Bottom and top reconstruction as benchmarks for the ILD Detector

M.S. Amjad[a], A. Irles[b], V. Lohezic[c], Y. Okugawa[d], R. Pöschl[b], F. Richard[b], H. Yamamoto[d], R. Yonamine[d]

[a] *National University of Technology, Islamabad, Pakistan*
[b] *Laboratoire de l'Accélerateur Linéaire, CNRS/IN2P3 et Université de Paris-Sud, Centre Scientifique d'Orsay Bâtiment 200, BP 34, 91898 Orsay CEDEX, France*
[c] *Université Paris-Sud, 15 Rue Georges Clemenceau, 91400 Orsay, France*
[d] *Tohoku University, 6-3 Aoba Aramaki-aza Aoba-ku, Sendai 980-8578, Japan*

Abstract

A comprehensive study of bottom-quark pair and top-quark pair production using the semi-leptonic decay channel at $\sqrt{s} = 500\,\text{GeV}$ is presented that benchmark the performance of the current so-called *large* and *small* models of the ILD Detector. The event reconstruction exploits distinguished features of the detector such as lepton identification, vertex charge reconstruction and particle ID with the central TPC of the ILD Detector. With these techniques the final state leptons, the b-quark and $\bar{b}$-quark and the W can be unambiguously reconstructed. Both detector models perform similarly well with a slight advantage of the large detector model. In case of top quark pair production the selection efficiency is between 30% and 60% for semi-leptonic events. For this channel the note presents an update of the perspective of the physics potential w.r.t. earlier studies for an integrated luminosity of $3200\,\text{fb}^{-1}$. The results for $e^+e^- \to b\bar{b}$ demonstrate that also this channel can accurately measured at $\sqrt{s} = 500\,\text{GeV}$.

1. Introduction

Heavy quarks may be messengers of new physics of primary importance [1]. Their large mass compared with other fermions can be explained in Randall Sundrum models [2, 3] featuring warped extra dimensions that are dual to model, which assume that the heavy quarks are composite objects [4]. High precision e^+e^- collisions with polarised beams around the TeV scale are ideally suited to detect new physics effects [5, 6]. Precise measurements of the electroweak couplings of third-generation quarks require superb detector performance in terms of flavour tagging including the event by event determination of the charge of the final state jets to avoid for example migrations in polar angle spectra and/or to reconstruct events in which the heavy quark charge is the only handle to distinguish between particles and anit-particles. The charge determination happens mainly by a combination of the determination of the summed charge of tracks pointing to a secondary vertex or by the identification of the charge of a final state Kaon. This is turn requires a successful particle identification by the detector. Therefore processes with heavy quark final states, i.e. $e^+e^-\text{ø}b\bar{b}$ and $e^+e^- \to t\bar{t}$ are highly relevant for the benchmarking of the detector performance. In short one can test the following detector capacities.

- Track finding efficiency
- Stringent test of (secondary) vertexing
- Particle ID

In case of $e^+e^- \to t\bar{t}$ leptonic and semi-leptonic decays of the t-quark pair provide an important additional handle for the accurate measurement of the final state. The analysis presented in this

note focuses on the semi-leptonic decay mode of the top-quark pair. The analysis of $t\bar{t}$ production and $b\bar{b}$ production share a number of commonalities. Therefore these two analyses are joined in this note. The analyses presented in this note start out from the PhD thesis of Sviatoslav Bilokin that are based on the DBD samples and software versions [7]. This work has in part been published in Ref [6]. The analyses are ported to the large, IDR-L, and small, IDR-S, detector models, respectively, of the ILD detector [8] for the International Linear Collider [9]. The research program includes high-statistics running at $\sqrt{s} = 250\,\text{GeV}$ and $\sqrt{s} = 500\,\text{GeV}$. For further details of the operation scenarios see Ref. [10]. For the process $e^+e^- \to b\bar{b}$ an analysis at $\sqrt{s} = 500\,\text{GeV}$ is presented instead of $\sqrt{s} = 250\,\text{GeV}$ as in Ref. [6]. The results here benefit from a refined analysis strategy for the ILD paper that is under review in ILD.

2. Methods, tools and Monte Carlo samples

For the event reconstruction we use the `ILCSoft` version `v02-00-02`. The software allows for a full detector simulation and event reconstruction including digitisation steps. For the analyses presented we use the versions `ILD_s5_o1_v02` and `ILD_l5_o1_v02` of the detector model.

For the event reconstruction we use the following methods

- 'Core tools'
  - In case of $e^+e^- \to t\bar{t}$ we use the standard Durham algorithm for jet clustering and the Valencia algorithm for $\gamma\gamma$ background removal.
  - In case of $e^+e^- \to b\bar{b}$ we use the Valencia jet algorithm implemented in the `LCFIPlus` package [11] as available in `ILCSoft`. We apply the option in which the algorithm does not break secondary vertices. In this algorithm the distance between two objects is calculated as
$$d_{ij} = 2\min(E_i^{2\beta}, E_j^{2\beta})(1-\cos\theta_{ij})/R^2 \tag{1}$$
The distance of a particle i to the beam is calculated according to.
$$d_{iB} = E^{2\beta}\sin^{2\gamma}\theta_{iB} \tag{2}$$
The jet algorithm is run with the following settings: $\alpha = \beta = \gamma = 1$, $R = 1.4$
  - We use the standard LeptonFinder of ILD to identify isolated electrons and muons in semi-leptonic $t\bar{t}$ events.
- Tools developed for the study
  - We identify reconstructed tracks that have not been associated to secondary vertices from B-Meson decays but belongs to this decay according to the Monte Carlo Truth information. It then recovers the 'lost' tracks by means of the impact parameters d_0 (transversal) and z_0 (longitundinal). In this present note the recovery uses only the impact parameter d_0 since the algorithms needs to be adapted for the vertex smearing present in the simulation for the IDR.
  - We identify Kaons by means of the dE/dx measured in the TPC of ILD. It selects a strip in the dE/dx-momentum plane with a high kaon concentration. The efficiency and the purity of the Kaon selection vary as a function of the width of this strip.
  - We finally calculate the jet charge and the polar angle of the bottom and top quark pair, respectively.
- The following method combines the results of the two (nearly independent) charge measurements on the b and $\bar{b}$ jet into a robust charge determination.

$$\left.\begin{aligned} N_{acc} &= Np^2 + Nq^2 \\ N_{rej} &= 2Npq \\ 1 &= p + q \end{aligned}\right\} \quad N_{corr} = N_{acc} \cdot \frac{p^2}{p^2+q^2} \tag{3}$$

where N is total number of events, N_{acc} and N_{rej} are the number of events that were accepted and rejected, respectively. The p and q values represent the probabilities for a correct or an incorrect reconstruction of the b-quark charge, respectively. Solving this equation allows to correct for migrations caused by imperfect reconstruction or $B^0 - \bar{B}^0$ oscillations in the resulting polar angle spectrum. The correction has been applied to the $b\bar{b}$ studies but not (yet) for $t\bar{t}$. For the latter the selection scheme in $t\bar{t}$ is more involved rendering its application less straight forward since e.g. the b-quarks are not necessarily back-to-back.

2.1. Monte Carlo samples

Samples generated with WHIZARD 1.95. Top quark pair production is the dominant process in the $e^+e^- \to b\bar{b}\ell\nu q\bar{q}$ sample, but it contains also single t and WWZ. In case of the t quark study we have analysed samples for the beam polarisations $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = \pm 1, \mp 1$. In total we analyse about $2200\,\mathrm{fb}^{-1}$ for each of the beam polarisations but we discard final states for which $W \to \tau\nu$. In case of the b quark study we have studied only the configuration $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = \mp 1$ since this study is just considered as a supporting study for the t quark analysis and the chosen polarisation configuration is the more challenging for the detector performance due to larger migration effects in the polar angle reconstruction. $e^+e^- \to b\bar{b}$: The $b\bar{b}$ events are extracted from a sample of simulated $q\bar{q}$. The total integrated luminosity is about $46\,\mathrm{fb}^{-1}$. The genuine $e^+e^- \to b\bar{b}$ at 500 GeV is about 4% of the total cross section.

