

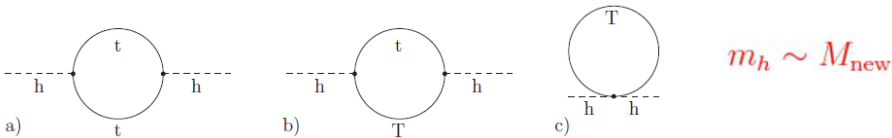
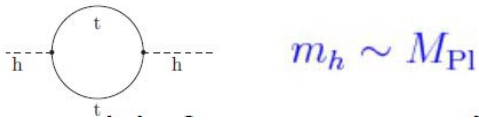
Higgs Coupling Precision at the ILC Using the EFT Formalism

Tim Barklow (SLAC)
ILC Project Meeting, DESY
September 15, 2017

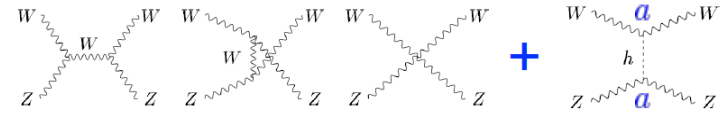
This is a talk about achieving percent level precision for many Higgs couplings at the ILC.

Why is % level Higgs coupling accuracy so interesting?

The Spin-less Higgs



Higgs and Unitarity



$$\Lambda \sim \frac{4\pi v}{\sqrt{1-a^2}}$$

$$\Delta g(hWW) = 10\% \quad \longrightarrow \quad \Lambda \sim 8 \text{ TeV}$$

$$\Delta g(hWW) = 1\% \quad \longrightarrow \quad \Lambda \sim 23 \text{ TeV}$$

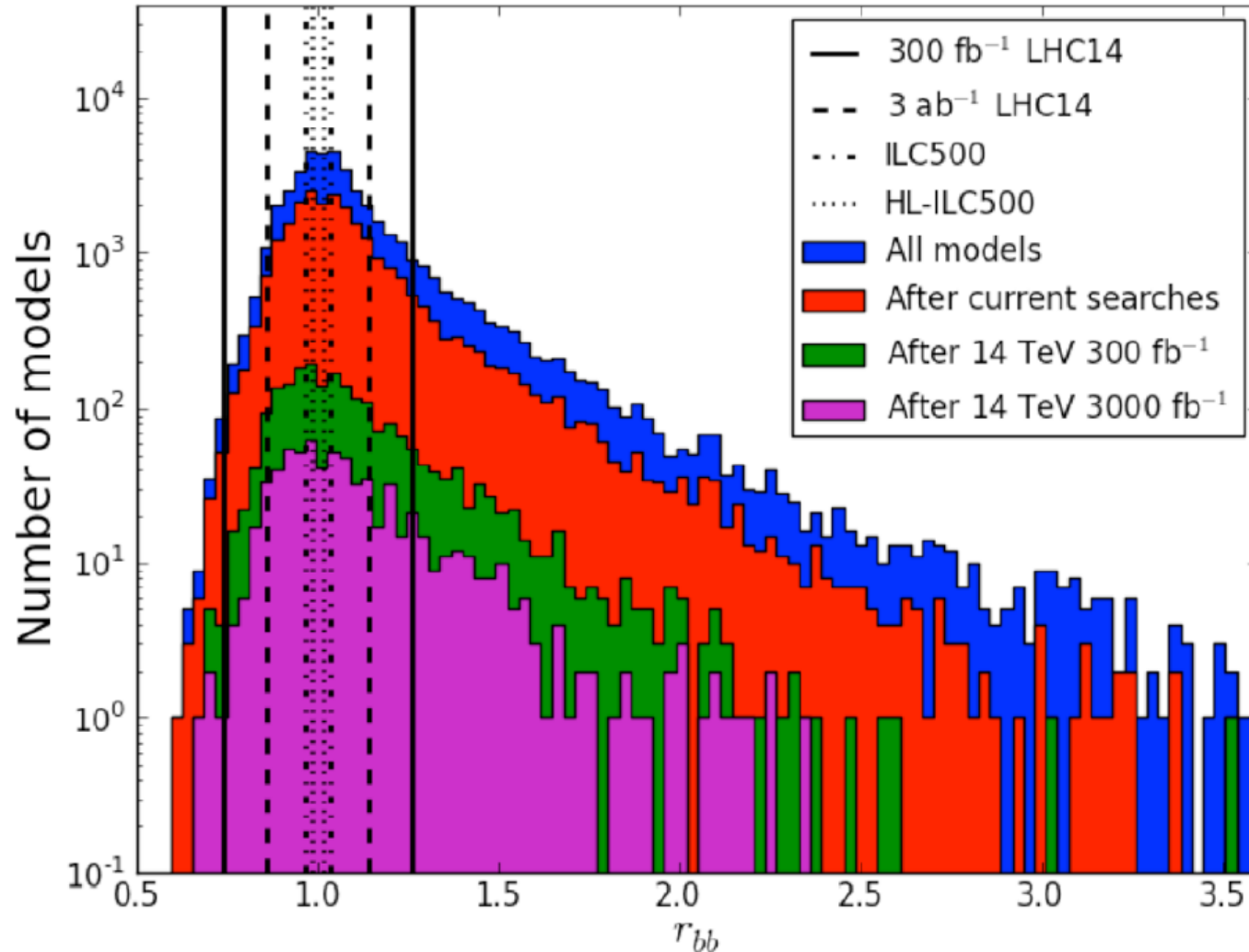
$$\Delta g(hWW) = 0.3\% \quad \longrightarrow \quad \Lambda \sim 42 \text{ TeV}$$

Electroweak Phase Transition

- Electroweak Baryogenesis only works if EWPT is 1st order
- In the SM, EWPT is 2nd order \Rightarrow no EW Baryogenesis
- New particles, coupled to the Higgs could lead to 1st order EWPT
- Almost all 1st order EWPT models predict large shift in Higgs self coupling
- In many 1st order EWPT models the Higgs couplings to gluons, γ 's, W/Z are shifted by 1-5%

Direct LHC Searches and e^+e^- Precision Higgs Couplings Measurements are Complementary. For example, for SUSY:

