

International Workshop on Future Linear Colliders

LCWS2019

Sendai

October 28 – November 1

ANATOMY OF THE TTHH PHYSICS **AT HL-LHC**

Tao Liu

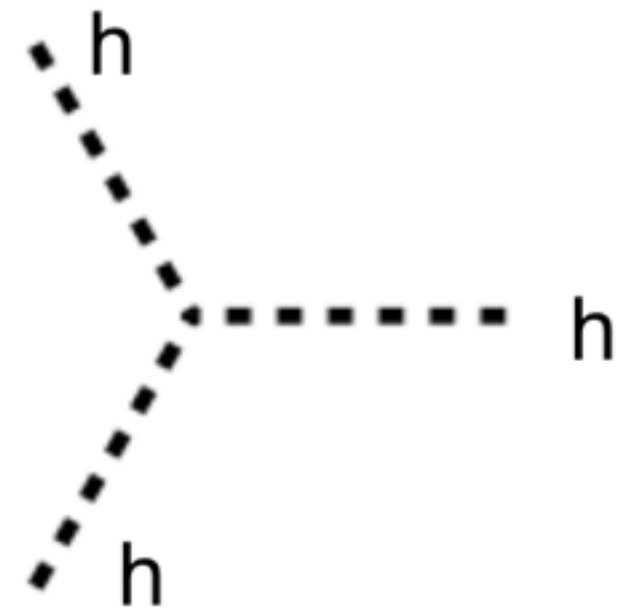
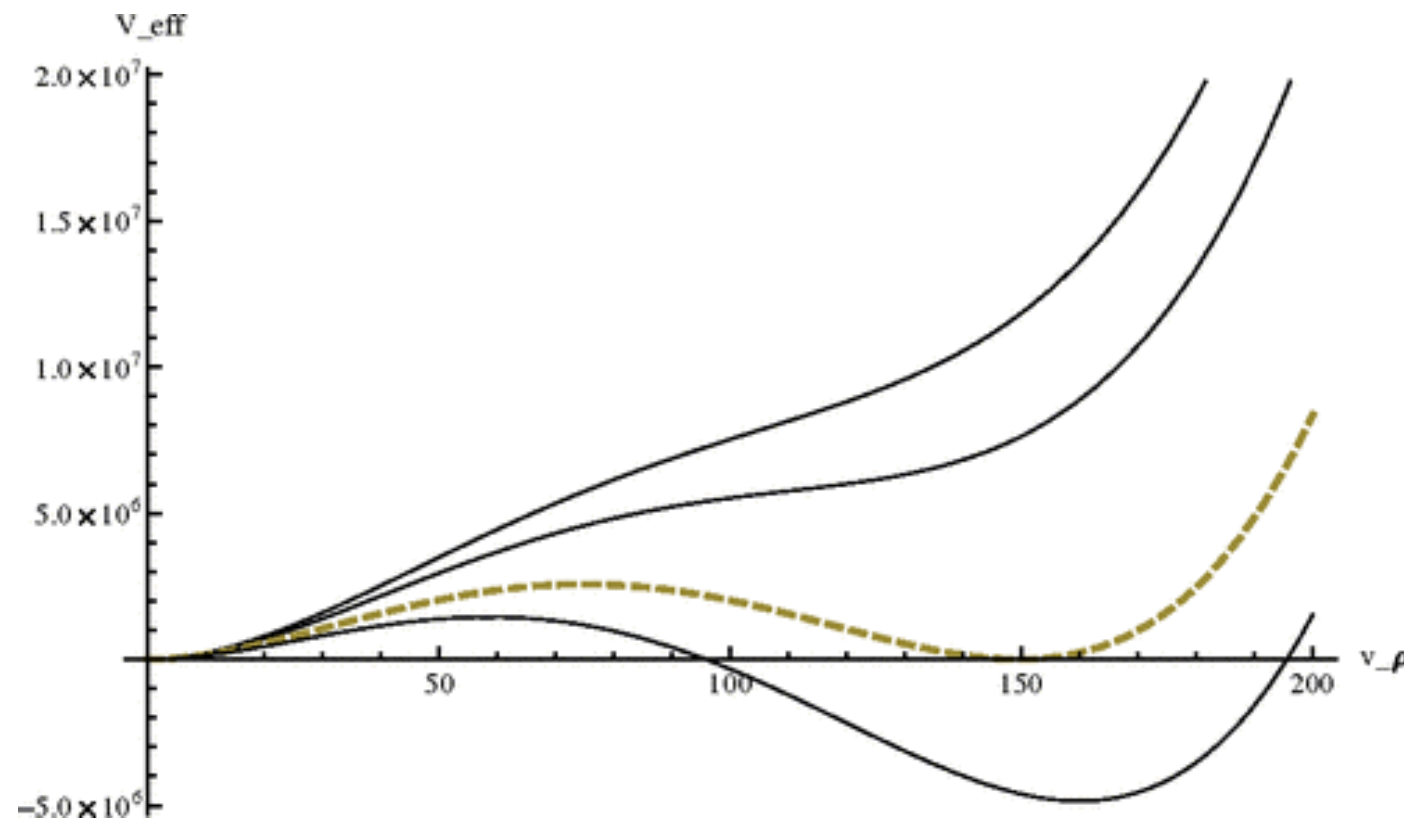
The Hong Kong University of Science and Technology

Based on [arXiv:1905.03772](https://arxiv.org/abs/1905.03772)

in collaboration with Lingfeng Li and Ying-Ying Li



Non-Resonant $t\bar{t}h$ - Cubic Higgs Coupling

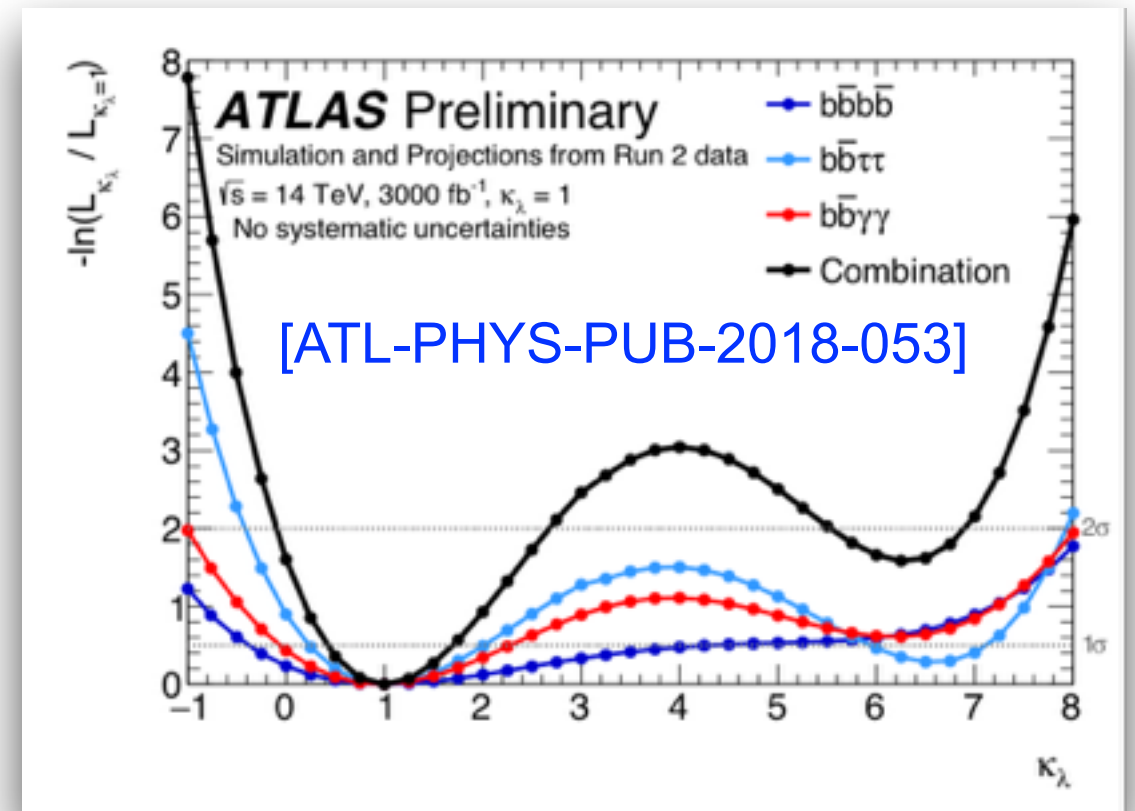
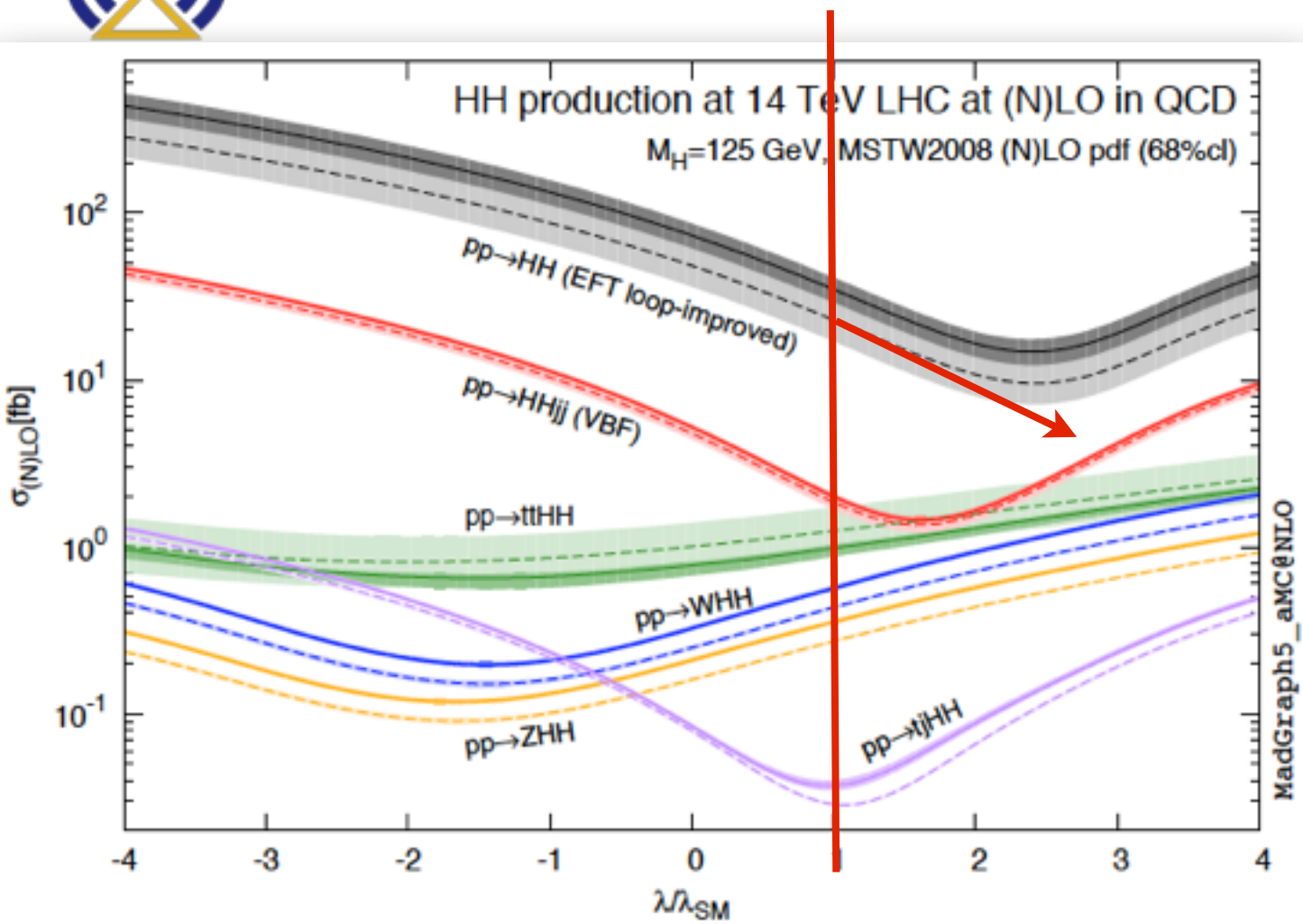


Known for some time that there exists a correlation between the nature of electroweak phase transition and the cubic Higgs coupling.

