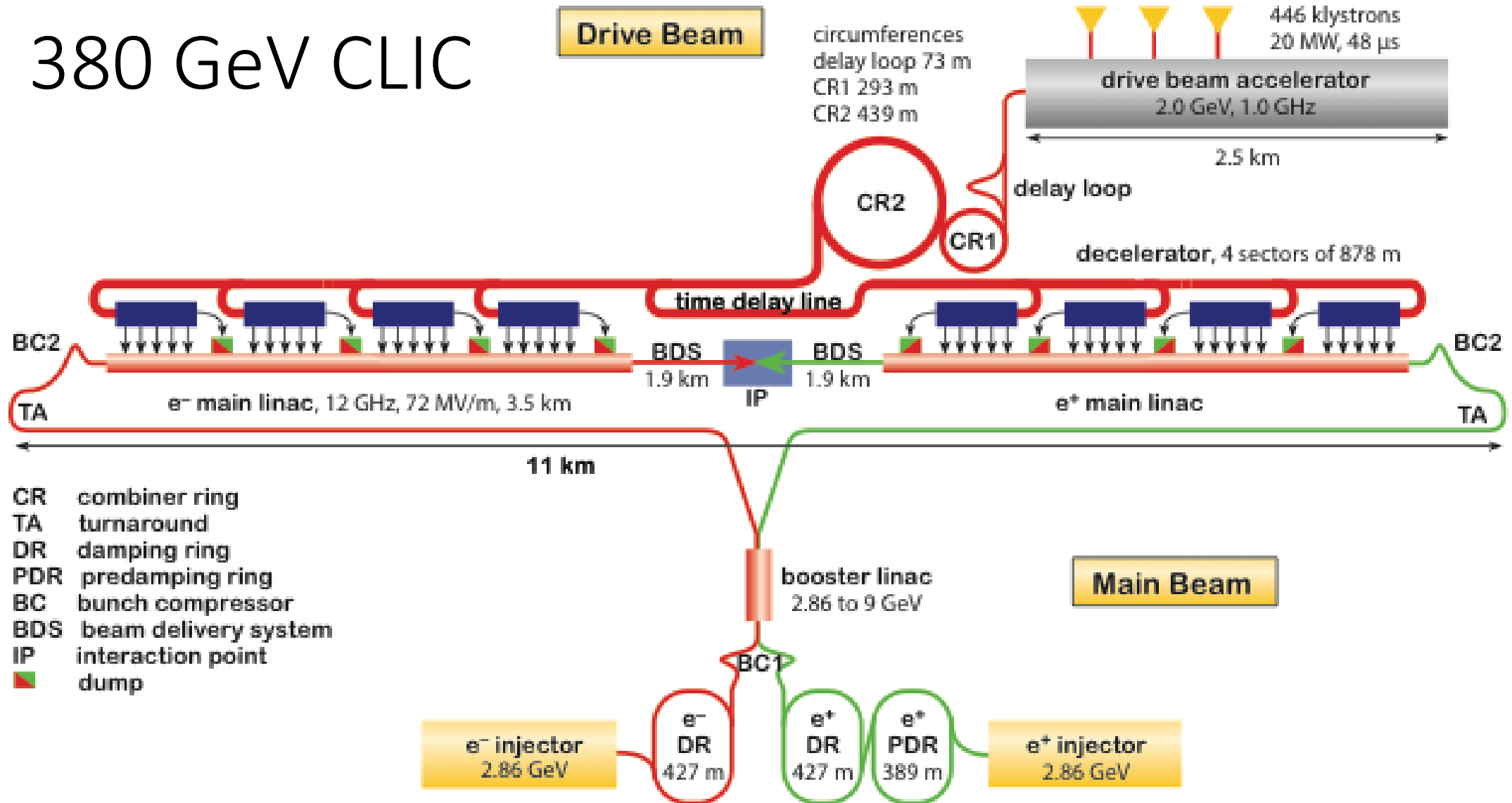


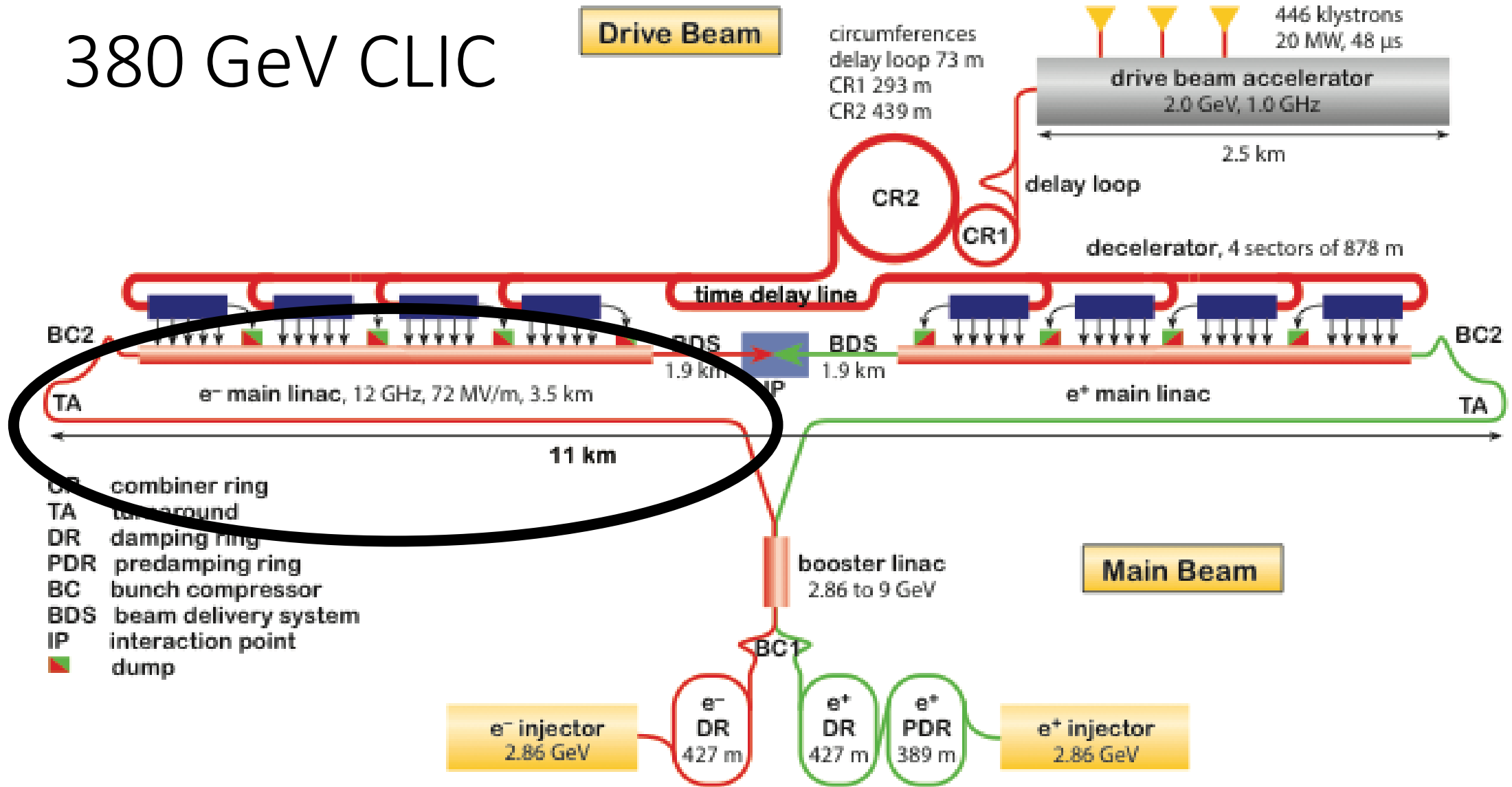
# 380 GeV CLIC - RTML Stray Field Simulations