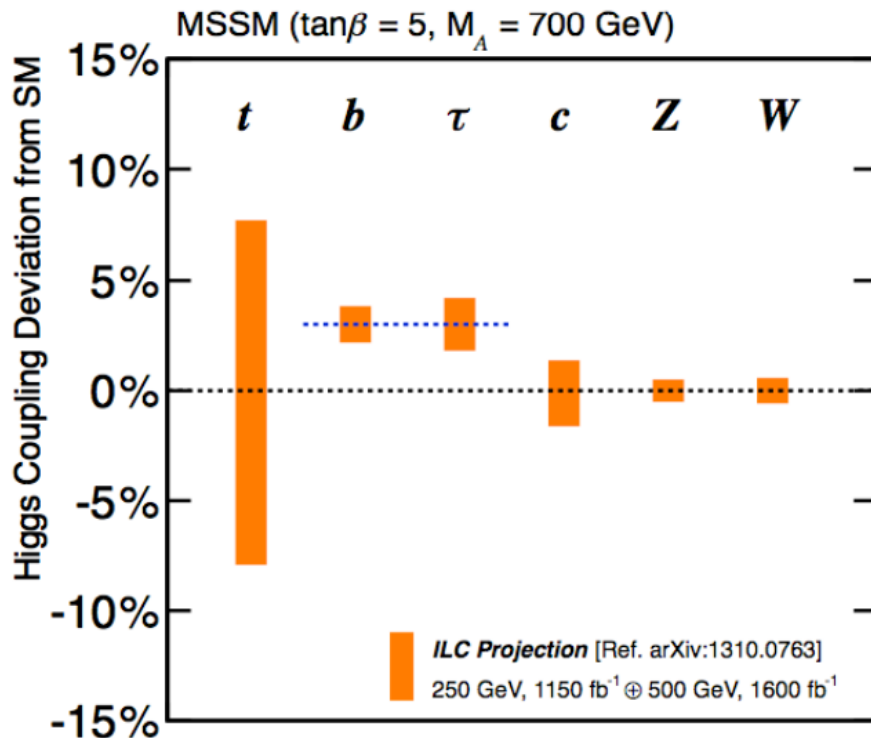
$$\Gamma(h \rightarrow bb)/(SM)$$



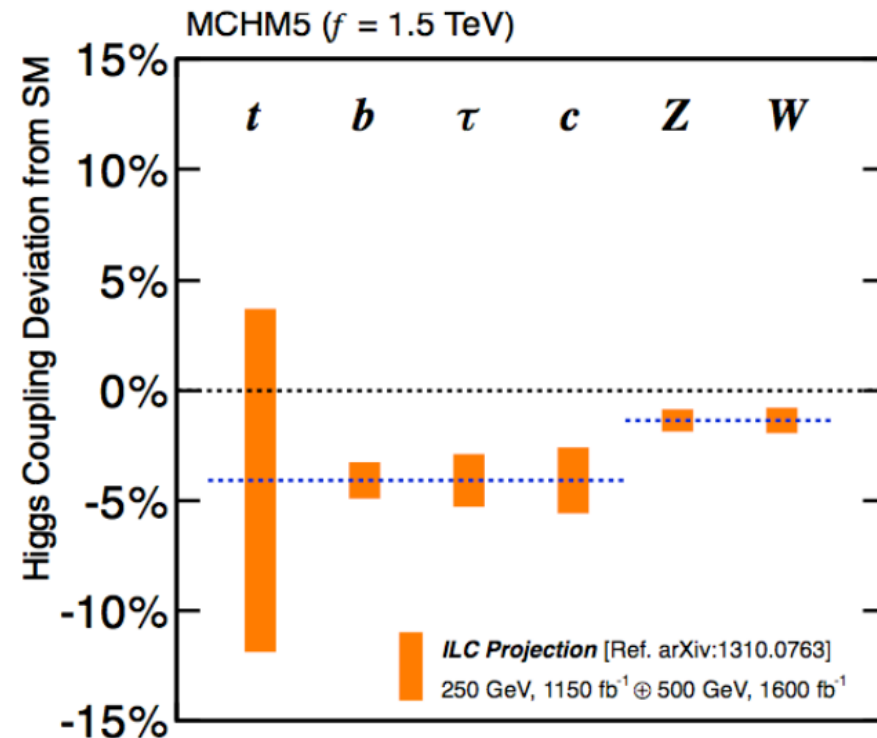
Cahill-Rowley, Hewett, Ismail, Rizzo

Excellent Model Discrimination with e^+e^- Precision Measurements of Many Higgs Couplings

SUSY



Composite Higgs



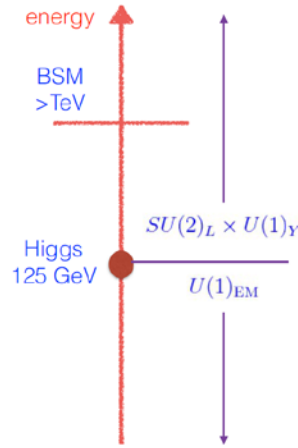
Kanemura, Tsumura, Yagyu, Yokoya

Higgs Effective Field Theory

- LHC results strongly suggest that there is a significant mass gap between the Higgs and BSM particles
- In this situation, BSM corrections to Higgs properties are parametrically small:

$$\delta O \sim m_h^2/M_{\text{BSM}}^2$$

- Moreover, BSM physics must respect the full gauge symmetry of the SM
- Effective Field Theory (EFT) gives a systematic way to parametrize correction to Higgs properties under these conditions, by adding “effective operator” terms to the SM Lagrangian



General $SU(2) \times U(1)$ gauge invariant Lagrangian with dimension-6 operators in addition to the SM.

$$\begin{aligned} \Delta \mathcal{L} = & \frac{c_H}{2v^2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) + \frac{c_T}{2v^2} (\Phi^\dagger \overleftrightarrow{D}^\mu \Phi) (\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) - \frac{c_6 \lambda}{v^2} (\Phi^\dagger \Phi)^3 \\ & + \frac{g^2 c_{WW}}{m_W^2} \Phi^\dagger \Phi W_{\mu\nu}^a W^{a\mu\nu} + \frac{4gg' c_{WB}}{m_W^2} \Phi^\dagger t^a \Phi W_{\mu\nu}^a B^{\mu\nu} \\ & + \frac{g'^2 c_{BB}}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g^3 c_{3W}}{m_W^2} \epsilon_{abc} W_{\mu\nu}^a W^{b\nu\rho} W^{c\rho\mu} \\ & + i \frac{c_{HL}}{v^2} (\Phi^\dagger \overleftrightarrow{D}^\mu \Phi) (\bar{L} \gamma_\mu L) + 4i \frac{c'_{HL}}{v^2} (\Phi^\dagger t^a \overleftrightarrow{D}^\mu \Phi) (\bar{L} \gamma_\mu t^a L) \\ & + i \frac{c_{HE}}{v^2} (\Phi^\dagger \overleftrightarrow{D}^\mu \Phi) (\bar{e} \gamma_\mu e) - c_{\tau\Phi} \frac{y_\tau}{v^2} (\Phi^\dagger \Phi) \bar{L}_3 \cdot \Phi \tau_R + h.c. \end{aligned}$$

After EWSB

$$\begin{aligned}
\Delta\mathcal{L}_h = & \frac{1}{2}\partial_\mu h\partial^\mu h - \frac{1}{2}m_h^2 h^2 - (1 + \eta_h)\bar{\lambda}vh^3 + \frac{\theta_h}{v}h\partial_\mu h\partial^\mu h \\
& + (1 + \eta_W)\frac{2m_W^2}{v}W_\mu^+W^{-\mu}h + (1 + \eta_{WW})\frac{m_W^2}{v^2}W_\mu^+W^{-\mu}h^2 \\
& + (1 + \eta_Z)\frac{m_Z^2}{v}Z_\mu Z^\mu h + \frac{1}{2}(1 + \eta_{ZZ})\frac{m_Z^2}{v^2}Z_\mu Z^\mu h^2 \\
& + \zeta_W\hat{W}_{\mu\nu}^+\hat{W}^{-\mu\nu}\left(\frac{h}{v} + \frac{1}{2}\frac{h^2}{v^2}\right) + \frac{1}{2}\zeta_Z\hat{Z}_{\mu\nu}\hat{Z}^{\mu\nu}\left(\frac{h}{v} + \frac{1}{2}\frac{h^2}{v^2}\right) \\
& + \frac{1}{2}\zeta_A\hat{A}_{\mu\nu}\hat{A}^{\mu\nu}\left(\frac{h}{v} + \frac{1}{2}\frac{h^2}{v^2}\right) + \zeta_{AZ}\hat{A}_{\mu\nu}\hat{Z}^{\mu\nu}\left(\frac{h}{v} + \frac{1}{2}\frac{h^2}{v^2}\right).
\end{aligned}$$

$$\begin{aligned}
\Delta\mathcal{L}_{ehZ} = & -\frac{g}{2c_w}(c_{HL} - c'_{HL})(\bar{\nu}_L\gamma_\mu\nu_L)Z^\mu\left(1 + 2\frac{h}{v} + \frac{h^2}{v^2}\right) \\
& -\frac{g}{2c_w}(c_{HL} + c'_{HL})(\bar{e}_L\gamma_\mu e_L)Z^\mu\left(1 + 2\frac{h}{v} + \frac{h^2}{v^2}\right) \\
& -\frac{g}{2c_w}(c_{HE})(\bar{e}_R\gamma_\mu e_R)Z^\mu\left(1 + 2\frac{h}{v} + \frac{h^2}{v^2}\right) \\
& \frac{g}{\sqrt{2}}(c'_{HL})(\bar{e}_L\gamma_\mu\nu_L W^{-\mu} + \bar{\nu}_L\gamma_\mu e_L W^{+\mu})\left(1 + 2\frac{h}{v} + \frac{h^2}{v^2}\right)
\end{aligned}$$

$$\begin{aligned}
\Delta\mathcal{L}_{TGC} = & ig_V\left\{V^\mu(\hat{W}_{\mu\nu}^-W^{+\nu} - \hat{W}_{\mu\nu}^+W^{-\nu}) + \kappa_V W_\mu^+W_\nu^-\hat{V}^{\mu\nu}\right. \\
& \left. + \frac{\lambda_V}{m_W^2}\hat{W}_{\mu}^{-\rho}\hat{W}_{\rho\nu}^+\hat{V}^{\mu\nu}\right\},
\end{aligned}$$

similar qqhW, qqhZ contact terms
for hadronic W,Z decays in $h \rightarrow WW, ZZ$

The couplings η_x ζ_x as well as TGC's & EWPO's are functions of the dim 6 operator coefficients c_j .

