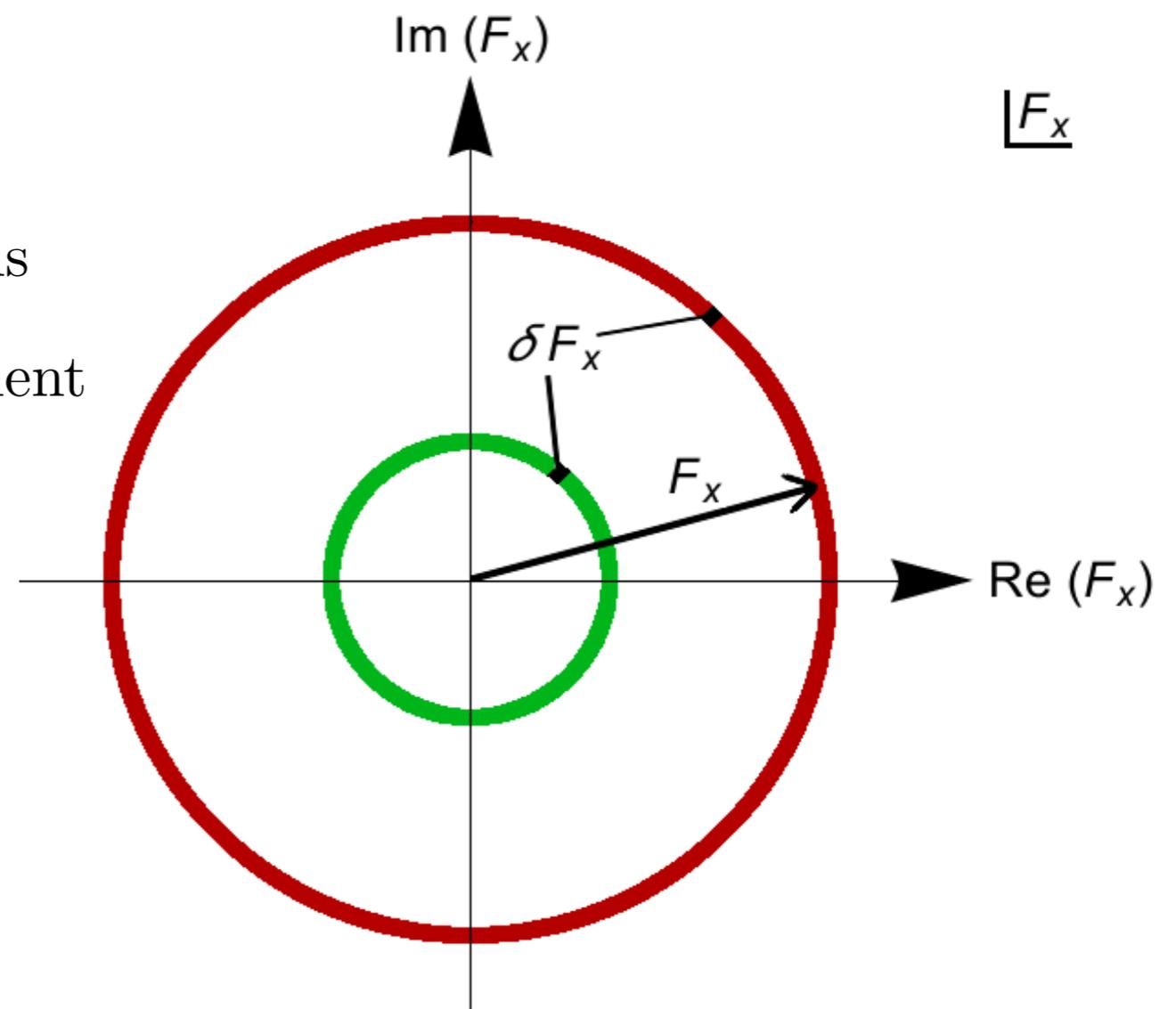
Large $m_{3/2}$ generates small $\mu \sim 100 - 200$ GeV!

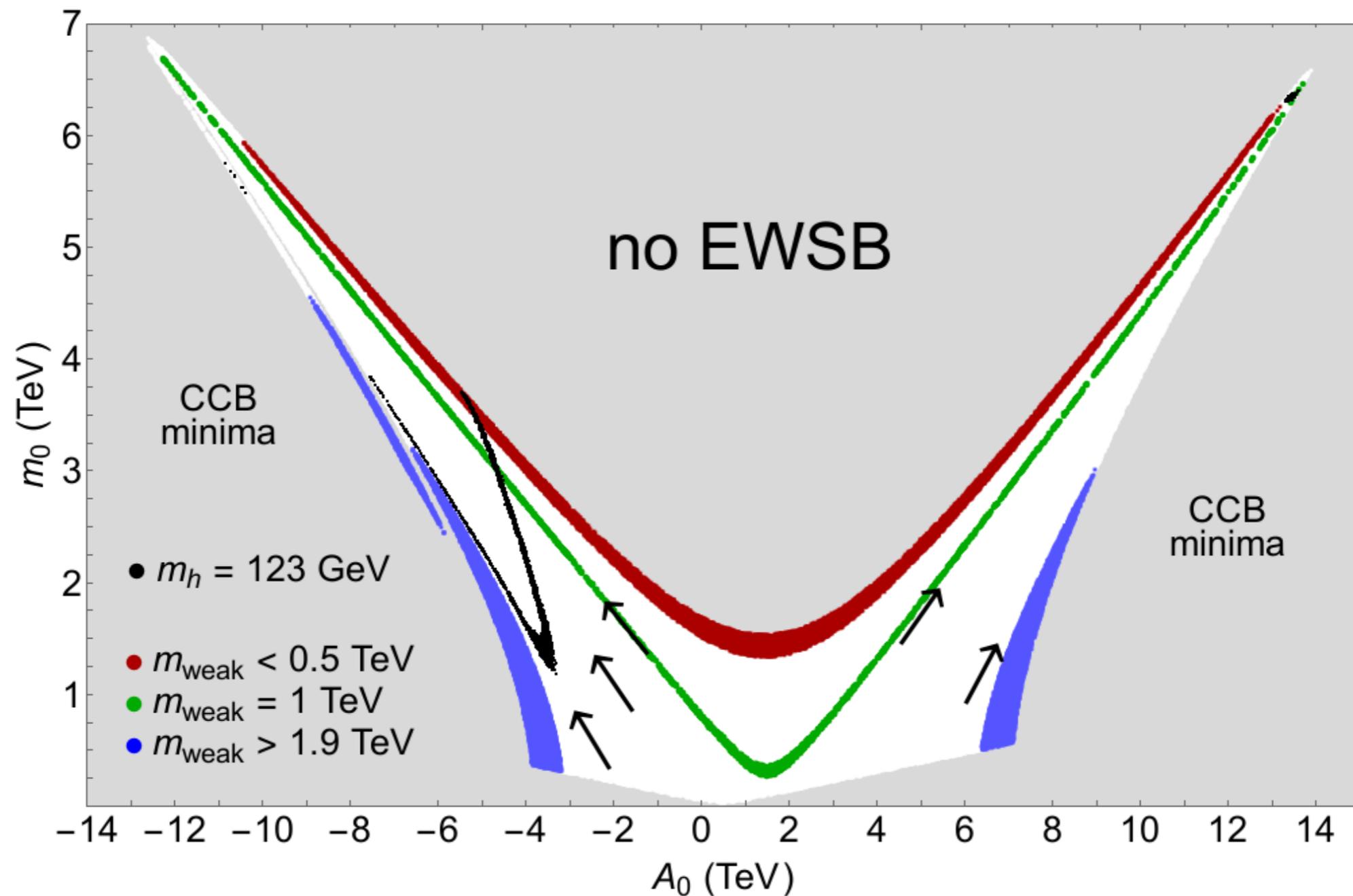
Why do soft terms take on values needed for natural (barely-broken) EWSB? string theory landscape?

- assume model like MSY/CCK where $\mu \sim 100$ GeV
- then $m(\text{weak})^2 \sim |m_{H_u}^2|$
- If all values of SUSY breaking field $\langle F_X \rangle$ equally likely, then mild (linear) statistical draw towards large soft terms
- This is balanced by anthropic requirement of weak scale $m_{\text{weak}} \sim 100$ GeV

Agrawal, Barr, Donoghue and Seckel

Anthropic selection of $m_{\text{weak}} \sim 100$ GeV:
If m_W too large, then weak interactions $\sim (1/m_W^4)$ too weak
weak decays, fusion reactions suppressed
elements not as we know them