[A. Noble and M. Perelstein, '07]

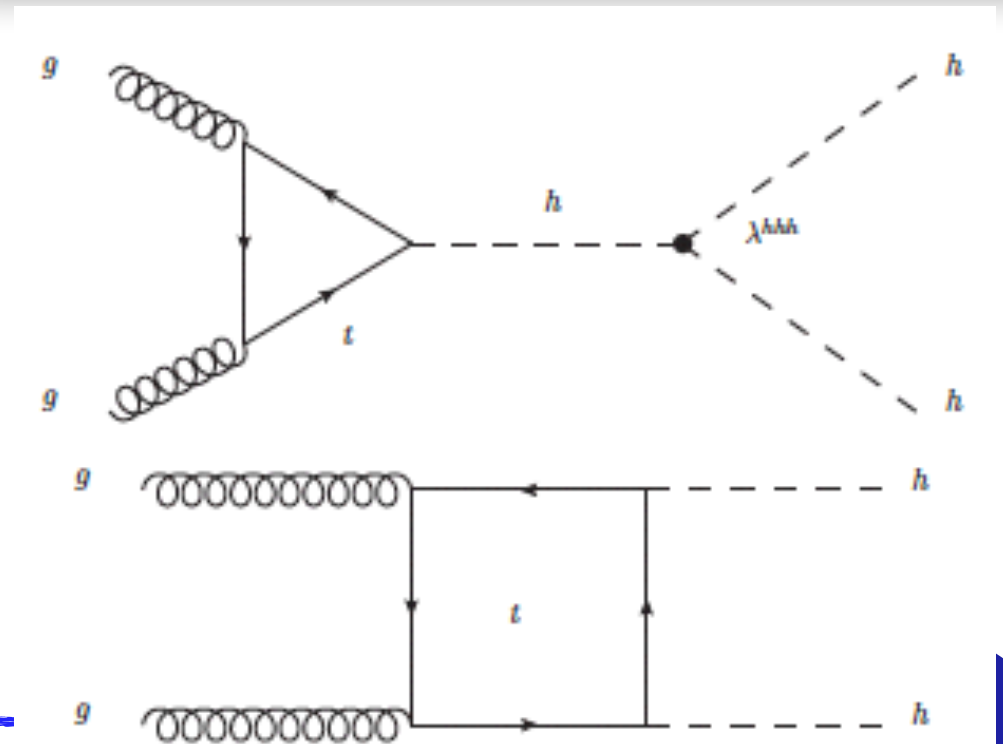


Non-Resonant tthh - Cubic Higgs Coupling



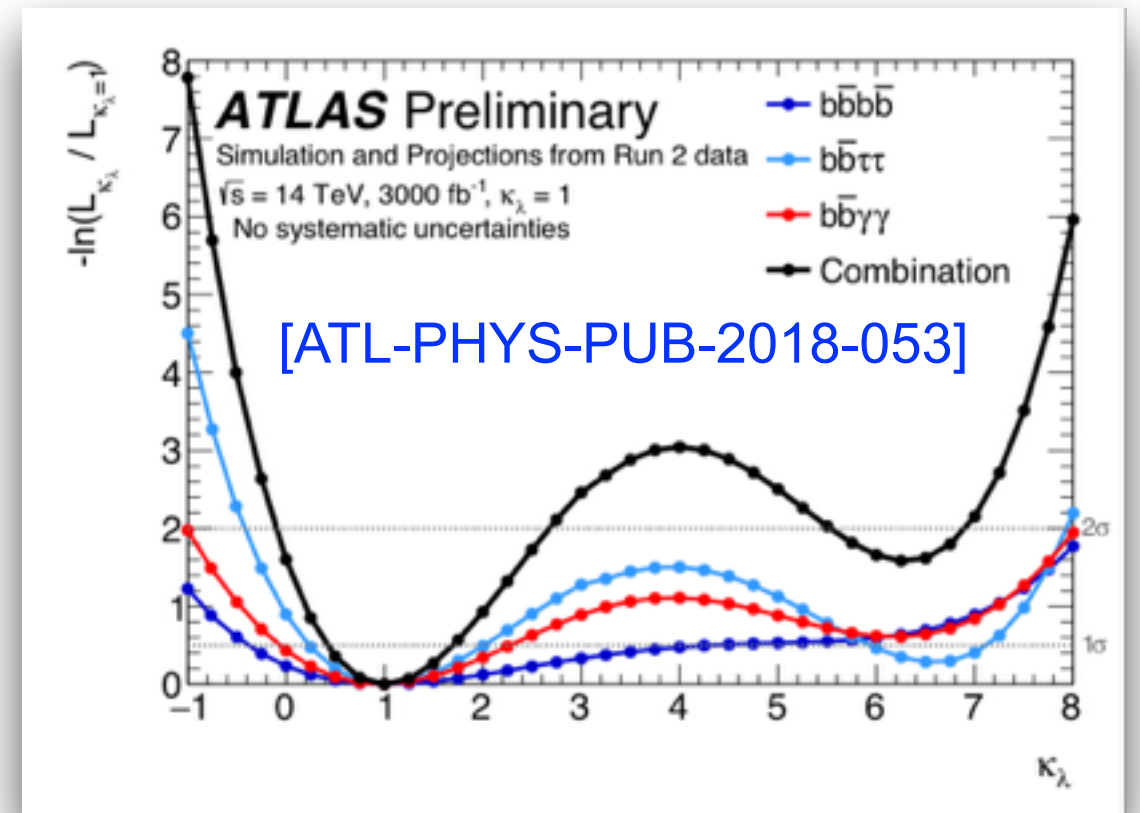
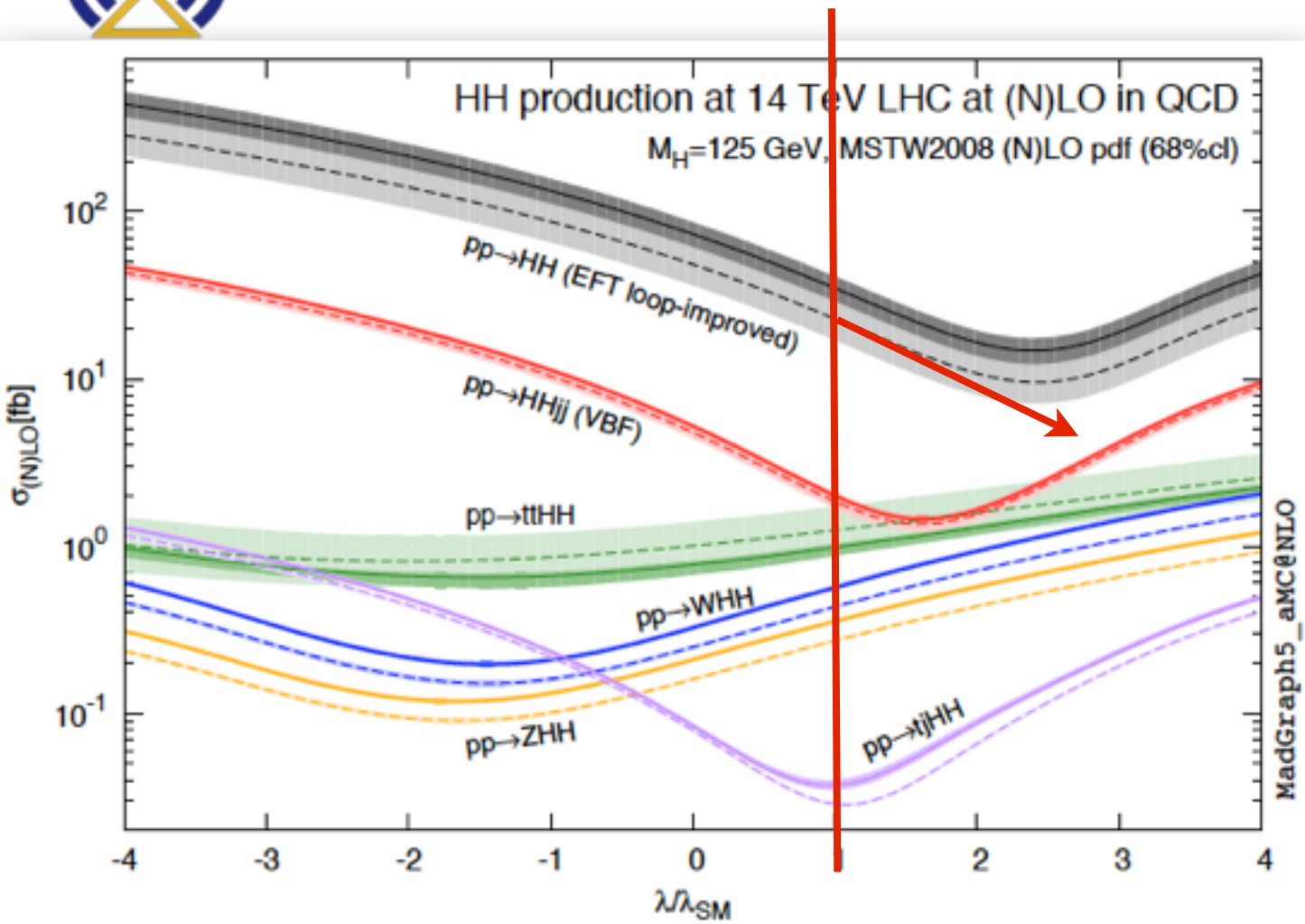
[R. Frederix et. al., arXiv:1401.7340]