Some examples:

Higgs couplings

$$\eta_Z = \left(-c_T - \frac{1}{2}c_H - c'_{HL}\right)$$

$$\eta_{2Z} = \left(-5c_T - c_H - 2c'_{HL}\right)$$

$$\eta_W = \left(-\frac{1}{2}c_H - c'_{HL}\right)$$

$$\eta_{2W} = \left(-c_H - c'_{HL}\right) .$$

$$\zeta_W = \zeta_{2W} = 8(c_{WW})$$

$$\zeta_Z = \zeta_{2W} = 8\left(c_0^2 c_{WW} + 2s_0^2 c_{WB} + \frac{s_0^4}{c_0^2} c_{BB}\right)$$

TGC's

$$\Delta_g = \frac{1}{c_0^2 - s_0^2} \left(\frac{1}{2}c_T - c'_{HL} - 8\frac{s_0^2}{c_0^2} c_{WB}\right)$$

$$\Delta_\kappa = +8c_{WB}$$

$$\Delta_\lambda = -6\frac{e_0^2}{s_0^2} c_{3W}$$

EWPO's

$$m_W^2/m_Z^2 = c_0^2 + \frac{c_0^2}{c_0^2 - s_0^2} (c_0^2 c_T - 2s_0^2 (c'_{HL} + 8c_{WB}))$$

$$s_*^2 = s_0^2 + \frac{s_0^2}{c_0^2 - s_0^2} (c'_{HL} + 8c_{WB} - c_0^2 c_T) - \frac{1}{2}c_{HE} - s_0^2 (c_{HL} - c_{HE})$$

The κ framework for parameterizing BSM Physics in Higgs Couplings is not model independent

In the κ framework SM Higgs couplings are scaled by factor κ_i so that

$$\begin{aligned}\sigma(e^+e^- \rightarrow Zh) &\sim \kappa_Z^2 \\ \Gamma(h \rightarrow ZZ^*) &\sim \kappa_Z^2\end{aligned}$$

Relations such as the following that are used to calculate the total Higgs width remain valid for $\kappa_Z \neq 1$:

$$\frac{\sigma(e^+e^- \rightarrow Zh)}{BR(h \rightarrow ZZ^*)} = \frac{\sigma(e^+e^- \rightarrow Zh)}{\Gamma(h \rightarrow ZZ^*)/\Gamma_h} \sim \Gamma_h$$

In the dim 6 EFT framework the Lorentz structure for the hZZ vertex includes the momentum-dependent Z field strength tensor:

$$\delta\mathcal{L} = (1 + \eta_Z) \frac{m_Z^2}{v} h Z_\mu Z^\mu + \zeta_Z \frac{h}{2v} Z_{\mu\nu} Z^{\mu\nu}$$

Integration over phase space gives different dependencies on ζ_Z for cross sections and partial widths:

$$\begin{aligned}\sigma(e^+e^- \rightarrow Zh) &= (SM) \cdot (1 + 2\eta_Z + (5.5)\zeta_Z) \\ \Gamma(h \rightarrow ZZ^*) &= (SM) \cdot (1 + 2\eta_Z - (0.50)\zeta_Z)\end{aligned}$$

and we can't use $\frac{\sigma(e^+e^- \rightarrow Zh)}{BR(h \rightarrow ZZ^*)}$ to extract Γ_h

In practice, since $BR(h \rightarrow ZZ^*)$ is so small, the WW fusion process is used in combination with Higgsstrahlung BR measurements to obtain Γ_h within the κ framework, even at $\sqrt{s} = 250$ GeV

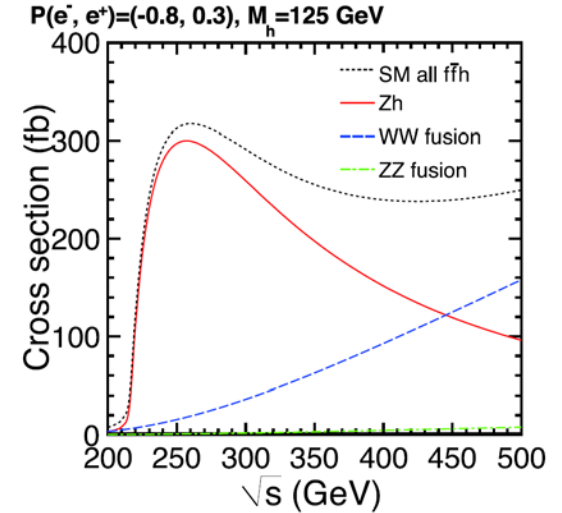
κ framework using WW fusion:

$$\sigma(e^+e^- \rightarrow \nu_e \bar{\nu}_e h) \sim \kappa_W^2$$

$$\Gamma(h \rightarrow WW^*) \sim \kappa_W^2$$

$$\frac{[\sigma(e^+e^- \rightarrow \nu_e \bar{\nu}_e h) \cdot BR(h \rightarrow b\bar{b})]}{BR(h \rightarrow b\bar{b})BR(h \rightarrow WW^*)} = \frac{\sigma(e^+e^- \rightarrow \nu_e \bar{\nu}_e h)}{\Gamma(h \rightarrow WW^*) / \Gamma_h} \sim \Gamma_h$$

This gives a more accurate estimate of Γ_h than $\sigma(e^+e^- \rightarrow Zh)$ and $BR(h \rightarrow ZZ^*)$ at $\sqrt{s} = 250$ GeV, but an even more accurate estimate is made if WW fusion data is collected at $\sqrt{s} = 350$ or 500 GeV.



Within the κ framework it is impossible to achieve better than 2% accuracy on any Higgs coupling (other than hZZ) by running solely at $\sqrt{s} = 250$ GeV with a luminosity 2 ab^{-1} . For this reason it was long advocated that ILC running at $\sqrt{s} = 350$ GeV or 500 GeV was essential to obtaining interesting Higgs coupling accuracy. As we shall see this is not the case in the EFT framework. Many couplings can be measured with $O(1\%)$ running solely at $\sqrt{s} = 250$ GeV with 2 ab^{-1} luminosity. This is due to the relationship between the hZZ and hWW couplings in the $SU(2) \times U(1)$ invariant EFT framework:

$$\eta_Z = \left(-c_T - \frac{1}{2}c_H - c'_{HL}\right) \approx \eta_W = \left(-\frac{1}{2}c_H - c'_{HL}\right)$$

(c_T is tightly constrained by precision electroweak constraints)

In order to estimate the Higgs coupling accuracy within the EFT framework a linear least squares fit of 20 parameters is performed using EWPO's, LHC measurements of ratios of Higgs partial widths, and ILC measurements of Higgs cross sections, Higgs cross section times branching ratios, and TGC's.