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(01-09-17)

# 380 GeV CLIC

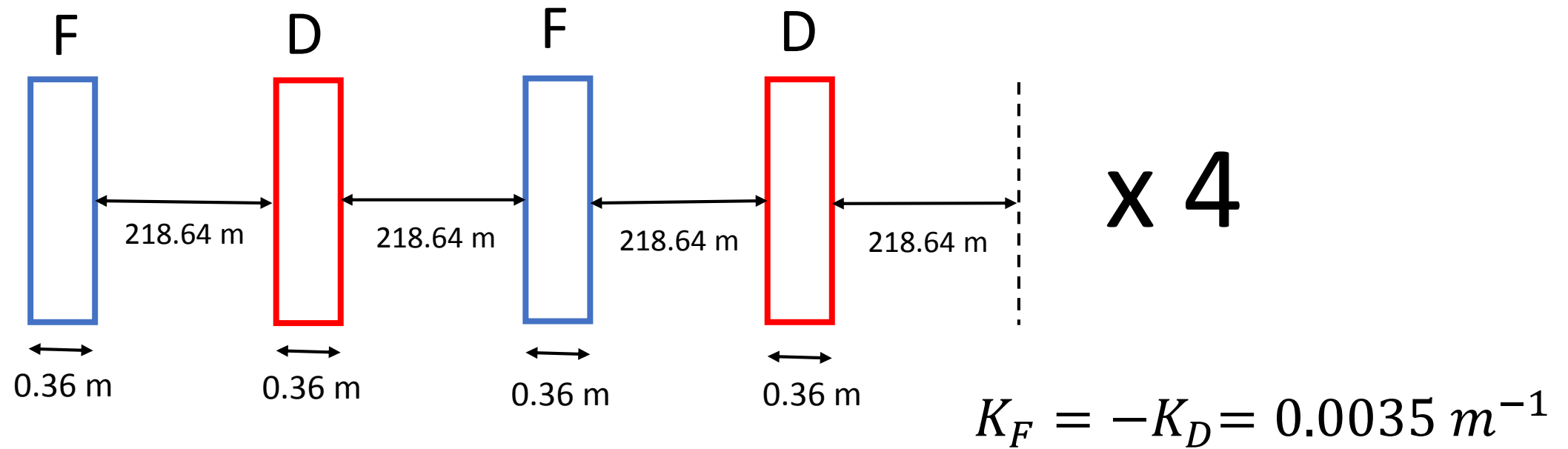


# 380 GeV CLIC



# RTML Transfer Line

- Length = 3,504 m
- Consists of one cell repeated four times:



# Stray Fields

$$\begin{aligned}\mathbf{B}(\mathbf{s}, t) &= B_0 \sin(\mathbf{k} \cdot \mathbf{s} - \omega t) \hat{\mathbf{x}} \\ &= B_0 \sin(\mathbf{k} \cdot \mathbf{s}) \cos(\omega t) \hat{\mathbf{x}} + B_0 \sin(\omega t) \cos(\mathbf{k} \cdot \mathbf{s}) \hat{\mathbf{x}} \\ &\approx B_0 \sin(\mathbf{k} \cdot \mathbf{s}) \hat{\mathbf{x}}\end{aligned}$$

- Model the stray field as a standing wave in the horizontal direction.
- This will kick the beam in the vertical direction:

$$\delta(s) = A \sin(ks) = A \sin\left(\frac{2\pi s}{\lambda}\right), \quad \text{where} \quad A = \frac{B_0 ec}{E}$$

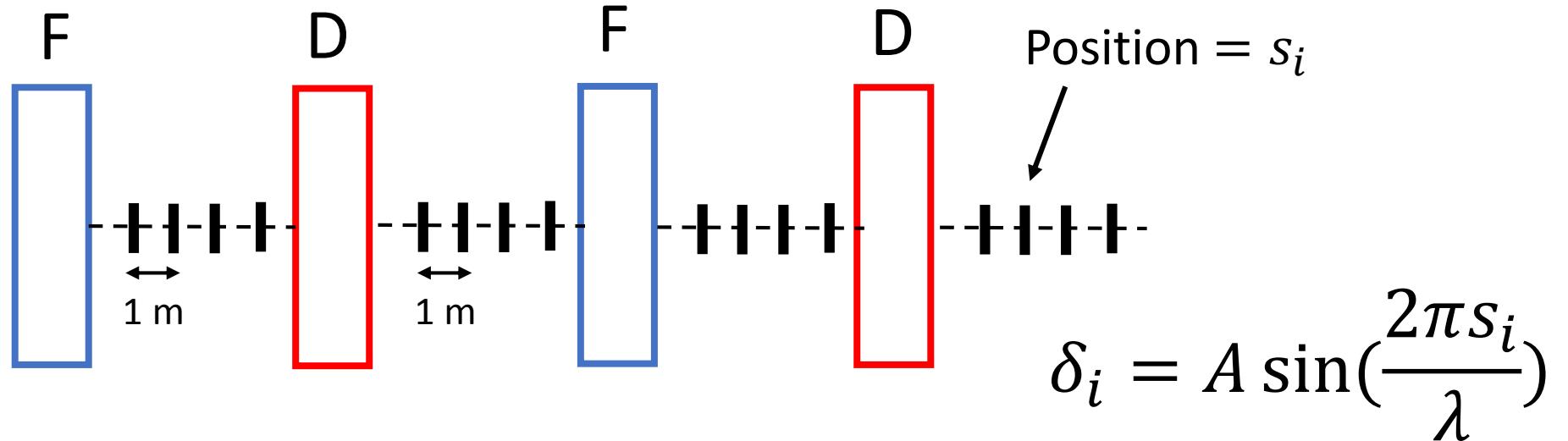
# Simulation of Stray Fields

- Simulations were done with particle tracking code PLACET.
- This tracks a number of particles in a Gaussian beam.
- It computes the emittance at the end of a beamline using the final positions and divergence of each particle in the beam.
- 10,000 particles were used in the simulations.
- Simulation of the RTML with no stray fields gives final emittances:

$$\varepsilon_{x,0} = 7.74 \times 10^{-7} \text{ m}, \quad \varepsilon_{y,0} = 0.057 \times 10^{-7} \text{ m}$$