Note that the analysis of the $e^+_R e^-_L \to t\bar{t}$ came quite late in the study since it is generally considered that the $e^+_L e^-_R \to t\bar{t}$ is the more challenging channel. Therefore, the corresponding control plots in Secs. 3 and 4.2 are only given for the latter in this note.

3. Efficiencies and Control plots

Figure 1 demonstrates that the studies presented in this note test the detector performance for very different momenta of the final state b quark.

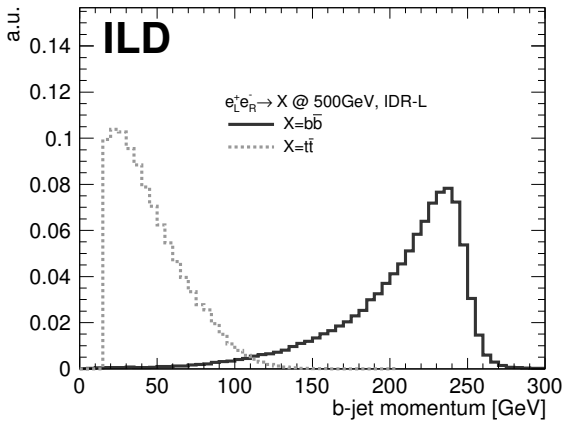

Figure 1: *Momentum of the b-jet with cheated identification for $e^+e^- \to b\bar{b}$ and $e^+e^- \to t\bar{t}$ processes.*

The Figs. 2 and 3 show the missed tracks before and after vertex recovery for the $e^+e^- \to b\bar{b}$ and $e^+e^- \to t\bar{t}$ analyses, respectively. Both figures suggest a systematic improvement in the assignment of secondary vertices.

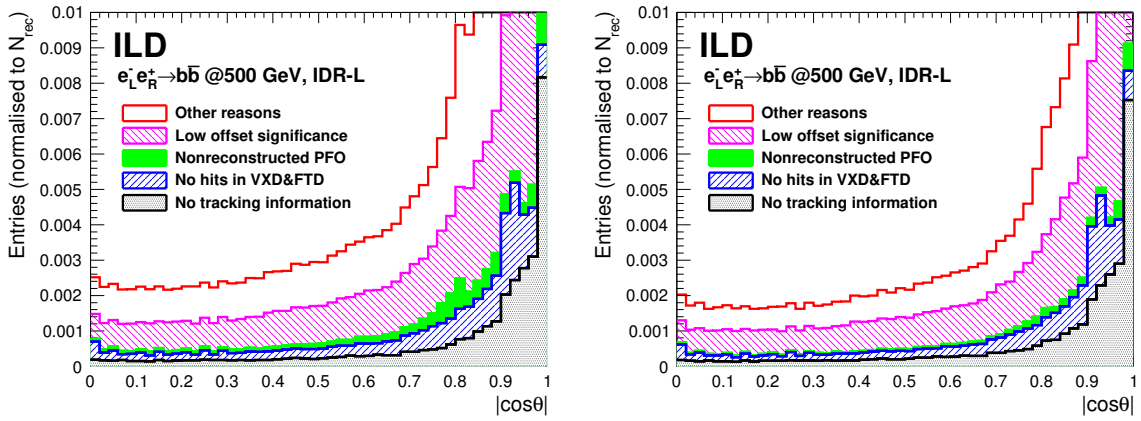

Figure 2: *Polar angle of missed tracks before (left) and after (right) vertex recovery in case of the $e^+e^- \to b\bar{b}$ process.*

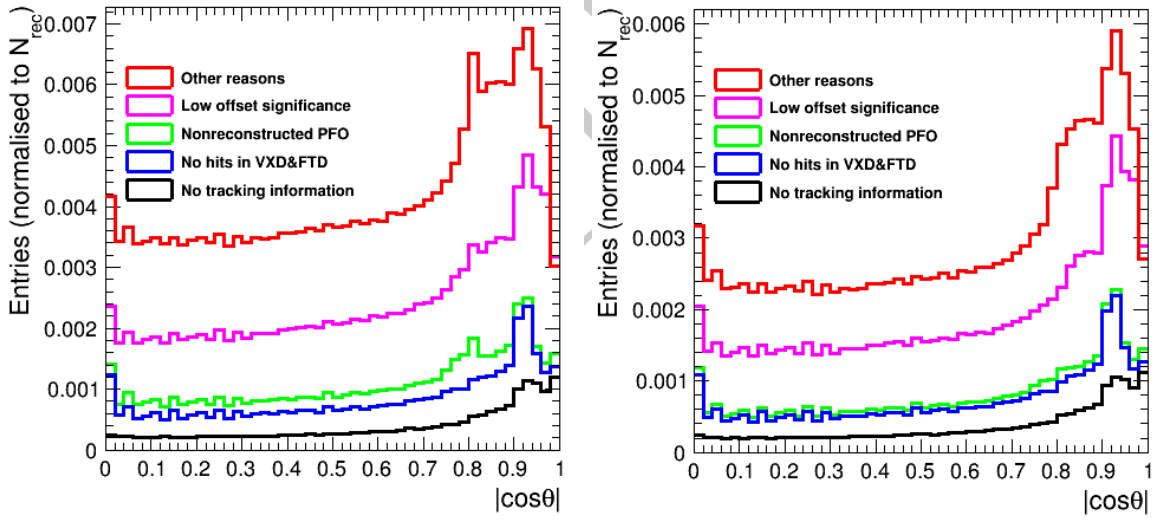

Figure 3: *Polar angle of missed tracks before (left) and after (right) vertex recovery in case of the $e^+e^- \to t\bar{t}$ process.*

This improvement is quantified in Figs. 4 and 5 where the purity of the b-charge reconstruction is shown as a function of the $b-tag$ value, the reconstructed b-momentum $|p_{had}|$ the number of reconstructed tracks assigned to a secondary vertex N_{rec} and finally the polar angle of the b-hadron. here denoted as $|\cos\theta|$. The b-charge purity is defined as

$$p_b = N_{correct}/N_{jet,tot.} \tag{4}$$

with $N_{correct}$ being the number of b-jets with correctly reconstructed b quark charge. This value normalised to the total number of b-jets $N_{b-jet,tot.}$ for which a charge assignment according to e.g. Table 2 can be made.

The improvement is is larger for the process $e^+e^- \to t\bar{t}$ than for $e^+e^- \to b\bar{b}$. Qualitatively this is expected since as a consequence of the different b-jet momenta, see Fig. 1, also the tracks produced in the decay of the b-hadron are softer in case of top-pair production. In case of $e^+e^- \to t\bar{t}$ the improvement is 10% over a large range in $|\cos\theta|$ and mainly driven by three to five prong decays. Both results will further improve once the vertex recovery takes also the the impact parameter z_0 into account. All results shown so far in this section have been obtained for the large detector model. The conclusions for the small detector model are similar.

The lower right panels of Figs. 4 and 5 show a drop in purity for large values of $|\cos\theta|$. This is compatible with the drop in acceptance that is shown in Fig. 6 for the case $e^+e^- \to b\bar{b}$ as a function of the polar angle of the reconstructed b-jet $|\cos\theta_b|$. Within statistical errors the results are the same for the large and the small detector model. However, towards large values of $|\cos\theta_b|$ the large detector performs systematically better than the small detector.