- Most studies so far were focused on gluon fusion di-Higgs production
- Destructive interference => Sensitivity degeneracy w.r.t. $\kappa_\lambda = \lambda/\lambda_{\text{SM}}$ at HL-LHC



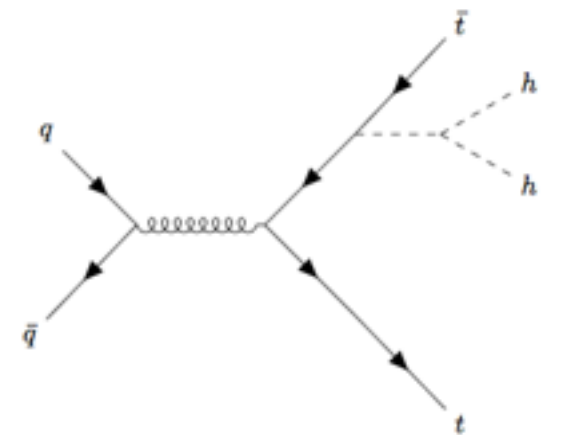
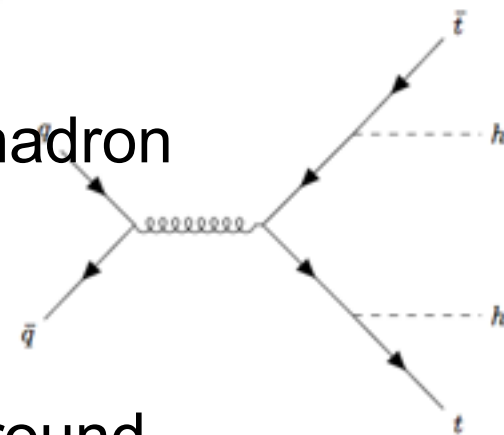


Non-Resonant tthh - Cubic Higgs Coupling



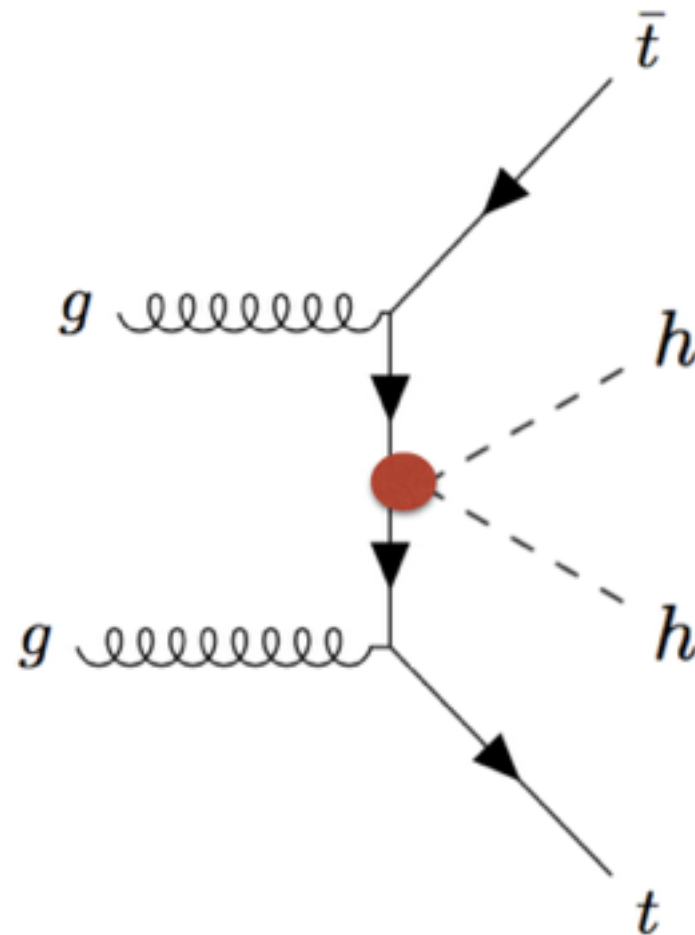
[R. Frederix et. al., arXiv:1401.7340]

- Can be broken using kinematics at future hadron colliders
- Alternatively, tthh could help at HL-LHC: constructive interference + reduced background





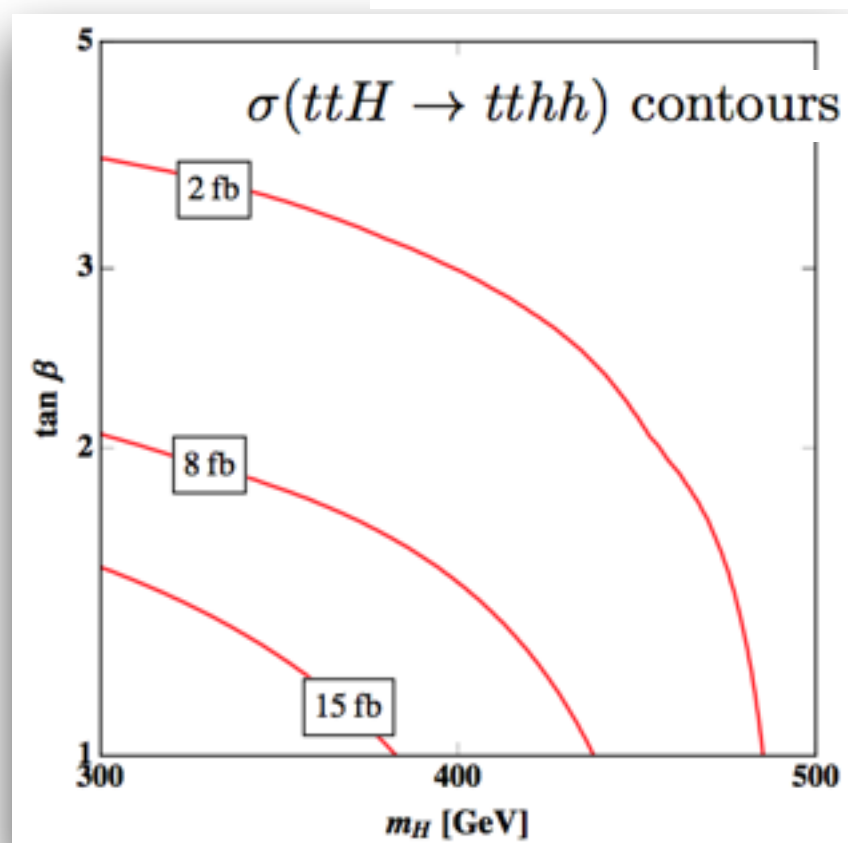
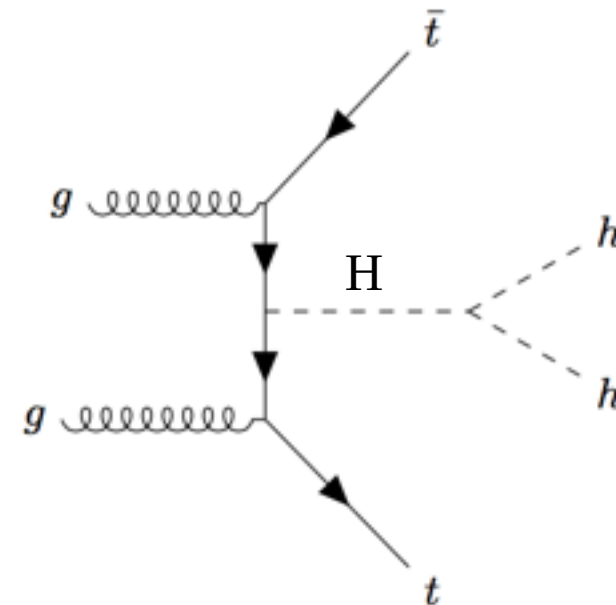
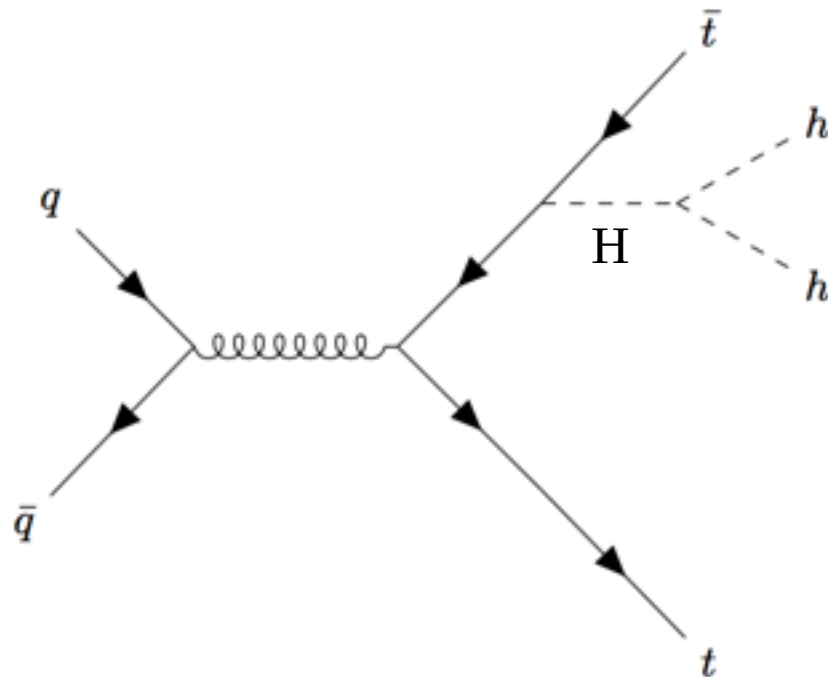
Non-Resonant tthh - Contact Interaction



- tthh contact interaction may arise from, e.g., 6D operator in CHM: $\frac{H^\dagger H Q_L H t_R}{\Lambda^2}$
- Been applied for electroweak baryogenesis (EWBG), as new CP-violating source - minimal EWBG [S. Huber, M. Pospelov and A. Ritz, arXiv:hep-ph/0610003]



