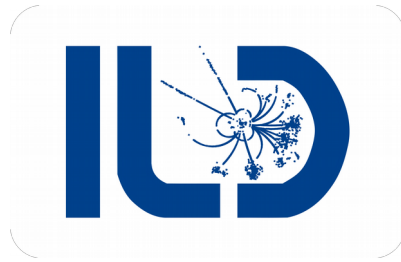


Heavy quark production in high energy electron positron collisions

Adrián Irlés^{*†}, R. Pöschl^{*}, F. Richard^{*}
on behalf of the **ILD Collaboration**



* **IJCLab** CNRS/IN2P3

† **IFIC** (UV / CSIC)

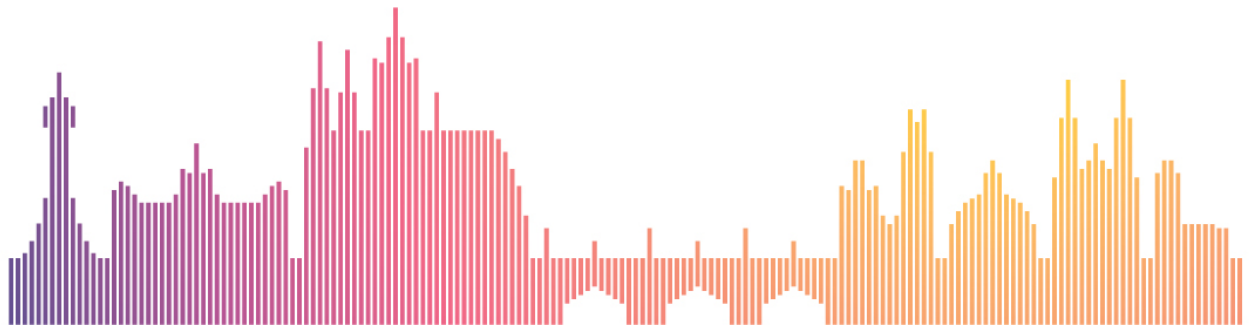
ICHEP 2020 | PRAGUE

40th INTERNATIONAL CONFERENCE
ON HIGH ENERGY PHYSICS

**VIRTUAL
CONFERENCE**

28 JULY - 6 AUGUST 2020

PRAGUE, CZECH REPUBLIC



Outline of the talk

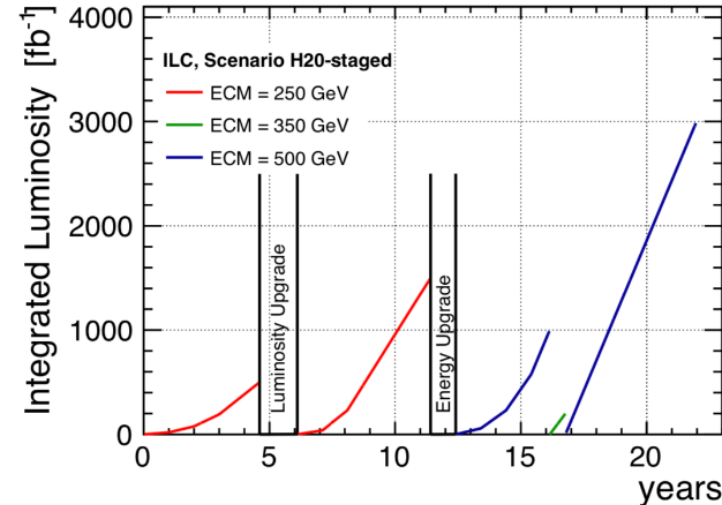
- ▶ What do we want to measure and why?
- ▶ Where? International Linear Collider, **ILC250**, and the International Large Detector **ILD**
- ▶ **top/b/c-quark electroweak couplings** extracted from **differential cross section** measurements
 - Experimental prospects based on full simulation including a comprehensive study of the systematic uncertainties



ILC physics program

- ▶ ILC is a Linear Collider Project, to be hosted in Japan.
 - **Matured technology: TDR since 2013.**
- ▶ **All Standard Model particles within reach of planned linear colliders**
- ▶ High precision tests of Standard Model over wide range to detect onset of New Physics
- ▶ **Machine settings can be “tailored” for specific processes**
 - **Centre-of-Mass energy**
 - **Beam polarisation** (straightforward at linear colliders)
- ▶ Background free searches for BSM through beam polarisation

current ILC run plan: (basis of projections)



250 GeV: 2 ab⁻¹, 500 GeV: 4ab⁻¹, 350 GeV: 0.2 ab⁻¹

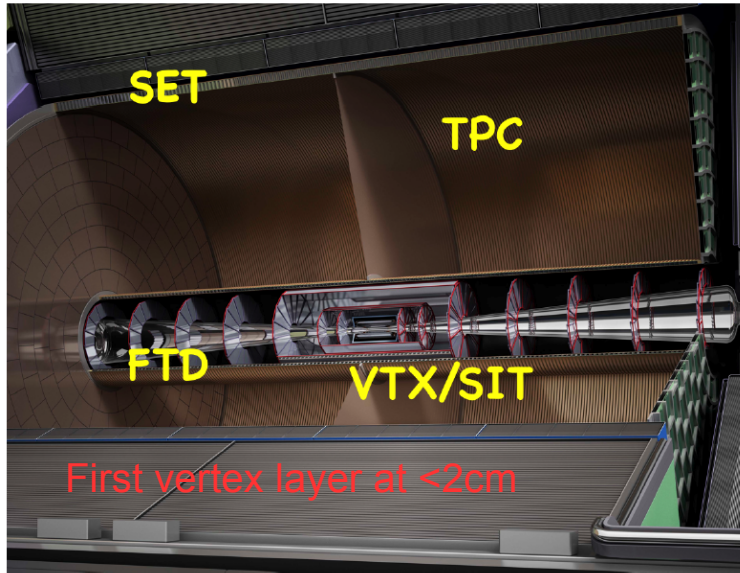
also, runs at 91 GeV (5B Z's) and 1000 GeV (8 ab⁻¹)

L upgrade: 5 Hz → 10 Hz; **E upgrade:** extend the linac

M. Peskin Snowmass (EF Workshop 21st July 2020)

ILD highlights

▶ ILD snapshot

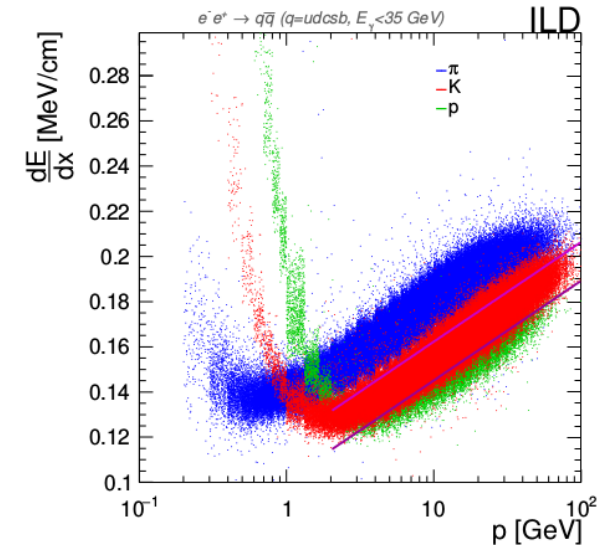
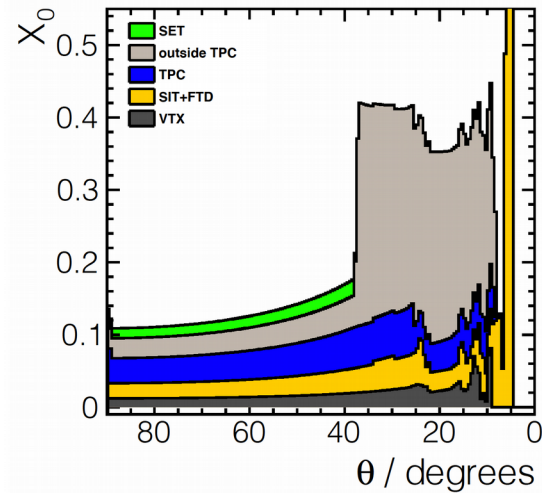


- ▶ High angular coverage with minimum material budget and PID (TPC)

- ▶ ILC experiments, as the ILD, will provide excellent:

- Beam IP constraint
- Secondary vertex separation and excellent flavour tagging
- Tracking efficiency (>99%)

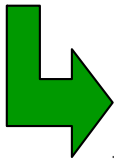
- ▶ Particle Flow optimized detector with high granularity calorimeters (> 10^8 cells!)