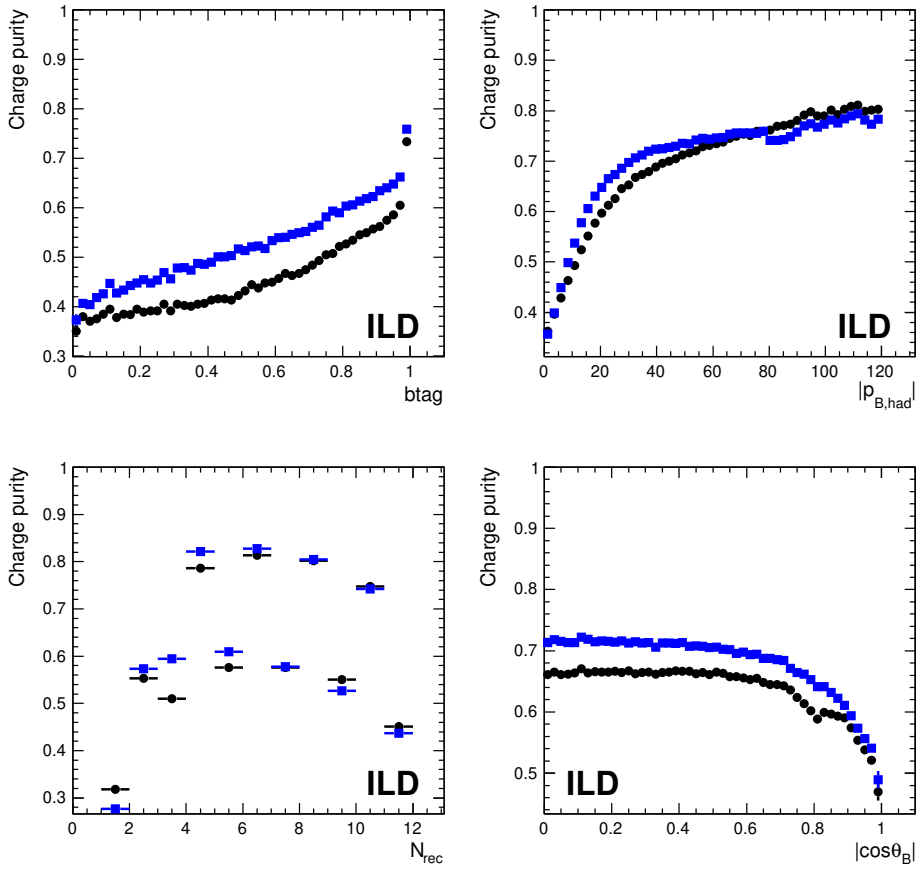

Figure 4: *Purity before and after vertex recovery in case of the $e^+e^- \to t\bar{t}$ process for different observables.*

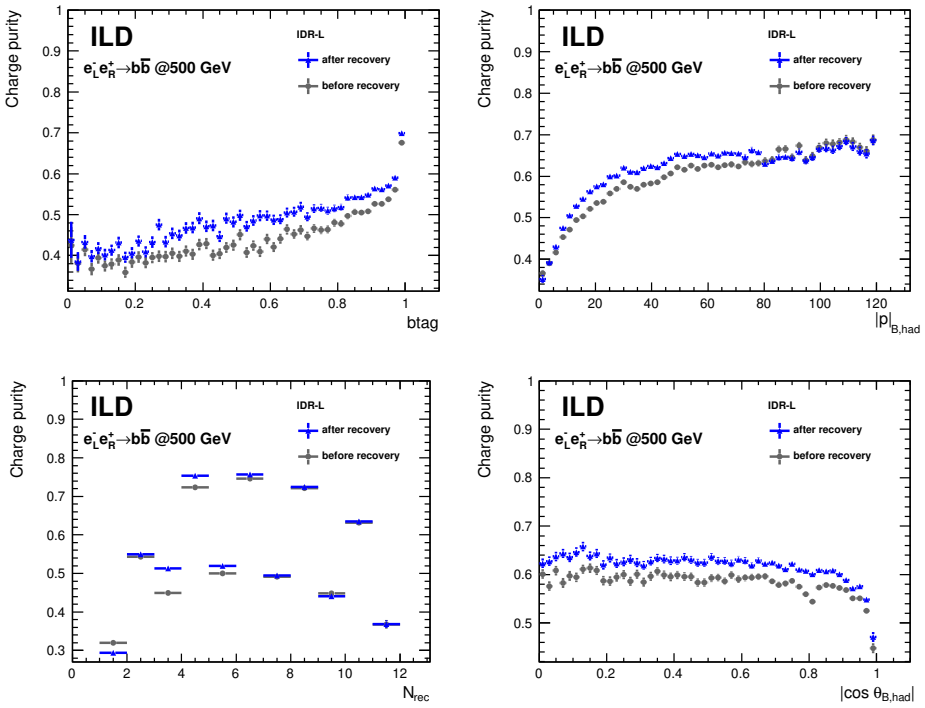

Figure 5: *Purity before and after vertex recovery in case of the $e^+e^- \to t\bar{t}$ process for different observables.*

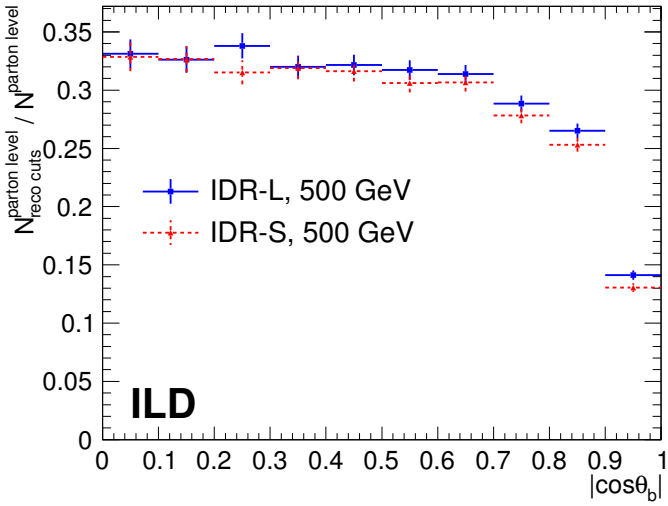

Figure 6: Detector acceptance distribution for b-tagged jets. Note the figure shows the acceptance after the full selection given in Tab. 1.

A component that distinguishes the ILD Detector from other proposals for e^+e^- colliders is the TPC as the central tracking system. Beside the precise momentum measurement the dE/dx measurement in the gaseous medium allows for a particle identification. Since around 87% of B-Mesons (neutral or charged) contain a charged Kaon among their decay products the particle ID can support greatly the

charge determination of the b quark.

The left part of Fig. 7 shows the dE/dx values obtained in simulation for different particle types as a function of the particle momentum. The lines indicate a strip with an accumulation of signals produced by Kaons. A minimum momentum of 2 GeV is required for the selection of Kaons. The right part of Fig. 7 shows the variation of the purity as as a function of the Kaon selection efficiency that corresponds to a variation of the width of the strip in the previous figure. A closer look into the separation power in different momentum ranges is given in App. A.