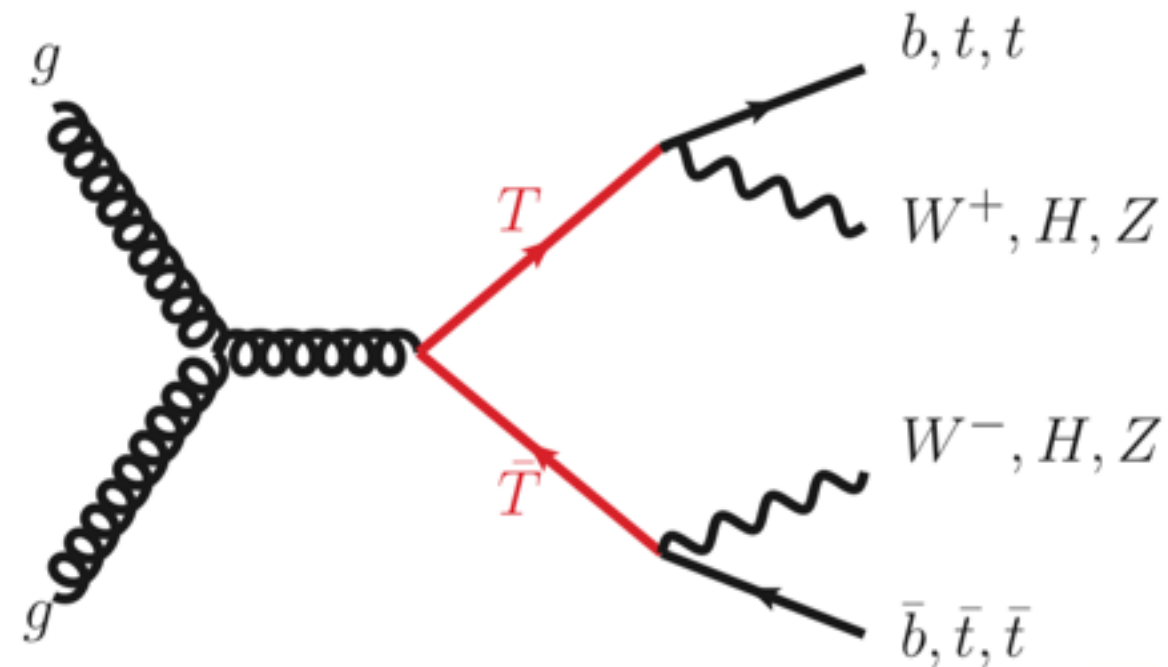
Resonant $t\bar{t}h$ - Heavy Higgs Siblings



- BSM physics with extended Higgs sector such as two Higgs doublet model (THDM)
- Especially important for Higgs sibling search, e.g., in the nearly decoupling limit $2m_h < m_H < 2m_t$ with low tan_beta in type II THDM or MSSM:
 - On-shell $H \rightarrow t\bar{t}$ turned off
 - H coupling to gauge bosons suppressed



Resonant tthh - Fermionic Top Partners



- One of the most important modes for searching pair production of top partners in CHMs. $\text{Br}(T \rightarrow t\bar{t})$ typically varies from 25% to 50%, depending on the EW charge carried by T

[R. Contino, et al. arXiv:hep-ph/0612190]



Early Sensitivity Analyses

	signal		backgrounds					
	$\xi = 1$	$\xi = 4$	$t\bar{t}b\bar{b}b\bar{b}$	$t\bar{t}h\bar{b}b$	$t\bar{t}hZ$	$t\bar{t}Z\bar{b}b$	$t\bar{t}ZZ$	$Wb\bar{b}b\bar{b}$
trigger	0.10	0.23	4.75	1.38	0.64	1.37	1.36×10^{-2}	1.33
jet cuts	7.40×10^{-2}	0.17	1.44	0.76	0.40	0.65	8.74×10^{-3}	7.46×10^{-2}
5 b tags	1.23×10^{-2}	2.83×10^{-2}	4.46×10^{-2}	6.19×10^{-2}	7.24×10^{-3}	4.43×10^{-2}	1.25×10^{-3}	5.35×10^{-4}
$2 \times h \rightarrow b\bar{b}$	7.33×10^{-3}	1.69×10^{-2}	1.59×10^{-2}	2.71×10^{-2}	3.41×10^{-3}	1.56×10^{-2}	4.28×10^{-4}	$< 1 \times 10^{-4}$
lep./had. t	5.04×10^{-3}	1.12×10^{-2}	9.50×10^{-3}	1.66×10^{-2}	2.29×10^{-3}	9.42×10^{-3}	2.69×10^{-4}	$< 1 \times 10^{-4}$
lep. t only	2.33×10^{-3}	5.29×10^{-3}	5.03×10^{-3}	9.36×10^{-3}	1.14×10^{-3}	4.90×10^{-3}	1.39×10^{-4}	$< 1 \times 10^{-4}$
had. t only	2.71×10^{-3}	5.93×10^{-3}	4.47×10^{-3}	7.20×10^{-3}	1.16×10^{-3}	4.44×10^{-3}	1.30×10^{-4}	$< 1 \times 10^{-4}$
6 b tags	2.21×10^{-3}	4.97×10^{-3}	3.80					
$2 \times h \rightarrow b\bar{b}$	1.81×10^{-3}	5.94×10^{-3}	2.01					

[C. Englert, F. Krauss, M. Spannowsky, J. Thompson, arXiv:1409.8074]

$\sqrt{s} = 14$ TeV	$t\bar{t}hh$	$t\bar{t}b\bar{b}b\bar{b}$	$t\bar{t}b\bar{b}c\bar{c}$	$t\bar{t}h\bar{b}b$	$t\bar{t}Z\bar{b}b$	$t\bar{t}hc\bar{c}$
Preselection	39.0	390.6	353.1	222.7	126.8	98.2
Di-Higgs rec.	33.0	269.3	242.1	171.0	93.5	76.8
Top rec.	19.5	160.7	149.0	102.8	54.6	47.1

[TL, H. Zhang, arXiv:1410.1855]

- Luminosity of $3/\text{ab} \Rightarrow S/\sqrt{B} \sim 1 - 2\sigma$
- A subsequent analysis [ATL-PHYS-PUB-2016-023] $\Rightarrow S/\sqrt{B} \sim 0.35\sigma$

[TL, talk in TOP2018]



Early Sensitivity Analyses

[C. Englert, F. Krauss, M. Spannowsky, J. Thompson, arXiv:1409.8074]

[TL, H. Zhang, arXiv:1410.1855]

- [ATL-PHYS-PUB-2016-023]

Yet, these analyses were not optimized.

- Handle I: clean channels such as same-sign di-lepton and multiple leptons, etc.
- Handle II: advanced tools such as BDT
 - All of these analyses were cut-based
 - Incapable for addressing complicated topology and kinematics: jet multiplicity, combinatorial backgrounds, etc.
- Most importantly, BSM physics may significantly modify the kinematics and signal rates of the SM $t\bar{t}h$, yielding valuable sensitivities at HL-LHC



Strategies of New Analysis

Analysis	$5b1\ell$	$5b2\ell$	$SS2\ell$	Multi- ℓ	$\tau\tau$
Triggers	$\cup\{1, 4, 5\}$	$\cup\{1, 2, 4, 5\}$	$\cup\{2\}$	$\cup\{1, 2, 3, 4, 5\}$	$\cup\{1, 4, 5, 6, 7\}$

5(or more)b+1 lepton

tt(semi-leptonic)bbbb

5(or more)b + OS dilepton

tt(leptonic)bbbb

SS dilepton (w/ ≥ 4 b jets)

tt(semi)VV*bb

Multi-lepton (w/ ≥ 4 b jets)

tt(semi)VV*bb

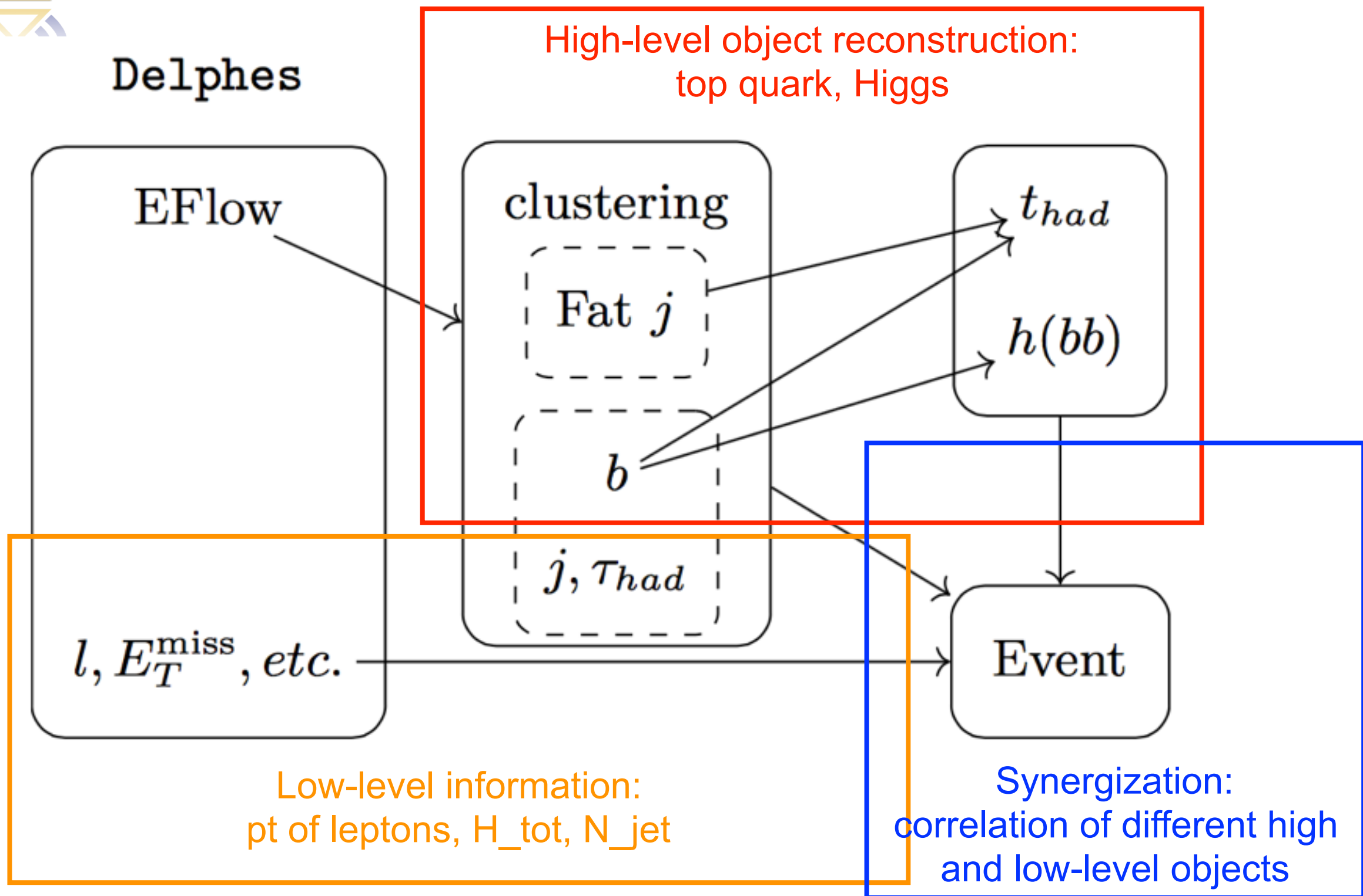
2 τ jets (w/ ≥ 4 b jets & 1 lepton)

tt(semi)bb $\tau\tau$



Strategies for New Analysis

Delphes





Strategies for New Analysis

- Analysis I: non-resonant $t\bar{t}h$ (SM)
- Analysis II: non-resonant $t\bar{t}h$ with anomalous cubic Higgs coupling and $t\bar{t}h$ contact interaction

$$\mathcal{L} \supset -y \frac{m_t}{v} t\bar{t}h - \kappa \frac{1}{3!} \frac{3m_h^2}{v} h^3 - c_t \frac{1}{2!} \frac{m_t}{v^2} t\bar{t}hh$$