20 parameters:

9 operators modifying h, γ, W, Z interactions

5 operators modifying h coupling to b, c, τ, μ, g

2 parameters to account for invisible and exotic Higgs decays

4 SM parameters g, g', v, λ that get shifted by the dim 6 operator coefficients

(also some $qqhW$ & $qqhZ$ contact terms that are controlled with precision measurements of the W & Z boson widths)

Measureables:

EWPO's:

$$\alpha(m_Z), G_F, m_W, m_Z, m_h, A_{LR}(l), \Gamma(Z \rightarrow l^+ l^-)$$

LHC:

$$BR(h \rightarrow \gamma\gamma) / BR(h \rightarrow ZZ)$$

ILC TGC's:

$$g_{1Z}, \kappa_\gamma, \lambda_\gamma$$

ILC polarized $\sigma(e^+ e^- \rightarrow Zh)$

$$\sigma = \frac{2}{3} \frac{\pi \alpha_w^2}{c_w^4} \frac{m_Z^2}{(s - m_Z^2)} \frac{2k_Z}{\sqrt{s}} \left(2 + \frac{E_Z^2}{m_Z^2}\right) \cdot Q_Z^2 \cdot \left[1 + 2a + 2 \frac{3\sqrt{s} E_Z / m_Z^2}{(2 + E_Z^2 / m_Z^2)} b\right]$$

$$Q_{ZL} = \left(\frac{1}{2} - s_w^2\right), \quad a_L = -c_H/2$$

$$b_L = c_w^2 \left(1 + \frac{s_w^2}{1/2 - s_w^2} \frac{s - m_Z^2}{s}\right) (8c_{WW})$$

$$Q_{ZR} = (-s_w^2), \quad a_R = -c_H/2$$

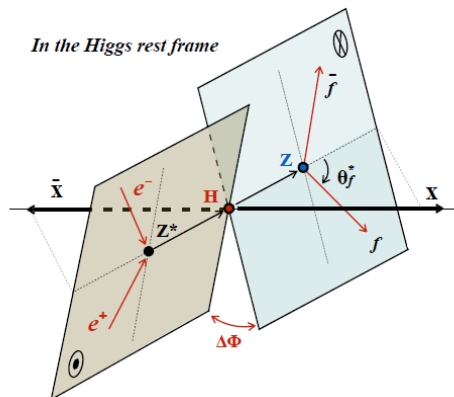
$$b_R = c_w^2 \left(1 - \frac{s - m_Z^2}{s}\right) (8c_{WW}) .$$

Measureables:

ILC $\sigma \times \text{BR}$

| | 250 GeV | | 500 GeV | |
|-------------------------------|---------|-----------------|---------|-----------------|
| | Zh | $\nu\bar{\nu}h$ | Zh | $\nu\bar{\nu}h$ |
| $h \rightarrow \text{invis.}$ | 0.9 | | 3.4 | |
| $h \rightarrow b\bar{b}$ | 1.2 | 10.5 | 2.54 | 0.99 |
| $h \rightarrow c\bar{c}$ | 8.3 | | 18.4 | 8.8 |
| $h \rightarrow gg$ | 7.0 | | 15.6 | 5.8 |
| $h \rightarrow WW$ | 6.4 | | 13.0 | 3.4 |
| $h \rightarrow \tau\tau$ | 3.2 | | 7.6 | 12.7 |
| $h \rightarrow ZZ$ | 19 | | 35 | 11.6 |
| $h \rightarrow \gamma\gamma$ | 34 | | 48 | 27 |
| $h \rightarrow \mu\mu$ | 72 | | 124 | 102 |

ILC Angular Analysis of $e^+e^- \rightarrow Zh$



Angular Asymmetry derived from the new structures

The Lorentz structure

$$\mathcal{L}_{ZZH} = M_Z^2 \left(\frac{1}{v} + \frac{a_Z}{\Lambda} \right) Z_\mu Z^\mu H + \frac{b_Z}{2\Lambda} \hat{Z}_{\mu\nu} \hat{Z}^{\mu\nu} H + \frac{\tilde{b}_Z}{2\Lambda} \hat{Z}_{\mu\nu} \tilde{\hat{Z}}^{\mu\nu} H$$

- “ a_Z ” is a simple normalization parameter which affects the overall cross section of processes. (just rescales the SM-coupling)
- “ b_Z ” has a different tensor structure which affects momentum spectra and changes angular/spin correlations.
- “ \tilde{b}_Z ” is a CP-violating parameter which affects angular/spin correlations.

Evolution of EFT parameter measurements

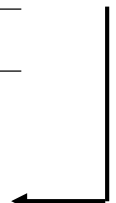
250 GeV

| c_I | prec. EW | + WW | + LHC | + Zh | ILC 250 |
|-----------------|----------|-------|-------|-------|---------|
| c_T | 0.011 | 0.051 | 0.051 | 0.048 | 0.052 |
| c_{HE} | 0.043 | 0.026 | 0.085 | 0.047 | 0.055 |
| c_{HL} | 0.042 | 0.035 | 0.035 | 0.032 | 0.039 |
| c'_{HL} | — | 0.028 | 0.028 | 0.028 | 0.047 |
| δc_{WB} | — | 0.078 | 0.080 | 0.076 | 0.090 |
| δc_{BB} | — | — | 0.20 | 0.16 | 0.11 |
| δc_{WW} | — | — | 0.21 | 0.13 | 0.13 |
| c_H | — | — | — | 1.12 | 1.20 |

500 GeV

@ 500 GeV Higgs meas. give better
precision on c_{HE} and c_{HL} than EWPO's

| c_I | prec. EW | + WW | + LHC | + Zh | ILC 500 | 250+500 |
|-----------------|----------|-------|-------|-------|---------|---------|
| c_T | 0.011 | 0.046 | 0.047 | 0.041 | 0.037 | 0.030 |
| c_{HE} | 0.043 | 0.015 | 0.077 | 0.040 | 0.010 | 0.009 |
| c_{HL} | 0.042 | 0.030 | 0.030 | 0.027 | 0.016 | 0.013 |
| c'_{HL} | — | 0.027 | 0.028 | 0.026 | 0.014 | 0.011 |
| δc_{WB} | — | 0.070 | 0.072 | 0.067 | 0.052 | 0.041 |
| δc_{BB} | — | — | 0.20 | 0.15 | 0.088 | 0.062 |
| δc_{WW} | — | — | 0.21 | 0.11 | 0.044 | 0.039 |
| c_H | — | — | — | 4.78 | 1.24 | 0.65 |



Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ κ fit |
|---------------------------------|--|
| $g(hb\bar{b})$ | 2.1 |
| $g(hc\bar{c})$ | 2.7 |
| $g(hgg)$ | 2.4 |
| $g(hWW)$ | 1.9 |
| $g(h\tau\tau)$ | 2.3 |
| $g(hZZ)$ | 0.36 |
| $g(h\gamma\gamma)$ | 7.4 |
| $g(h\mu\mu)$ | 14 |
| $g(hbb) / g(hWW)$ | 0.85 |
| $g(hWW) / g(hZZ)$ | 3.29 |
| Γ_h | 4.4 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.72 |
| $BR(h \rightarrow inv)$ | 0.39 |
| $BR(h \rightarrow other)$ | 1.6 |
| $g(hhh)$ | – |

Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ κ fit | ILC 250 GeV 2 ab ⁻¹ EFT |
|---------------------------------|--|---|
| $g(hb\bar{b})$ | 2.1 | 1.0 |
| $g(hc\bar{c})$ | 2.7 | 1.8 |
| $g(hgg)$ | 2.4 | 1.6 |
| $g(hWW)$ | 1.9 | 0.65 |
| $g(h\tau\tau)$ | 2.3 | 1.2 |
| $g(hZZ)$ | 0.36 | 0.66 |
| $g(h\gamma\gamma)$ | 7.4 | 1.2 |
| $g(h\mu\mu)$ | 14 | 5.5 |
| $g(hbb) / g(hWW)$ | 0.85 | 0.82 |
| $g(hWW) / g(hZZ)$ | 3.29 | 0.07 |
| Γ_h | 4.4 | 2.4 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.72 | 0.70 |
| $BR(h \rightarrow inv)$ | 0.39 | 0.30 |
| $BR(h \rightarrow other)$ | 1.6 | 1.5 |
| $g(hhh)$ | – | – |

Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ κ fit | ILC 250 GeV 2 ab ⁻¹ EFT | ILC 250 + 500 GeV 2 ab ⁻¹ + 4 ab ⁻¹ κ fit |
|---------------------------------|--|---|---|
| $g(hb\bar{b})$ | 2.1 | 1.0 | 0.70 |
| $g(hc\bar{c})$ | 2.7 | 1.8 | 1.2 |
| $g(hgg)$ | 2.4 | 1.6 | 1.0 |
| $g(hWW)$ | 1.9 | 0.65 | 0.42 |
| $g(h\tau\tau)$ | 2.3 | 1.2 | 0.90 |
| $g(hZZ)$ | 0.36 | 0.66 | 0.31 |
| $g(h\gamma\gamma)$ | 7.4 | 1.2 | 3.4 |
| $g(h\mu\mu)$ | 14 | 5.5 | 5.0 |
| $g(hbb) / g(hWW)$ | 0.85 | 0.82 | – |
| $g(hWW) / g(hZZ)$ | 3.29 | 0.07 | – |
| Γ_h | 4.4 | 2.4 | 1.8 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.72 | 0.70 | 0.62 |
| $BR(h \rightarrow inv)$ | 0.39 | 0.30 | 0.29 |
| $BR(h \rightarrow other)$ | 1.6 | 1.5 | – |
| $g(hhh)$ | – | – | 27 |

Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ EFT | CEPC 250 GeV 5 ab ⁻¹ EFT |
|---------------------------------|---|--|
| $g(hb\bar{b})$ | 1.0 | 0.98 |
| $g(hc\bar{c})$ | 1.8 | 1.4 |
| $g(hgg)$ | 1.6 | 1.3 |
| $g(hWW)$ | 0.65 | 0.80 |
| $g(h\tau\tau)$ | 1.2 | 1.1 |
| $g(hZZ)$ | 0.66 | 0.80 |
| $g(h\gamma\gamma)$ | 1.2 | 1.3 |
| $g(h\mu\mu)$ | 5.5 | 5.1 |
| $g(hbb) / g(hWW)$ | 0.82 | 0.58 |
| $g(hWW) / g(hZZ)$ | 0.07 | 0.07 |
| Γ_h | 2.4 | 2.1 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.70 | 0.50 |
| $BR(h \rightarrow inv)$ | 0.30 | 0.30 |
| $BR(h \rightarrow other)$ | 1.5 | 1.1 |
| $g(hhh)$ | – | – |

Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ EFT | FCC-ee 250 + 350 GeV 5 ab ⁻¹ + 1.5 ab ⁻¹ EFT | ILC 250 + 500 GeV 2 ab ⁻¹ + 4 ab ⁻¹ EFT |
|---------------------------------|---|---|--|
| $g(hb\bar{b})$ | 1.0 | 0.66 | 0.55 |
| $g(hc\bar{c})$ | 1.8 | 1.2 | 1.1 |
| $g(hgg)$ | 1.6 | 0.99 | 0.89 |
| $g(hWW)$ | 0.65 | 0.42 | 0.34 |
| $g(h\tau\tau)$ | 1.2 | 0.75 | 0.71 |
| $g(hZZ)$ | 0.66 | 0.42 | 0.34 |
| $g(h\gamma\gamma)$ | 1.2 | 1.0 | 1.0 |
| $g(h\mu\mu)$ | 5.5 | 4.9 | 5.0 |
| $g(hbb) / g(hWW)$ | 0.82 | 0.51 | 0.43 |
| $g(hWW) / g(hZZ)$ | 0.07 | 0.06 | 0.05 |
| Γ_h | 2.4 | 1.5 | 1.5 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.70 | 0.22 | 0.61 |
| $BR(h \rightarrow inv)$ | 0.30 | 0.27 | 0.28 |
| $BR(h \rightarrow other)$ | 1.5 | 0.94 | 1.2 |
| $g(hhh)$ | – | – | 27 |

Higgs Coupling Precision (%)

| | ILC 250 GeV 2 ab ⁻¹ EFT | ILC 250 + 500 GeV 2 ab ⁻¹ + 4 ab ⁻¹ EFT |
|---------------------------------|---|--|
| $g(hb\bar{b})$ | 1.0 | 0.55 |
| $g(hc\bar{c})$ | 1.8 | 1.1 |
| $g(hgg)$ | 1.6 | 0.89 |
| $g(hWW)$ | 0.65 | 0.34 |
| $g(h\tau\tau)$ | 1.2 | 0.71 |
| $g(hZZ)$ | 0.66 | 0.34 |
| $g(h\gamma\gamma)$ | 1.2 | 1.0 |
| $g(h\mu\mu)$ | 5.5 | 5.0 |
| $g(hbb) / g(hWW)$ | 0.82 | 0.43 |
| $g(hWW) / g(hZZ)$ | 0.07 | 0.05 |
| Γ_h | 2.4 | 1.5 |
| $\sigma(e^+e^- \rightarrow Zh)$ | 0.70 | 0.61 |
| $BR(h \rightarrow inv)$ | 0.30 | 0.28 |
| $BR(h \rightarrow other)$ | 1.5 | 1.2 |
| $g(hhh)$ | – | 27 |

- Selection of 9 models with all new particles outside of projected reach of direct searches at HL-LHC

| Model | $b\bar{b}$ | $c\bar{c}$ | gg | WW | $\tau\tau$ | ZZ | $\gamma\gamma$ | $\mu\mu$ |
|---------------------------------|------------|------------|------|------|------------|------|----------------|----------|
| 1 MSSM [34] | -0.2 | -0.2 | -0.2 | 0.0 | +0.4 | -0.5 | +0.1 | +0.3 |
| 2 Type II 2HD [36] | +10.1 | -0.2 | -0.2 | 0.0 | +9.8 | 0.0 | +0.1 | +9.8 |
| 3 Type X 2HD [36] | -0.2 | -0.2 | -0.2 | 0.0 | +7.8 | 0.0 | 0.0 | +7.8 |
| 4 Type Y 2HD [36] | +10.1 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.1 | -0.2 |
| 5 Composite Higgs [38] | -6.4 | -6.4 | -6.4 | -2.1 | -6.4 | -2.1 | -2.1 | -6.4 |
| 6 Little Higgs w. T-parity [39] | 0.0 | 0.0 | -6.1 | -2.5 | 0.0 | -2.5 | -1.5 | 0.0 |
| 7 Little Higgs w. T-parity [40] | -7.8 | -4.6 | -3.5 | -1.5 | -7.8 | -1.5 | -1.0 | -7.8 |
| 8 Higgs-Radion [41] | -1.5 | -1.5 | -1.5 | -1.5 | -1.5 | -1.5 | -1.0 | -1.5 |
| 9 Higgs Singlet [42] | -3.5 | -3.5 | -3.5 | -3.5 | -3.5 | -3.5 | -3.5 | -3.5 |