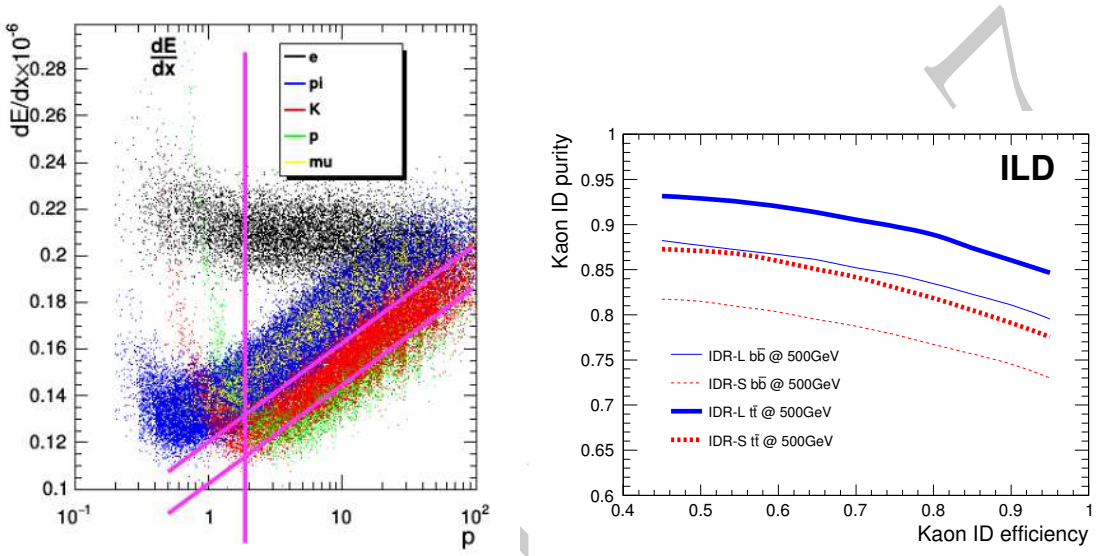

Figure 7: Left: Simulated dE/dx spectrum for different particle types. The lines indicate the cuts for the Kaon selection. Right:Purity of the Kaon selection as a function of the Kaon selection efficiency.

4. Event selection

In the following the event selection of the two final states under study will be presented. The different complexities of the final states require different set of cuts. The t quark is composed from its decay products, the b quark and the W boson and the final state le

4.1. Analysis details specific to the $e_L^- e_R^+ \to b\bar{b}$ analysis

Table 1 shows the selection efficiencies for the $e_L^- e_R^+ \to b\bar{b}$ analysis. In this case events that are subject to the radiative return to the Z, implying an energetic final state photon, have to be suppressed. Therefore cuts on the sum of the masses of the two jets and a cut on the photon energy are introduced. The overall efficiency after selection of events with consistent b quark charge is with around 28% to 29% similar for both detector models. For the b-charge measurement opposite charges in opposite jets are required. The charges are either derived from the tracks pointing to the secondary vertex or from the Kaon charge or from a combination of both. An event is selected if there is one combination with a consistent result. The efficiencies for the different methods are given in Tab. 2. The purity of the different methods is shown in Fig. 8. In both cases there is no large difference between the two detector models although the large detector seems to perform slightly better for the double Kaon method. This suggests that the smaller outer radius of the TPC puts a, however minor, on the dE/dx measurement.

$e_L^- e_R^+ \to b\bar{b}$ at 500 GeV

	IDR-L			IDR-S		
	Signal	$B_{q\bar{q}}/S$	$B_{rad.Z}/S$	Signal	$B_{q\bar{q}}/S$	$B_{rad.Z}/S$
Full sample	100.0%	1800.5%	359.1%	100.0%	1800.6%	359.0%
$b_{tag}(jet_1) > 0.9$ and $b_{tag}(jet_2) > 0.2$	70.2%	2.3%	147.7%	69.9%	2.3%	149.0%
$m_{jet_1+jet_2} > 200GeV$	68.2%	1.4%	6.7%	67.8%	1.2%	6.7%
$E_{photon} < 100GeV$	64.8%	1.3%	1.7%	64.3%	1.2%	1.6%
double jet-charge measurement	28.9%	1.0%	1.0%	27.9%	0.9%	1.0%

Table 1: Selection efficiency and B/S rejection for some bkg sources

$e_L^- e_R^+ \to b\bar{b}$ at 500 GeV

	IDR-L	IDR-S
Vtx+Vtx	12.9%	12.8%
K+K	4.4%	4.0%
Vtx+K (diff. jets)	3.9%	3.7%
Vtx+K (same jet)	7.7%	7.4%

Table 2: Final selection efficiency, after double jet-charge measurement

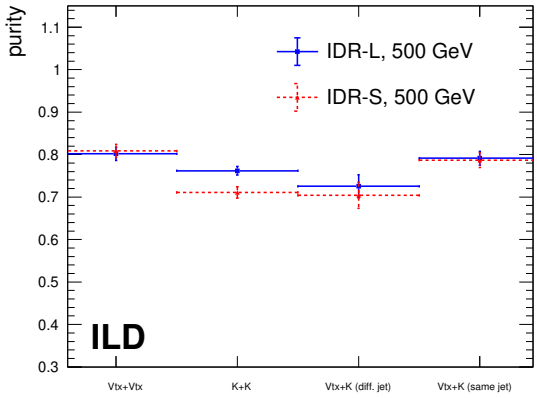

Figure 8: Purity of the methods listed in Tab. 2 used for the reconstruction of the vertex charge in the $e^-_L e^+_R \to b\bar{b}$.

4.2. Analysis details specific to the $e^+e^- \to t\bar{t}$ analysis

The t quark is composed from its decay products, the b quark and the W boson and the charge of the lepton is a measure to distinguish the t from the $\bar{t}$ quark.

Figure 9 shows in the left panel the energy distribution of the isolated lepton in the laboratory frame. The distribution features a maximum at around 30 GeV and a tail towards higher energies. The right panel shows the polar angle spectrum of the isolated lepton. The distribution decreases slightly with a sharp drop at the acceptance limit of the detector. The distribution reveals also acceptance drops at $\cos\theta_\ell = 0$ and $\cos\theta_\ell = 0.8$ that correspond to the position of the TPC anode plate and the barrel-endcap transition region, respectively.

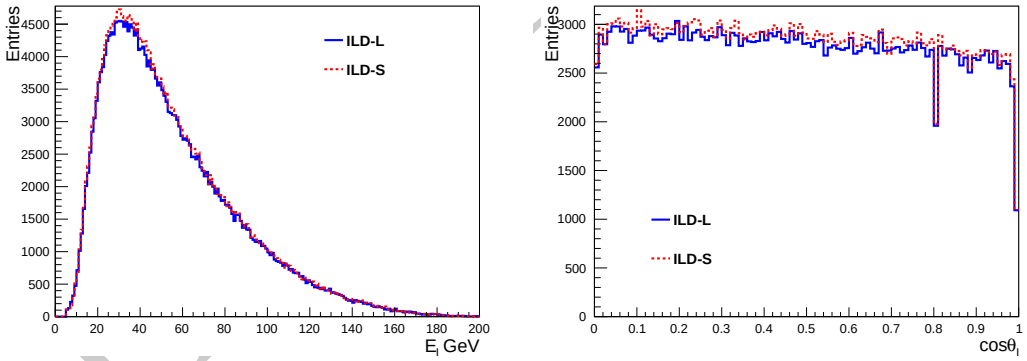

Figure 9: *Left: Energy of the isolated lepton in $e^+_L + e^-_R \to t\bar{t}$. Right: Polar angle distribution of the isolated lepton in $e^+_L e^-_R \to t\bar{t}$.*

For completeness Fig, 10 shows the mass distribution of the hadronic W and the hadronic t quark.

