

# Prospects for determining the top quark Yukawa coupling at future $e^+e^-$ colliders

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Based on: Boselli, Hunter, Mitov arXiv:1805.12027

# Intro: why the top-Yukawa coupling?

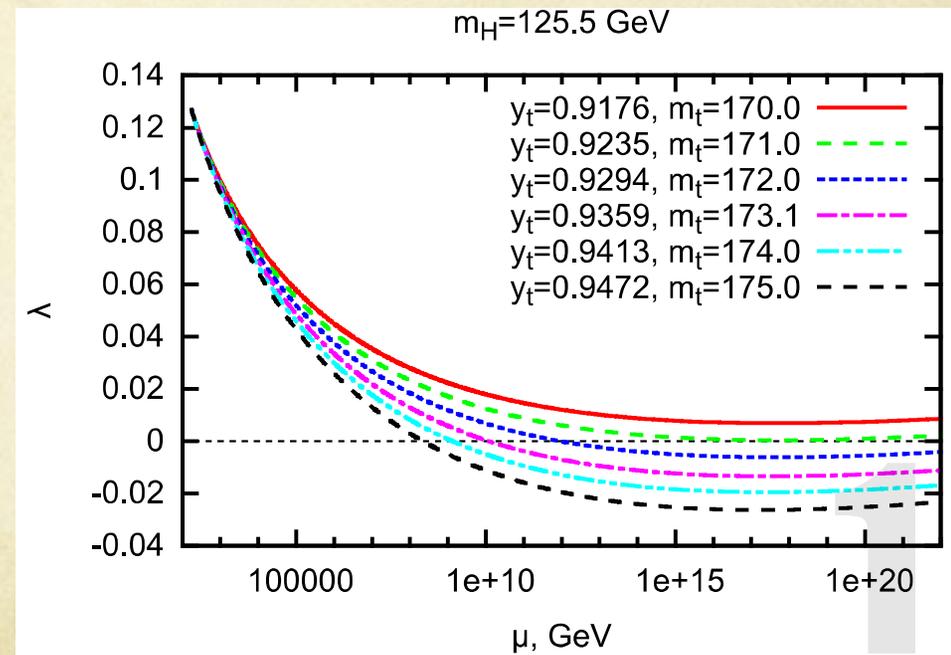
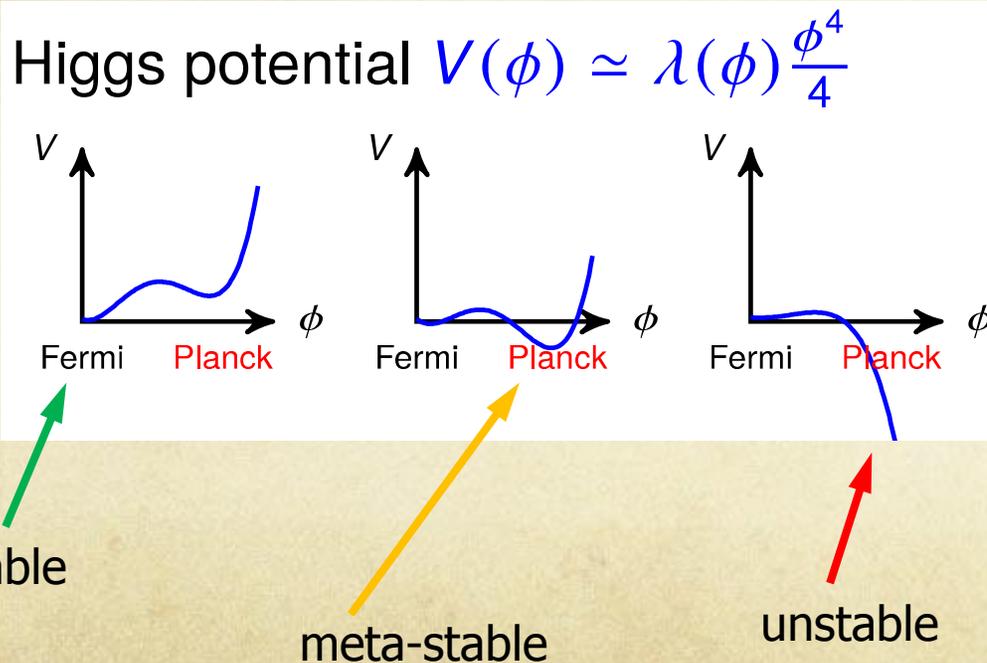
- ✓  $y_t$  is a derived parameter in the SM. It is known fairly accurately:

$$y_t(\mu = m_t) \approx 0.9369 + 0.0056 \left( \frac{\Delta m_t}{1\text{GeV}} \right) - 0.0006 \left( \frac{\Delta \alpha_s}{0.001} \right) \pm \dots$$

$$y_t(\mu = m_{\text{PL}}) \approx 0.3825 + 0.0051 \left( \frac{\Delta m_t}{1\text{GeV}} \right) - 0.003 \left( \frac{\Delta \alpha_s}{0.001} \right) \pm \dots$$

- ✓ The dominant uncertainty is  $m_{\text{top}}$  but, currently, only at the sub-percent level!
- ✓ It is very relevant for understanding the stability of EW vacuum

Plots courtesy of F. Bezrukov



# Intro: why the top-Yukawa coupling?

- ✓ We would like to measure  $y_t$  directly and verify its SM value
- ✓ It is quite likely that  $y_t$  is modified if BSM physics is present

$$y_t = y_t^{\text{SM}} + \Delta y_t$$

- ✓ How to measure  $\Delta y_t$ ?
- ✓ And here is the puzzle:
  - ✓ At the LHC we may be able to measure  $y_t$  with 10% precision (at HL-LHC)
  - ✓ A 100 TeV hadron collider can measure  $y_t$  with 1% precision  
Mangano, Plehn, Reimitz, Schell, Shao '15
- ✓ What about  $e^+e^-$  colliders?
  - ✓ Usual wisdom: obtain  $y_t$  from  $t\bar{t}h$  final states.
    - ✓ This offers clean(er) interpretation of the measurement
    - ✓ However, we need a 500GeV c.m. energy to produce  $t\bar{t}h$ !
  - ✓ Accessible only at CLIC and ILC (among all proposed colliders)
  - ✓ Existing studies show that  $y_t$  can be measured with 4.3% at CLIC and 5% at ILC
  - ✓ Such a prospect is a bit underwhelming, isn't it?

See talk by Yixuan Zhang on Tue

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# Intro: why the top-Yukawa coupling?

- ✓ Why is the precision from  $t\bar{t}h$  so low?
- ✓ Answer: luminosity is low, despite the very good sensitivity of the x-section w/r to  $y_t$ .
- ✓ In this work we ask the question: how can one do better (if possible at all)?
- ✓ Clearly, one has to look at different observables; ideally ones with high expected event yields.
- ✓ We consider events with a single Higgs in the final state but no top quarks.
  - ✓  $t\bar{t}$  final states also have some sensitivity to  $y_t$ . It is low; has been studied in the context of  $m_{\text{top}}$  determination
- ✓ We consider all proposed colliders (CEPC, CLIC, FCC-ee, ILC)
- ✓ A great benefit from using single Higgs final states: they can be produced in (relative) abundance at all energies and at all colliders!
- ✓ Where does the  $y_t$  sensitivity come from in such processes?
- ✓ From coupling of the Higgs to top quarks in loops (working in the NWA for the Higgs)