# Simulation of Stray Fields

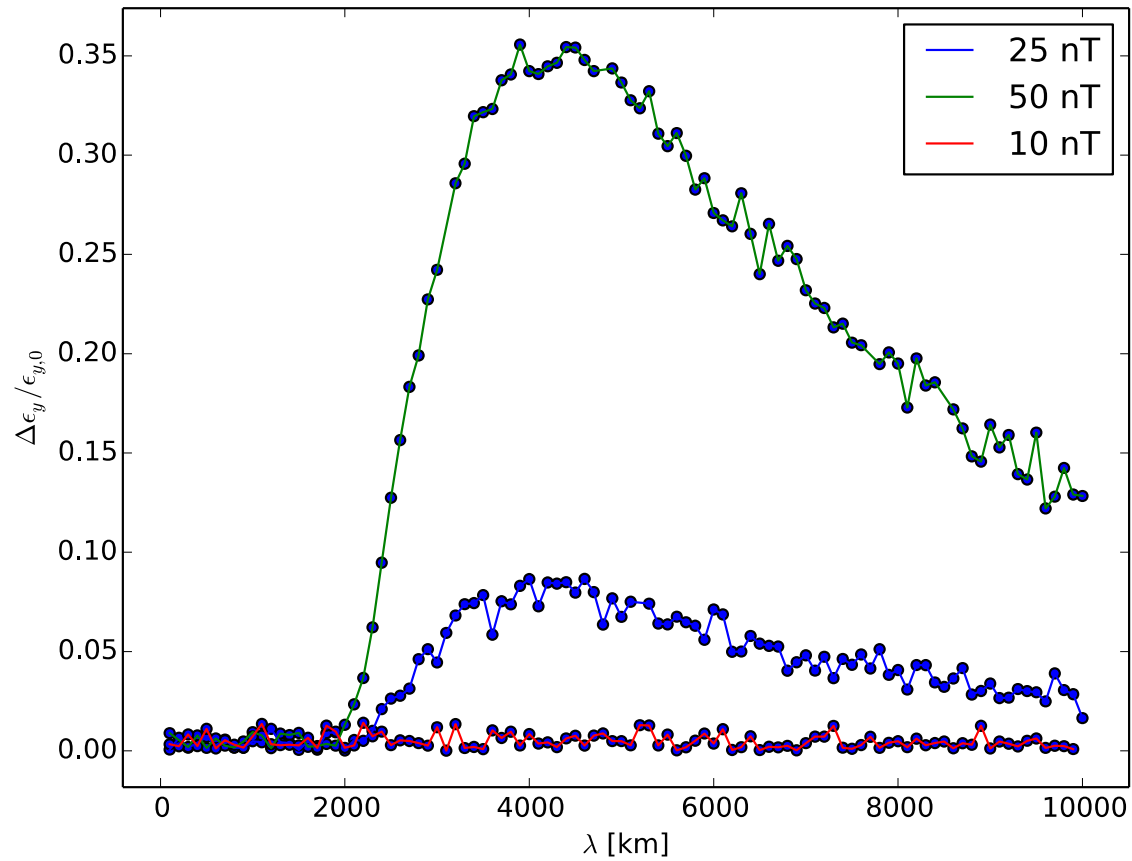
- Kicks are applied at discrete locations using dipole elements of zero length:



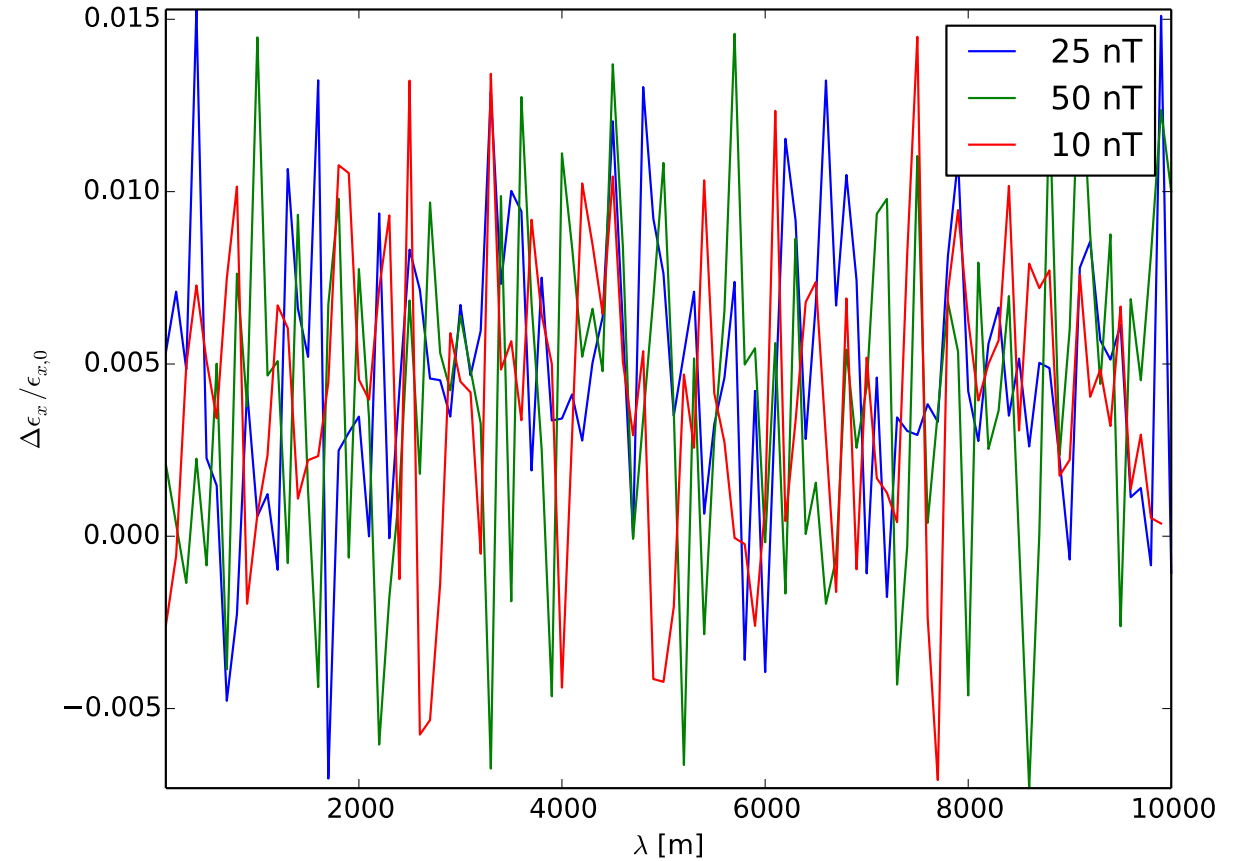
- We need to specify  $A$  and  $\lambda$ .
- Dipole are spaced 1 m apart, therefore can't simulate  $\lambda < 1$  m.