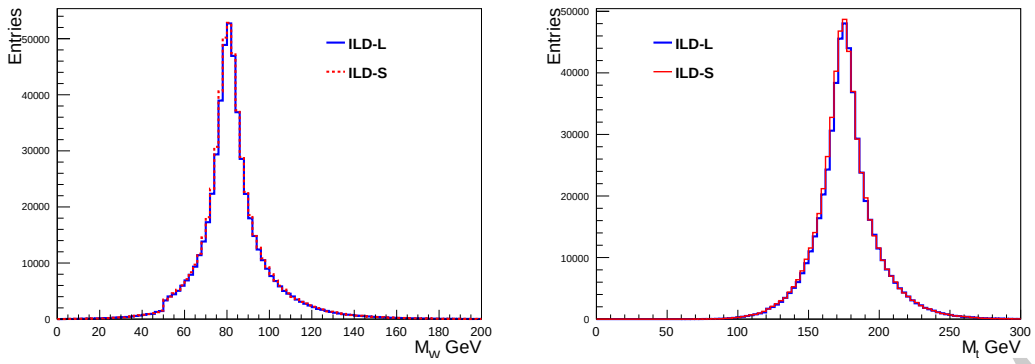

Figure 10: *Left:* Mass distribution of the hadronic W in $e_L^+ + e_R^- \to t\bar{t}$. *Right:* Mass distribution of the hadronic t quark in $e_L^+ e_R^- \to t\bar{t}$.

Tables 3 and 5 give the efficiencies after each cut applied for the selection of $t\bar{t}$ events. The first part lists *General selection cuts* that were already used in Ref. [5].

The polar angle of the t quark $\cos\theta_t$ is reconstructed from the hadronically decaying t quark. For the polar angle spectrum the charge of the t quark has to be determined and the b quark and the W-boson have to be correctly associated. This is more involved in the $e_L^- e_R^+ \to t\bar{t}$ case than in the $e_R^- e_L^+ \to t\bar{t}$ case due to the different kinematics provoked by the $V-A$ interaction of the t quark decay. The different steps for an accurate reconstruction of the polar angle spectrum are listed in the following[1].

- In a first step further cuts on the sum of the Lorentz factor of the two tops of $\gamma_t^{had} + \gamma_t^{\ell} > 2.4$ is applied. Here γ_t^{had} is the Lorentz factor of the hadronically decaying t quark and γ_t^{ℓ} the Lorentz factor of the leptonically decaying t quark. In case of $e_L^- e_R^+ \to t\bar{t}$ a cut on the B-hadron momentum of $p_{B,had.} > 15\,\text{GeV}$ is applied in addition.

- The semi-leptonic decay of the t quark gives powerful information for the event reconstruction giving rise to the variable L_{cut}, which means the charge of the isolated lepton plus a cut on the event quality of $\chi 2 < 15$ that is motivated in Ref. [5].

- For the b quark charge determination the used methods are very much similar to those in 2. The vertex charge is supported by the requirements of $btag > 0.8$ and a minimal hadron momentum of 25 GeV as motivated in Ref. [6][2].

- The various methods of measuring the b quark charge are also combined with the charge of the isolated lepton L. In this case an additional cut on $\gamma_t^{had} > 1.23$ is applied.

- The final decision on the t quark charge is obtained from the sum of the charges associated to the different methods. If the sum is smaller (greater) than zero then the hadronically decaying t quark candidate is said to be the t quark ($\bar{t}$ quark).

Table 3 gives the final selection efficiency for the case $e_L^- e_R^+ \to t\bar{t}$ after the inclusion of the respective cuts. Table 4 shows the *accumulated* efficiencies for $t\bar{t}$ identification after progressive application of the various methods described in the list of items above. The addition of methods other than L_{cut} increases the efficiency by around 38%. Figure 11 shows for completeness the purity of the selection

[1] We are aware that the set of cuts does not look straight forward and needs revision in the Post-IDR phase.

[2] In the $e^+e^- \to b\bar{b}$ analysis this additional requirement was removed. It would have to be investigated whether this requirement can be omitted in case of $e_L^- e_R^+ \to t\bar{t}$, too

$e_L^- e_R^+ \to t\bar{t}$ at 500 GeV

General selection cuts	IDR-L	IDR-S
Isolated Lepton	92.1%	92.1%
$btag_1 > 0.8$ or $btag_2 > 0.3$	81.2%	81.1%
Thrust < 0.9	81.2%	81.1%
Hadronic mass	78.2%	78.2%
Reconstructed m_W and m_t	73.4%	73.4%
t quark polar angle spectrum		
$\gamma_t^{had.} + \gamma_t^{\ell} > 2.4$	62.2%	61.8%
$\lvert p_{B,had} \rvert > 15\,\text{GeV}$	34.5%	33.9%
"$t\bar{t}$ identification"	30.6%	30.2%
b quark polar angle spectrum		
No additional cuts		

Table 3: Event selection efficiencies after pre-selection and reconstruction of the polar angle spectrum of the t quark and that of the underlying b quark for the process $e_L^- e_R^+ \to t\bar{t}$.

Methods	IDR-L	IDR-S
1 L_{cut}	22.1%	21.9%
2 L+Vtx	28.6%	28.4%
3 L+K	29.6%	29.3%
4 Vtx+Vtx	30.1%	29.7%
5 K+K	30.3%	29.9%
6 Vtx+K (same jet)	30.5%	30.1%
7 Vtx+K (different jet)	30.6%	30.2%

Table 4: Efficiency increase after the progressive application of the various methods introduced to ensure a correct association of the W boson and the b quark in case of $e_L^- e_R^+ \to t\bar{t}$. The final efficiency in the last line of this table corresponds to the efficiency given in the line "$t\bar{t}$ identification" of Tab. 3.

for those cases in which the information from the t quark and the $\bar{t}$ quark decay have been combined. As in case of $e^+e^- \to b\bar{b}$ differences between the large and the small detector are observed for those combinations that include Kaons with the biggest difference for the pure Kaon combination $K + K$. The Kaon measurement is the domain of the TPC and the two models feature different outer TPC radii.

As a supplementary remark please note that methods available also for the fully hadronic final state yield efficiencies of 14.4% when using Methods 4-7 of Tab. 4 and 6.9% when using only Method 4, respectively.

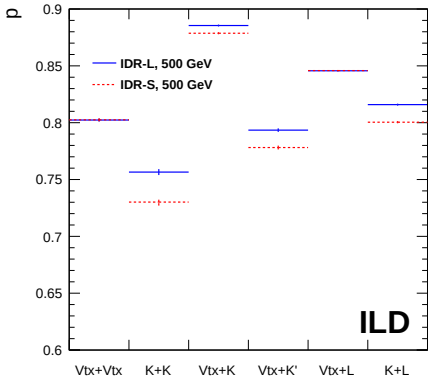

Figure 11: Purities of the various combinations to distinguish between the t and the $\bar{t}$ quark.

$e_R^- e_L^+ \to t\bar{t}$ at 500 GeV

General selection cuts	IDR-L	IDR-S
Isolated Lepton	94.1%	94.0%
$btag_1 > 0.8$ or $btag_2 > 0.3$	84.9%	84.8%
Thrust < 0.9	84.9%	84.8%
Hadronic mass	82.2%	82.3%
Reconstructed m_W and m_t	77.6%	77.5%
t quark polar angle spectrum		
$\gamma_t^{had.} + \gamma_t^{\ell} > 2.4$	64.1%	64.1%
b quark polar angle spectrum		
Vtx+Vtx	10.8%	10.3%

Table 5: Event selection efficiencies after pre-selection and reconstruction of the polar angle spectrum of the t quark and that of the underlying b for the process $e_R^- e_L^+ \to t\bar{t}$

Table 5 shows the selection efficiencies for the case $e_R^- e_L^+ \to t\bar{t}$. The cut scenario for the reconstruction of the polar angle of the t quark is much simpler than in case of $e_L^- e_R^+ \to t\bar{t}$. The reason is that the t quark direction is in first approach given by the W boson such that a wrong association of W boson and b quark doesn't alter the t quark direction. For the polar angle spectrum of the underlying b quark the analysis is restricted to the the combination of the vertex charge. This is further discussed in Sec. 5.

5. Results

Figure 12 shows the spectrum of the polar angle $\cos\theta_b$ after the selection given in Tabs. 1 and 2 and the application of Eq. 3 for $e^-_L e^+_R \to b\bar{b}$. Large and small detector agree within statistical uncertainties. It seems however that there is larger migration for the small detector.