See talks by  
Yuichiro Kiyo  
Frank Simon  
Pablo Lopez

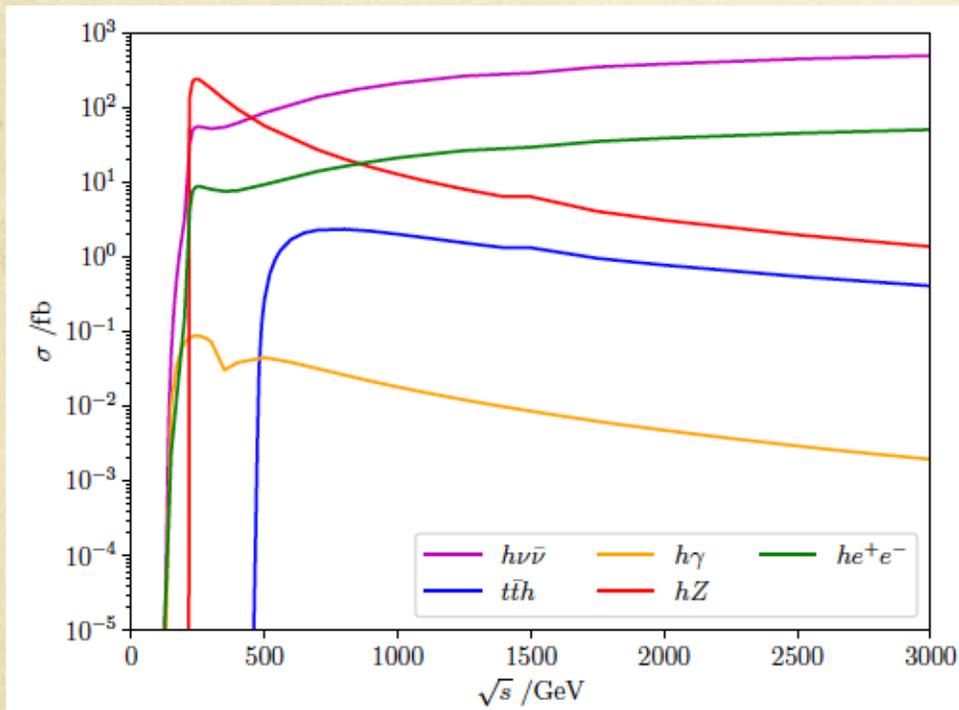
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# Our approach

✓ We have identified 3 such loop-induced processes:

- ✓  $e^+e^- \rightarrow h\gamma$  (with  $h \rightarrow bb$ )
- ✓  $h \rightarrow \gamma\gamma$  (from  $e^+e^- \rightarrow hZ/h\nu\nu/hee$ )
- ✓  $h \rightarrow gg$  (from  $e^+e^- \rightarrow hZ/h\nu\nu/hee$ )

✓ Here is a LO estimate of x-sections and event yields



	FCC-ee		CEPC
$\sqrt{s}$ (GeV)	240	350	240
$\mathcal{L}_{\text{int.}}$ ( $\text{fb}^{-1}$ )	$1.0 \cdot 10^4$	$2.6 \cdot 10^3$	$5.0 \cdot 10^3$
$\sigma_{hZ}$ (fb)	240	130	240
$\mathcal{N}_{hZ}$	$2.4 \cdot 10^6$	$3.38 \cdot 10^5$	$1.2 \cdot 10^6$
$\sigma_{\nu\bar{\nu}h}$ (fb)	54.4	54.7	54.4
$\mathcal{N}_{\nu\bar{\nu}h}$	$5.44 \cdot 10^5$	$1.42 \cdot 10^5$	$2.72 \cdot 10^5$
$\sigma_{eeh}$ (fb)	7.9	7.13	7.9
$\mathcal{N}_{eeh}$	$7.9 \cdot 10^4$	$1.85 \cdot 10^4$	$3.95 \cdot 10^4$
$\sigma_{h\gamma}$ (fb)	$8.96 \cdot 10^{-2}$	$3.18 \cdot 10^{-2}$	$8.96 \cdot 10^{-2}$
$\mathcal{N}_{h\gamma}$	896	82	448

	CLIC			ILC	
$\sqrt{s}$ (GeV)	350	1400	3000	250	500
$\mathcal{L}_{\text{int.}}$ ( $\text{fb}^{-1}$ )	$5.0 \cdot 10^2$	$1.5 \cdot 10^3$	$2.0 \cdot 10^3$	$2.0 \cdot 10^3$	$4.0 \cdot 10^3$
$\sigma_{hZ}$ (fb)	130	6.42	1.37	240	57.2
$\mathcal{N}_{hZ}$	$6.50 \cdot 10^4$	$9.6 \cdot 10^3$	$2.74 \cdot 10^3$	$4.80 \cdot 10^5$	$2.29 \cdot 10^5$
$\sigma_{\nu\bar{\nu}h}$ (fb)	54.4	293	498	55.0	85.2
$\mathcal{N}_{\nu\bar{\nu}h}$	$2.73 \cdot 10^4$	$4.39 \cdot 10^5$	$9.96 \cdot 10^5$	$1.10 \cdot 10^5$	$3.41 \cdot 10^5$
$\sigma_{eeh}$ (fb)	7.13	28.3	49.1	8.2	8.7
$\mathcal{N}_{eeh}$	$3.56 \cdot 10^3$	$4.24 \cdot 10^4$	$9.82 \cdot 10^4$	$1.64 \cdot 10^4$	$3.48 \cdot 10^4$
$\sigma_{t\bar{t}h}$ (fb)	-	1.33	0.41	-	0.27
$\mathcal{N}_{t\bar{t}h}$	-	1995	820	-	$1.08 \cdot 10^3$
$\sigma_{h\gamma}$ (fb)	$3.18 \cdot 10^{-2}$	$1.20 \cdot 10^{-2}$	$3.08 \cdot 10^{-3}$	$8.97 \cdot 10^{-2}$	$4.74 \cdot 10^{-2}$
$\mathcal{N}_{h\gamma}$	16	18	6	179	189