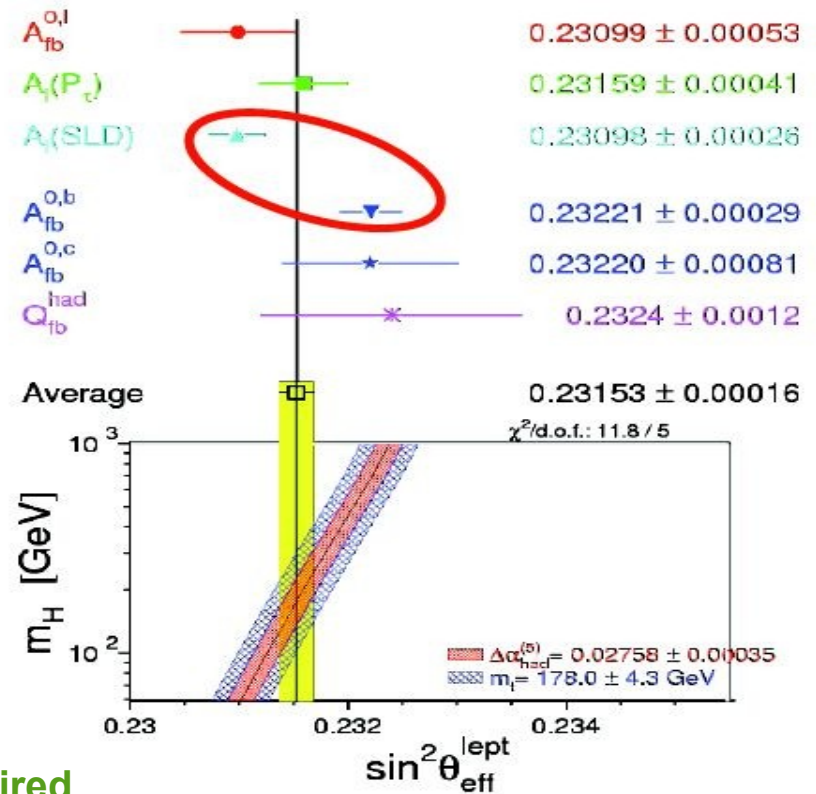
Motivation: LEP/SLC tension

- ▶ Current LEP & SLC best $\sin^2\theta_{eff}^l$ measurements show tension
 - This measurement is the one with **largest tension with the SM fit.**
 - Most precise single Individual determination of $\sin^2\theta_{eff}^l$ from SLC → Left-right asymmetry of leptons
 - Most precise single Individual determination of $\sin^2\theta_{eff}^l$ from LEP → forward backward asymmetry (b-quark)
- ▶ Heavy quark effect, effect on all quarks/fermions, no effect at all?

The **resolution** of this issue requires improving the the measurements precission an order of magnitude



Per mil level of experimental precision is required



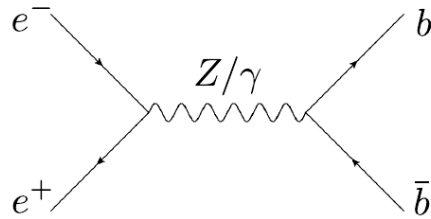
Motivation: Two fermion processes

- ▶ Many **BSM scenarios** (i.e. Randal Sundrum, compositeness, Higgs unification models...) predict heavy resonances coupling to the (t,b) doublet and also lighter fermions (i.e. c/s quarks)
 - **BSM resonances** tend to **couple** to the **right components**.
 - Only (t,b) doublet
 - Peskin, Yoon arxiv:1811.07877
 - Djouadi et al arxiv:hep-ph/0610173
 - All fermions
 - Hosotani et al arxiv:1705.05282 arxiv:2006.02157
 - For an EFT review see M. Perelló's talk and arxiv:1907.10619
- ▶ **How do we probe these BSM scenarios ?**

Probe such BSM require at least **per mil level of experimental precision**
tt/**bb**/cc... (ss?) **Can we do it?**
(this talk)

Motivation: Two fermion processes

- ▶ Differential cross section for (relativistic) di-fermion production



$$\frac{d\sigma}{d\cos\theta}(e_L^- e_R^+ \rightarrow f \bar{f}) = \Sigma_{LL}(1 + \cos\theta)^2 + \Sigma_{LR}(1 - \cos\theta)^2$$

$$\frac{d\sigma}{d\cos\theta}(e_R^- e_L^+ \rightarrow f \bar{f}) = \Sigma_{RR}(1 + \cos\theta)^2 + \Sigma_{RL}(1 - \cos\theta)^2$$

- The helicity amplitudes Σ_{IJ} , contain the couplings g_L/g_R (or Form factors or EFT factors)
- Left/right asymmetries (characteristic for each fermion)

- ▶ **Only beam polarisation allows inspection of the 4 helicity amplitudes for all fermions**

- ▶ **These processes have been deeply studied at LEP/SLC at the Z-pole**

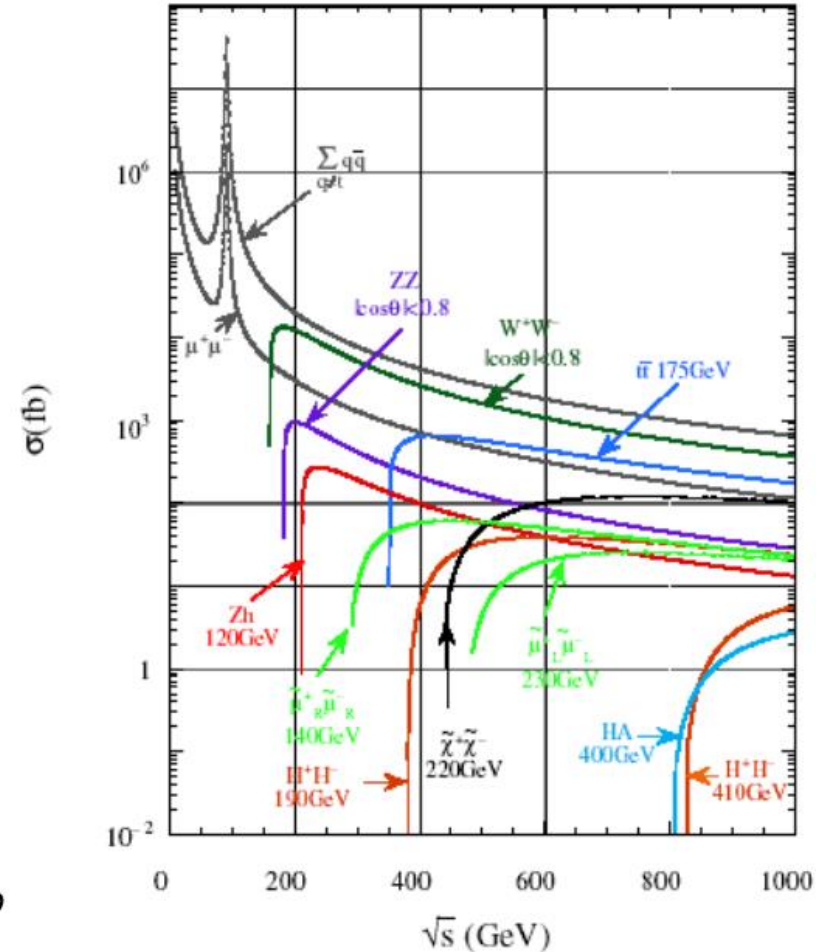
- no access to the γ or Z/γ interferences
- Moderated quark tagging or charge measurements capabilities.
- Also moderated angular acceptance of the detectors