# Fractional Emittance Growth

## Vertical



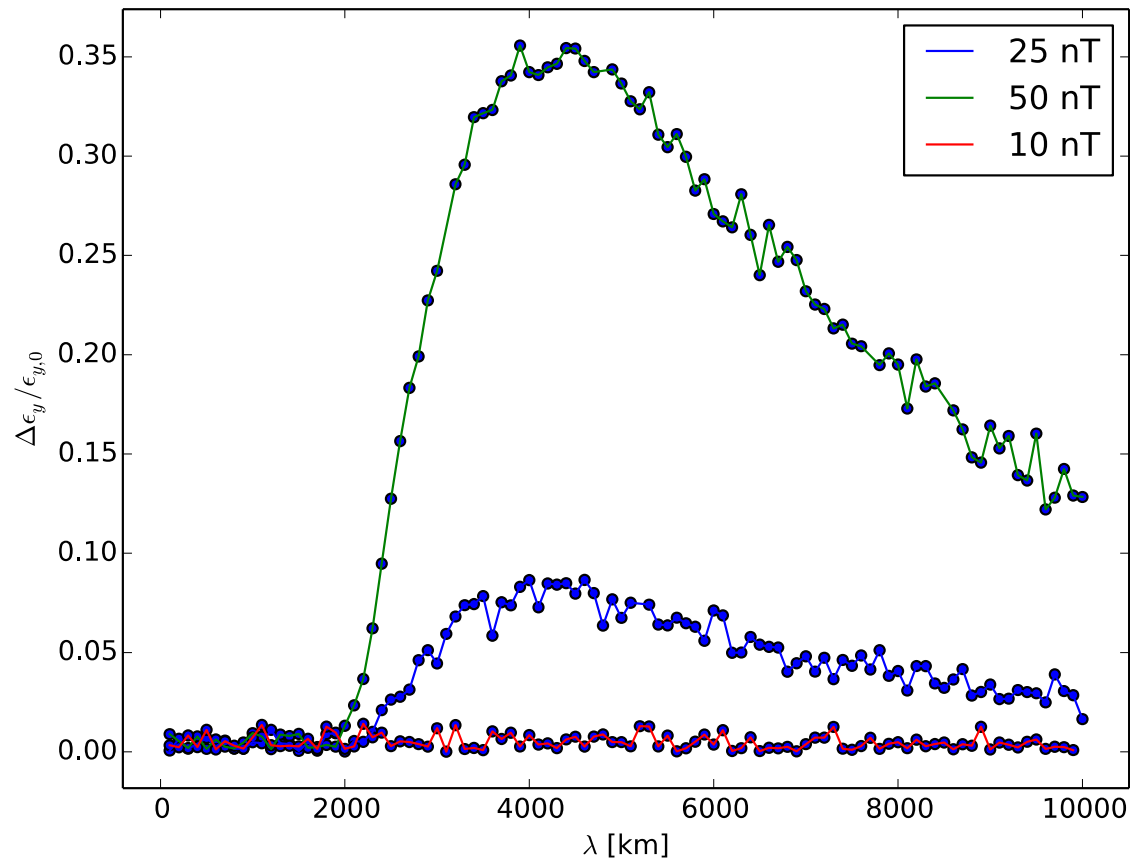
## Horizontal





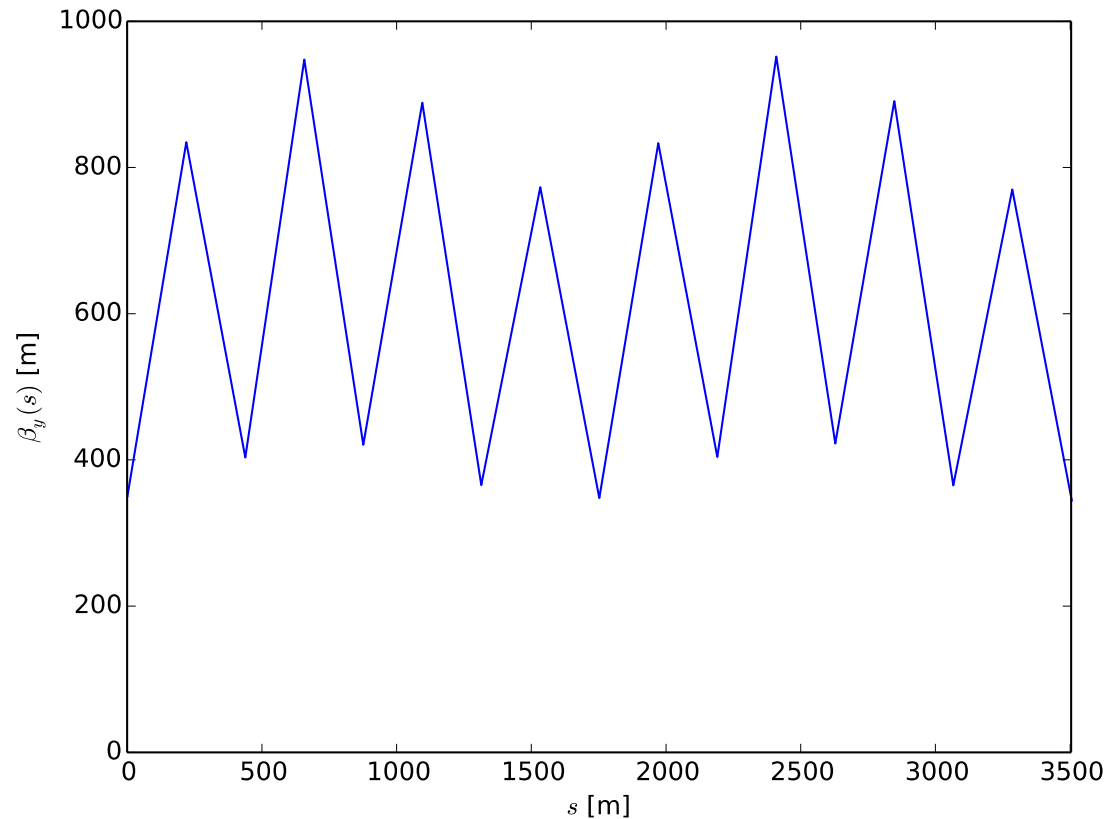
# Fractional Emittance Growth

Vertical



- Sensitivity occurs for  $B_0 > 10$  nT.
- Resonance around  $\lambda = 4000$  m.
- Emittance growth scales with  $A^2$ .

