# EWPOs using radiative return 

- motivation
- AlR $_{\text {(electron) }}$
- $A_{f}(f=b / c / \mu / \tau)$
- $\mathrm{R}_{\mathrm{f}}$
- $g_{L} \& g_{R}$

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recap 1: Higgs couplings are related to EW couplings (EWPOs)

$$
4 i \frac{c_{H L}^{\prime}}{v^{2}}\left(\Phi^{\dagger} t^{a} \overleftrightarrow{D}{ }^{\mu} \Phi\right)\left(\bar{L} \gamma_{\mu} t^{a} L\right)
$$

$$
i \frac{c_{H E}}{v^{2}}\left(\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi\right)\left(\bar{e} \gamma_{\mu} e\right)
$$


$e^{+} e^{-}->Z h h$

$\mathrm{e}^{+} \mathrm{e}^{-}->$Zh


Z-pole

- contact interactions from снн/снц'/сне in Higgs processes can be constrained by EWPOs at Z-pole: $\boldsymbol{A}_{L R}, \boldsymbol{\Gamma}_{\boldsymbol{I}}$


## a gift from ISR: radiative return @ ILC250



- ISR is mostly collinear
- asymmetric collision at Z-pole
- ISR (QED) retains initial e-/e+ chirality


## \# of radiative return events @ ILC250



- ~108 events at ILC250 with $2 \mathrm{ab}^{-1}$
- > 5 (100) times than all $Z$ at LEP (SLC)
- and now all with beam polarizations
- potentially much better $\boldsymbol{A}_{\boldsymbol{f}}$ and $\boldsymbol{R}_{\boldsymbol{f}}$ measurements


## study of e+e- -> $\mathrm{y} Z$ @ ILC250



- reconstruction method from LEP 2: Z mass can be determined by only directions of two fermions
- shortly after our SMEFT studies in 2017, I proposed to use this process at ILC250 for improving ALR (T.Barklow@AWLC17)
- SiD performed a fast simulation (T.Ueno@LCWS18)
- ILD full simulation ongoing (T.Mizuno)
- following are some expectations


## inputs and assumptions for $\mathrm{A}_{\mathrm{f}}$ and $\mathrm{R}_{\mathrm{f}}$

|  | efficiency | systematics |
| :---: | :---: | :---: |
| Z->hadrons | $73 \%{ }^{[1]}$ | 0 |
| Z-> $\mu \mu / \mathrm{ee}$ | 88\% ${ }^{[2]}$ | 0 |
| Z-> $T$ T (Rf) | 80\% ${ }^{[3]}$ | 0.1\% |
| Z->t $\tau$ (Af) | 80\% ${ }^{[3]}$ | 0 |
| Z->bb (Rf) | $73 \% \times 80 \%{ }^{[4]}$ | 0.1\% |
| Z->bb (Af) | 73\% $\times 40 \%{ }^{[4]}$ | 0 |
| Z-> cc (Rf) | $73 \% \times 30 \%{ }^{[4]}$ | 0.5\% |
| Z-> cc (Af) | $73 \% \times 10 \%{ }^{[4]}$ | 0 |

[1] Takayuki Ueno, Master Thesis
[2] T.Suehara et al; [3] D.Jeans; [4] R.Poeschl et al

## result: $A_{\text {LR }}$ using e+e- -> $\gamma Z\left(A_{e}\right)$

$$
\begin{aligned}
A_{L R}=\frac{\sigma_{L}-\sigma_{R}}{\sigma_{L}+\sigma_{R}}=A_{e} \quad & \Delta A_{e}=\frac{1}{\sqrt{N}} \sqrt{K^{2}-A_{e}^{2}} \\
& K=\frac{1+\left|P_{e^{-}} P_{e^{+}}\right|}{\left|P_{e^{-}}\right|+\left|P_{e^{+}}\right|} \quad \mathrm{N}: \text { \# of sig. events }
\end{aligned}
$$

- to measure: cross sections for ( $-0.8,+0.3$ ) and (+0.8,-0.3), using all hadronic and leptonic channels

| ILC250 | $\mathrm{N}_{\mathrm{L}}$ | $\mathrm{N}_{\mathrm{R}}$ | $\Delta \mathrm{A}_{\mathrm{LR}}$ |
| :---: | :---: | :---: | :---: |
| hadronic | 46 M | 31 M | 0.00015 |
| leptonic | 7.2 M | 4.9 M | 0.00035 |
| combined | 54 M | 36 M | 0.00014 |

- main systematics would be uncertainty of $K$ factor (effective polarization), which would be determined using WW data -> c.f. R.Karl's thesis, and further study would be needed