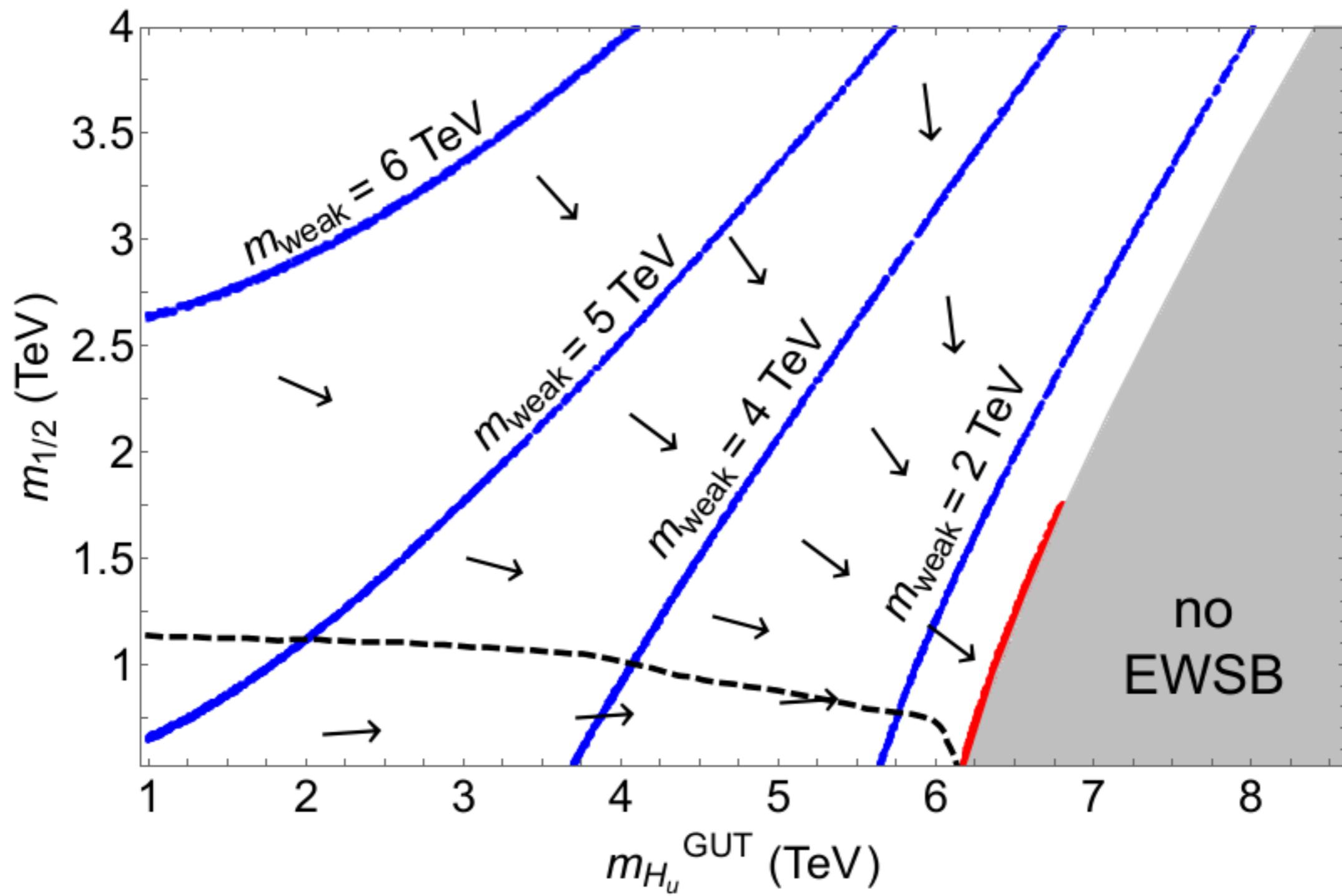




statistical draw to large soft terms balanced by anthropic draw toward red ($m(\text{weak}) \sim 100$ GeV): then $m(\text{Higgs}) \sim 125$ GeV and natural SUSY spectrum!

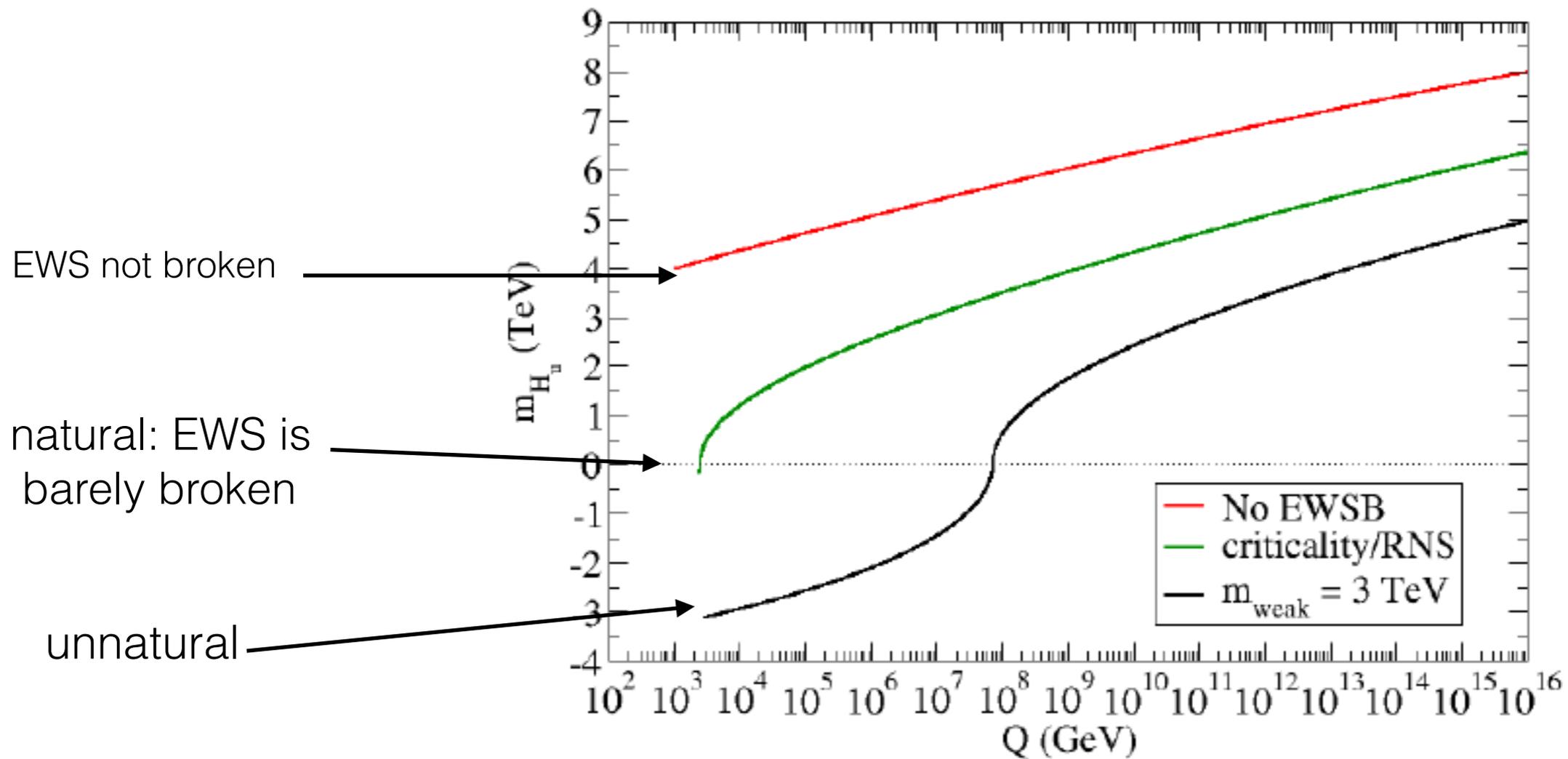
Giudice, Rattazzi, 2006

HB, Barger, Savoy, Serce, PLB758 (2016) 113



statistical/anthropic draw toward FP-like region

radiative corrections drive $m_{H_u}^2$ from unnatural GUT scale values to naturalness at weak scale:
radiatively-driven naturalness



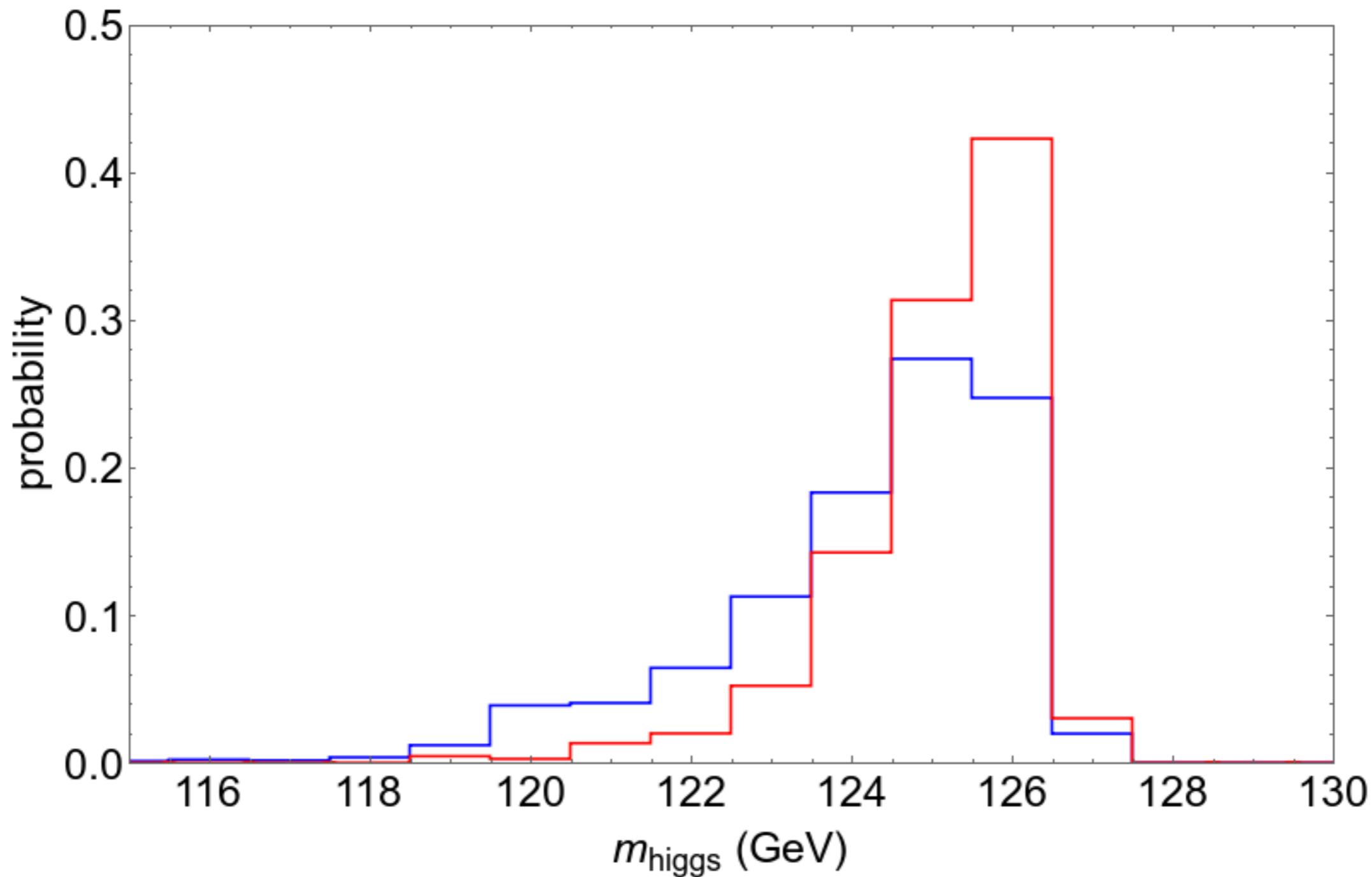
Evolution of the soft SUSY breaking mass squared term $sign(m_{H_u}^2)\sqrt{|m_{H_u}^2|}$ vs. Q