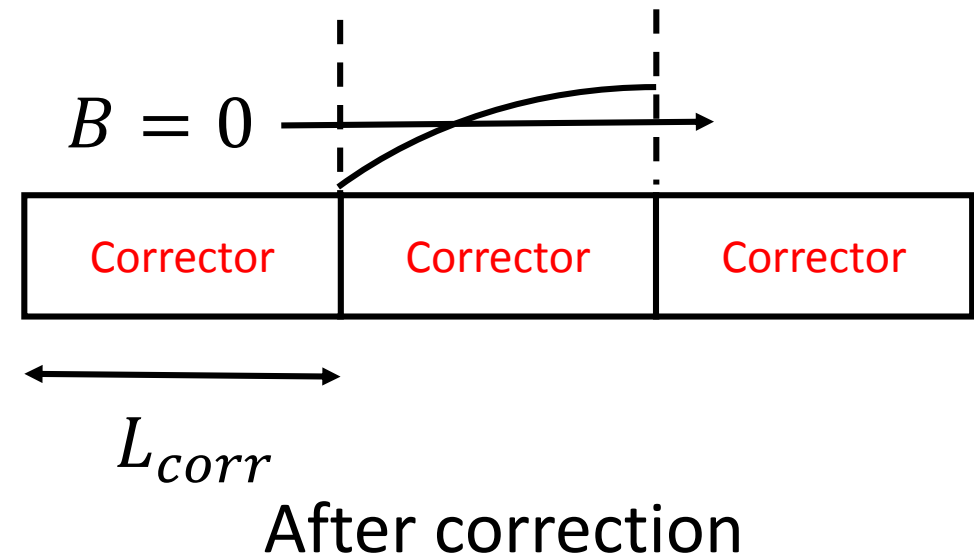
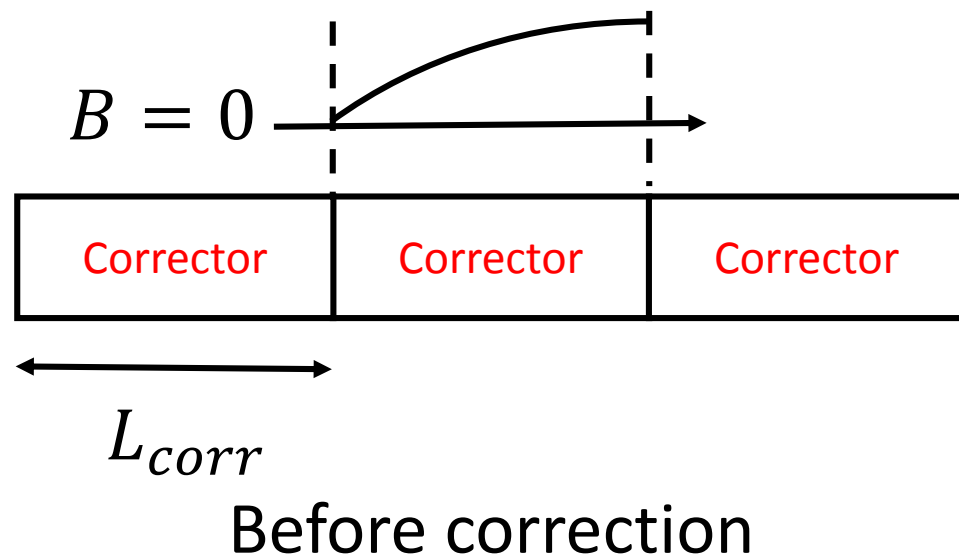
# Beta Function in the Transfer Line



- Calculated at the centres of each quadrupole.
- Shows some beta beating.
- The betatron wavelength is
$$\lambda_\beta = 2\pi \langle \beta \rangle \approx 3600 \text{ m.}$$
- Emittance growth occurs when  $\lambda$  approaches  $\lambda_\beta$ .

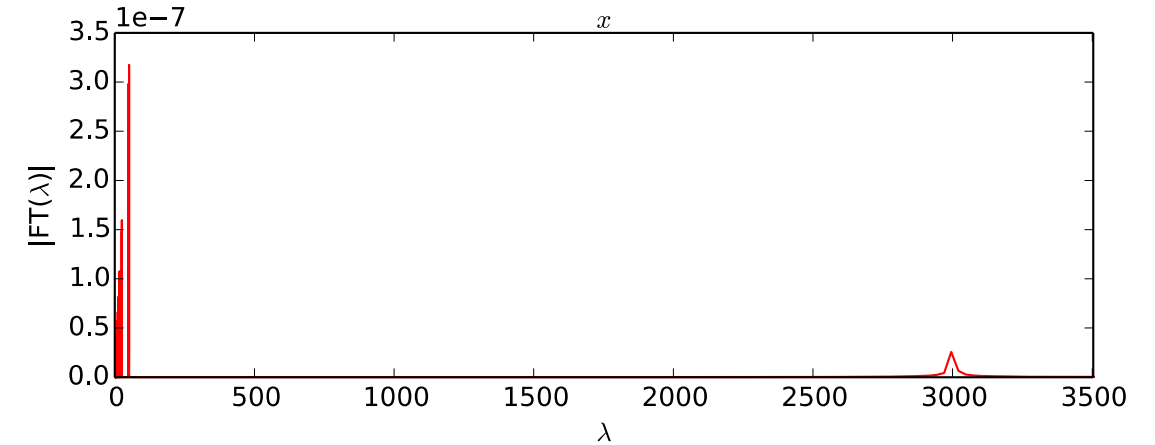
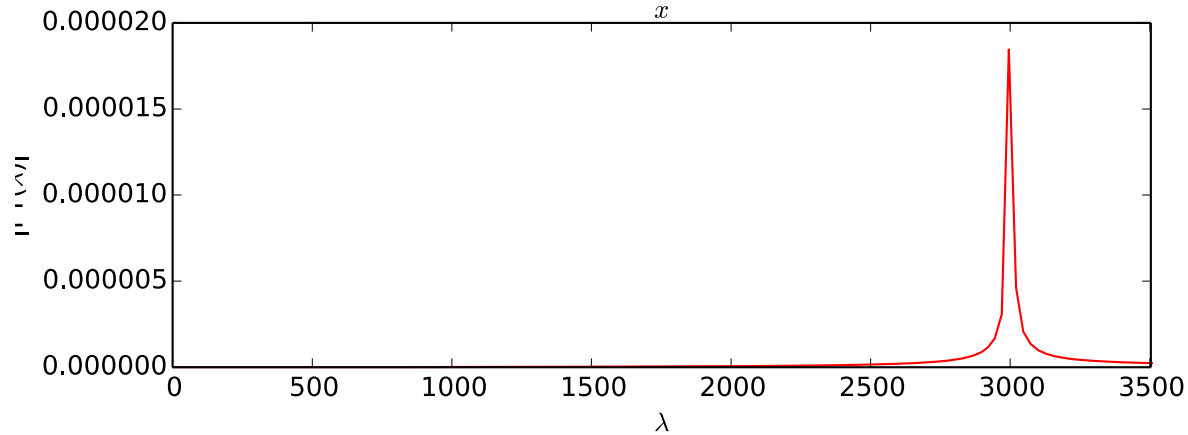
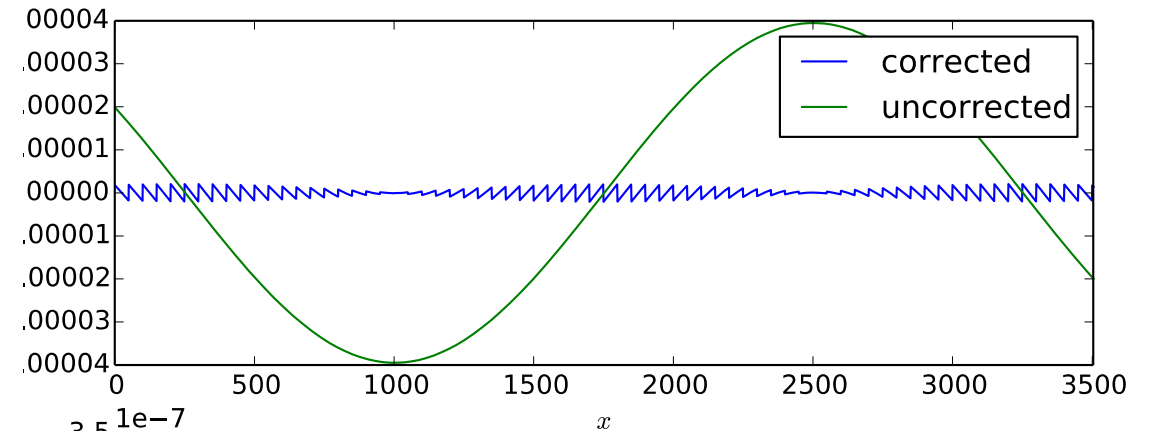
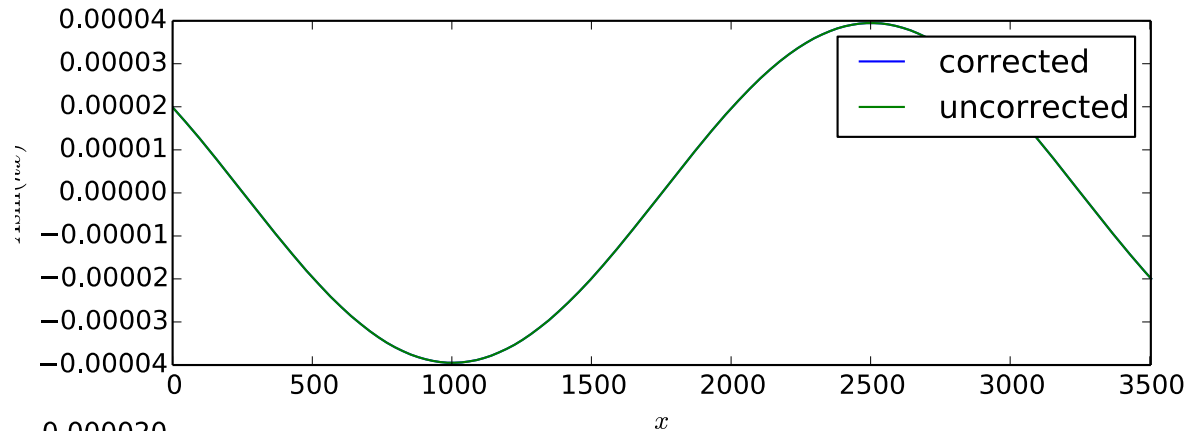
# Corrections