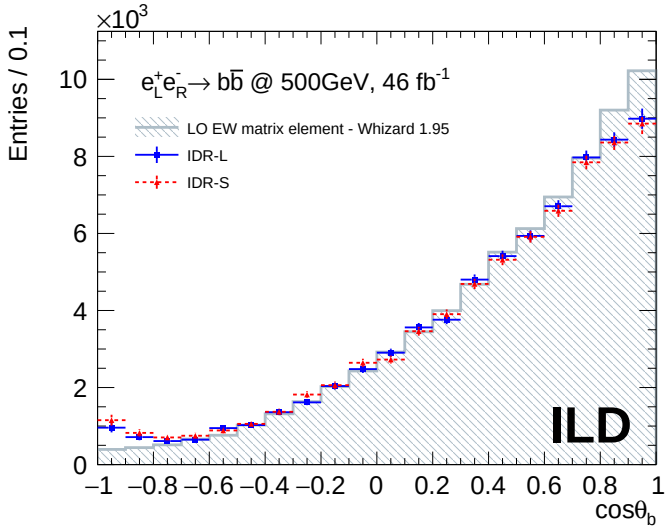

Figure 12: Polar angle spectrum for $e^-_L e^+_R \to b\bar{b}$ at $\sqrt{s} = 500\,\text{GeV}$ after correction for acceptance and migrations due to charge mismeasurement, see Eq. 3, for the two ILD Detector models.

The left part of Fig. 13 shows the polar angle distribution of $t\bar{t}$ of the generated and reconstructed data for the large and the small detector models. For this all cuts and methods given in Tabs. 3 and 4 have been applied. Overlaid and here and in the following figures, mainly to guide the eye, is a polynomial of second degree motivated by the approximate $d\sigma/d\cos\theta_t \sim S\cos^2\theta_t + A\cos\theta_t$ dependence of the differential cross section. The right part shows the polar angle distribution of the underlying b quark for the same set of cuts. The polar spectrum can be accurately reconstructed over the entire polar angle. Acceptance drops at large absolute values of the polar angle become visible in the polar angle spectrum of the b quark (In this case no attempt was made to correct for acceptance as in Fig. 12). However, in the range $-0.8 < \cos\theta_b < 0.8$ also the polar angle of the b quark can be accurately reconstructed.

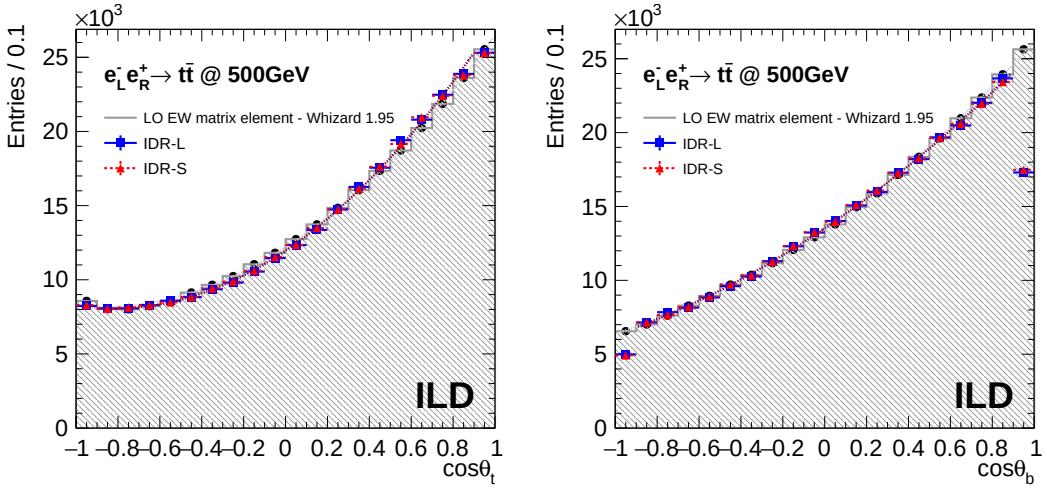

Figure 13: *Results for $e^-_L e^+_R \to t\bar{t}$ at $\sqrt{s} = 500\,GeV$ for the two ILD Detector models. Left: Polar angle distribution for t quark. Right: Polar angle distribution for the b quark that is issue of the t quark decay. The distributions for IDR-S is normalised to the one for IDR-L so that both histograms will be on the same level. For details on the selection, see text and Tabs. 3 and 4.*

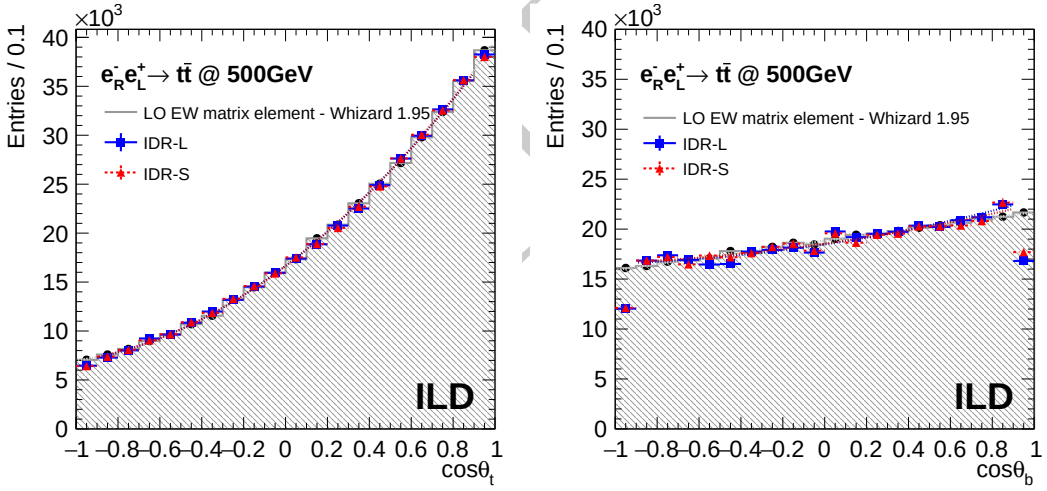

Figure 14: *Same as Fig. 13 for $e^-_R e^+_L \to t\bar{t}$. For details on the selection, see text and Tab. 5.*

The left part of Fig. 14 shows the polar angle spectrum for the case $e^-_R e^+_L \to t\bar{t}$ after application of the cuts introduced in Tab. 5. The generated spectrum can be very precisely reconstructed. The right hand part of Fig. 14 shows the polar angle spectrum of the underlying b quark. Here only events with consistent vertex-charge measurements have been included. The polar angle of the b quark can also in this case be very well reconstructed. However, the efficiency drops to 10% as already quantified in the lower part of Tab. 5. The inclusion of the other methods is subject to further studies. Preliminary results show that taking into account the isolated lepton "swamps" the polar angle spectrum with events in which the b quark direction is not constrained. One reason is certainly that in the case of $e^-_R e^+_L \to t\bar{t}$ the b quark is on average softer than in the case of $e^-_L e^+_R \to t\bar{t}$.

Comparing the spectra of the underlying b quarks demonstrates more clearly than the actual t quark polar angle spectra the different polarisations of the t quarks projected out by the flight direction of

the b quark. In case of $e^-_L e^+_R \to t\bar{t}$ the final state is enriched with left-handed t quarks. In this case the b quark is preferably emitted in the direction of the t quark. Therefore the two polar angle spectra in Fig. 13 look similar to each other. In case of $e^-_R e^+_L \to t\bar{t}$ the final state is enriched with right-handed t quarks. In this case the b quark is preferably emitted opposite to the direction of the t-quark. The polar angle of the b quark is a consequence of the Jacobian peak in the vicinity of $\cos\theta_{Wb} = 0$, with $\cos\theta_{Wb}$ being the opening angle between the b quark and the W boson, generated by the boost of the back-to-back configuration in the centre-of-mass frame of the decaying t quark into the laboratory frame.