Cross sections

$$\sigma_{e^-e^+ \rightarrow q\bar{q}}$$

	Channel	σ_{unpol} [fb]	$\sigma_{-,+}$ [fb]	$\sigma_{+,-}$ [fb]
500 GeV	q=t	572	1564	724
	q=b	372	1212	276
	q=u+d+s+c	2208	6032	2793
250 GeV	q=t	--	--	--
	q=b	1756	5677	1283
	q=c	3020	8518	3565
	q=u+d+s	6750	18407	5463

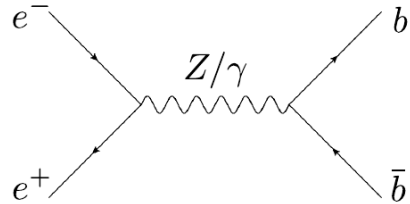
► Beam polarisation also enhances the cross section values



In this talk I concentrate on b-quark pair production at 250 GeV

Observables

- ▶ Quark (fermion) **electroweak couplings** can be **inferred from cross section, R_q** and forward backward asymmetry **AFB** observables.



$$R_q^0 = \Gamma_{q\bar{q}} / \Gamma_{had}(Z-pole)$$

$$\rightarrow R_q^{cont.} = \sigma_{q\bar{q}} / \sigma_{had}(s > Z-pole)$$

Quark identification. No need to measure an angular distribution, a priori.

$$\frac{d\sigma}{d\cos\theta}$$

Angular Distribution.

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

Quark ID + charge measurement (quark – antiquark disentangling)

Gives access to all left/right couplings.

Flavor tagging and charge measurement

► Flavor tagging

- Indispensable for analysis with final state quarks

► Quark charge measurements

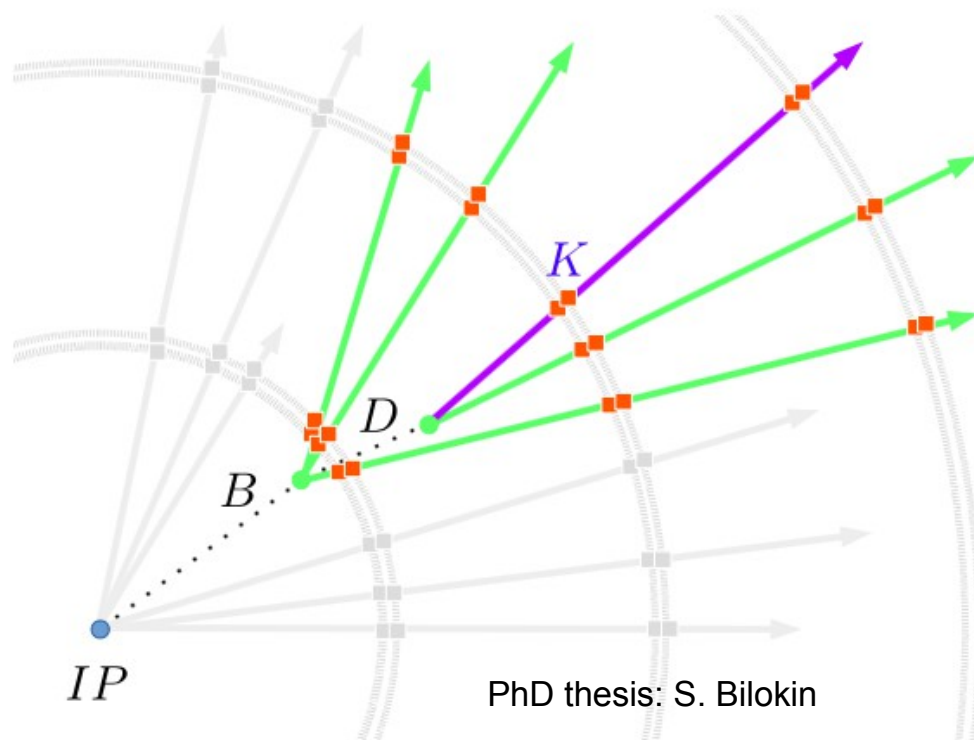
- Important for top-quark studies but Indispensable for $ee \rightarrow bb/cc/ss...$

► Charge measurements:

- Vtx charge and **Kaon Identification**
- High purity \rightarrow control of the migrations (**double tagging**)

► Future detectors can base their entire measurements on double Tagging and vertex charge

- LEP/SLC had to include single tags and semi-leptonic events



b/c-quarks: reconstruction efficiencies

Arxiv:1709.04289, PoS(EPS-HEP2019)624
ILD Note in progress

- ▶ 2 back-to back jets topology
- ▶ Main source of systematics in LEP/SLC:
 - Uncertainties related to tagging efficiency
 - The tagging efficiency needs to be measured (not MC estimated) to reach the per mil level of accuracy.
- ▶ New systematics sources for LC operating polarised beams far from the Z-pole
 - Beam polarisation
 - Event selection → **backgrounds from radiative return** events and **WW/ZZ/HZ**

Polarization	$\sigma_{e^-e^+ \rightarrow q\bar{q}}(E_\gamma < 35 \text{ GeV})[\text{fb}]$			$\sigma_{e^-e^+ \rightarrow q\bar{q}}(E_\gamma > 35 \text{ GeV})[\text{fb}]$		
	$b\bar{b}$	$c\bar{c}$	$q\bar{q} (q = uds)$	$b\bar{b}$	$c\bar{c}$	$q\bar{q} (q = uds)$
$e_L^-e_R^+$	5677.2	8518.1	18407.3	20531.4	18363.8	57651.3
$e_R^-e_L^+$	1283.2	3565.0	5643.5	12790.8	11810.8	36179.5

q \bar{q} signal

Rad. Ret. BKG

Up to x10 signal

Double Tag Method

- ▶ Needed to reach the per mil precision
- ▶ The sample consisted on events made of two hadronic jets (qqbar)
 - The LEP/SLC preselection consisted on a “simple” veto of $Z \rightarrow$ leptons events
- ▶ The method is based on the comparison of single vs double tagged samples
 - f_1 = ratio of number jets that are tagged as b-jets
 - f_2 = ratio of events in which both jets are tagged as b-jets

To remove modelling dependence on the efficiency of b-tagging

$$f_1 = \epsilon_b R_b + \epsilon_c R_c + \epsilon_{uds} R_{uds}$$

$$f_2 = \epsilon_b^2 (1 + \rho_b) R_b + \epsilon_c^2 R_c + \epsilon_{uds}^2 R_{uds}$$

- R_b and ϵ_b are measured simultaneously.