## result: $A_{L R}$ using e+e- -> $\gamma Z\left(A_{f}, f=\mu / \tau / b / c\right)$

$$
A_{L R F B}=\frac{\sigma_{L F}-\sigma_{L B}-\sigma_{R F}+\sigma_{R B}}{\sigma_{L F}+\sigma_{L B}+\sigma_{R F}+\sigma_{R B}}=\frac{3}{4} A_{f} \quad \Delta A_{f}=\frac{1}{\sqrt{N}} \sqrt{\frac{16}{9} K^{2}-A_{f}^{2}}
$$

- to measure: + cross sections in the forward/backward regions

|  | $\Delta \mathrm{Ab}$ | $\Delta \mathrm{Ac}$ | $\Delta \mathrm{A} \mu$ | $\Delta \mathrm{AT}$ |
| :---: | :---: | :---: | :---: | :---: |
| ILC250 | 0.00053 | 0.0014 | 0.00080 | 0.00083 |

- systematics: + uncertainties from charge identification, momentum directions -> c.f. study by Adrain for bb/cc, need further study here due to the boost of $Z$ (more forward jets); also need to study tau channel the about the momentum directions


## result: BR(Z->ff) using e+e- -> $\gamma Z\left(R_{f}, f=\mu / \tau / b / c\right)$

$$
R_{b / c}=\frac{\Gamma(Z \rightarrow b \bar{b} / c \bar{c})}{\sum_{q} \Gamma(Z \rightarrow q \bar{q})} \quad R_{\mu / \tau}=\frac{\sum_{q} \Gamma(Z \rightarrow q \bar{q})}{\Gamma\left(Z \rightarrow \mu^{+} \mu^{-} / \tau^{+} \tau^{-}\right)}
$$

- to measure: $\sigma x B R$, relative simple

|  | $\Delta R b$ | $\Delta R c$ | $\Delta R \mu$ | $\Delta R \tau$ | $\Delta R e$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| ILC250 | 0.00023 | 0.00087 | 0.011 | 0.024 | 0.024 |

- systematics: uncertainty on various efficiencies


## result: Zff couplings (gL/gR) using e+e--> $\mathrm{Y} Z$

$$
\begin{gathered}
\delta g_{L}=\frac{A_{f}}{2\left(A_{f}+1\right)} \delta A_{f} \oplus \frac{1}{2} \delta R_{f} \oplus \frac{1}{2} \delta \Gamma_{\text {had }} \\
\delta g_{R}=\frac{A_{f}}{2\left(A_{f}-1\right)} \delta A_{f} \oplus \frac{1}{2} \delta R_{f} \oplus \frac{1}{2} \delta \Gamma_{\text {had }} \\
\uparrow
\end{gathered}
$$