Simple landscape prediction for $m(h)$:

blue: $n=1$

red: $n=2$

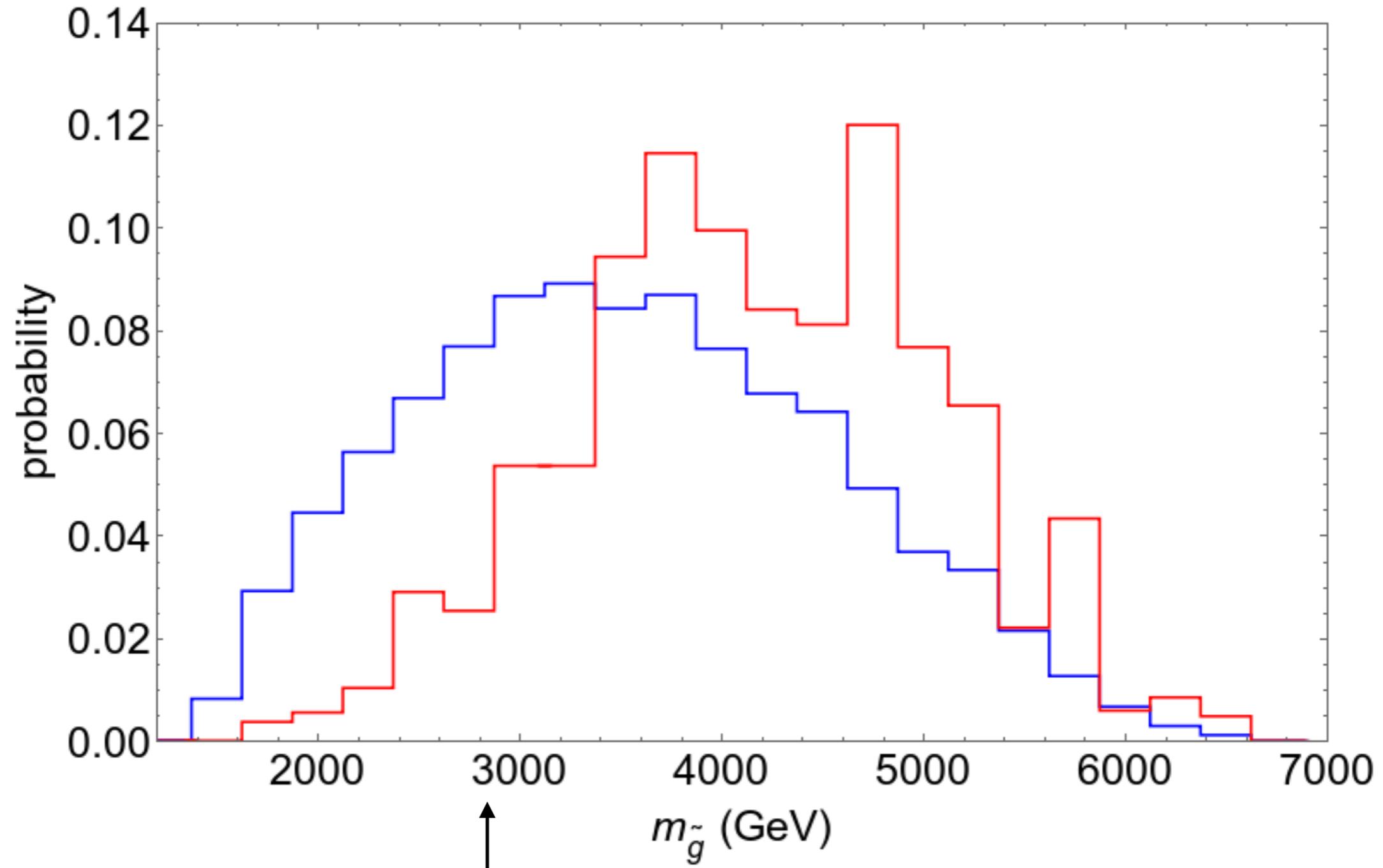
$$P(m_{SUSY}) \sim m_{SUSY}^n$$



Simple landscape prediction for $m(\text{gluino})$

blue: $n=1$

red: $n=2$

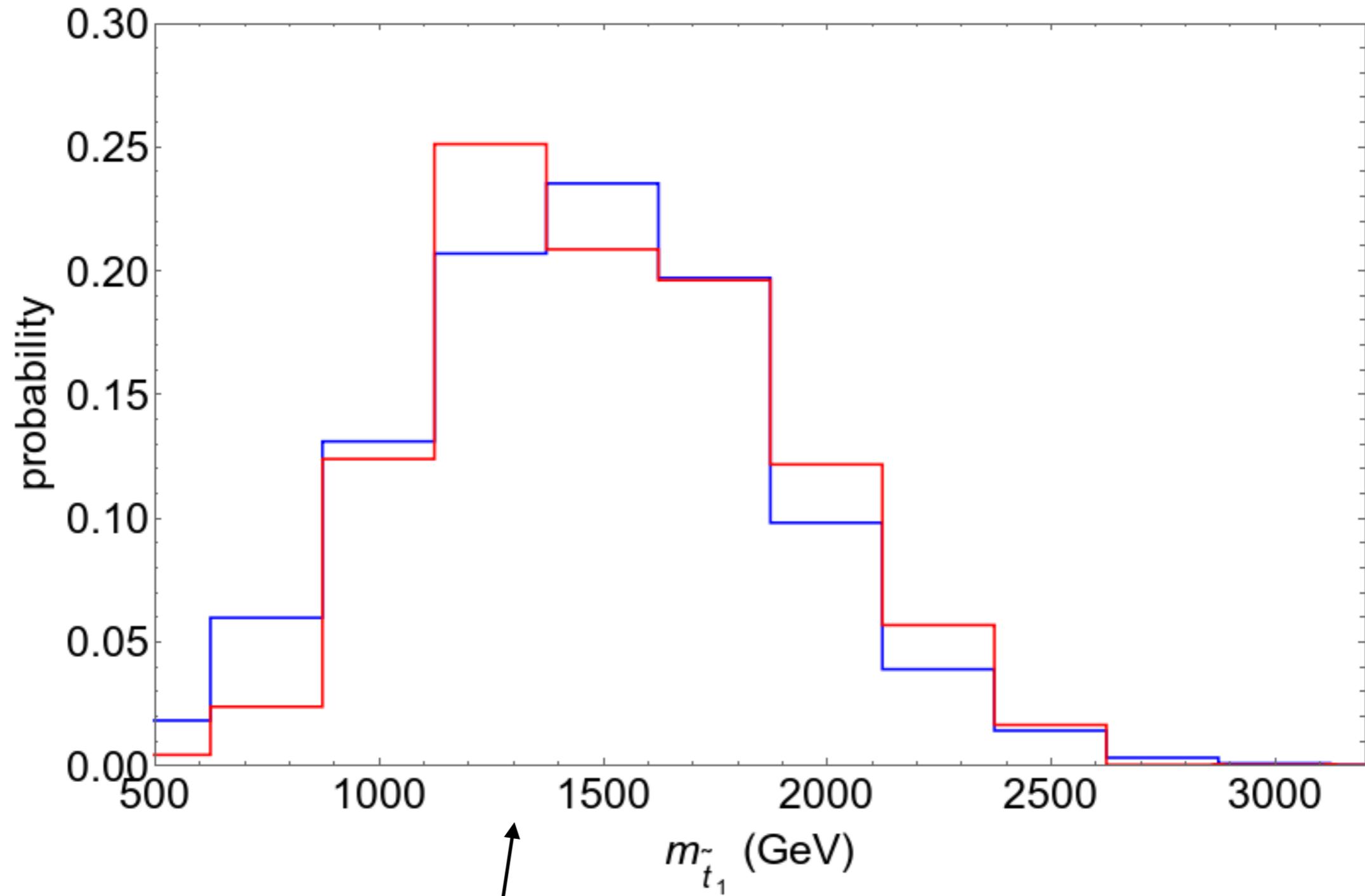


HL-LHC reach is to $m(\text{gluino}) \sim 2.8 \text{ TeV}$

Simple landscape prediction for $m(t_1)$

blue: $n=1$

red: $n=2$



HL-LHC reach is to $m(t_1) \sim 1.2-1.4$ TeV

Mirage mediation: comparable moduli- & anomaly-mediation

Choi, Falkowski, Nilles, Olechowski, Pokorski

Generalized mirage mediation model:

HB, Barger, Serce, Tata: arXiv:1610.06205

$$M_a = (\alpha + b_a g_a^2) m_{3/2} / 16\pi^2, \quad (10)$$

$$A_\tau = (-a_3 \alpha + \gamma_{L_3} + \gamma_{H_d} + \gamma_{E_3}) m_{3/2} / 16\pi^2, \quad (11)$$

$$A_b = (-a_3 \alpha + \gamma_{Q_3} + \gamma_{H_d} + \gamma_{D_3}) m_{3/2} / 16\pi^2, \quad (12)$$

$$A_t = (-a_3 \alpha + \gamma_{Q_3} + \gamma_{H_u} + \gamma_{U_3}) m_{3/2} / 16\pi^2, \quad (13)$$

$$m_i^2(1,2) = (c_m \alpha^2 + 4\alpha \xi_i - \dot{\gamma}_i) (m_{3/2} / 16\pi^2)^2, \quad (14)$$

$$m_j^2(3) = (c_{m3} \alpha^2 + 4\alpha \xi_j - \dot{\gamma}_j) (m_{3/2} / 16\pi^2)^2, \quad (15)$$

$$m_{H_u}^2 = (c_{H_u} \alpha^2 + 4\alpha \xi_{H_u} - \dot{\gamma}_{H_u}) (m_{3/2} / 16\pi^2)^2, \quad (16)$$

$$m_{H_d}^2 = (c_{H_d} \alpha^2 + 4\alpha \xi_{H_d} - \dot{\gamma}_{H_d}) (m_{3/2} / 16\pi^2)^2, \quad (17)$$

elevate $a_3, c_m, c_{m3}, c_{H_u}, c_{H_d}$ from discrete to continuous:
soft terms depend on location of fields in compactified manifold!

p-space: $\alpha, m_{3/2}, c_m, c_{m3}, a_3, c_{H_u}, c_{H_d}, \tan \beta$ (GMM)

$\alpha, m_{3/2}, c_m, c_{m3}, a_3, \tan \beta, \mu, m_A$ (GMM'). \Leftarrow

allows for natural mirage mediation

Allows to generate **mini-landscape** spectra:
heterotic string compactified on Z_6 -II orbifold

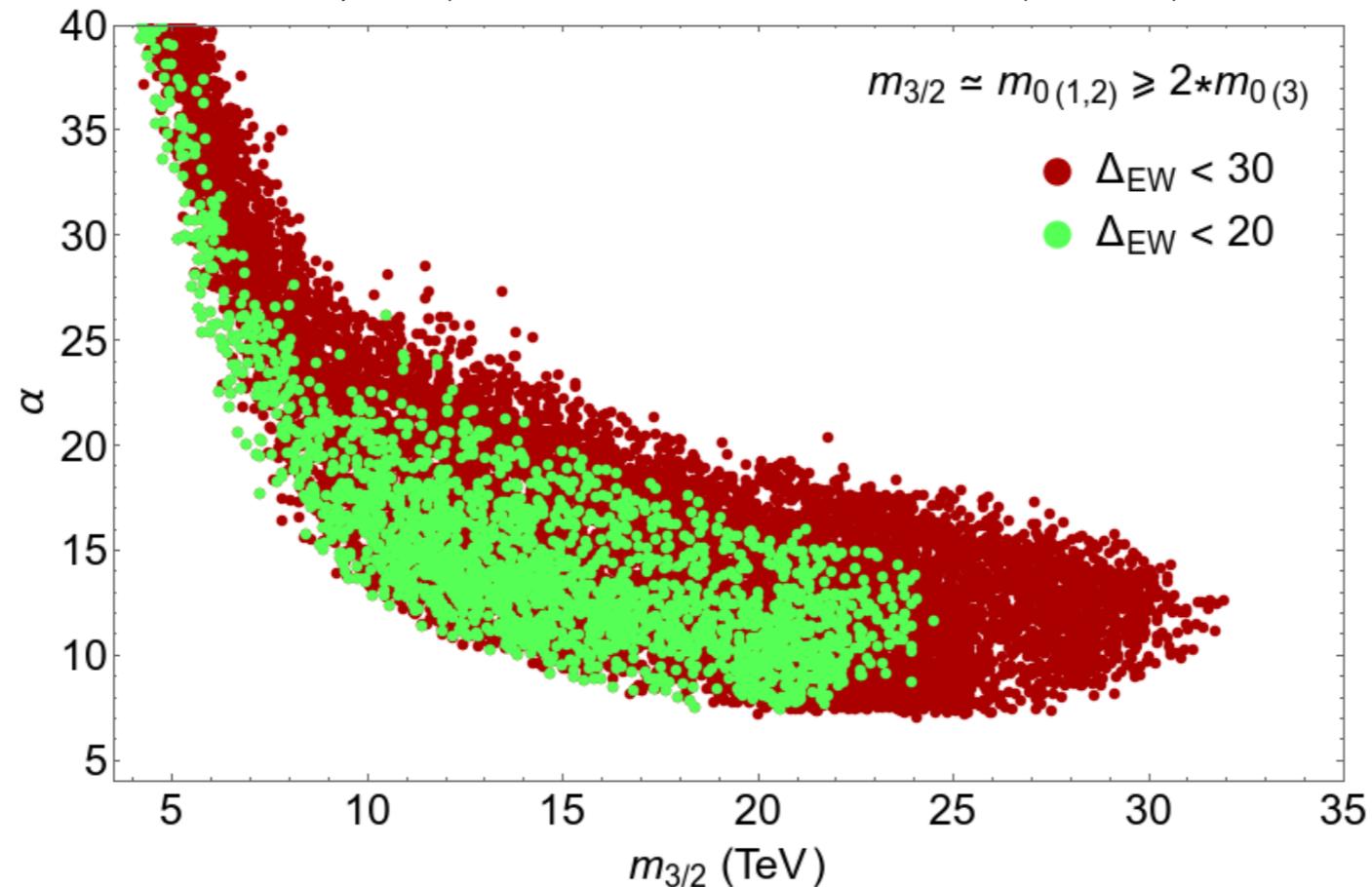
Lebedev, Nilles, Raby, Ramos-Sanches, Ratz, Vaudrevange
but with radiatively-driven naturalness

HB, Barger, Savoy, Serce, Tata, arXiv:1705.01578

- Look for fertile patch of landscape giving MSSM
- 1,2 gen lives on orbifold fixed points/tori: in 16 of $SO(10)$
- 3rd gen, Higgs, gauge live more in bulk: split multiplets
- $m(1,2) \sim m(3/2) \sim 10-30$ TeV
- $m(3) \sim m(H) \sim A's \sim m(\text{inos}) \sim 1-3$ TeV
- soft terms that of mirage mediation
- programmed Isajet 7.86

To generate mini-landscape, take:

$$c_m = (16\pi^2/\alpha)^2 \text{ so that } m_0(1,2) \simeq m_{3/2}$$



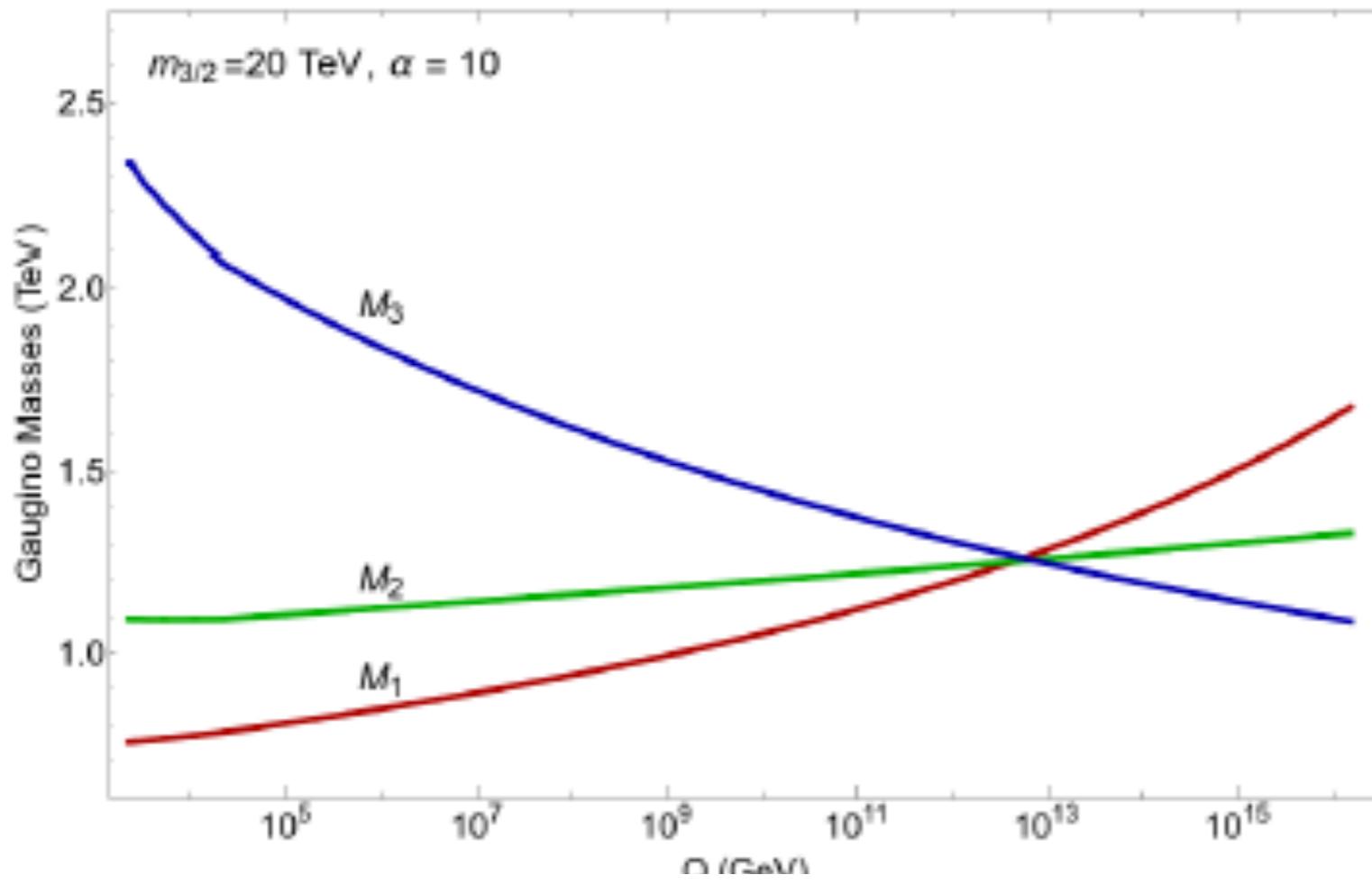
Then get upper bound $m_{3/2} < 25 - 30$ TeV and $\alpha > 7$
else too large $m_0(1,2)$ drives 3rd generation tachyonic

Martin, Vaughn, 2-loop RGEs

Increased upper bound on $m(\text{gluino}) < 6$ TeV

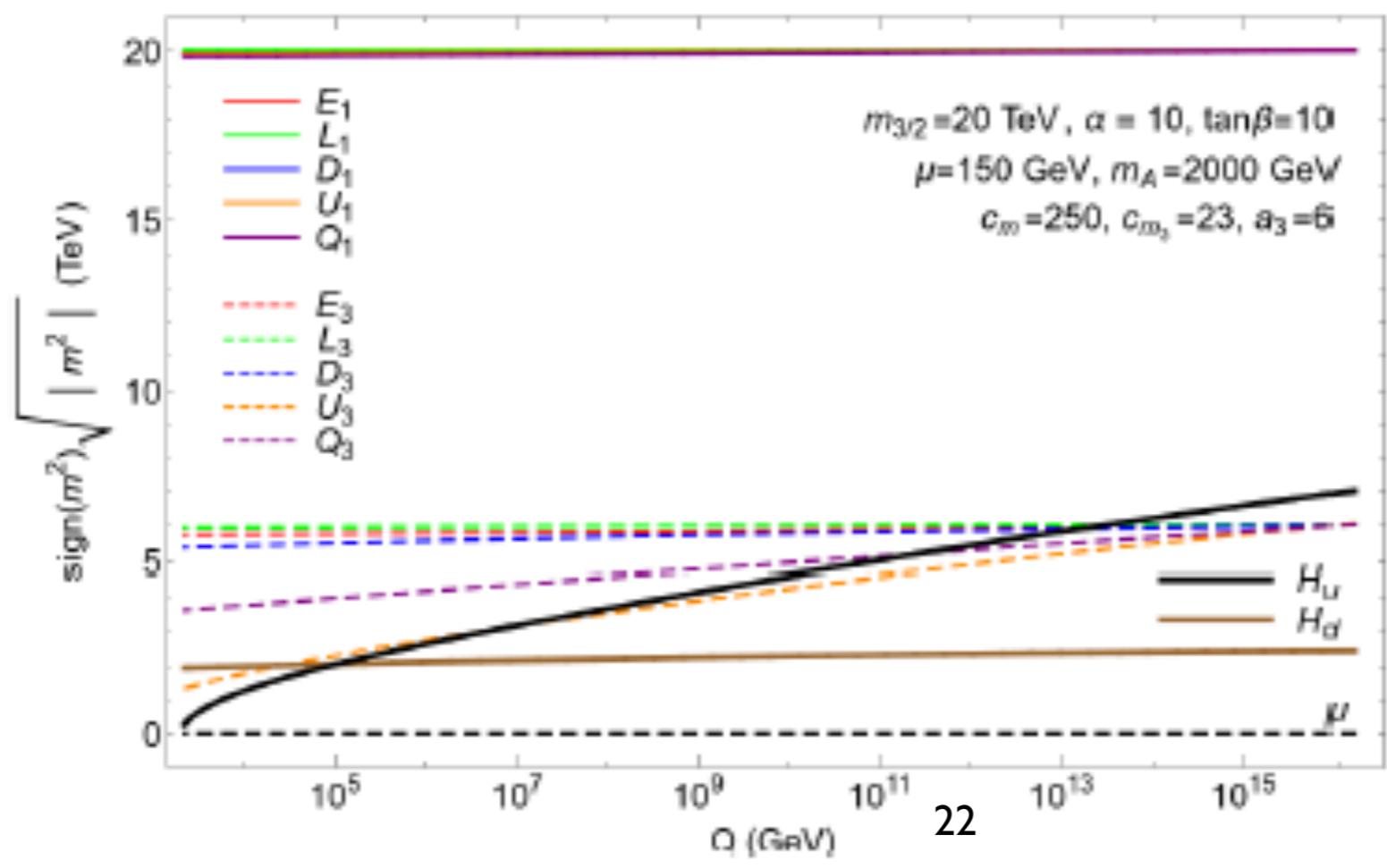
Alpha bound \Rightarrow mirage unif scale $> 10^{11}$ GeV

(not too much compression of inos)