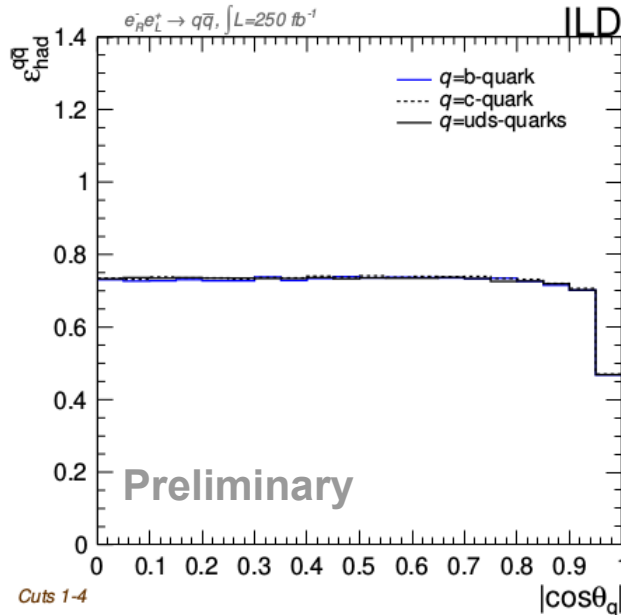
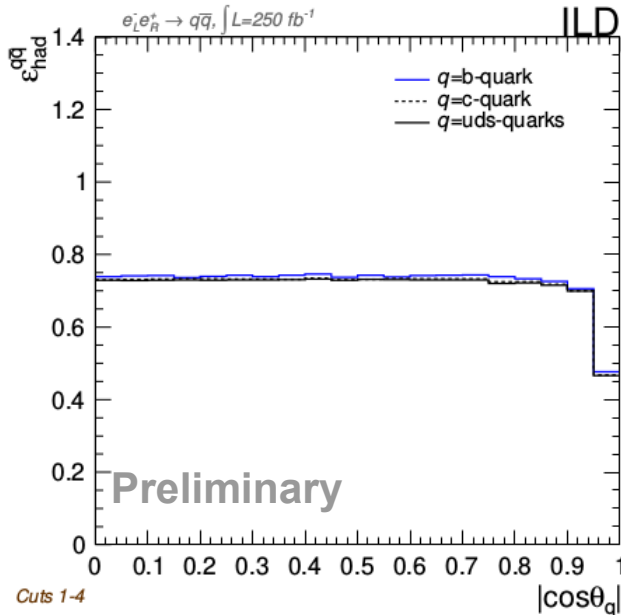
- ϵ_b = b-tagging efficiency
- ρ_b = b-tagging correlation factor
- ϵ_c = probability of tagging a c-quark jet as b-jet
- ϵ_{uds} = probability of tagging an uds-quark jet as b-jet

These values must be as small as possible and with small uncertainties

to not spoil our accuracy (not covered in this talk)

Double Tag Method

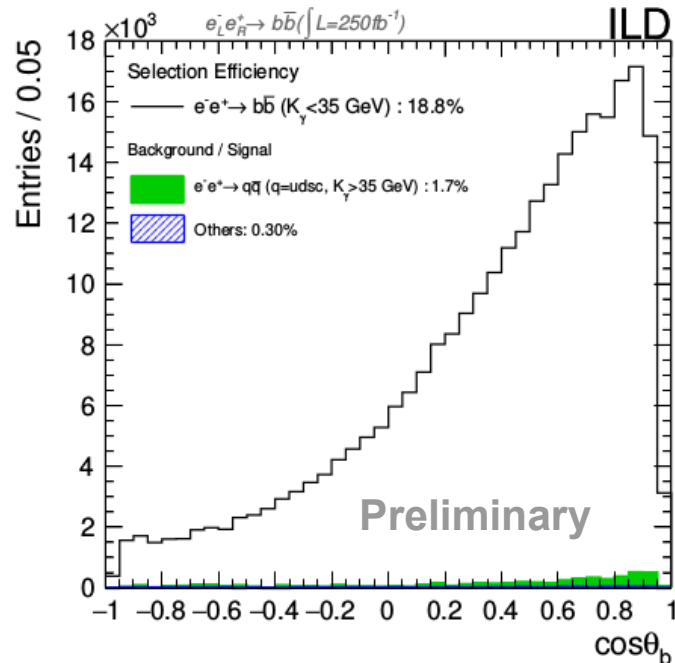
- ▶ **This method requires** (to minimise modelling uncertainties)
 - **Preselection with similar efficiency for all quark flavours**
 - Preselection that reduces to the minimum the main backgrounds
 - High quark tagging efficiencies with minimal mis-tagging efficiencies



Double Tag Method

► This method requires (to minimise modelling uncertainties)

- Preselection with similar efficiency for all quark flavours
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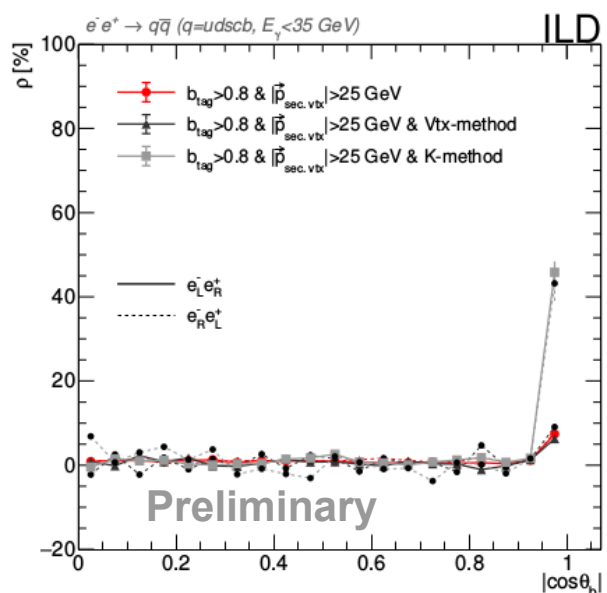
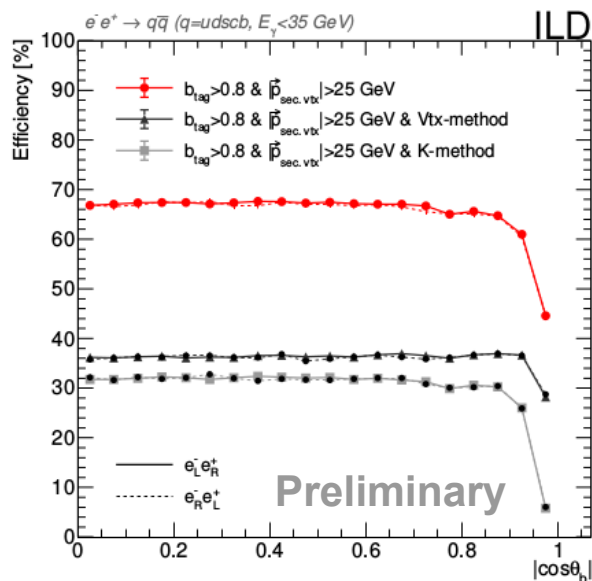


- Main bkg $ee \rightarrow Z\gamma(\text{ISR})$
 - $\sim x10$ larger than signal
 - For **$\sim 90\%$ of such ISR photons are lost in the beam pipe** \rightarrow events filtered by energy (angular) conservation arguments
 - The remaining **$\sim 10\%$ are filtered by identifying photons** in the detector (efficiency of $\sim 90\%$)
- Very small B/S $\sim 2\%$

Double Tag Method

► This method requires (to minimise modelling uncertainties)

- Preselection with similar efficiency for all quark flavours
- Preselection that reduces to the minimum the main backgrounds
- High quark tagging efficiencies with minimal mis-tagging efficiencies



► Excellent prospects for b-tagging (or c-tagging) with very low correlation factor ~ 0 ($\sim 2\%$ at LEP)

► Differential measurements!

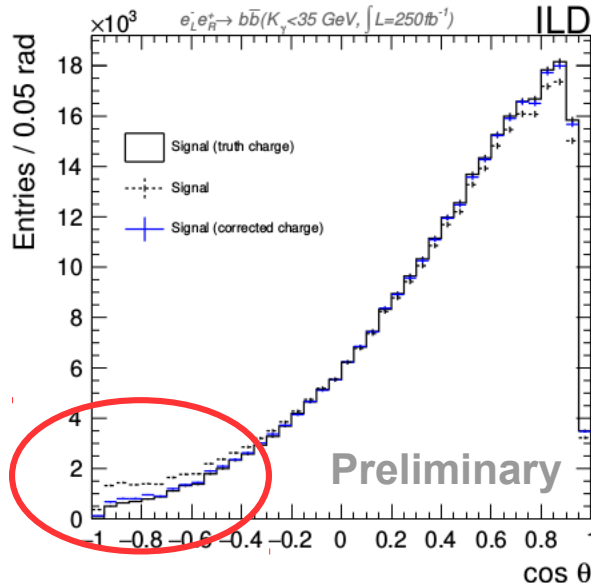
- Constant values for most of the angles
- Drop of acceptance the very forward region \rightarrow optimizations are under consideration