# Fit methodology

- ✓ Extract  $y_t$  from a  $\chi^2$  fit (assuming this is the only parameter to be fit; more later)

$$\chi^2(\Delta y_t) = \sum_{i=1}^{N_p} \sum_{j=1}^{N_d} \frac{[\mu_{ij}(\Delta y_t) - 1]^2}{\delta_{ij}^2}$$

Sums over these pairs of channels

Collider	$\sqrt{s}$ (GeV)	$\mathcal{L}$ ( $\text{fb}^{-1}$ )	$h \rightarrow gg$		$h \rightarrow \gamma\gamma$		$h \rightarrow b\bar{b}$	
			$hZ$	$\nu\bar{\nu}h$	$hZ$	$\nu\bar{\nu}h$	$h\gamma$	$t\bar{t}h$
FCC- $ee$	240	$1.0 \cdot 10^4$	1.4%	-	3.0%	-	4.4%	-
	350	$2.6 \cdot 10^3$	3.1%	4.7%	14%	21%	14%	-
CEPC	240	$5.0 \cdot 10^3$	1.2%	-	9.0%	-	6.2%	-
CLIC	350	$5.0 \cdot 10^2$	6.1%	10%	-	-	-	-
	1400	$1.5 \cdot 10^3$	-	5.0%	-	15%	-	8.0%
	3000	$2.0 \cdot 10^3$	-	4.3%	-	10%	-	12.5%
ILC	250	$2.0 \cdot 10^3$	2.5%	-	12%	-	10%	-
	500	$4.0 \cdot 10^3$	3.9%	1.4%	12%	6.7%	9.8%	9.9%

where: 
$$\mu_{ij} = \left( \frac{\sigma_i}{\sigma_i^{\text{SM}}} \right) \left( \frac{\Gamma_j}{\Gamma_j^{\text{SM}}} \right) \left( \frac{\Gamma_h}{\Gamma_h^{\text{SM}}} \right)^{-1}$$

- ✓ SM above means, basically, that  $\Delta y_t = 0$
- ✓ One-sigma uncertainties  $\delta_{ij}$  are taken from the literature.
  - ✓ An exception is  $e^+e^- \rightarrow h\gamma$  which is estimated by us based purely on the expected number of events (see previous slide). Likely to be optimistic.

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# Signal-strengths

✓ Here are the needed signal-strengths

$$\mu_{h\gamma} = \begin{pmatrix} \sqrt{s} = 240 \text{ GeV} \\ \sqrt{s} = 250 \text{ GeV} \\ \sqrt{s} = 350 \text{ GeV} \\ \sqrt{s} = 500 \text{ GeV} \end{pmatrix} = \frac{\sigma_{h\gamma}}{\sigma_{h\gamma}^{\text{SM}}} = 1 - \begin{pmatrix} 0.43 \\ 0.45 \\ 0.73 \\ 0.13 \end{pmatrix} \Delta y_t$$

$$\mu_{t\bar{t}h} = \begin{pmatrix} \sqrt{s} = 500 \text{ GeV} \\ \sqrt{s} = 1400 \text{ GeV} \\ \sqrt{s} = 3000 \text{ GeV} \end{pmatrix} = \frac{\sigma_{t\bar{t}h}}{\sigma_{t\bar{t}h}^{\text{SM}}} = 1 + \begin{pmatrix} 1.99 \\ 1.83 \\ 1.71 \end{pmatrix} \Delta y_t,$$

$$\mu_{h \rightarrow gg} = \frac{\Gamma_{h \rightarrow gg}}{\Gamma_{h \rightarrow gg}^{\text{SM}}} = 1 + 2\Delta y_t,$$

$$\mu_{h \rightarrow \gamma\gamma} = \frac{\Gamma_{h \rightarrow \gamma\gamma}}{\Gamma_{h \rightarrow \gamma\gamma}^{\text{SM}}} = 1 - 0.56\Delta y_t$$