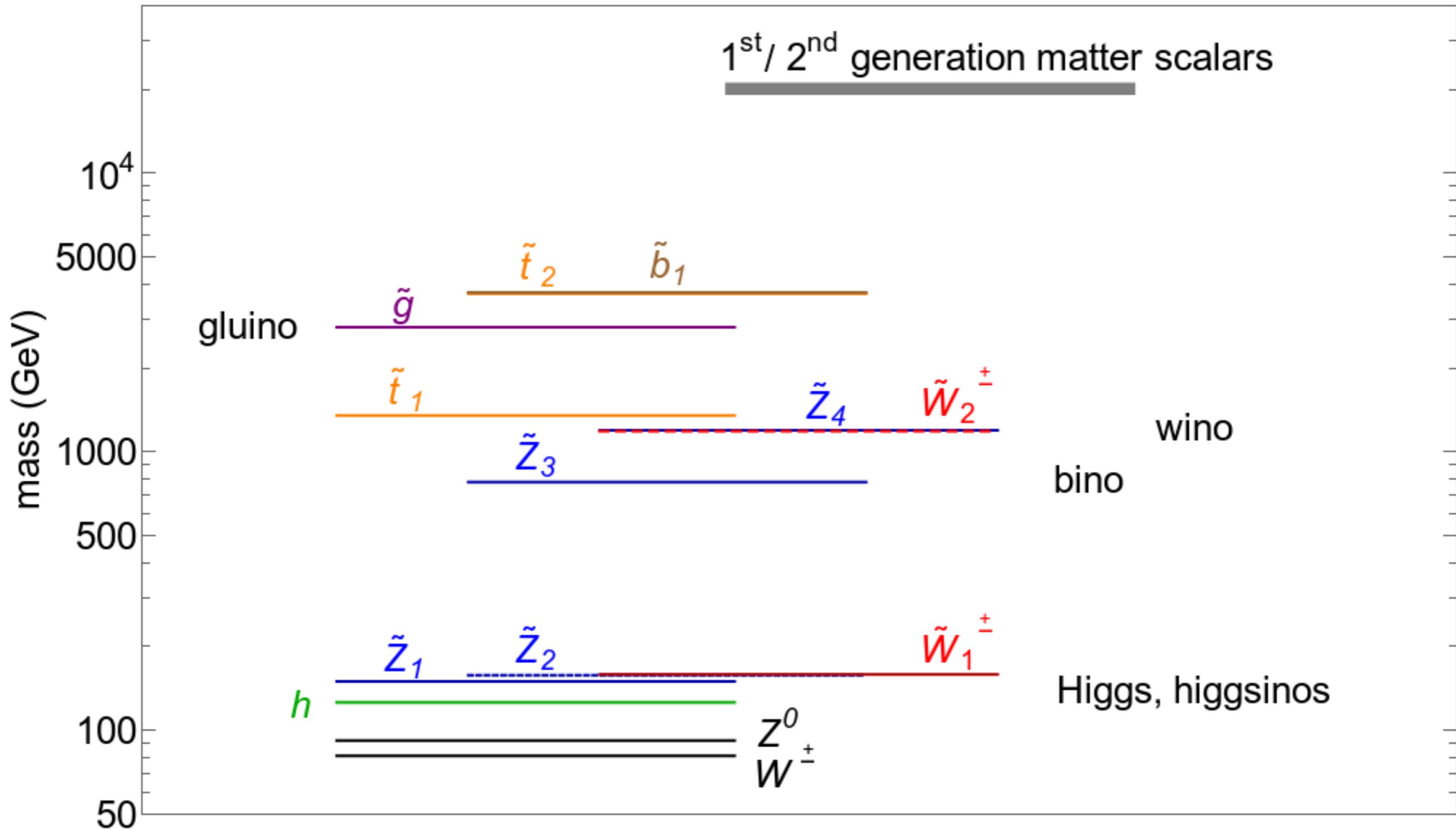


$$\Delta_{EW} = 17.6$$

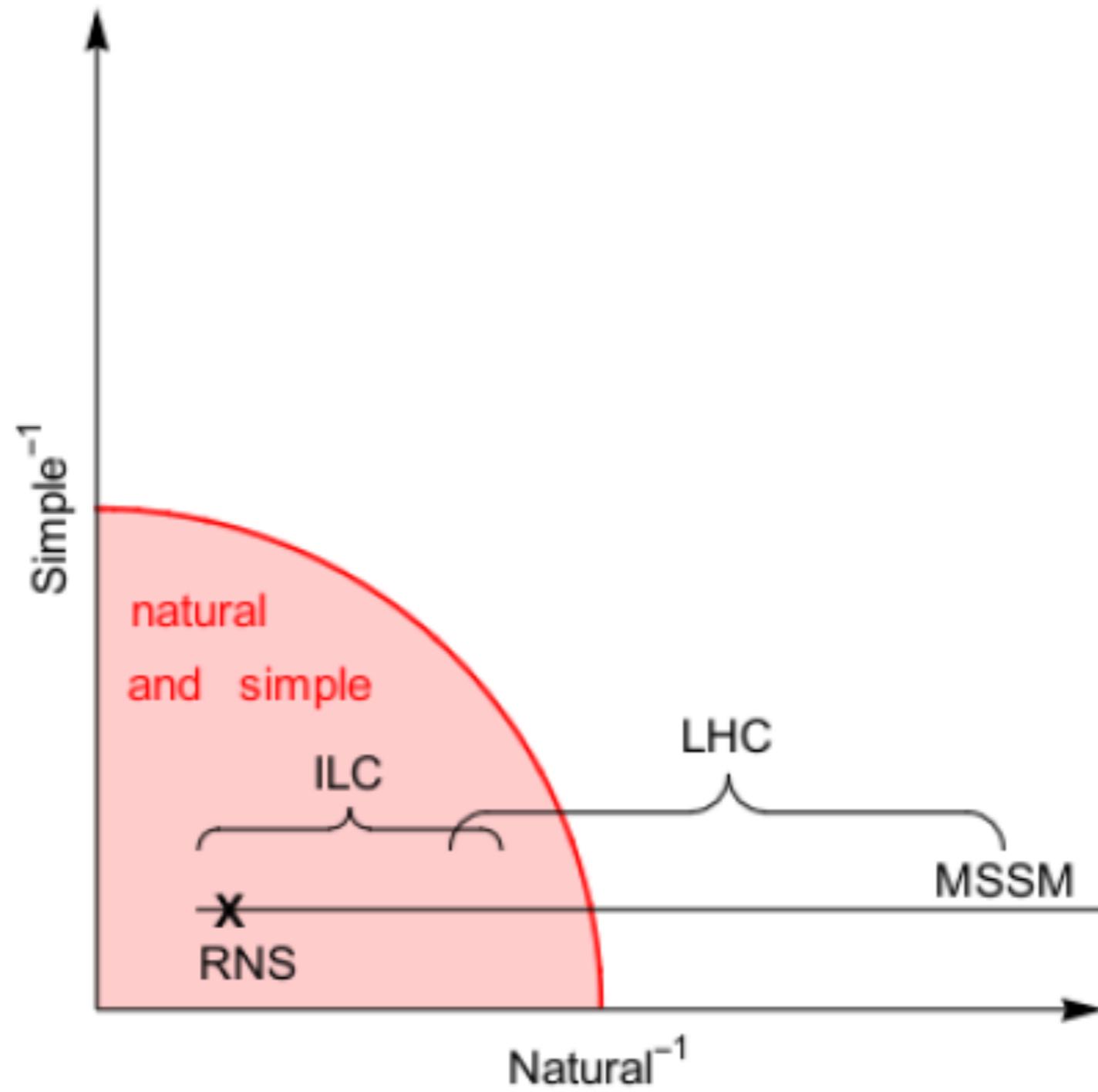
$$m_h \simeq 125 \text{ GeV}$$



mass spectrum for mini2 benchmark point



SUSY may not be accessible to HL-LHC;
 must have ILC to discover required light higgsinos



Conclusions

- ILC250 has strong motivation for discovering new physics via virtual effects
- But case for direct production of superparticles (higgsinos) at ILC highly motivated by naturalness: **this is most plausible rendition of SUSY and in dire need of ILC!**
- small μ related to PQ mechanism: unification of intermediate scales: SUSYbreaking, PQ, Majorana
- string landscape favors large soft terms/barely broken EWS: get mh right but favor high sparticle masses except $\mu \sim 100\text{-}200$ GeV
- UV completion for MSSM within heterotic string compactified on Z6-II orbifold: mini-landscape picture with localized SUSY GUTs and natural with light higgsinos
- imprint of higgsino mass splitting tells whether gaugino masses unify or else mirage unification!
- (dark matter: in this case, likely axion-higgsino-like WIMP admixture)

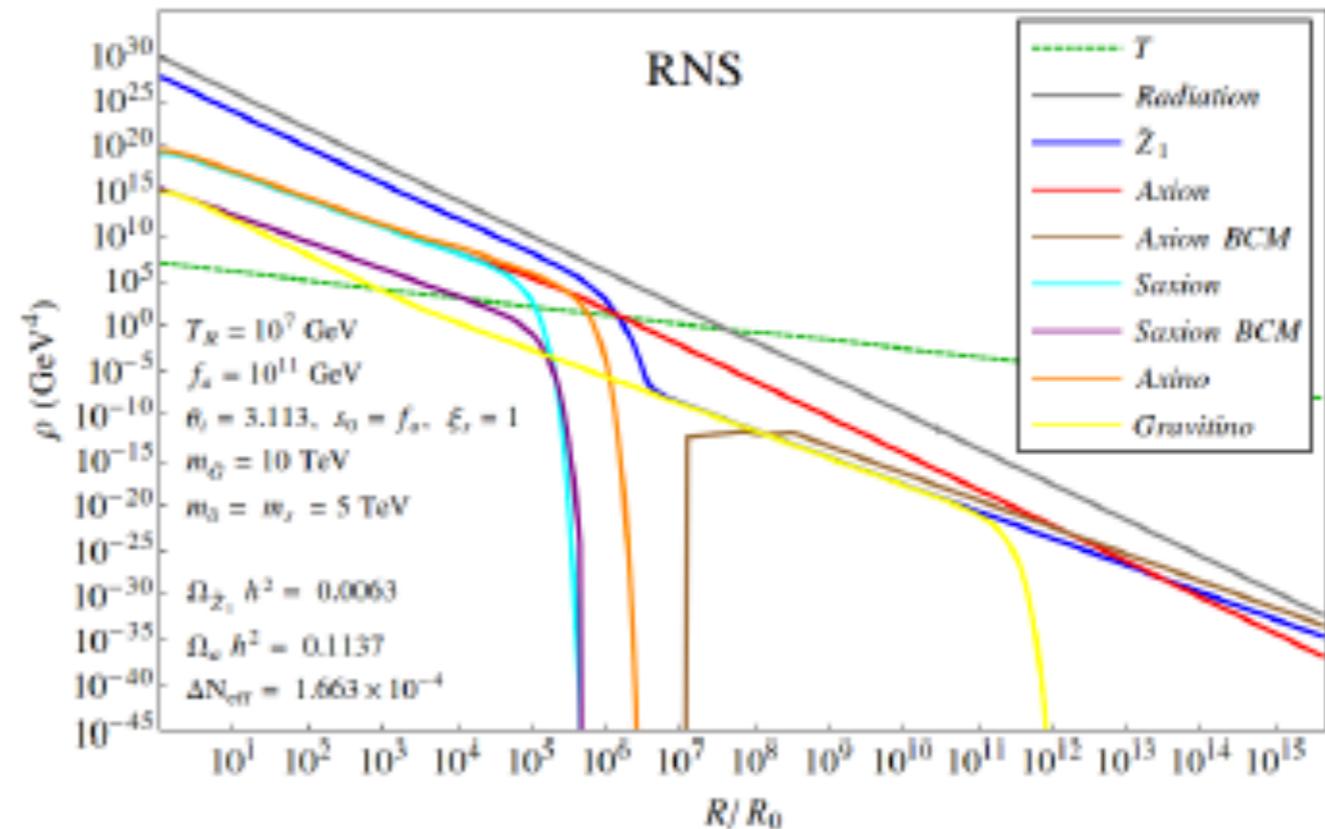
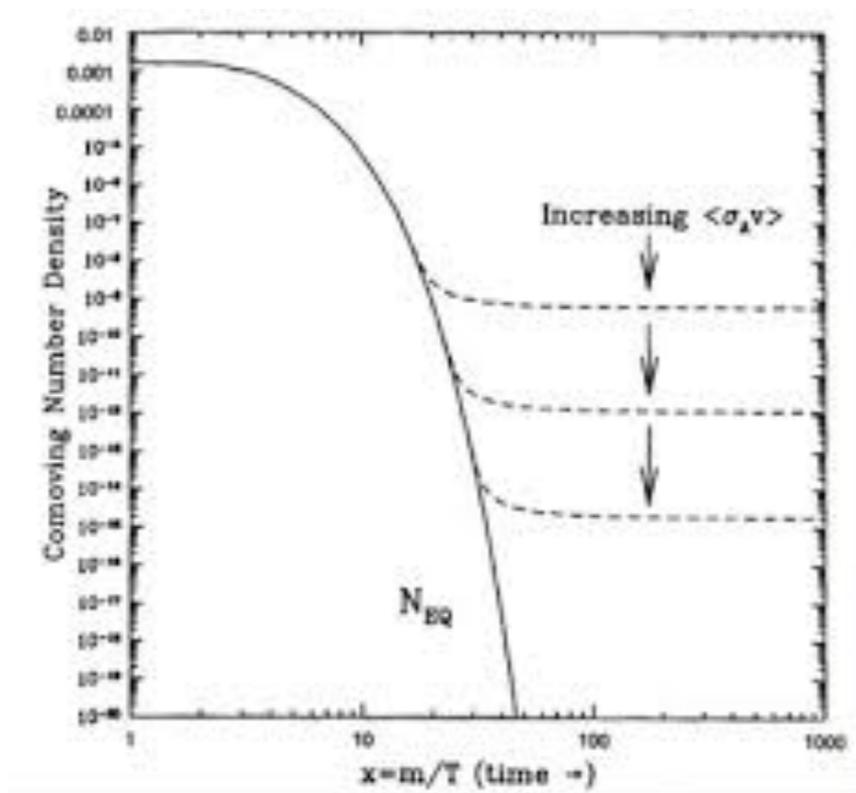
What happens to SUSY WIMP dark matter?