- Analysis III: resonant $t\bar{t}h$ with heavy Higgs boson
- Analysis IV: resonant $t\bar{t}h$ with fermionic top partners



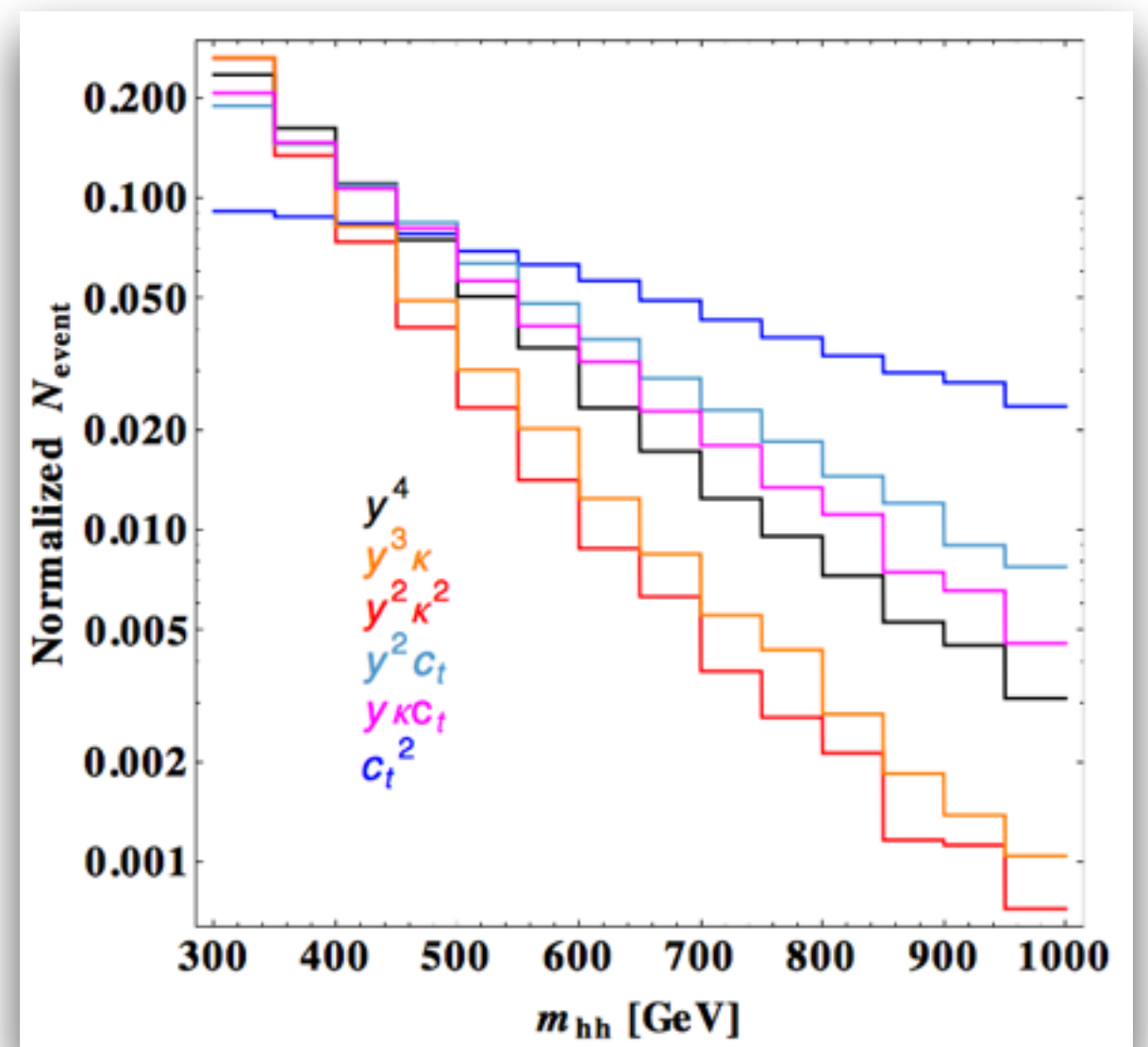
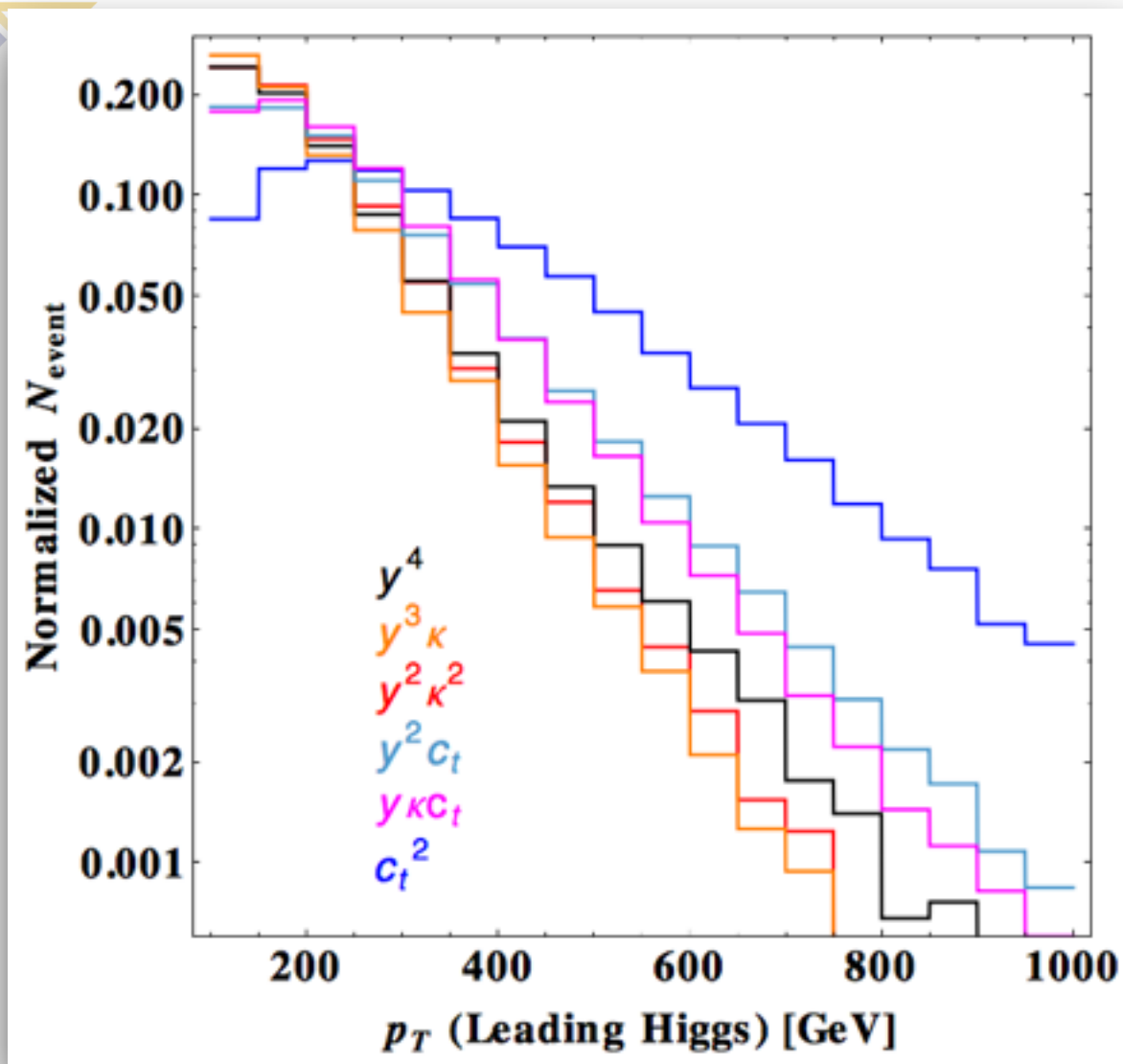
Analysis I: The SM $t\bar{t}h\bar{h}$ Production

	No cut	Preselection	$5b1\ell$	$5b2\ell$	$SS2\ell$	Multi- ℓ	$\tau\tau$
$t\bar{t}h\bar{h}$	2.9e3	7.37e2	50.9 (97.2)	6.1 (12.0)	14.6 (15.7)	8.6 (9.2)	3.6 (3.8)
$t\bar{t}4b$	1.1e6	1.79e5	6.56e3 (1.31e4)	664 (1.30e3)	212 (223)	115 (121)	94.1 (95.1)
$t\bar{t}2b2c$	3.1e5	4.28e4	621 (1.73e3)	59.4 (163)	38.0 (42.4)	24.1 (26.8)	43.6 (48.6)
$t\bar{t}VV$	4.4e4	3.64e3	20.7 (52.7)	3.5 (6.4)	51.8 (60.9)	32.4 (36.5)	3.1 (3.9)
$4t$	3.54e4	1.30e4	350 (804)	68.3 (152)	592 (635)	307 (324)	59.8 (64.2)
$t\bar{t}bbV$	8.29e4	1.54e4	353 (765)	47.8 (105)	114 (124)	203 (221)	22.2 (24.2)
$t\bar{t}bbh$	4.68e4	1.04e4	608 (1.15e3)	69.0 (136)	91.0 (98.0)	53.4 (56.2)	24.2 (25.9)
$t\bar{t}hZ$	4.65e3	881	28.1 (58.5)	4.1 (9.1)	8.8 (9.5)	18.5 (19.9)	2.3 (2.5)
Total	1.6e6	2.65e5	8.53e3 (1.76e4)	918 (1.88e3)	1.11e3 (1.19e3)	753 (806)	249 (265)
σ_{cut}			0.46 (0.62)	0.17 (0.23)	0.39 (0.40)	0.28 (0.29)	0.20 (0.20)
$(S/B)_{\text{cut}}(\%)$			0.42 (0.40)	0.47 (0.55)	1.1 (1.1)	0.9 (0.9)	1.1 (1.1)
σ_{BDT}			0.59 (0.79)	0.21 (0.30)	0.45 (0.46)	0.33 (0.35)	0.21 (0.21)
$(S/B)_{\text{BDT}}(\%)$			1.2 (1.0)	1.3 (1.6)	1.6 (1.6)	1.6 (1.9)	1.6 (1.6)
σ_{com}			0.86 (1.04)				

- $5b1\ell$ and $SS2\ell$ + multi- ℓ provide comparable sensitivities
- $5b1\ell$ gains more from BDT, in both statistical significance and S/B
- A combination yields ~ 0.9 sigma against background-only hypothesis