note the large factor for $b_{R}$ coupling
result: EWPOs from e+e--> $\gamma Z$

| Quantity | Value | $\begin{gathered} \text { current } \\ \delta\left[10^{-4}\right] \\ \hline \end{gathered}$ | $\begin{gathered} \text { GigaZ } \\ \delta_{\text {stat }}\left[10^{-4}\right] \end{gathered}$ | $\delta_{\text {sys }}\left[10^{-4}\right]$ | $\begin{gathered} 250 \mathrm{GeV} \\ \delta_{\text {stat }}\left[10^{-4}\right] \end{gathered}$ | $\delta_{\text {sys }}\left[10^{-4}\right]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| boson properties |  |  |  |  |  |  |
| $m_{W}$ | 80.379 | 1.5 | - | - |  | $0.3{ }^{\circ}$ |
| $m_{Z}$ | 91.1876 | 0.23 | - | - | - | - |
| $\Gamma_{Z}$ | 2.4952 | 9.4 |  | 3.2 | - | - |
| $\Gamma_{Z}($ had $)$ | 1.7444 | 11.5 |  | 3.2 | - | - |
| $Z$-e couplings |  |  |  |  |  |  |
| $1 / R_{e}$ | 0.0482 | 24. | 2. | $4^{\dagger}$ | 5.5 | $10^{+}$ |
| $A_{e}$ | 0.1515 | 139. | 0.1 | 5 * | 9.5 | 3 * |
| $g_{L}^{e}$ | -0.632 | 16. | 1. | 2.5 | 2.8 | 7.6 |
| $g_{R}^{e}$ | 0.551 | 18. | 1. | 2.5 | 2.9 | 7.6 |
| $Z-\ell$ couplings |  |  |  |  |  |  |
| $1 / R_{\mu}$ | 0.0482 | 16. | 2. | 2. ${ }^{\dagger}$ | 5.5 | $10^{+}$ |
| $1 / R_{\tau}$ | 0.0482 | 22. | 2. | 4. ${ }^{\dagger}$ | 5.7 | 10 + |
| $A_{\mu}$ | 0.1515 | 991. | 2. | 5 * | 54. | 3 * |
| $A_{\tau}$ | 0.1515 | 271. | 2. | 5 * | 57. | 3 * |
| $g_{L}^{\mu}$ | -0.632 | 66. | 1. | 1.8 | 4.5 | 7.6 |
| $g_{R}^{\mu}$ | 0.551 | 89. | 1. | 1.9 | 5.5 | 7.6 |
| $g_{L}^{\tau}$ | -0.632 | 22. | 1. | 2.5 | 4.7 | 7.6 |
| $g_{R}^{\tau}$ | 0.551 | 27. | 1. | 2.5 | 5.8 | 7.6 |
| $Z-b$ couplings |  |  |  |  |  |  |
| $R_{b}$ | 0.2163 | 31. | 0.4 | 7. \# | 3.5 | $10{ }^{+}$ |
| $A_{b}$ | 0.935 | 214. | 1. | 5. * | 5.7 | 3 * |
| $g_{L_{1}}^{b}$ | -0.999 | 54. | 0.31 | 4.0 | 2.2 | 7.6 |
| $g_{R}^{b}$ | 0.184 | 1540 | 7.2 | 36. | 41. | 23. |
| $Z-c$ couplings |  |  |  |  |  |  |
| $R_{c}$ | 0.1721 | 174. | 2. | $30^{\#,+}$ | 5.8 | 50 + |
| $A_{c}$ | 0.668 | 404. | 3. | 6. * | 21. | 3 * |
| $g_{L}^{c}$ | 0.816 | 119. | 1.2 | 15. | 5.1 | 26. |
| $g_{R}^{c}$ | -0.367 | 415. | 3.1 | 16. | 21. | 26. |

LCC Physics WG

## take-home message

if we work hard and be smart

ILC250 $=$ ILC250 $+100 \times$ LEP/SLC

## appendix

## impact of EWPOs



## experimental inputs and assumptions for section 1

|  | efficiency | systematics | $\cos \theta$ | bin width |
| :---: | :---: | :---: | :---: | :---: |
| $\mu \mu$ | $98 \%{ }^{[1]}$ | $0.1 \%$ | $[-0.95,0.95]$ | 0.1 |
| T T | $90 \%{ }^{[2]}$ | $0.2 \%$ | $[-0.95,0.95]$ | 0.1 |
| b b | $29 \%\left[^{[3]}\right.$ | $0.2 \%$ | $[-0.95,0.95]$ | 0.1 |
| c c | $7 \%{ }^{[3]}$ | $0.5 \%$ | $[-0.95,0.95]$ | 0.1 |
| e e | $97 \%{ }^{[1]}$ | $0.1 \%$ | $[-0.95,0.95]$ | 0.1 |

flat acceptance function within the fiducial volume
b- and c-charge identification efficiency applied
analytic differential cross section, w/o effects from beamstrahlung and ISR
[1] T.Suehara; [2] D.Jeans; [3] R.Poeschl

## computation for W and Y oblique parameters

(provided by M. Peskin \& M. Perelstein)