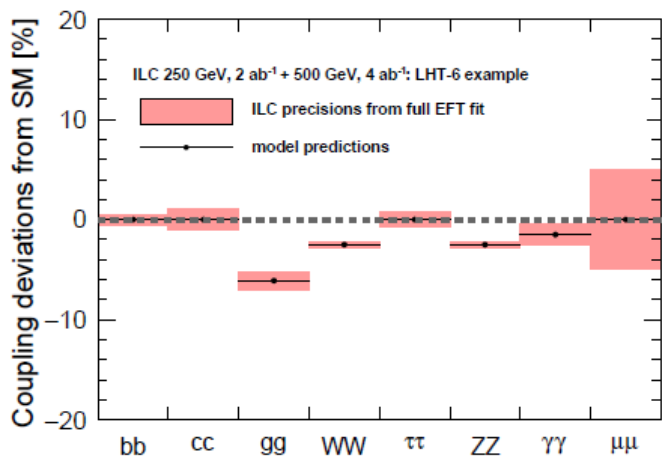
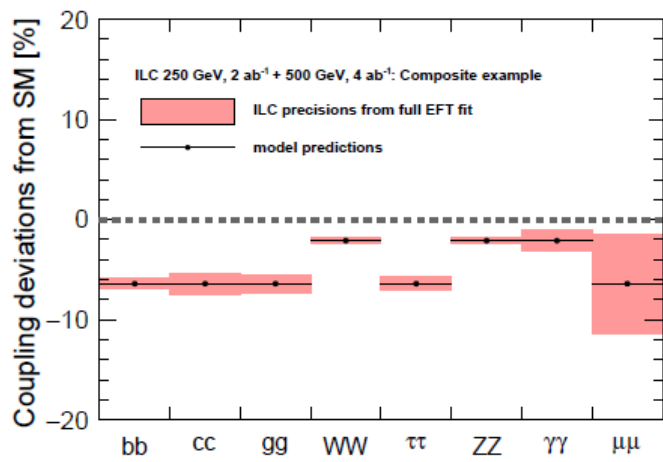
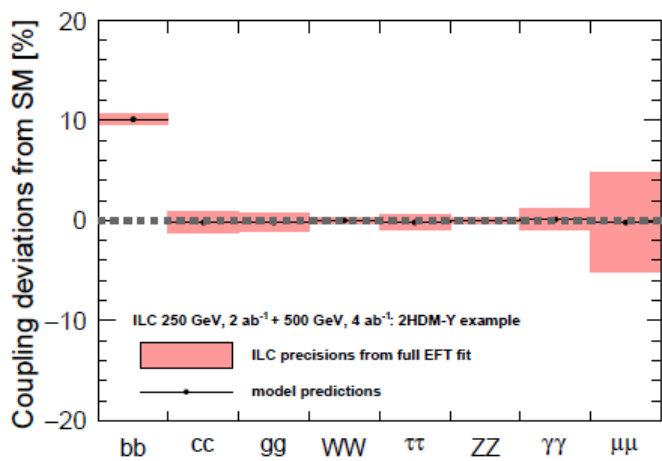
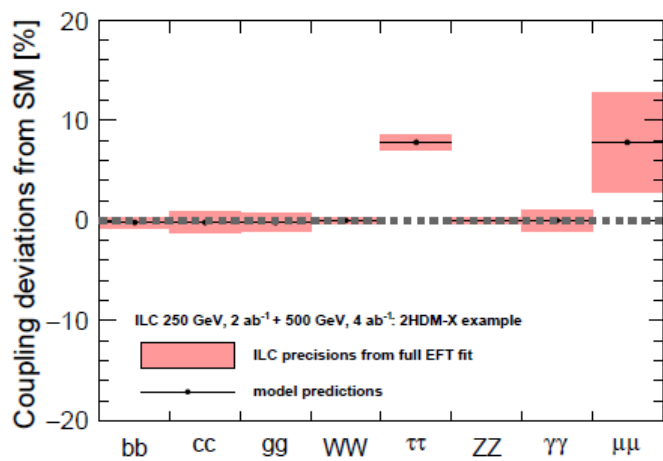
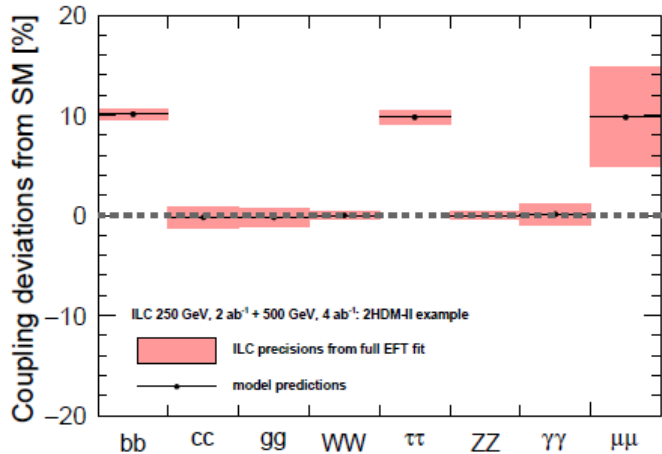
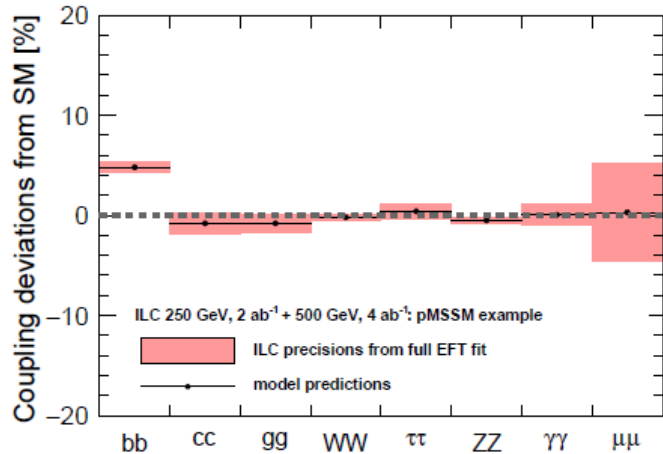
NATURALNESS

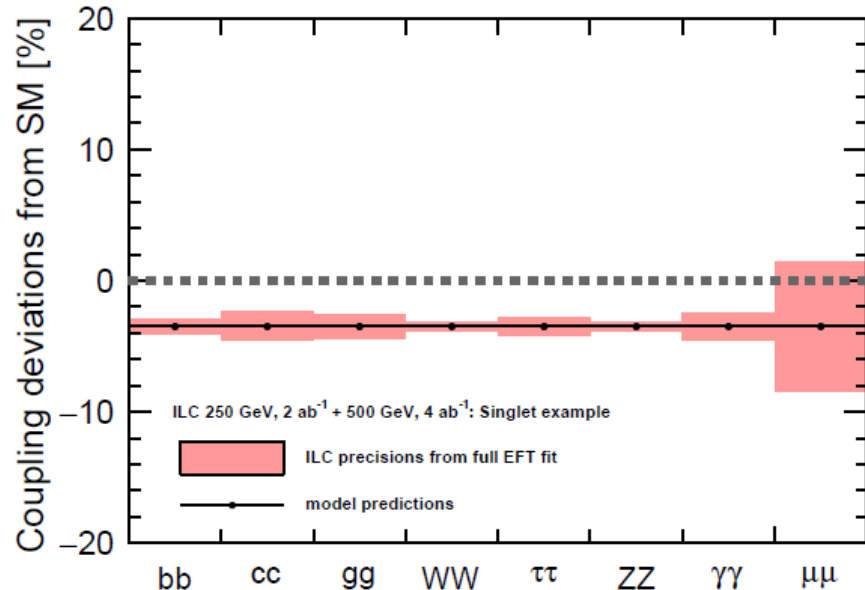
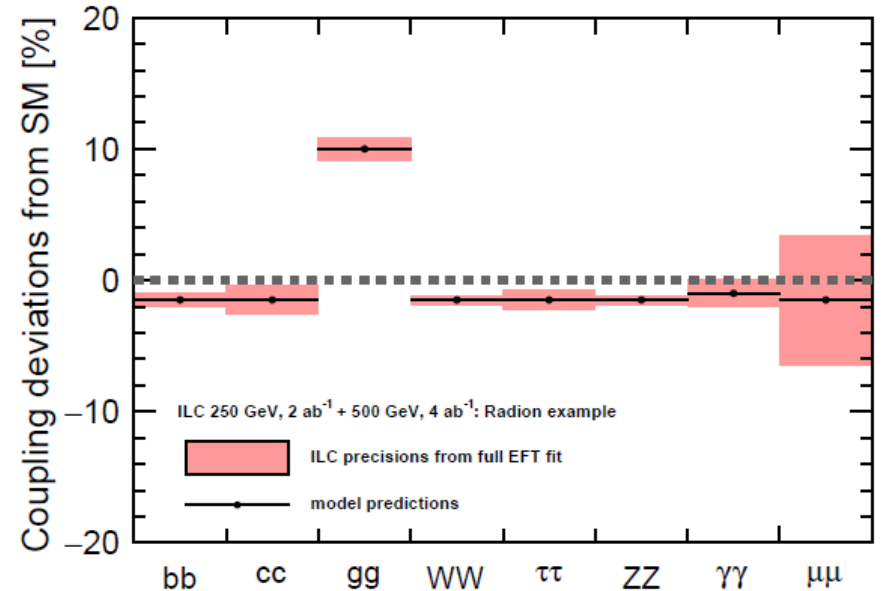
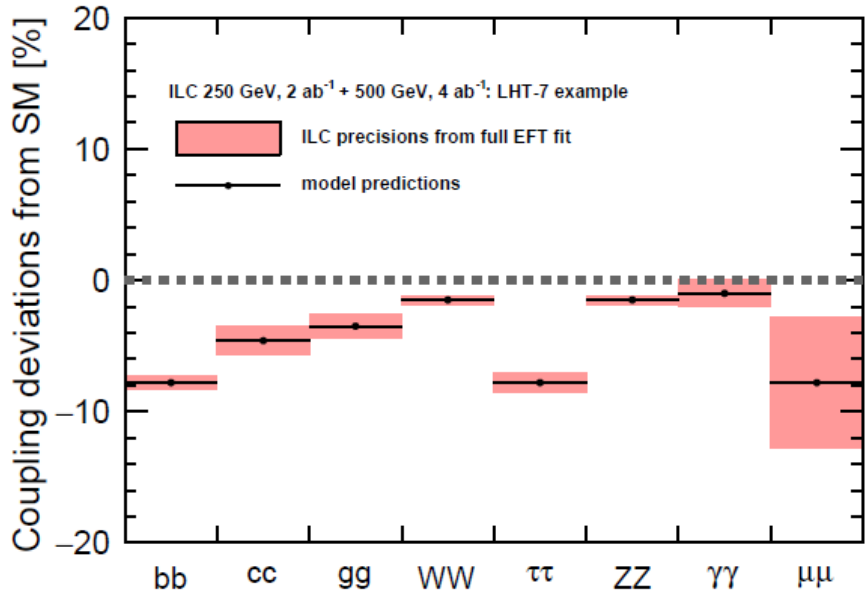
COMPOSITENESS