► Miss-efficiencies very small

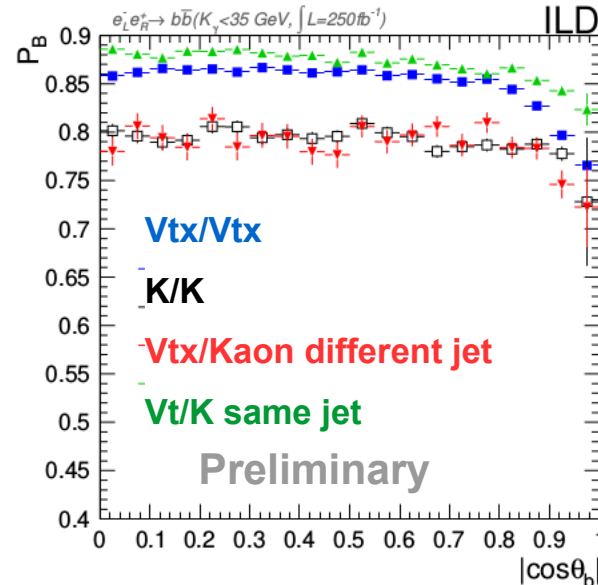
- $< 1\%$ for c-quark
- $\sim 0\%$ for uds

Charge measurement: migrations

- ▶ Mis-measurements of the jet charge produce a flip of the sign in the differential distribution: **migrations**.
 - Mistakes due to lost tracks, mis-identification of kaons...
- ▶ Migrations look as “new physics” → we need to correct them
 - Using data: double charge measurements with same and opposite charges (see back-up slides)
 - We measure the probability to reconstruct correctly the charge (P_B) and use it for correction
 - **DATA DRIVEN METHOD.**

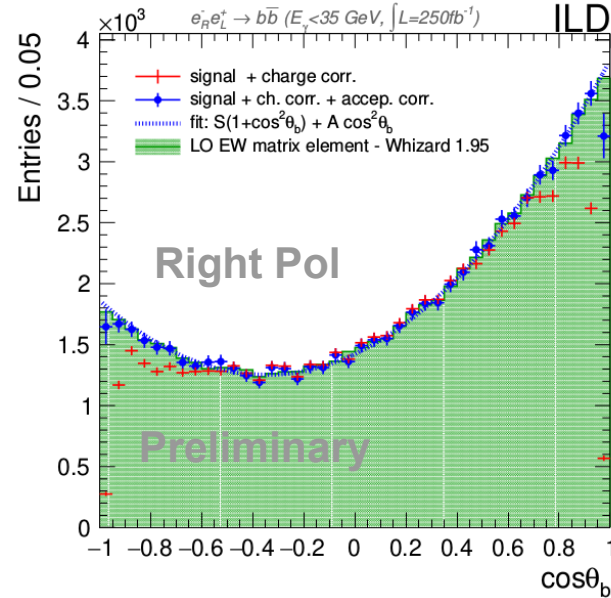
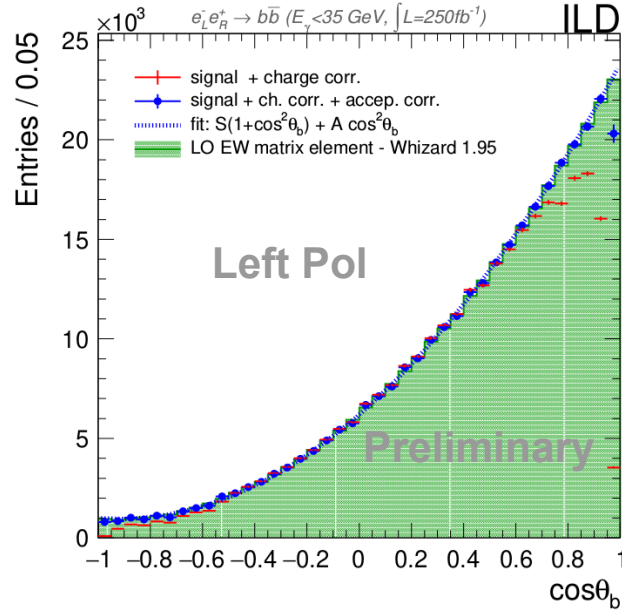


BSM or simple migrations?



▶ P_B limited by vertex reconstruction efficiency, Particle ID efficiency and B^0 oscillations.

Results (1)



Excellent agreement between predicted and reconstructed distributions

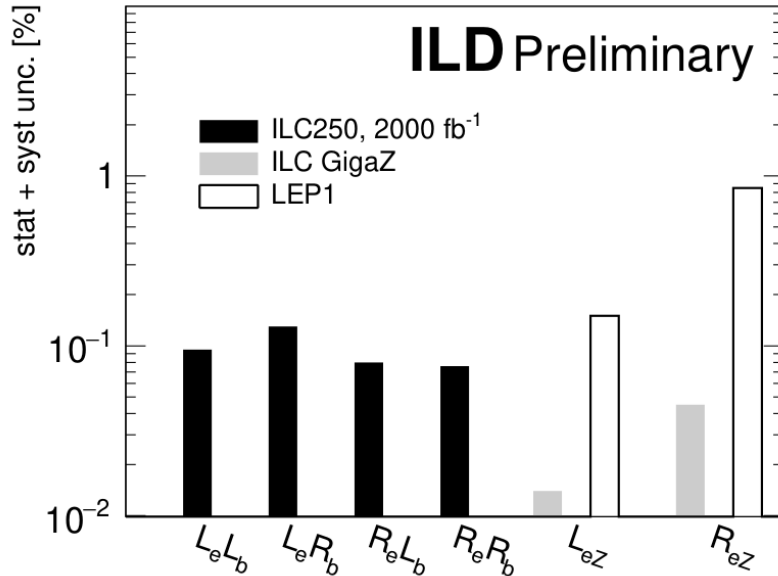
- ▶ Gap between red dots and green histogram = acceptance drop.
- ▶ Blue dots = corrected acceptance
- ▶ The fit is restricted to $|\cos\theta_b| < 0.8$
- *Minimal impact of the corrections*

	Beam Polarisation	
	(-+)	(+-)
$\Delta R_b^{cont.}$	0.12 (stat.) \pm 0.14 (syst.) %	0.15 (stat.) \pm 0.13 (syst.) %
$\Delta A_{FB}^{b\bar{b}}$	0.30 (stat.) \pm 0.05 (syst.) %	0.85 (stat.) \pm 0.10 (syst.) %

Stat unc (2000 fb⁻¹)

Syst unc.:

- Selection and background rejection
- quark tagging/mistagging (modélisation, QCD, correlations)
- Luminosity
- Polarisation



Couplings (notation for new resonances)

$$L_e L_b = Q_e Q_b + \frac{L_e Z L_b Z}{s^2 w c^2 w} B W Z + \sum_{Z'} \frac{L_e Z' L_b Z'}{s^2 w c^2 w} B W Z'$$