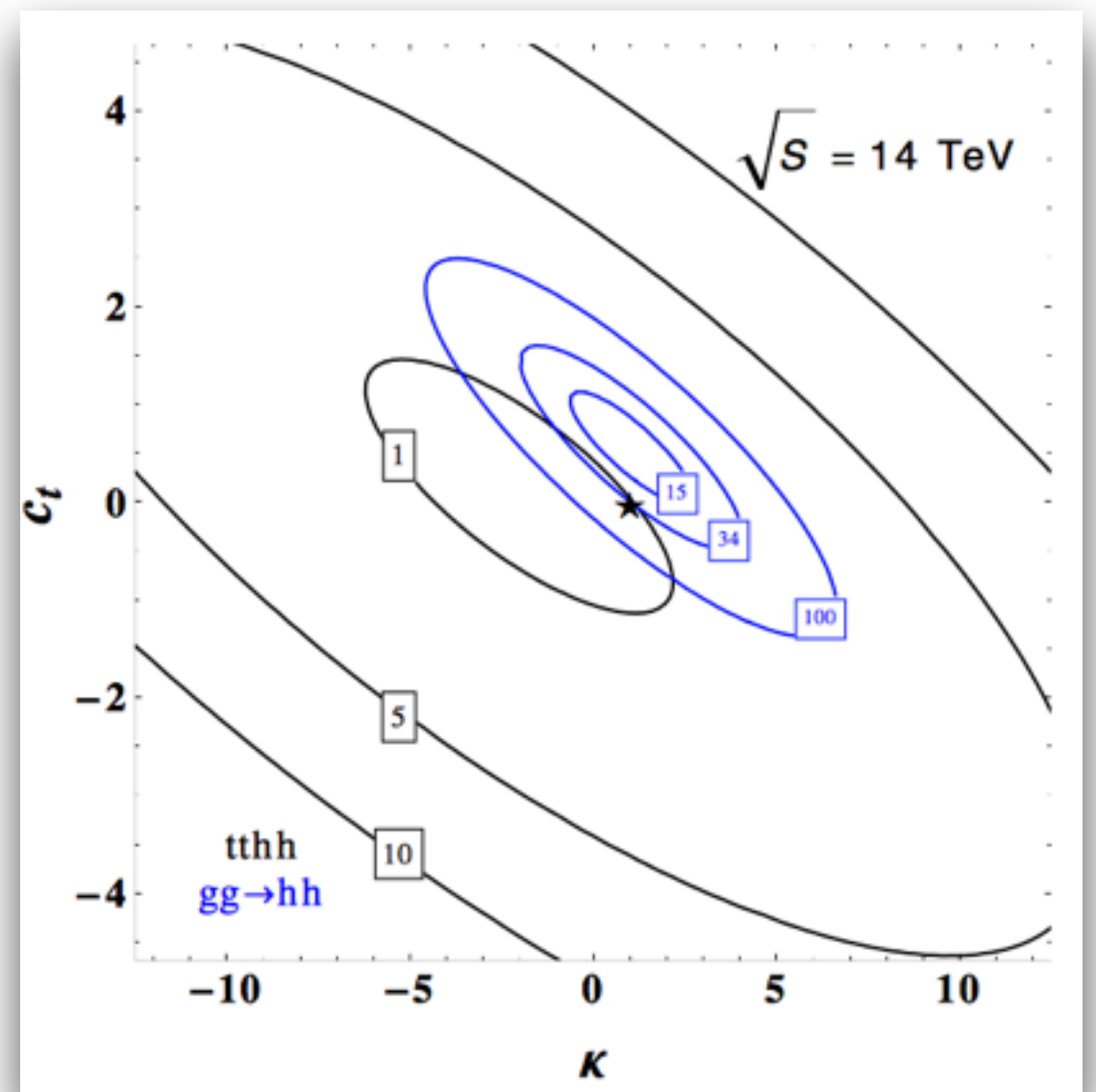
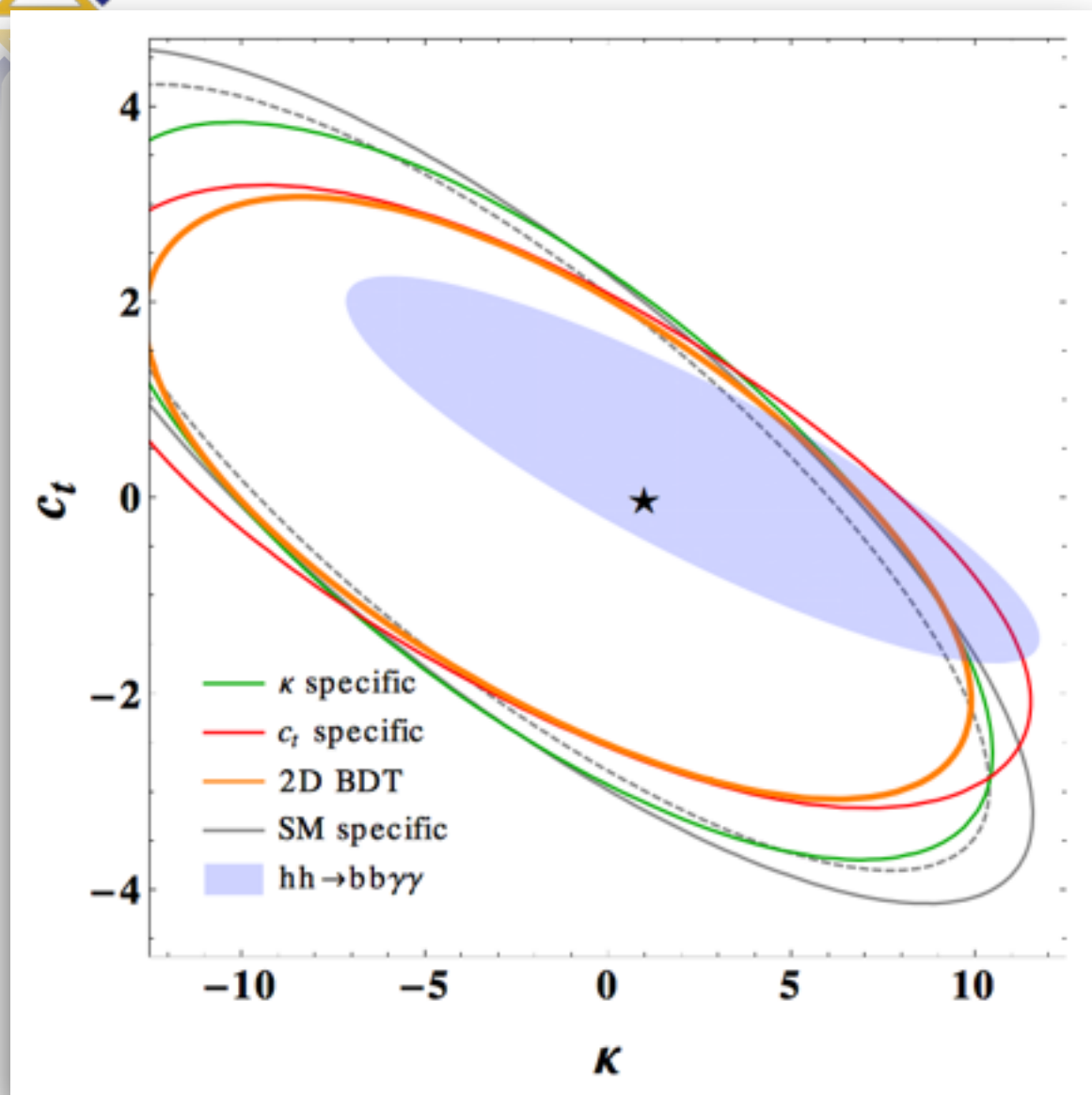
Analysis II: Anomalous $\kappa + c_t$



- Kinematics in opposite directions: soft (kappa) vs. hard (ct)
- 2-dim BDT, with $y^2 \kappa^2$ and c_t^2 kinematics being favored respectively by the two dimensions.



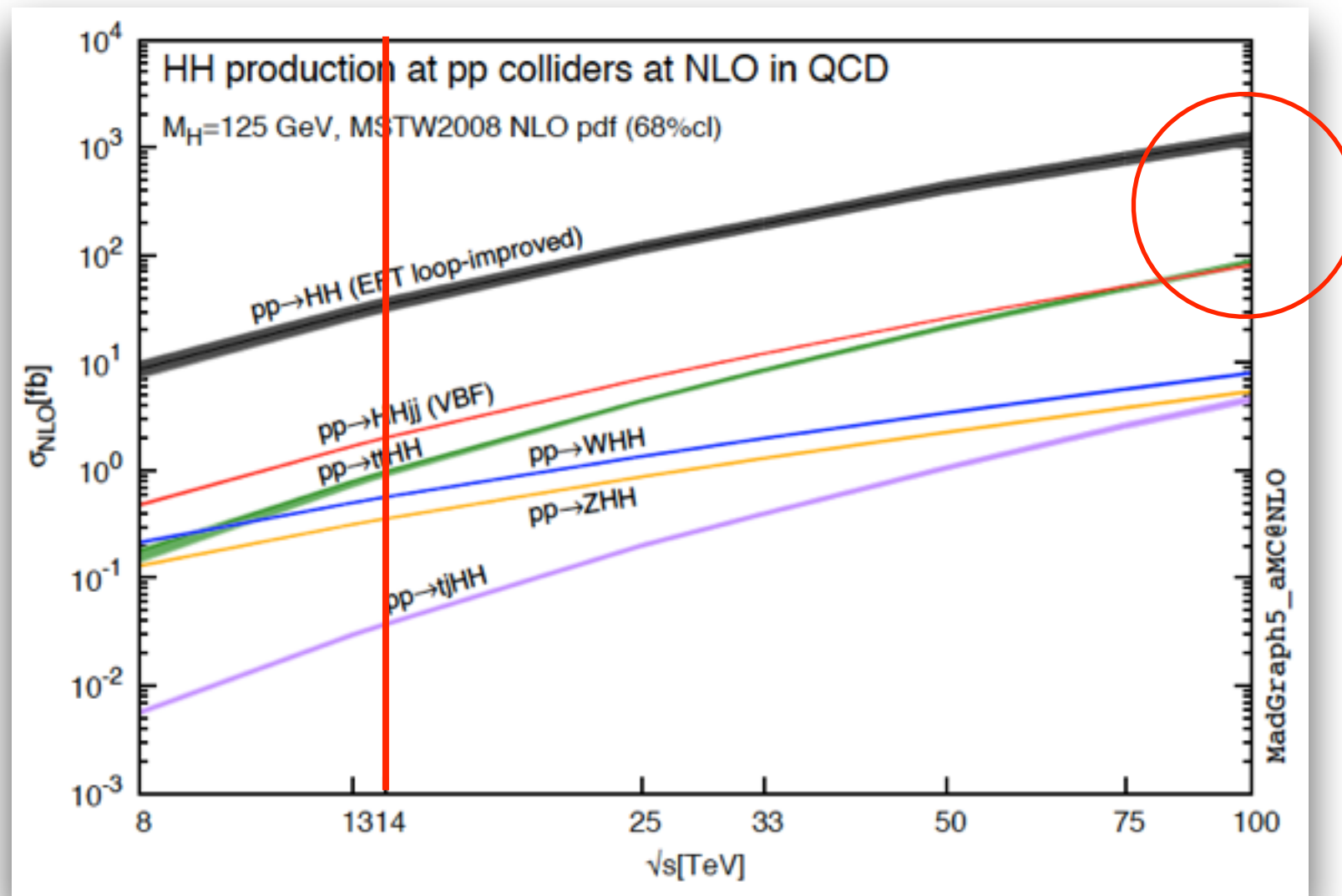
Analysis II: Anomalous $\kappa + c_t$



- Constrain κ to $\{6.9, -10\}$ for $c_t = 0$, and c_t to $\{-2.7, 1.8\}$ for $\kappa = 1$
- Partly breaks the sensitivity degeneracy w.r.t. the cubic Higgs self-coupling usually thought to exit in the gluon-fusion di-Higgs analysis at HL-LHC