- A series of correctors were placed over the entire length of the transfer line.
- A corrector has a finite length and will see the average stray field over its length:



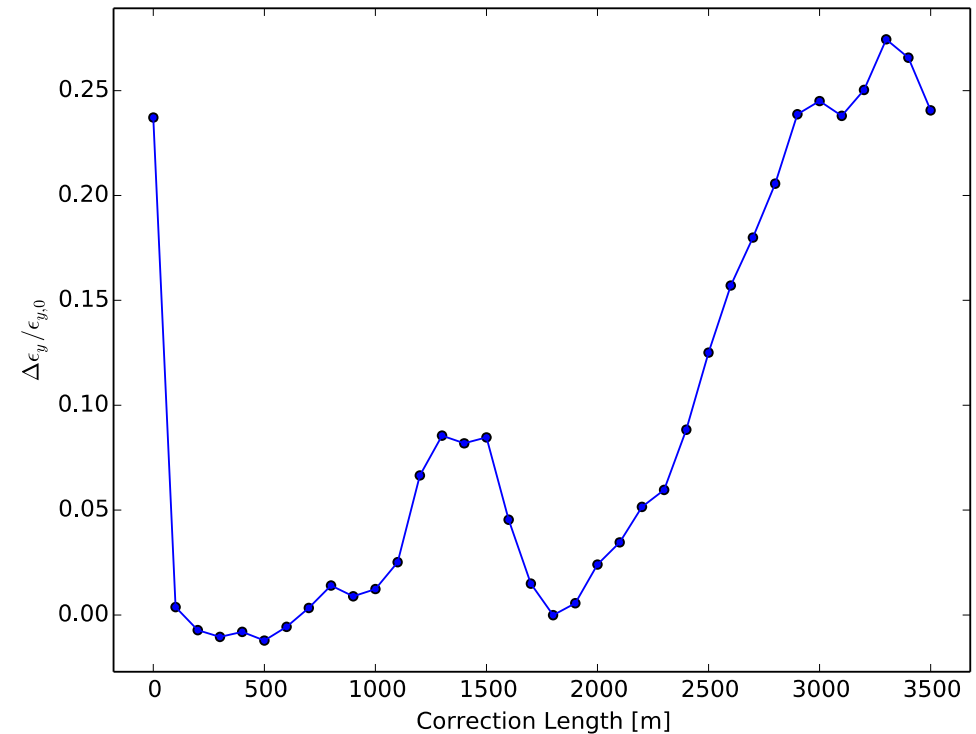
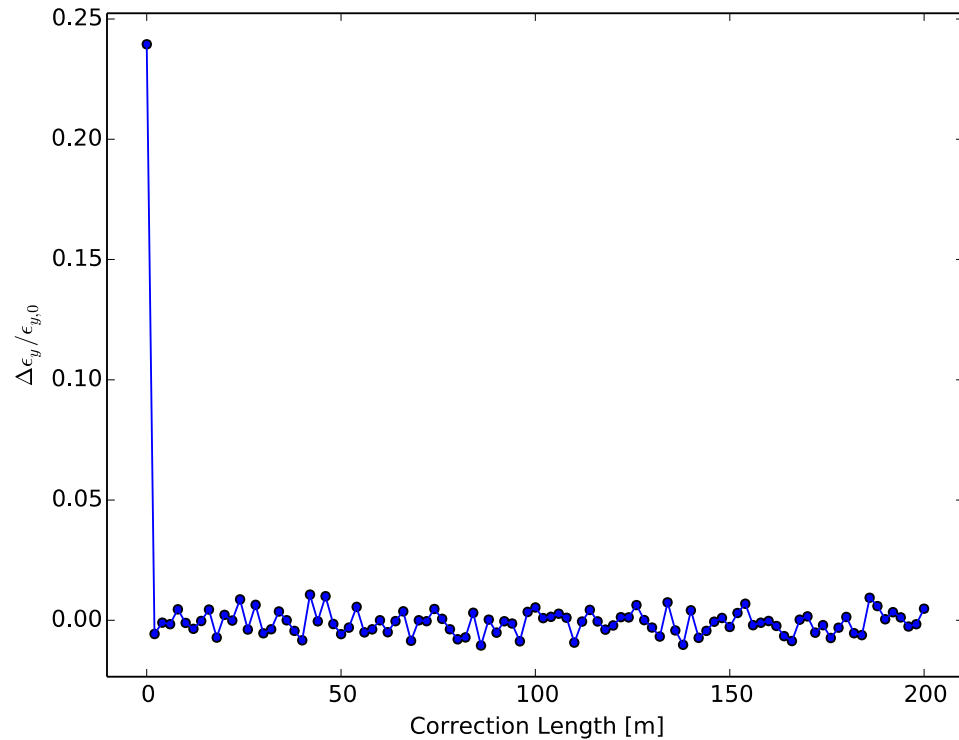
# Corrections