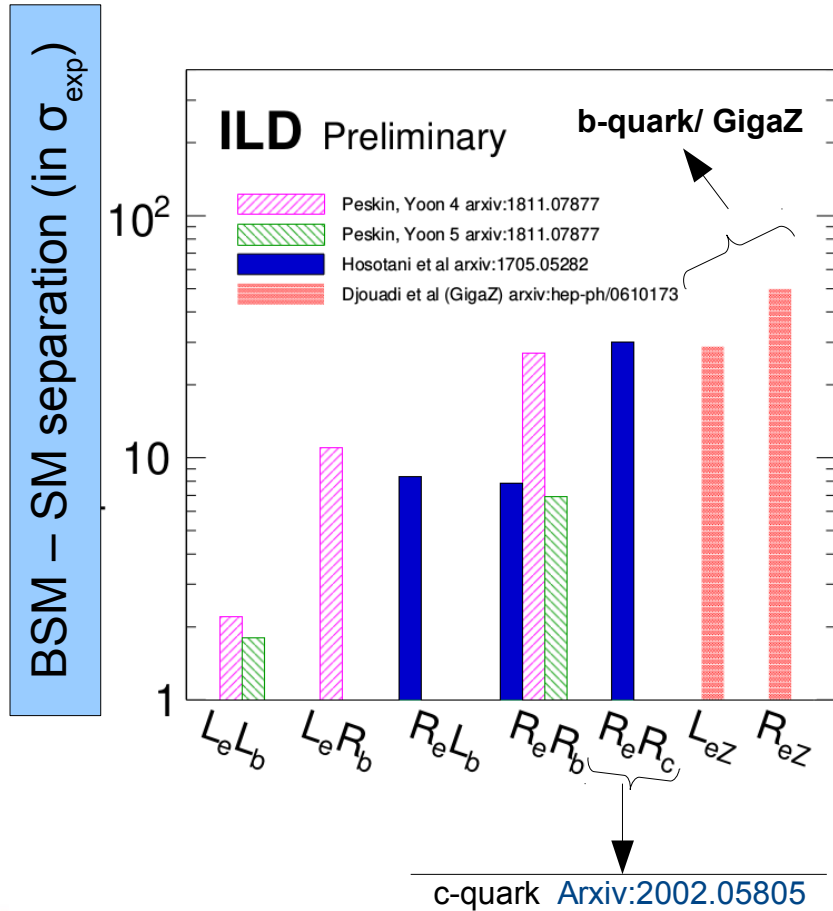
↓ ILC250
↓ SM
↓ GigaZ
↓ New resonances

↓
 Sensitive to Z-Z' mixing effects
 (that could explain AFB_b measurement of LEP?)

Prospects for couplings determination are order of magnitude better than at LEP

- ▶ Resolution of the LEP/SLC anomaly
- ▶ Full disentangling of helicity structure for all fermions only possible with polarised beams!!

Results (3) BSM benchmarks



► Many BSM predict deviations only for the right couplings

BEAM POLARISATION is crucial

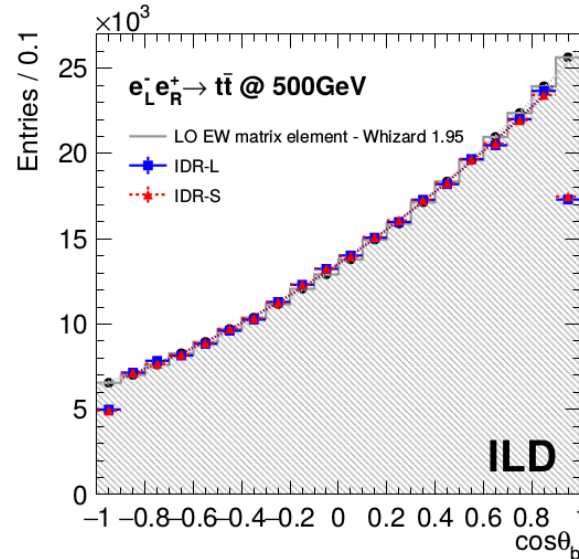
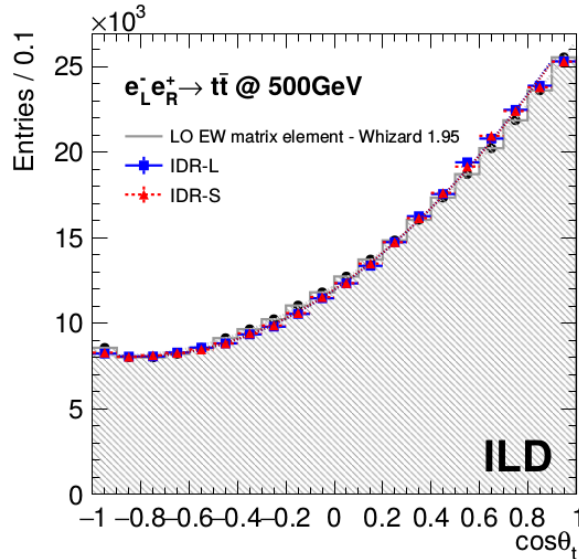
Expected number of standard deviations for different **RS/compositeness BSM scenarios** when determining the different EW couplings to c- and b-quark at **ILC250** (with GigaZ input).

- Models that predict multi-TeV Z' resonances
- With or without mixing at Z-pole
- See backup for more details on the models

Potential for discovery of new resonances $m_{Z'} \sim \mathcal{O}(20-30)$ TeV at ILC250

Top quark: results (1)

ILD-PHYS-PUB-2019-007



- ▶ Semi-leptonic channel
- ▶ Left polarisation plots
 - B-jet carries top direction information
 - Very useful for the hadronic channel!
- ▶ Right polarisation (not shown)
 - W-carries the top direction information → **lepton charge and c/s tagging become important**

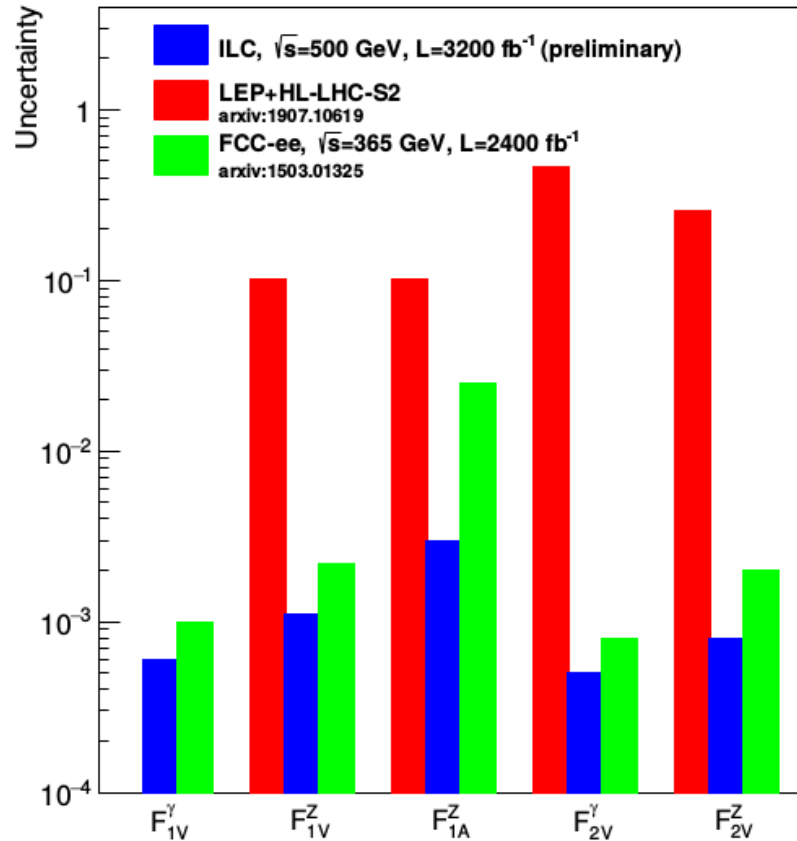
- ▶ Integrated Luminosity 4 fb^{-1}
- ▶ Thanks to the jet charge calculations capabilities, we could use all decay channels.
- ▶ Efficiencies of 75% (cross section) and 30% (differential cross section)
- ▶ Exact reproduction of generated spectra
 - Statistical precision on cross section: $\sim 0.1\%$ Statistical precision on A FB : $\sim 0.5\%$

Can expect that systematic errors will match statistical precision (but needs to be shown)

IDR-L/S
Are two detector
Concepts compared
In the ILD
Interim Design Report ILD
Arxiv:2003.01116

Top quark: results (2)

ILD-PHYS-PUB-2019-007



- ▶ e+e- collider way superior to LHC ($\sqrt{s} = 14$ TeV)
- ▶ Final state analysis at FCCee (polarisation)
 - Also possible at LC => Redundancy
- ▶ Two remarks:
 - 500 GeV is nicely away from QCD Matching regime
 - Less systematic uncertainties
 - The determination of axial form factors highly benefit from higher energies
- ▶ See M. Perelló's talk to interpret this plot in terms of EFT Wilson coefficients.