- higgsino-like WIMPs thermally underproduced
- 3 not four light pions \Rightarrow QCD theta vacuum
- EDM(neutron) \Rightarrow axions: no fine-tuning in QCD sector
- SUSY context: axion superfield, axinos and saxions
- DM= axion+higgsino-like WIMP admixture
- DFSZ SUSY axion: solves mu problem with $\mu \ll m_{3/2}$!
- ultimately detect both WIMP and axion?

usual picture

=>

mixed axion/WIMP



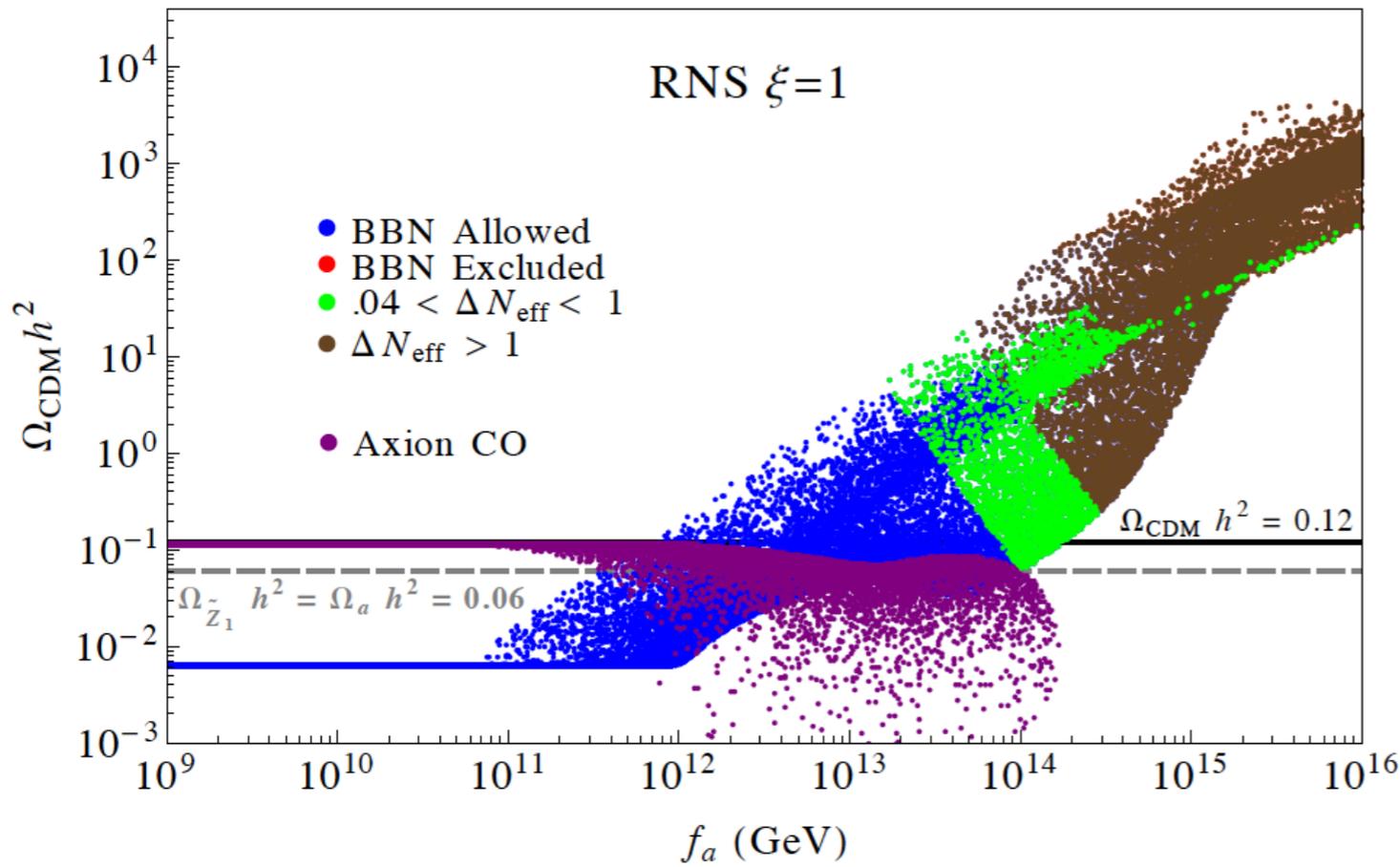
KJ Bae, HB, Lessa, Serce

much of parameter space is axion-dominated
with 10-15% WIMPs



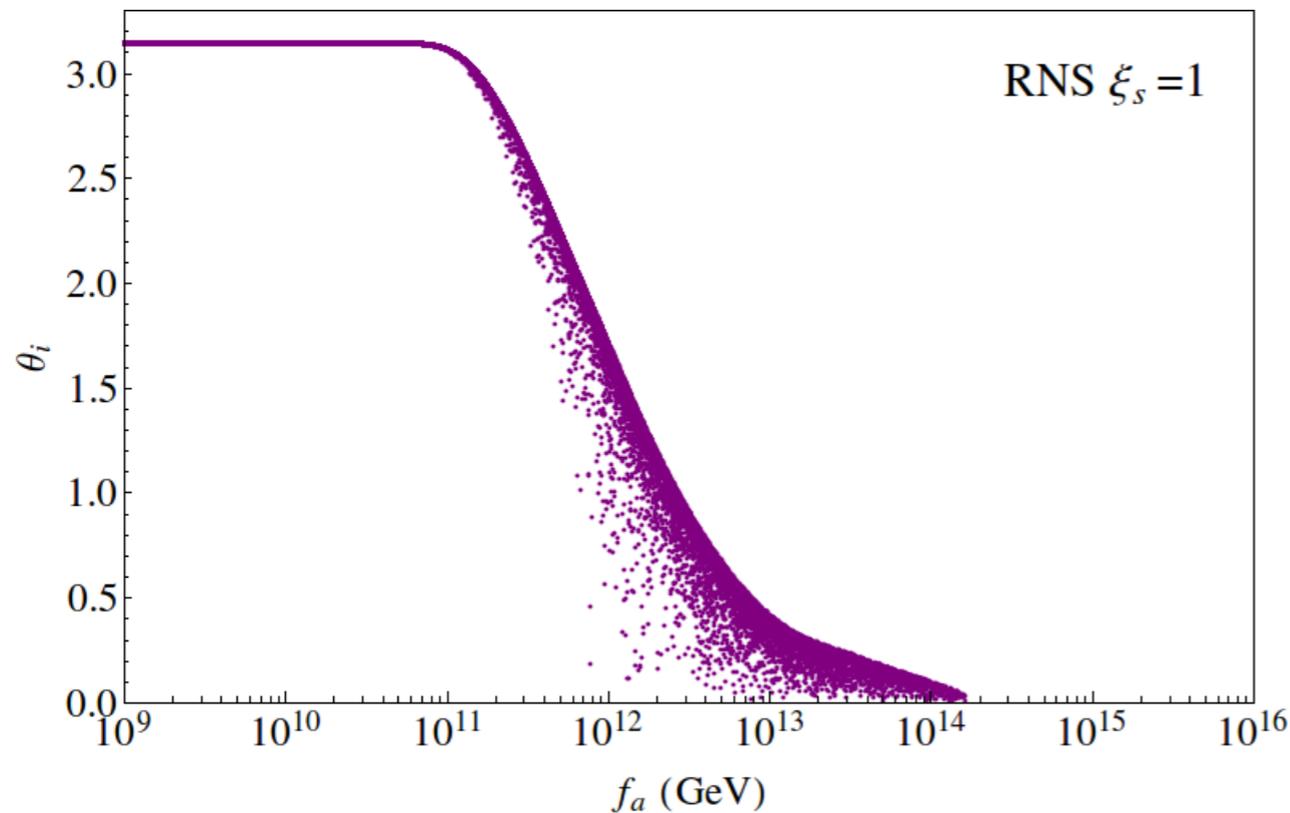
\Rightarrow





higgsino abundance

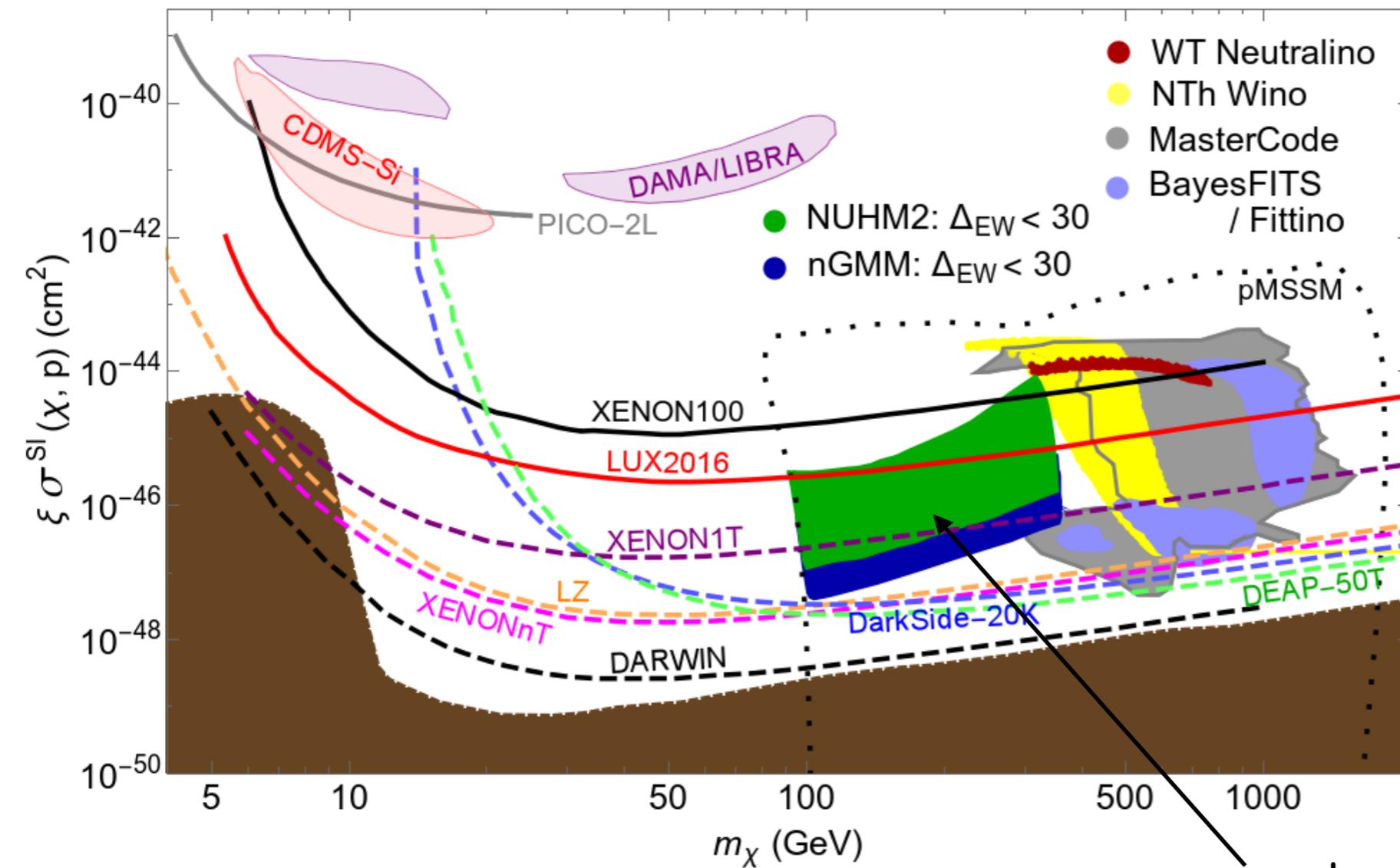
axion abundance