- Correcting like this can introduce discontinuities, which induce new small wavelengths: (below is  $L_{corr} = 50$  m)



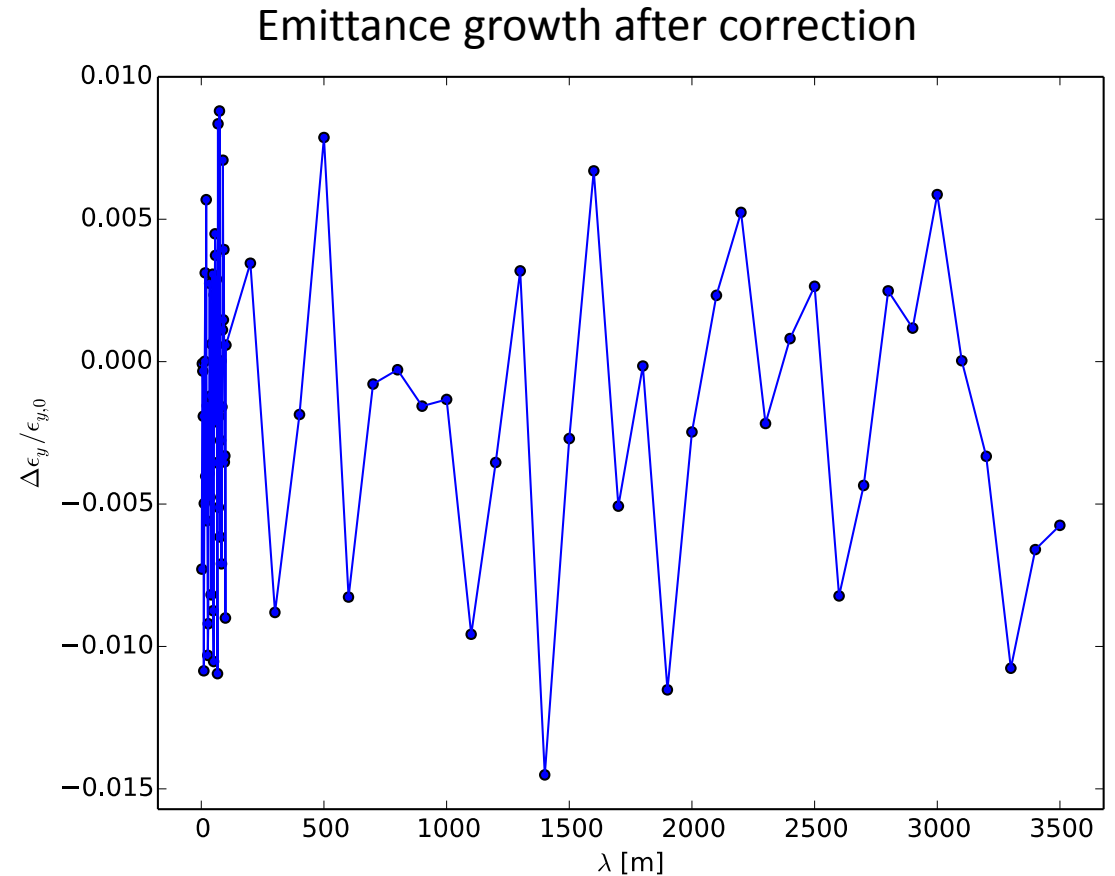
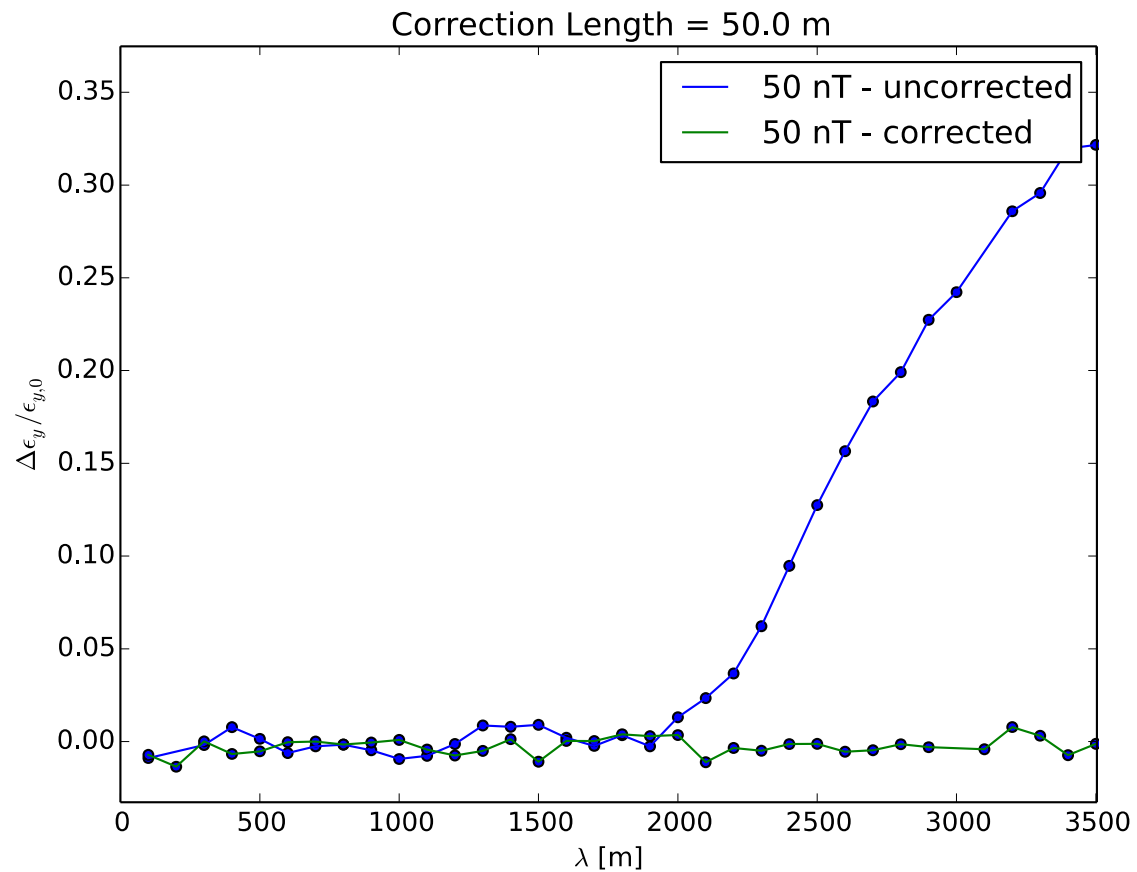
# Corrections – Varying $L_{corr}$

- Stray field with  $\lambda = 3000$  m and  $B_0 = 50$  nT, which causes emittance growth of  $\sim 0.24$ .