5.1. Interpretation of the results

Table 6 lists the generated and reconstructed value of the forward-backward asymmetry $A^t_{FB,reco}$ as an estimator for the quality of the reconstruction.

	$e^-_L e^+_R \to t\bar{t}$		$e^-_R e^+_L \to t\bar{t}$	
	IDR-L	IDR-S	IDR-L	IDR-S
$A^t_{FB,gen}$	0.329		0.430	
$A^t_{FB,reco}$	0.342	0.340	0.430	0.430

Table 6: *Selection efficiencies and resulting $A^t_{FB,reco}$ for both beam polarisations and the two detector models under study.*

So far the results have been presented for full beam polarisation. Using the known formula [12]

$$\sigma_{\mathcal{P}_{e^-},\mathcal{P}_{e^+}} = \frac{1}{4}\left[(1-\mathcal{P}_{e^-}\mathcal{P}_{e^+})(\sigma_{-,+}+\sigma_{+,-}) + (\mathcal{P}_{e^-}-\mathcal{P}_{e^+})(\sigma_{+,-}-\sigma_{-,+})\right], \tag{5}$$

with $\sigma_{-,+}$ and $\sigma_{+,-}$ being the fully polarised cross-sections, the results can be extrapolated to the realistic beam polarisations of $\mathcal{P}_{e^-},\mathcal{P}_{e^+} = \mp 0.8, \pm 0.3$. The resulting Born level cross sections are 1070 fb in case of $\mathcal{P}_{e^-},\mathcal{P}_{e^+} = -0.8, +0.3$ and 519 fb in case of $\mathcal{P}_{e^-},\mathcal{P}_{e^+} = +0.8, -0.3$. The resulting statistical errors for an integrated luminosity of $\mathcal{L} = 1600\,\text{fb}^{-1}$ at each of the two polarisation settings are given in Table 7.

	$\mathcal{P}_{e^-},\mathcal{P}_{e^+}$	$(\delta\sigma/\sigma)_{stat.}$ [%]	$(\delta A^t_{FB}/A^t_{FB})_{stat.}$ [%]
IDR-L	$-0.8,+0.3$	0.17	0.7
	$+0.8,-0.3$	0.25	0.53
IDR-S	$-0.8,+0.3$	0.17	0.7
	$+0.8,-0.3$	0.25	0.53

Table 7: *Statistical precisions expected for the cross sections and A^t_{FB} for the case $\mathcal{P}_{e^-},\mathcal{P}_{e^+} = -0.8, +0.3$ and the two detector models under study.*

For both, the cross section and the forward backward asymmetry it can be expected that even at full luminosity the statistical error has to be taken into account. However, the systematic errors need to be carefully estimated. For the present analysis it would have to be checked how much the sample is contaminated by events for which the semi-leptonic decay yields τ-leptons or the τ-leptons are taken into account as in [5]. The contamination by fully hadronic $t\bar{t}$ events can be expected to be small.

To put the results into context, the precisions on the cross-sections and the forward-backward asymmetries are translated into precisions on electromagnetic form factors of the t quark. Figure 15 shows the precisions at the $1\,\sigma$ level expected at ILC500. The precisions are compared with those expected after the full HL-LHC running and estimations produced for FCC-ee [13] at the same confidence level. For ILC500, the two sets F_1 and F_2 have been extracted separately but within each set the uncertainties have been extracted simultaneously. The projections for HL-LHC are derived from

the *individual* constraints of EFT Wilson coefficients presented in Tab. C2.3 of Ref. [14] (the most favorable scenario for HL-LHC). These figures demonstrate clearly the superiority of a linear e^+e^- collider with polarised beams operated at an adequate centre-of-mass energy.

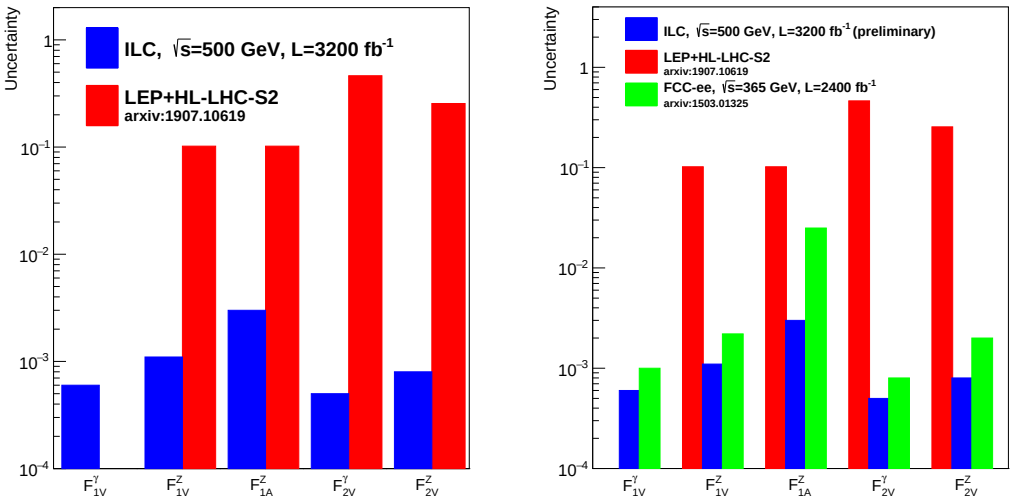

Figure 15: *Precisions the on electromagnetic t quark form factors expected after ILC500 compared with those expected after the full HL-LHC running and an estimation for FCC-ee after* $5000\,\mathrm{fb}^{-1}$. *The EFT does not include operators that map onto the* F_{1V}^{γ} *form factor. See text for further details.*

6. Summary

This note presents a study of the processes $e^+e^- \to t\bar{t}$ and $e^+e^- \to b\bar{b}$ with polarised beams at $\sqrt{s} = 500\,\mathrm{GeV}$. The events are simulated and reconstructed with the large and the small models of the ILD detector.

In case of t quark pair production the analysis focuses on the semi-leptonic decay channel in which the isolated lepton is available for the distinction between the t quark and the $\bar{t}$ quark. In case of the $e^+e^- \to b\bar{b}$ process the distinction has to be made by the measurement of the b quark charge, which helps also for a proper reconstruction of the $t\bar{t}$ quark pair. The charge of the b quark can be reconstructed with a purity of 80% using the combination of information available from the vertex charge, Kaons that have been measured in the TPC of ILD or isolated leptons in case of t-quark production.

The analysis shows that both, the large and the small detector model, are capable to provide a high precision measurement of the cross-section and the polar angle spectrum of semi-leptonic $t\bar{t}$ events with a mild advantage for the large detector. Assuming a total intergrated luminosity of $\mathcal{L} = 3200\,\mathrm{fb}^{-1}$ shared equally between the beam polarisations $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = \mp 0.8, \pm 0.3$, the cross sections of $t\bar{t}$ production can be measured to a statistical precision of about 0.2% and the forward backward asymmetry to a statistical precision of around 0.6%. The statistical precision on the cross section and the forward backward asymmetries are compatible with the scaling of the results found in [5].

For the first time the polar angle spectrum of the underlying b quark, issue of the t quark decay, is presented. This spectrum reveals more clearly the acceptance drop towards large polar angles. Still the polar angle of the b quark can be reconstructed accurately within $|\cos\theta < 0.8|$ for the two studied beam polarisations $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = \pm 1, \mp 1$. In case $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = +1, -1$ the efficiency drops however to 10% (compared with 30% for the case $\mathcal{P}_{e^-}, \mathcal{P}_{e^+} = +1, -1$). since only the vertex charge is used for this measurement. Here further work is clearly needed to improve the event yield. Both results allow

however already now for the perspective that in the future the fully hadronic final state can be taken into account for the analysis and that observables specific to the produced b quark can be addressed.