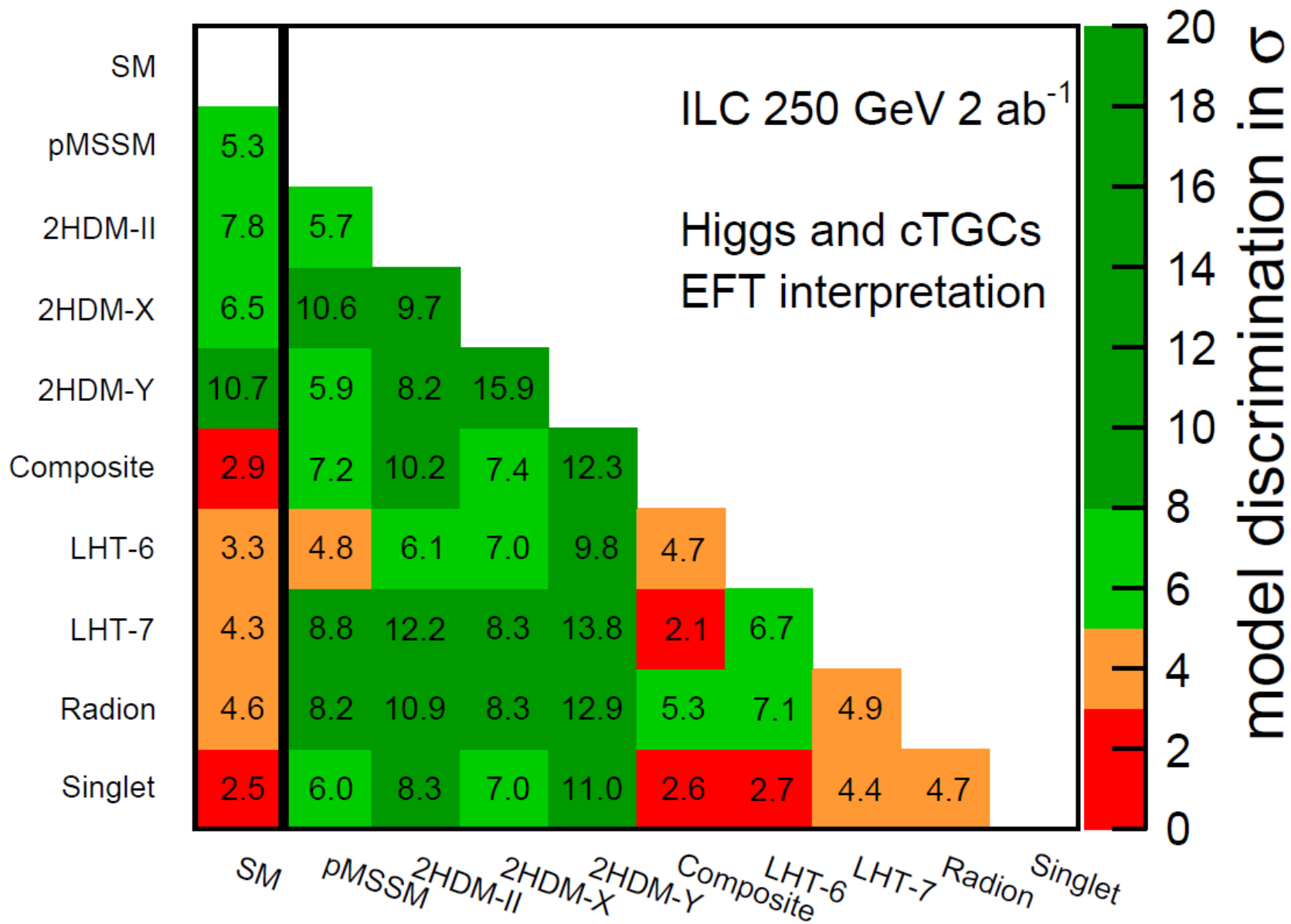
EW BARYOGENESIS

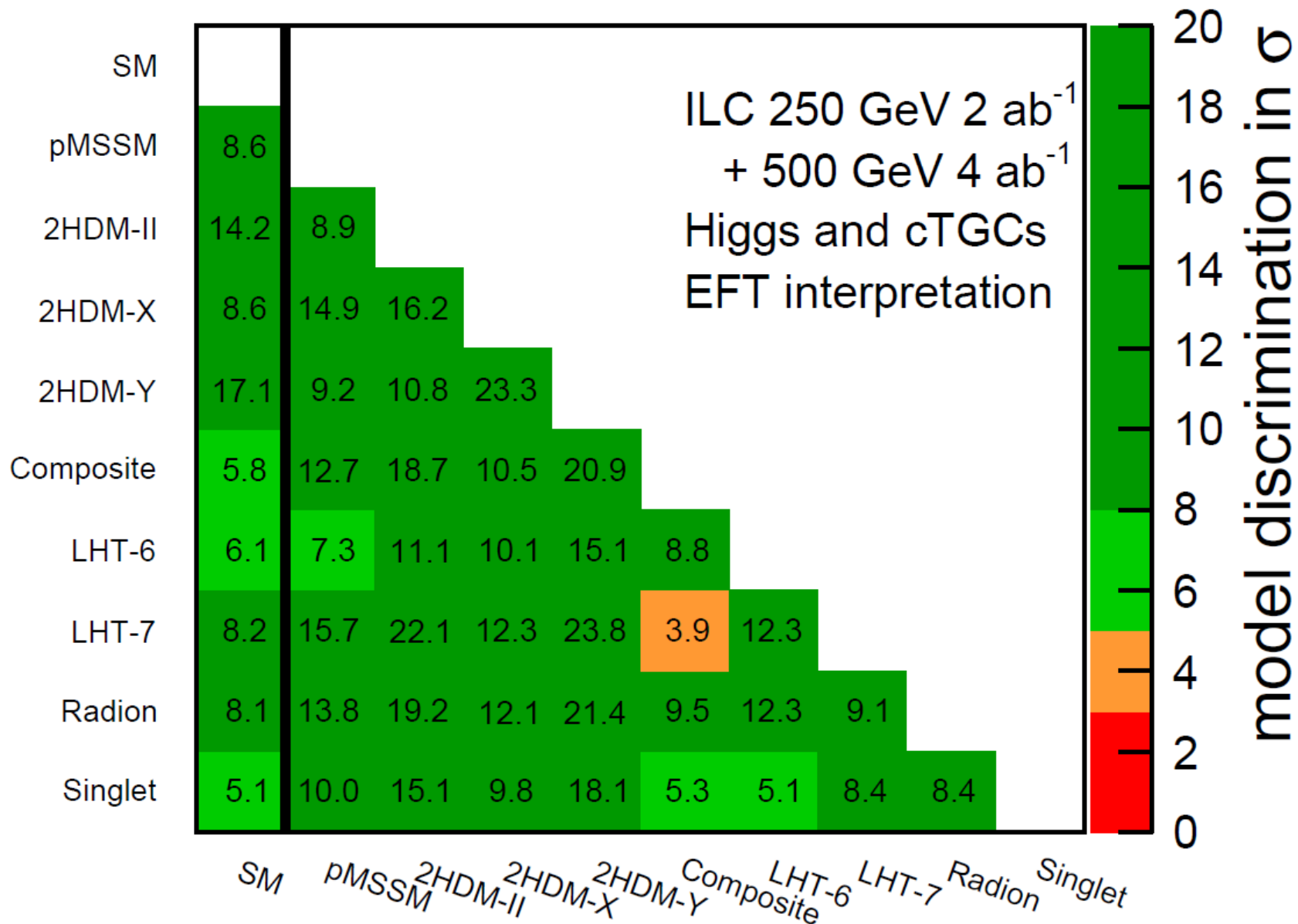
Quantify model discrimination using χ^2 deviation of Model A given Model b

$$(\chi^2)_{AB} = (g_A^T - g_B^T) [VCV^T]^{-1} (g_A - g_B)$$

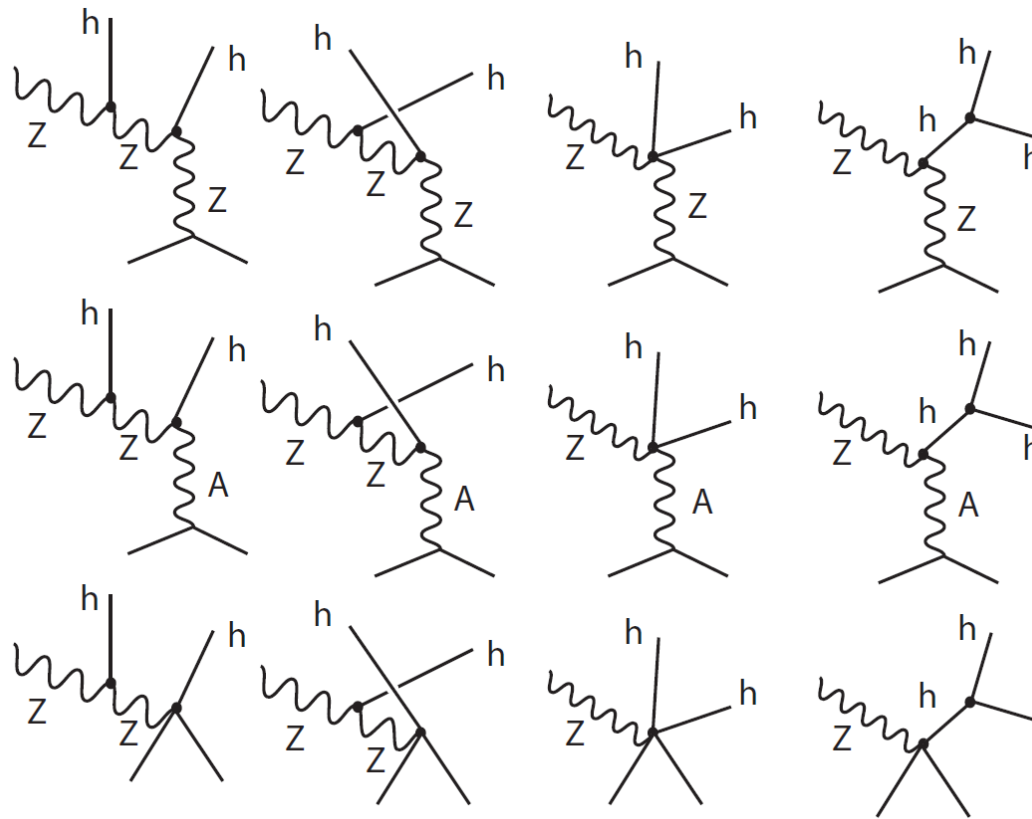
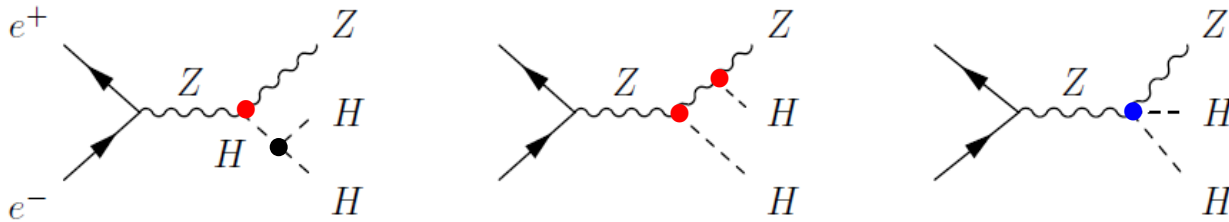








Applying EFT to ILC Triple Higgs Coupling Measurement



Applying EFT to ILC Triple Higgs Coupling Measurement

$$\begin{aligned} \sigma/(SM) = & 1 + 1.15\delta g_L + 0.85\delta g_R + 1.40\eta_Z + 1.02\eta_{ZZ} + 18.6\zeta_Z + 2.0\zeta_{AZ} \\ & + 0.56\eta_h - 1.58\theta_h + 62.1(c_{HL} + c'_{HL}) - 53.5c_{HE} \\ & - 3.9\delta m_h + 3.5\delta m_Z . \end{aligned}$$

$$\eta_h = \delta\bar{\lambda} + \delta v - \frac{3}{2}c_H + c_6$$

$$c_6 = \frac{1}{0.56} \left[\frac{\sigma_{Zhh}}{\sigma_{SM}} - 1 - \sum_i a_i c_i \right]$$

$$\Delta c_6 = \frac{1}{0.56} \left[\left(\frac{\Delta\sigma_{Zhh}}{\sigma_{SM}} \right)^2 + \sum_{i,j} a_i a_j (V_c)_{ij} \right]^{\frac{1}{2}}$$

Given the full ILC program of 2 ab^{-1} at 250 GeV and 4 ab^{-1} at 500 GeV

$$\left[\sum_{i,i} a_i a_j (V_c)_{ij} \right]^{\frac{1}{2}} = 0.04 \quad \ll \quad \frac{\Delta\sigma_{Zhh}}{\sigma_{SM}} = 0.168$$

Applying EFT to ILC Triple Higgs Coupling Measurement

Note that at a hadron collider:

- Many more unknown dim-6 op coefficients
- Fewer measurements to constrain coeff
- Leading production process $gg \rightarrow hh$ is loop level

ILC uniquely positioned to extract Higgs self coupling from double Higgs production