✓ Derived by us at LO (full one loop):

- ✓ Compute x-sections and decay widths for a number of values of  $\Delta y_t$ ,
- ✓ Fit this with a quadratic polynomial,
- ✓ Take the linear approximation for small  $\Delta y_t$ .
- ✓ Bottom contribution to  $h \rightarrow g g$  neglected.

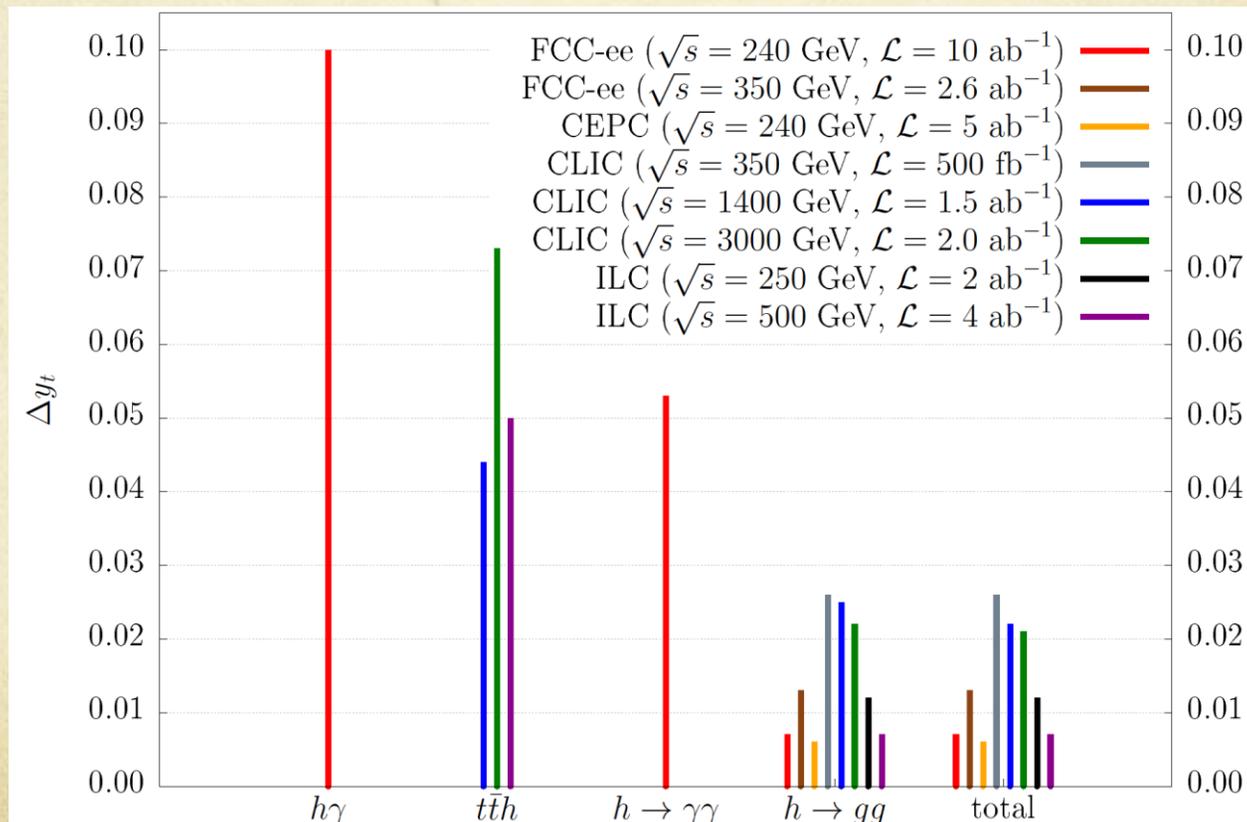
✓ Higher-order corrections in some cases have been included in the literature (CLIC 1.4 TeV).