Potential Gains at 100 TeV: Cross Section



[R. Frederix et. al., arXiv:1401.7340]

Cross section ratio (tthh/gluon fusion)
35 (14 TeV) vs. 14 (100 TeV)



Potential Gains at 100 TeV: Kinematics

The $t\bar{t}hh$ production becomes more sensitive to κ and c_t at non-linear level at 100 TeV

$$\frac{\sigma(gg \rightarrow hh \rightarrow bb\gamma\gamma)_{14}}{\sigma(gg \rightarrow hh \rightarrow bb\gamma\gamma)_{14}^{\text{SM}}} = 1.70 - 0.82\kappa + 0.12\kappa^2 - 3.79c_t + 0.98c_t\kappa + 2.68c_t^2$$

[A. Azatov et. al. arXiv:1502.00539]

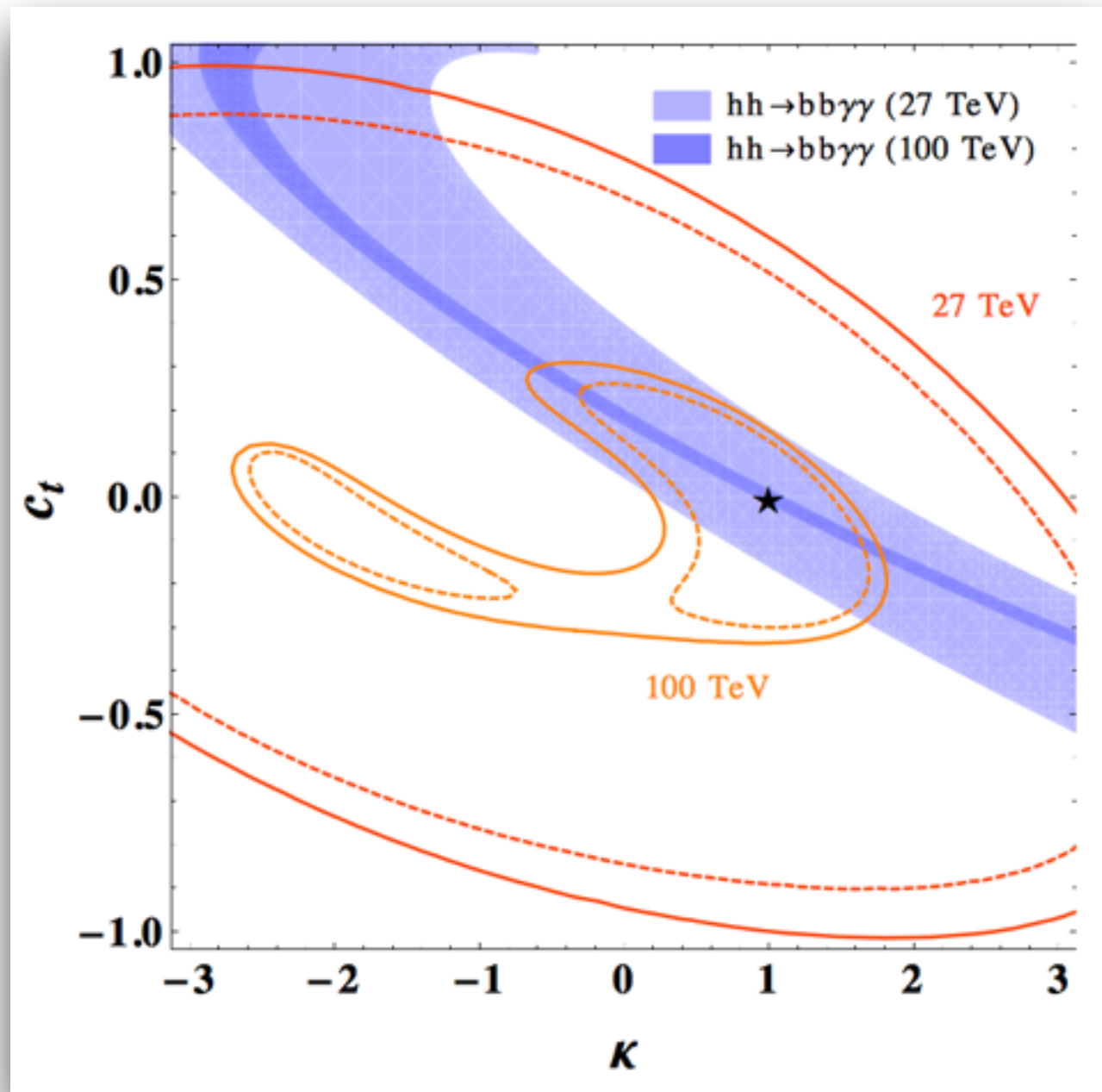
$$\frac{\sigma(gg \rightarrow hh \rightarrow bb\gamma\gamma)_{100}}{\sigma(gg \rightarrow hh \rightarrow bb\gamma\gamma)_{100}^{\text{SM}}} = 1.59 - 0.68\kappa + 0.09\kappa^2 - 3.83c_t + 0.92c_t\kappa + 3.20c_t^2$$

$$\frac{\sigma(t\bar{t}hh)_{14}}{\sigma(t\bar{t}hh)_{14}^{\text{SM}}} = 0.82 + 0.14\kappa + 0.04\kappa^2 + 0.28c_t + 0.21\kappa c_t + 0.44c_t^2$$

$$\frac{\sigma(t\bar{t}hh)_{100}}{\sigma(t\bar{t}hh)_{100}^{\text{SM}}} = 0.84 + 0.07\kappa + 0.09\kappa^2 + 0.15c_t + 0.41\kappa c_t + 1.73c_t^2$$



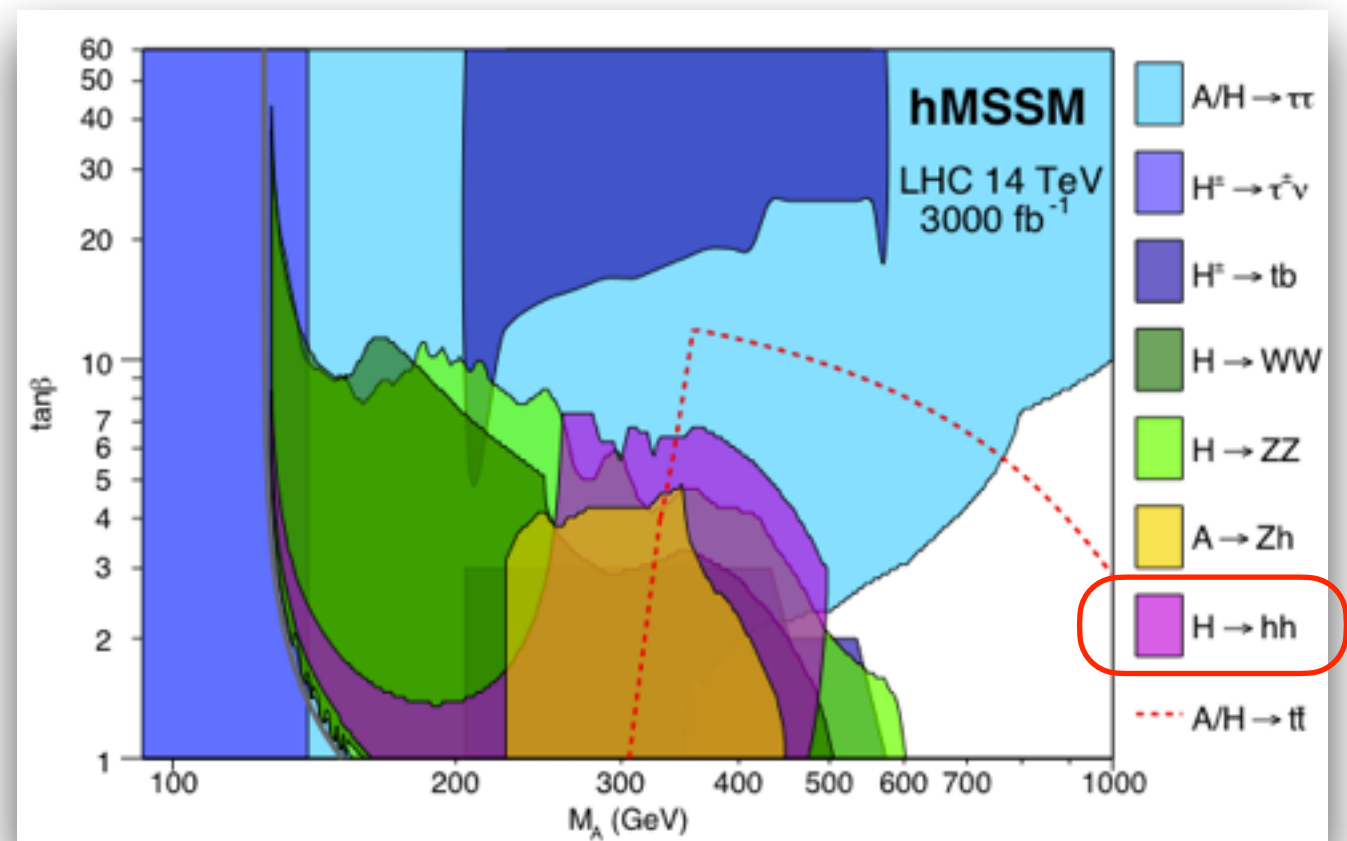
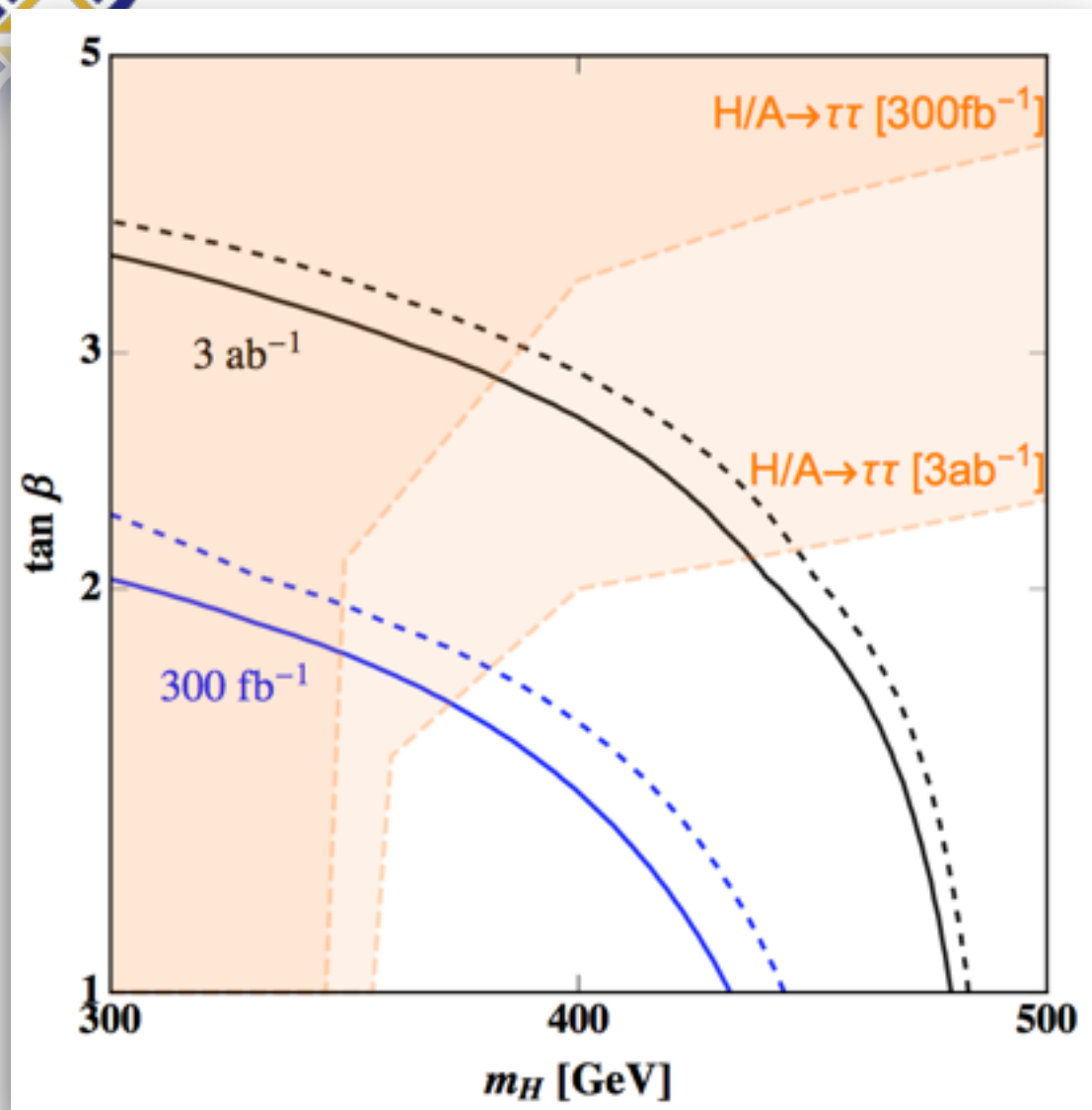
Sensitivity Projection at 100 TeV