mainly axion CDM
 for $f_a < \sim 10^{12}$ GeV;
 for higher f_a , then
 get increasing wimp
 abundance

Direct higgsino detection rescaled

for minimal local abundance $\xi \equiv \Omega_{\chi}^{TP} h^2 / 0.12$



Bae, HB, Barger, Savoy, Serce

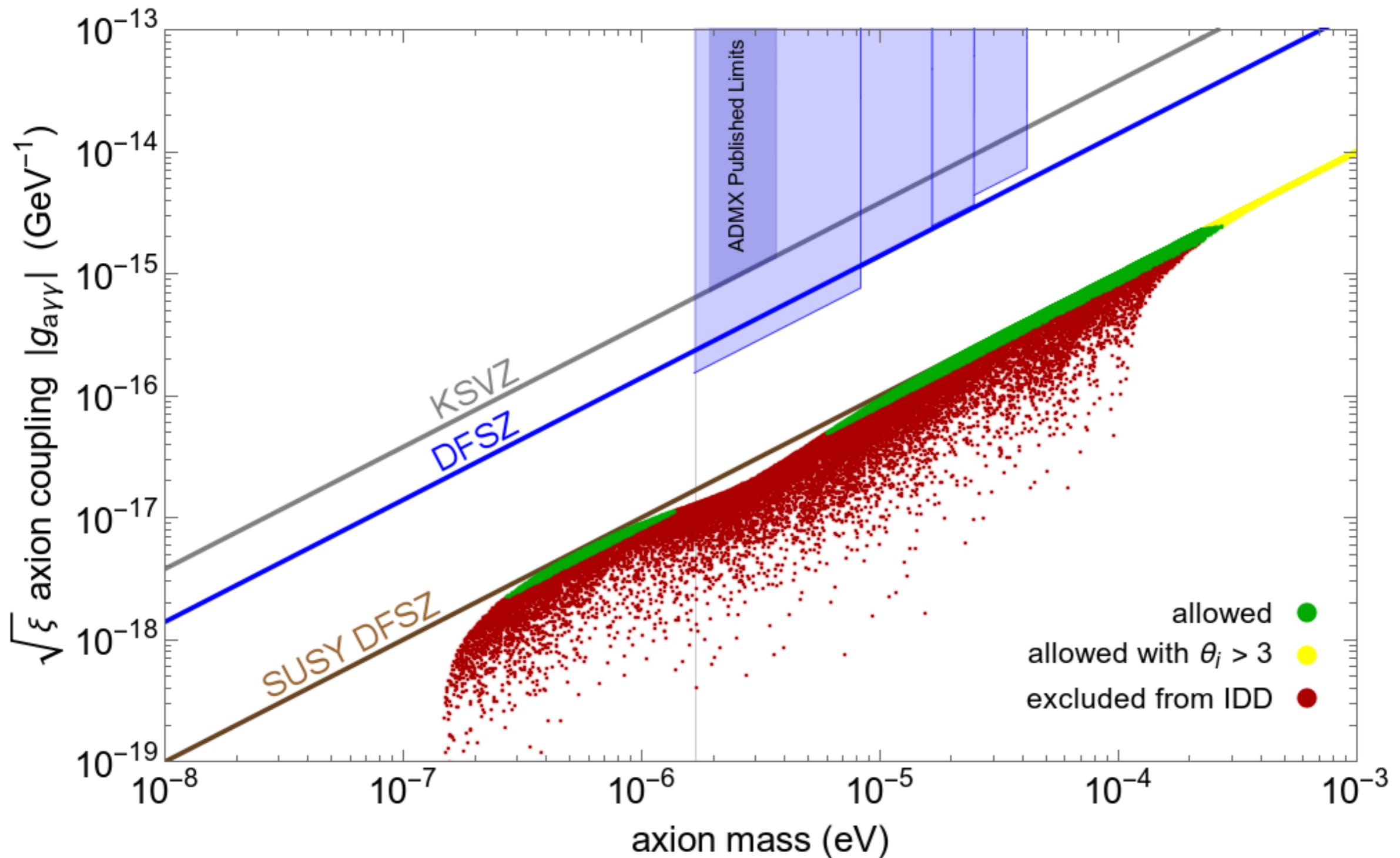
$$\mathcal{L} \ni -X_{11}^h \bar{\tilde{Z}}_1 \tilde{Z}_1 h$$

$$X_{11}^h = -\frac{1}{2} (v_2^{(1)} \sin \alpha - v_1^{(1)} \cos \alpha) (g v_3^{(1)} - g' v_4^{(1)})$$

Xe-1-ton
now operating!

natural SUSY

Can test completely with ton scale detector
or equivalent (subject to minor caveats)



SUSY DFSZ axion: large range in $m(a)$ but coupling reduced
 may need to probe broader and deeper!