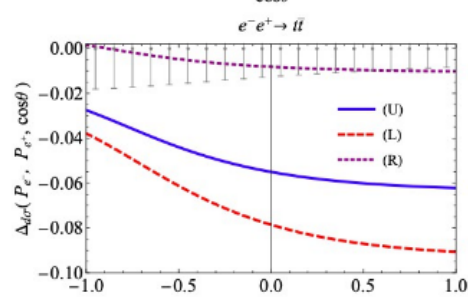
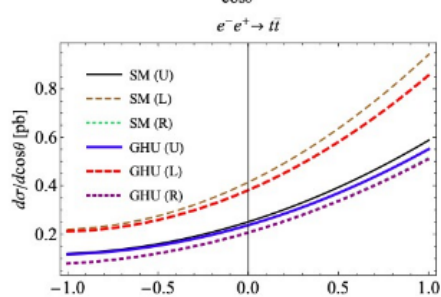
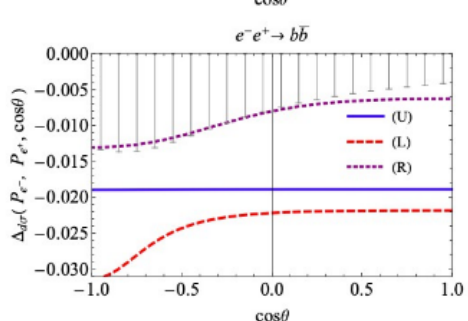
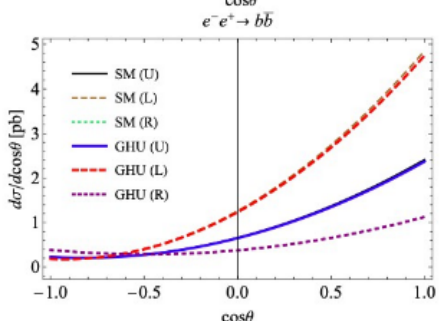
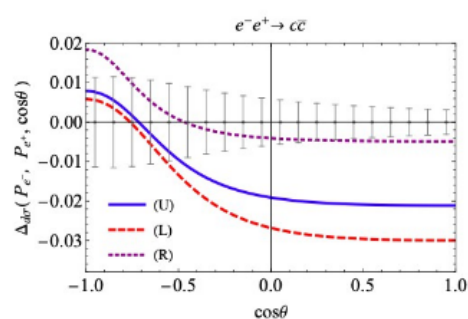
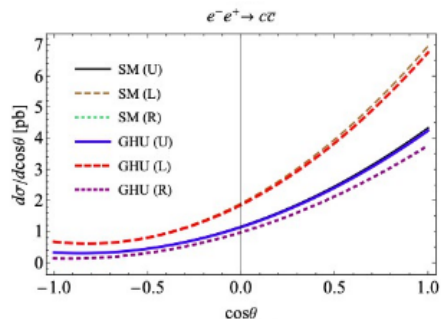
Summary / conclusions

- ▶ ILC is ideally suited for precision measurements of two-fermion final states
- ▶ **ILC will have the answer whether new physics acts on heavy doublet (t,b) only or on all fermions**
 - Will/would probe helicity structure of electroweak fermion couplings over at least one order of magnitude in energy (Z-Pole \rightarrow ~ 1 TeV)
- ▶ **Achievable experimental precisions $\sim 0.1 - 1\%$**
 - Demanding analysis requiring the full detector capabilities: Vertex charge and particle ID, PFO for final state jets, etc
 - **Comprehensive assessment of the systematic uncertainties done (b-quark)** or in progress (top and charm)
- ▶ **Effects may become already visible at 250 GeV stage for b quark and c quarks (and other light fermions)**
 - **Amplification of effects at higher energies** (studies at 500 GeV at preliminary stage)
 - **Clear and unique pattern thanks to polarised beams**
- ▶ Active phenomenological studies in terms of global analyses (EFT) and concrete models (not covered in this talk)
- ▶ Theory challenges (not covered in this talk)
 - Need at least NLO electroweak predictions (and MC programs) for correct interpretation of results

Thanks for your attention.

a BSM example: GUT Inspired Grand Higgs Unification Model

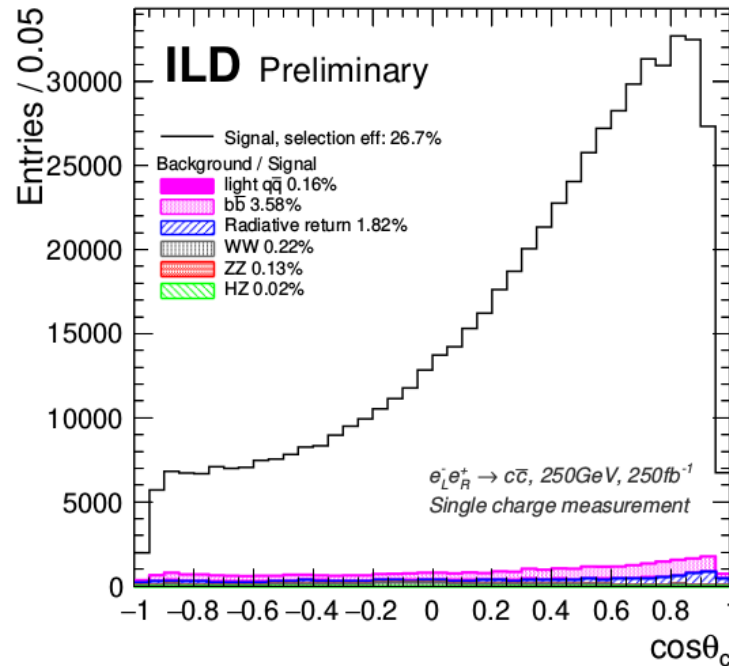
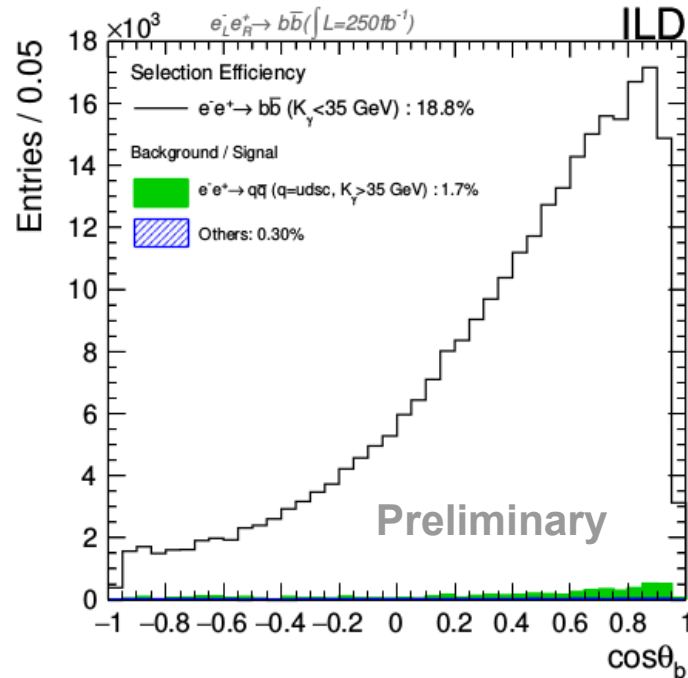
arxiv:2006.02157