The study of the $t\bar{t}$ production has been accompanied by the second benchmark study $e^+e^- \to b\bar{b}$. Since only $46\,\mathrm{fb}^{-1}$ are available for this channel it should be rather considered as an auxiliary study. However, it is shown that the polar angle spectrum can be very well reconstructed even for the hard b jets and that migrations can be controlled at a satisfactory level. It is therefore justified to conclude that ILD should be able to make precision measurements in this channel even at a centre-of-mass energy of 500 GeV. Although the results are similar for both the small and the large detector model, this analysis, more than the analysis of the $t\bar{t}$ process reveals a slight preference for the large model. This is most clearly visible in the purity of the charge measurement using Kaons, which may depend on the actual TPC radius.

Also in the future all heavy quark studies should be carried out in close cooperation with each other. As can be seen from the present study there are many common issues between the studies. In the future emphasis will be put on systematic uncertainties given e.g. by hemisphere correlations. These studies may be more involved in case of $t\bar{t}$ since in general the two b quarks are not back-to-back.

Acknowledgements

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A. Details on particle separation via dE/dx

In Fig. 8.6 of the IDR the separation power between different particle types is given. A useful supplementary information of this summary is the relativ frequency of the different particle types. The Fig. 16 displays the normalised dE/dx spectrum for different particles in different momentum ranges for the large and the small detector model. In both cases there is a clear separation of Kaons from pions. The latter are however much more abundant. There is only a small population of protons. Figure 17 shows the dE/dx spectra for the two processes under study.

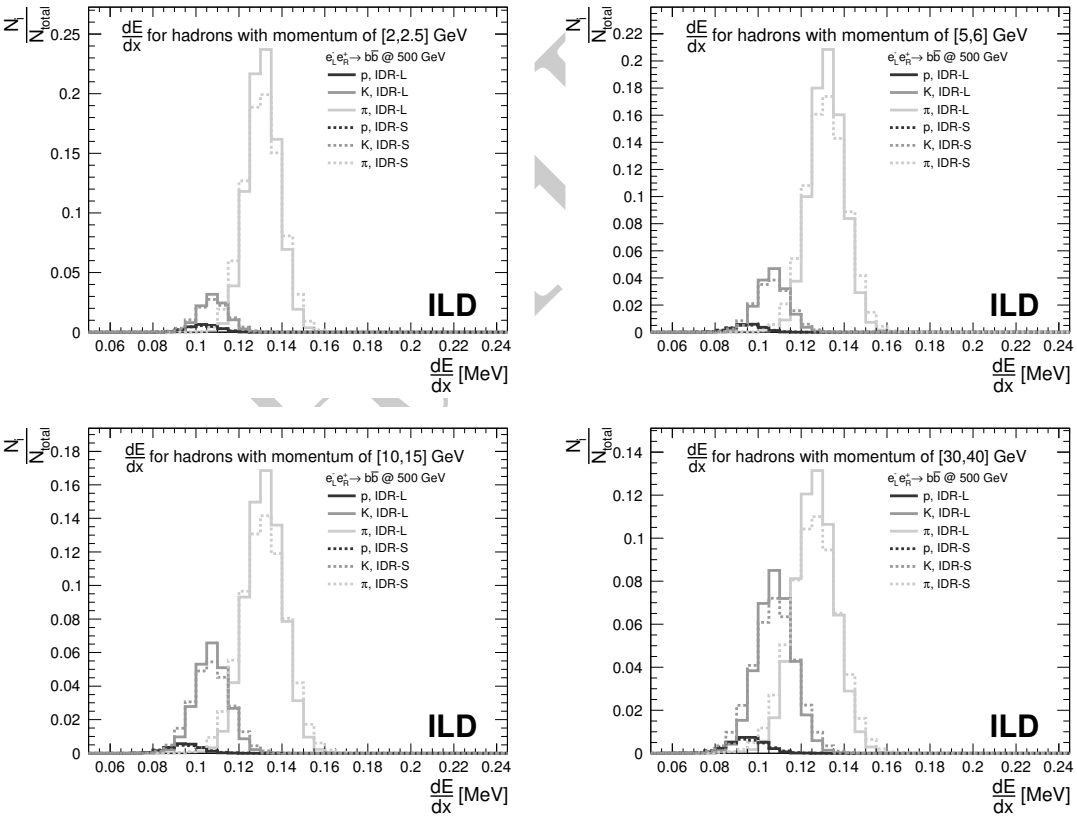

Figure 16: Projection of dE/dx for several momentum ranges. Comparison of hadron separation performance by different detector models in $b\bar{b}$ final states.

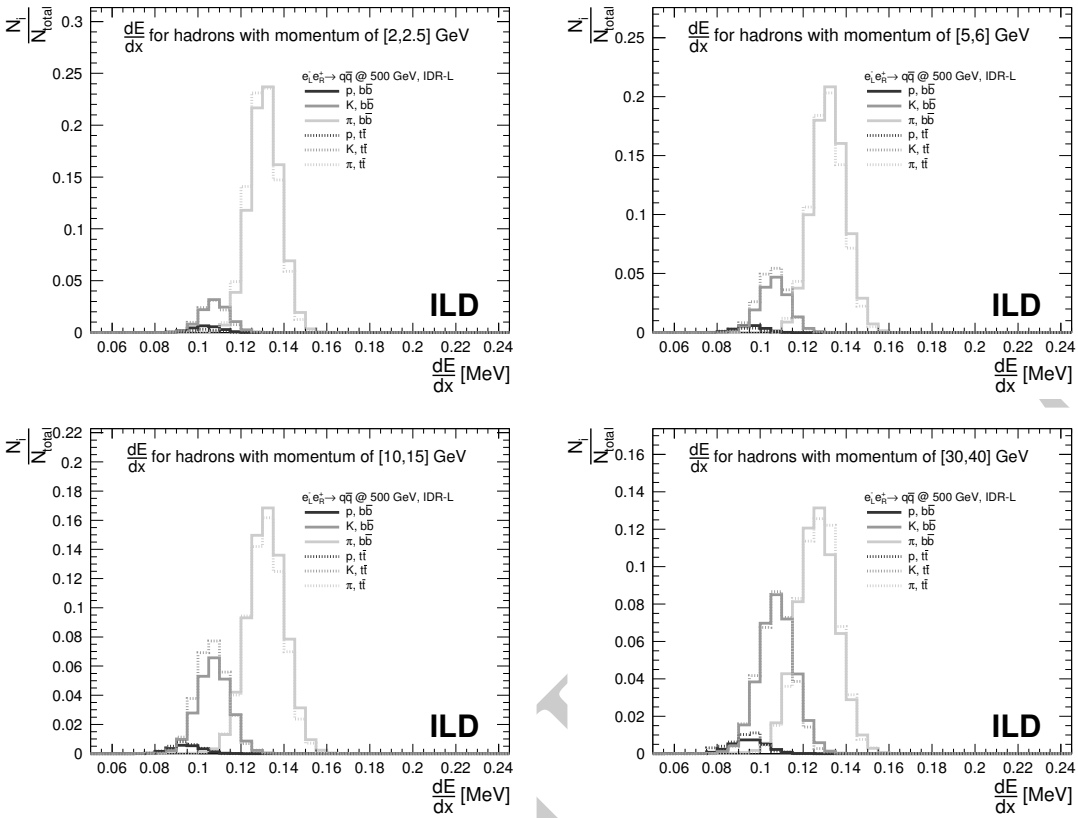

Figure 17: Projection of dE/dx for several momentum ranges. Comparison of hadron separation performance by the large model for different topologies.