Abramowicz et al., arXiv:1608.07538

- ✓ Slightly increases the expected precision

# Results: top-Yukawa precision prospects

Collider	$\sqrt{s}$ (GeV)	$\mathcal{L}$ ( $\text{fb}^{-1}$ )	$h \rightarrow gg$	$h \rightarrow \gamma\gamma$	$h\gamma$	$t\bar{t}h$	total
FCC- <i>ee</i>	240	$1.0 \cdot 10^4$	0.7%	5.3%	10%	-	0.7%
	350	$2.6 \cdot 10^3$	1.3%	21%	19%	-	1.3%
CEPC	240	$5.0 \cdot 10^3$	0.6%	16%	14%	-	0.6%
CLIC	350	$5.0 \cdot 10^2$	2.6%	-	-	-	2.6%
	1400	$1.5 \cdot 10^3$	2.5%	27%	-	4.4%	2.2%
	3000	$2.0 \cdot 10^3$	2.2%	18%	-	7.3%	2.1%
ILC	250	$2.0 \cdot 10^3$	1.2%	21%	23%	-	1.2%
	500	$4.0 \cdot 10^3$	0.7%	10%	75%	5.0%	0.7%



# Top-Yukawa precision prospects: few comments

- ✓  $h \rightarrow g g$  leads, by far, among all loop-induced processes
- ✓ This process offers potential  $y_t$  precision of about 0.6-0.7% at
  - ✓ 240 GeV CEPC and FCC-ee
  - ✓ 500 GeV ILC
- ✓  $h \rightarrow g g$  is better than  $t\bar{t}h$  for all energies and colliders by a factor of at least 2 (CLIC) and up to 7 (ILC)
- ✓  $e^+e^- \rightarrow h\gamma$  allows 10% determination. It is not great, but is comparable to HL-LHC
- ✓  $h \rightarrow \gamma\gamma$  allows about 5-6% precision at FCC-ee 240 GeV
- ✓ CLIC can measure  $y_t$  with precision of 2-2.5% (combining loop-induced and  $t\bar{t}h$ )
- ✓ ILC can measure  $y_t$  with precision of 1% or even better (combining loop-induced and  $t\bar{t}h$ )

# Limitations, assumptions and possible improvements

- ✓  $e^+e^- \rightarrow h\gamma$ : no detector simulation, efficiencies or background estimates. All done at LO.
- ✓  $m_{\text{top}}$ : we assume perfect knowledge of the top mass. In reality already after HL-LHC this error will be negligible
- ✓ Lack of proper EFT treatment:
  - ✓ We assume  $\Delta y_t$  is the only source of deviation from SM and so is the only parameter to fit
  - ✓ However, assuming BSM, no reason to have just one source of deviation from SM
  - ✓ Multiple Wilson coefficients will enter. This will dilute the expected precision on  $y_t$ .
  - ✓ However, after HL-LHC there will be many constraints on those coefficients.
- ✓ Assumed perfect knowledge of SM predictions.
  - ✓ In reality all is at LO (although fully one loop effects included)
  - ✓ NLO effects can be computed with some effort (2-loop amplitudes)
  - ✓ Realistic cuts imposed, etc.
- ✓ All of the above need to be done but we do not expect to change the picture qualitatively!

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# Conclusions

- This work tries to address the question: is it really not possible to measure  $y_t$  at a future  $e^+e^-$  collider with precision better than 4-5%?
- Such a prospect would be disheartening given we expect 10% from HL-LHC and 1% from a 100 TeV hadron collider
- This is an exploratory work. Its precision level is basic; still, we believe it is adequate in order to get a global picture about what is the ultimate possibility for measuring  $y_t$  at any one of the future  $e^+e^-$  colliders
- We consider *indirect* determination from loop-induced single Higgs processes
- Our findings are very promising. We find  $y_t$  can be measured
  - With precision as high as 0.6%
- This is almost an order of magnitude better than from purely  $t\bar{t}h$  final states.
- Such precision measurements can be done at any future  $e^+e^-$  colliders, especially at 240 GeV runs with  $hZ$  final states.
- Our work is very preliminary and can be made more precise in a number of ways
- We hope it provides useful input to the current discussion about which collider to build!

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