$$
\begin{aligned}
& \frac{d \sigma}{d \cos \theta}\left(e_{L}^{-} e_{R}^{+} \rightarrow f_{L} \bar{f}_{R}\right)=\frac{\pi \alpha^{2}}{2 s}\left|F_{L L}\right|^{2}(1+\cos \theta)^{2} \\
& \frac{d \sigma}{d \cos \theta}\left(e_{L}^{-} e_{R}^{+} \rightarrow f_{R} \bar{f}_{L}\right)=\frac{\pi \alpha^{2}}{2 s}\left|F_{L R}\right|^{2}(1-\cos \theta)^{2} \\
& \frac{d \sigma}{d \cos \theta}\left(e_{R}^{-} e_{L}^{+} \rightarrow f_{L} \bar{f}_{R}\right)=\frac{\pi \alpha^{2}}{2 s}\left|F_{R L}\right|^{2}(1-\cos \theta)^{2} \\
& \frac{d \sigma}{d \cos \theta}\left(e_{R}^{-} e_{L}^{+} \rightarrow f_{R} \bar{f}_{L}\right)=\frac{\pi \alpha^{2}}{2 s}\left|F_{R R}\right|^{2}(1+\cos \theta)^{2} \\
& F_{L L}(s)=-\left[Q_{f}+\frac{\left(\frac{1}{2}-s_{w}^{2}\right)\left(I_{f}^{3}-s_{w}^{2} Q_{f}\right)}{s_{w}^{2} c_{w}^{2}} \frac{s}{\left(s-m_{Z}^{2}\right)}-\frac{\frac{1}{2} I_{f}^{3}}{s_{w}^{2}} \frac{s}{m_{W}^{2}} \mathbf{W}-\frac{\frac{1}{2} Y_{f}}{c_{w}^{2}} \frac{s}{m_{W}^{2}} \mathbf{Y}\right] \\
& F_{L R}(s)=-\left[Q_{f}+\frac{\left(\frac{1}{2}-s_{w}^{2}\right)\left(-s_{w}^{2} Q_{f}\right)}{s_{w}^{2} c_{w}^{2}} \frac{s}{\left(s-m_{Z}^{2}\right)}-\frac{\frac{1}{2} Q_{f}}{c_{w}^{2}} \frac{s}{m_{W}^{2}} \mathbf{Y}\right] \\
& F_{R L}(s)=-\left[Q_{f}+\frac{\left(-s_{w}^{2}\right)\left(I_{f}^{3}-s_{w}^{2} Q_{f}\right)}{s_{w}^{2} c_{w}^{2}} \frac{s}{\left(s-m_{Z}^{2}\right)}-\frac{Y_{f}}{c_{w}^{2}} \frac{s}{m_{W}^{2}} \mathbf{Y}\right] \\
& F_{R R}(s)=-\left[Q_{f}+\frac{\left(-s_{w}^{2}\right)\left(-s_{w}^{2} Q_{f}\right)}{s_{w}^{2} c_{w}^{2}} \frac{s}{\left(s-m_{Z}^{2}\right)}-\frac{Q_{f}}{c_{w}^{2}} \frac{s}{m_{W}^{2}} \mathbf{Y}\right]
\end{aligned}
$$

## composite models

|  | $\eta_{L L}$ | $\eta_{R R}$ | $\eta_{L R}$ | $\eta_{R L}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\Lambda_{L L}^{+}$ | 1 | 0 | 0 | 0 |  |
| $\Lambda_{L L}^{-}$ | -1 | 0 | 0 | 0 | $\mathcal{L}_{L L}=\frac{g_{\text {contact }}^{2}}{2 \Lambda^{2}} \sum_{i, j} \eta_{L L}^{i j}\left(\bar{\psi}_{L}^{i} \gamma_{\mu} \psi_{L}^{i}\right)\left(\bar{\psi}_{L}^{j} \gamma^{\mu} \psi_{L}^{j}\right),$ |
| $\Lambda_{R R}^{+}$ | 0 | 1 | 0 | 0 | $\mathcal{L}_{R R}=\frac{g_{\text {contact }}^{2}}{2 \Lambda^{2}} \sum \eta_{R R}^{i j}\left(\bar{\psi}_{R}^{i} \gamma_{\mu} \psi_{R}^{i}\right)\left(\bar{\psi}_{R}^{j} \gamma^{\mu} \psi_{R}^{j}\right)$, |
| $\Lambda_{R R}^{-}$ | 0 | -1 | 0 | 0 |  |
| $\Lambda_{V V}^{+}$ | 1 | 1 | 1 | 1 | $R=\frac{2 \Lambda^{2}}{2 \Lambda_{i, j}} \sum_{L R}\left(\psi_{L}^{\ell} \gamma_{\mu} \psi_{L}^{t}\right)\left(\psi_{R}^{\prime} \gamma^{\mu} \psi_{R}^{j}\right),$ |
| $\Lambda_{V V}^{-}$ | -1 | -1 | -1 | -1 | $\mathcal{L}_{R L}=\frac{g_{\text {contact }}^{2}}{2 \Lambda^{2}} \sum_{i, j} \eta_{R L}^{i j}\left(\bar{\psi}_{R}^{i} \gamma_{\mu} \psi_{R}^{i}\right)\left(\bar{\psi}_{L}^{j} \gamma^{\mu} \psi_{L}^{j}\right),$ |
| $\Lambda_{A A}^{+}$ | 1 | 1 | -1 | -1 |  |
| $\Lambda_{A A}^{-}$ | -1 | -1 | 1 | 1 | Tanabashi, et al, |
| $\Lambda_{(V-A)}^{+}$ | 0 | 0 | 1 | 1 | PRD 98 (2018) 030001 |
| $\Lambda_{(V-A)}^{-}$ | 0 | 0 | -1 | -1 |  |