- Assuming the signal efficiency and background rejection at HL-LHC to be unchanged
- $t\bar{t}hh$ in SM: ~ 14 sigma with 30/ab@100 TeV
- Constrain κ to $\{-2.6, -1.6\} \cup \{0.2, 1.6\}$ for $c_t = 0$, and c_t to $\{-0.34, 0.16\}$ for $\kappa = 1$
- Note, the reference sensitivity of $hh \rightarrow bb\gamma\gamma$ was not optimized using kinematics



Analysis III: Heavy Higgs Bosons

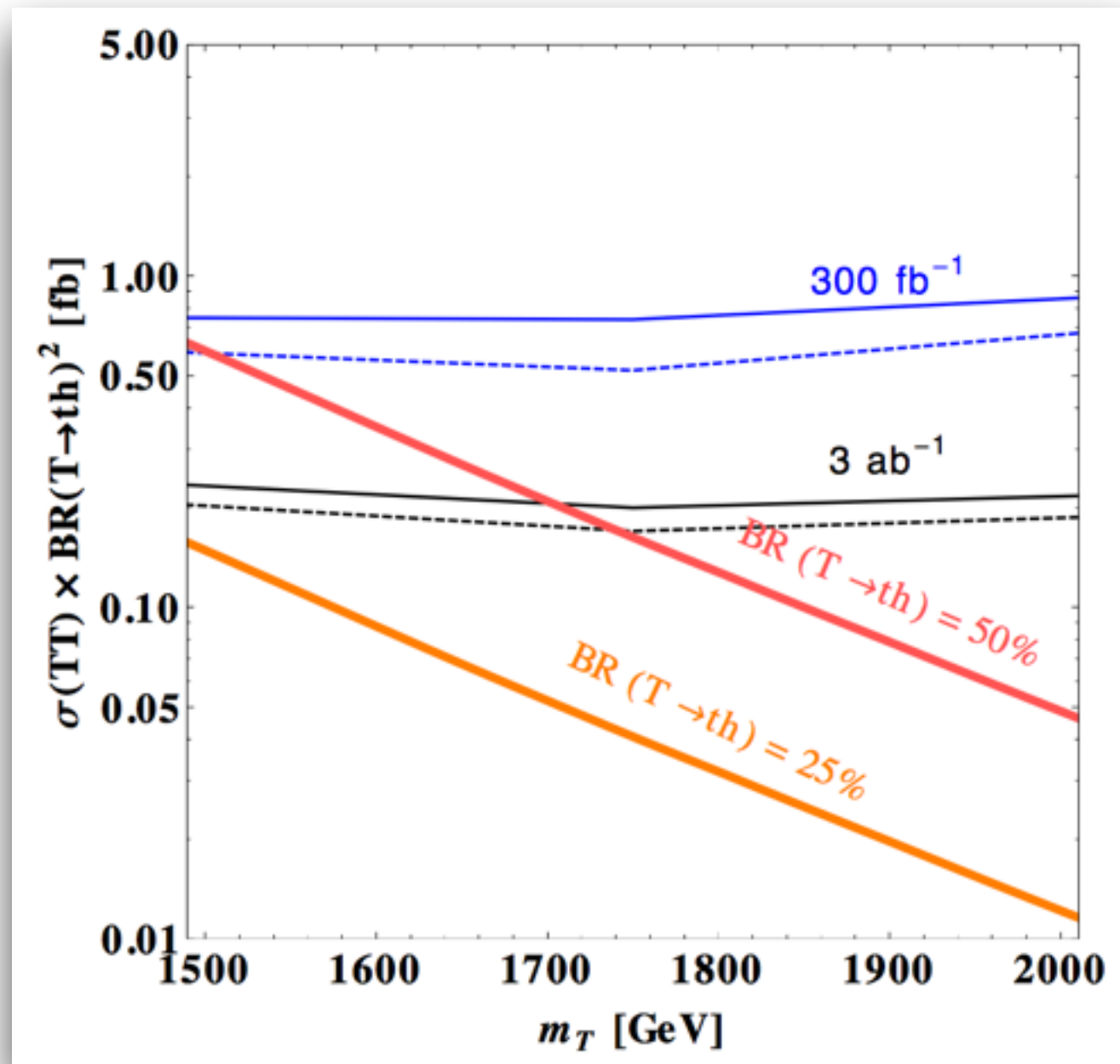


[Djouadi et. al.'arXiv:1502.05653]

- Complement the \tan_β favored channels well for $2m_h < m_H < 2m_t$ and even beyond
- Sensitivity is not bad compared to $gg \rightarrow H \rightarrow hh$: associated top pair ($ttH \rightarrow tthh$) compensates for relative smallness of the $tthh$ cross section



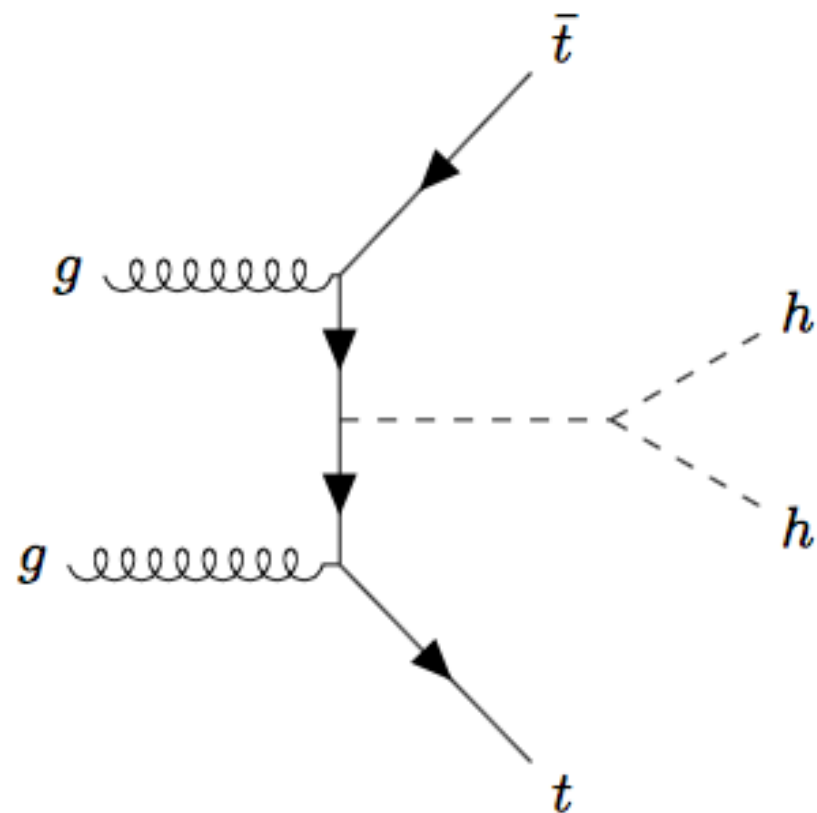
Analysis IV: Pair Production of Top Partners



Excludes the scenarios with $Br(T \rightarrow th) = 25\%$ and 50% up to 1.5 TeV and 1.7 TeV, respectively, at HL-LHC



Summary



- The $t\bar{t}h$ production can couple with BSM physics in multiple ways (deformed Higgs self-interaction, $t\bar{t}h$ contact interaction, new heavy resonances, etc.)
- Though its cross section in the SM is relatively low, the kinematics (e.g., resonance) and signal rate of the $t\bar{t}h$ could be significantly modified by the BSM physics, yielding valuable sensitivities at HL-LHC
- Given its potential role in probing BSM physics, a full exploration of its sensitivity at 100 TeV is important: still ongoing; will be presented somewhere else

Thank you!