- Correction like this works when the corrector length is small.

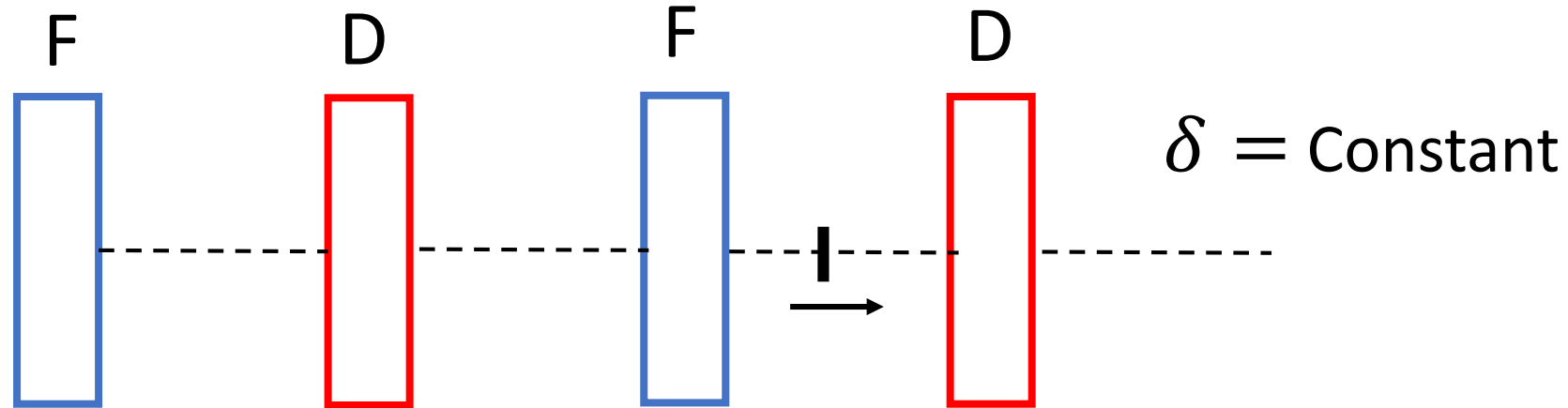
# Corrections - $L_{corr} = 50$ m



- Simulation with  $\lambda = 3000$  m,  $B_0 = 50$  nT.

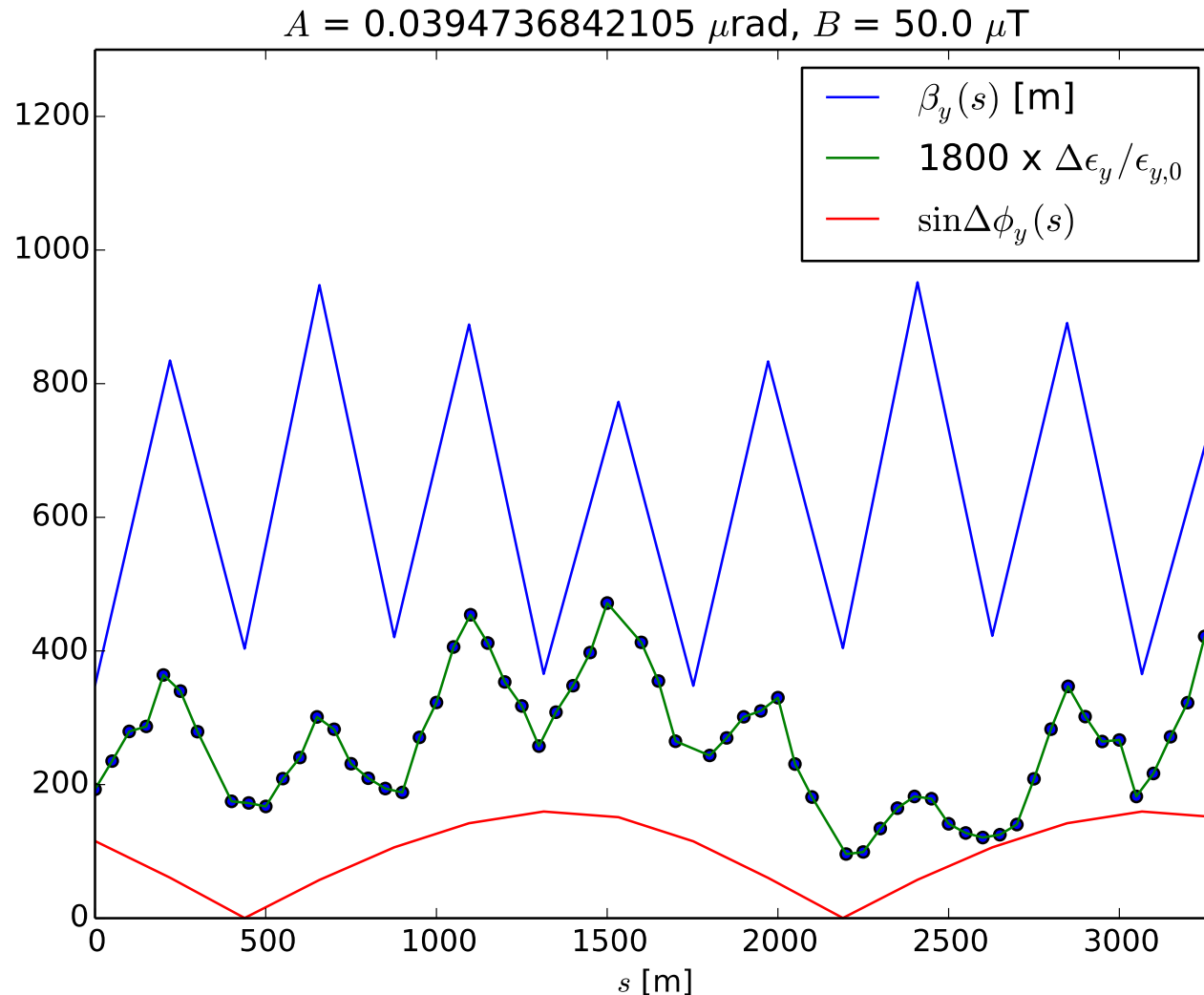
# Effect of a Single Kick

- The location of the kick matters.



- The equivalent kick of a  $50 \mu\text{T}$  magnetic field variation was placed at one location.
- This location for moved across the transfer line.

# Effect of a Single Kick



- Emittance growth is larger at locations of high beta.
- Locations that effect the beam most are high beta locations which also have a large phase advance.
- Looking into correcting only a fraction of the transfer line.