- Model parameter is Hosotani angle θ_H yielding the Higgs-Potential as consequence of Aharonov-Bohm Phase in 5th dimension
- Model defined in Randall-Sundrum warped extra dimensions
 - KK excitations of gauge bosons and new bosons modify fermion couplings
- Predictions for ILC
 - $m_{KK} = 13$ TeV and $\theta_H = 0.1$
- Deviations from SM of the order of a few %
 - Effects measurable already at 250 GeV
 - Effects amplified by beam polarisations
 - Effects for tt, bb and cc (and other light fermions)
- One concrete example for importance to measure full pattern of fermion couplings
- Full pattern only available with beam polarisation

b/c-quarks: reconstruction efficiencies

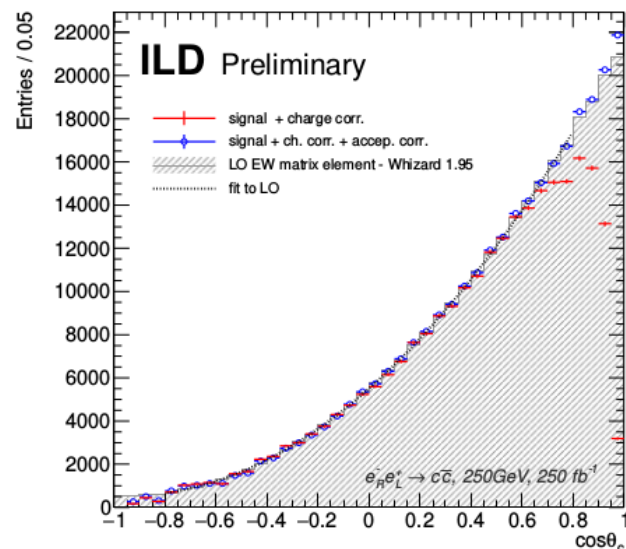
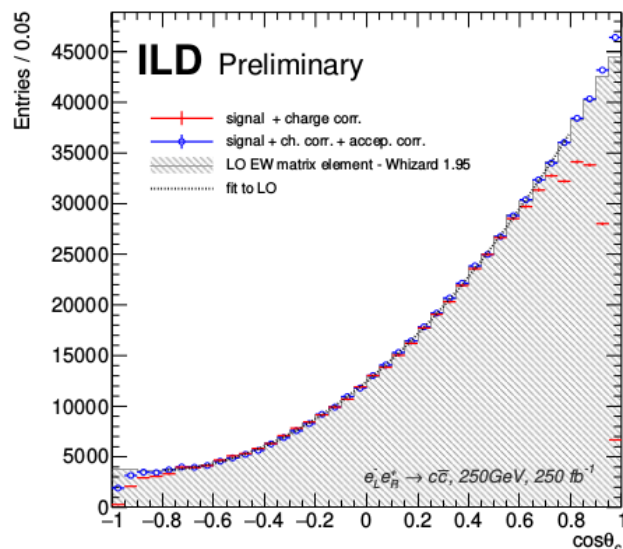
- ▶ After the preselection, we apply the b/c tagging including charge measurement for differential cross sections.
- Efficiencies for inclusive cross section are $\sim x2$ larger



▶ Background \sim free analysis!

Arxiv:2002.05805

b/c-quarks: Results (2)



c-quark case

- ▶ Similar precisions (work in progress)
- ▶ Lower tagging efficiency compensated by higher statistics for both polarisations.
- ▶ Kaon Identification becomes the most promising channel for the charge measurement

Top-quark: Reconstruction efficiencies

$$e_L^- e_R^+ \rightarrow t\bar{t} \text{ at } 500 \text{ GeV}$$

General selection cuts	IDR-L	IDR-S
Isolated Lepton	92.1%	92.1%
$btag_1 > 0.8$ or $btag_2 > 0.3$	81.2%	81.1%
Thrust < 0.9	81.2%	81.1%
Hadronic mass	78.2%	78.2%
Reconstructed m_W and m_t	73.4%	73.4%
<i>t</i> quark polar angle spectrum		
$\gamma_t^{had.} + \gamma_t^l > 2.4$	62.2%	61.8%
$ p_{R,had} > 15 \text{ GeV}$	34.5%	33.9%
" $t\bar{t}$ identification"	30.6%	30.2%
<i>b</i> quark polar angle spectrum		
No additional cuts		

$$e_R^- e_L^+ \rightarrow t\bar{t} \text{ at } 500 \text{ GeV}$$

General selection cuts	IDR-L	IDR-S
Isolated Lepton	94.1%	94.0%
$btag_1 > 0.8$ or $btag_2 > 0.3$	84.9%	84.8%
Thrust < 0.9	84.9%	84.8%
Hadronic mass	82.2%	82.3%
Reconstructed m_W and m_t	77.6%	77.5%
<i>t</i> quark polar angle spectrum		
$\gamma_t^{had.} + \gamma_t^l > 2.4$	64.1%	64.1%
<i>b</i> quark polar angle spectrum		
$Vtx+Vtx$	10.8%	10.3%

Total cross section

- ▶ Typical selection efficiencies for the 75%
- ▶ Independent of beam polarisation

Differential cross section

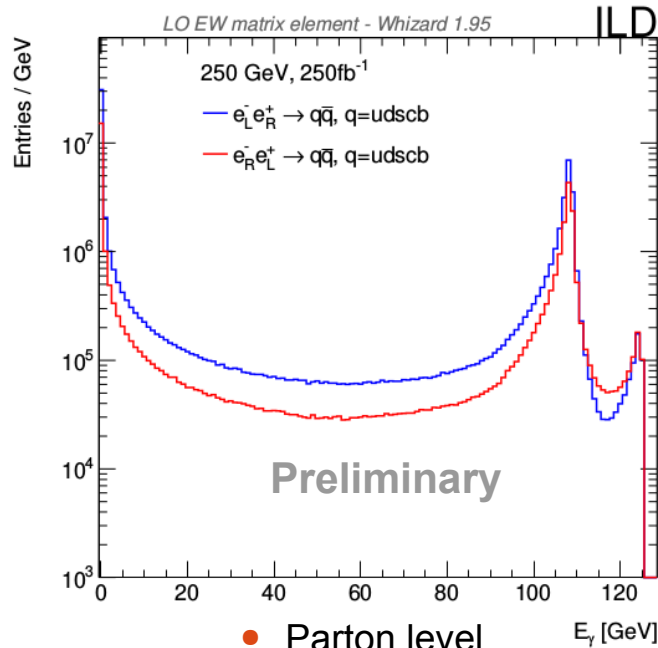
- ▶ Differences for beam polarisations
- ▶ Left hand polarisation more vulnerable to migrations
- ▶ Requires information from the hadronic state
- ▶ Vertex / Kaon as in the bb-case

Predictions (as a function of the ISR)

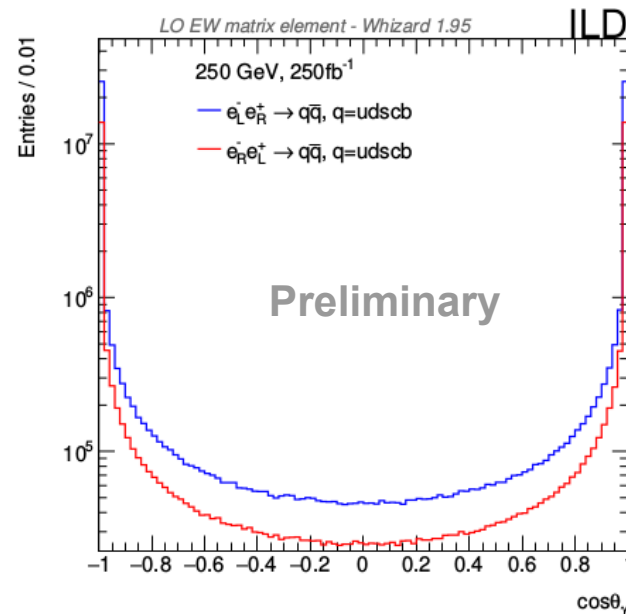
- ▶ The cross section depends on the “effective” center of mass energy
 - At which the Z/γ couple to the quark-antiquark pair

$$\frac{d\sigma_{e^-e^+ \rightarrow q\bar{q}}^{cont.}}{d\cos\theta_q}(s) \rightarrow$$

$$\rightarrow \frac{d\sigma_q^{cont. \bar{q}}}{d\cos\theta_q}(\hat{s} > s_{cut}) = \frac{d\sigma_{e^-e^+ \rightarrow q\bar{q}}^{cont.}}{d\cos\theta_q}(E_\gamma < K_{cut})$$



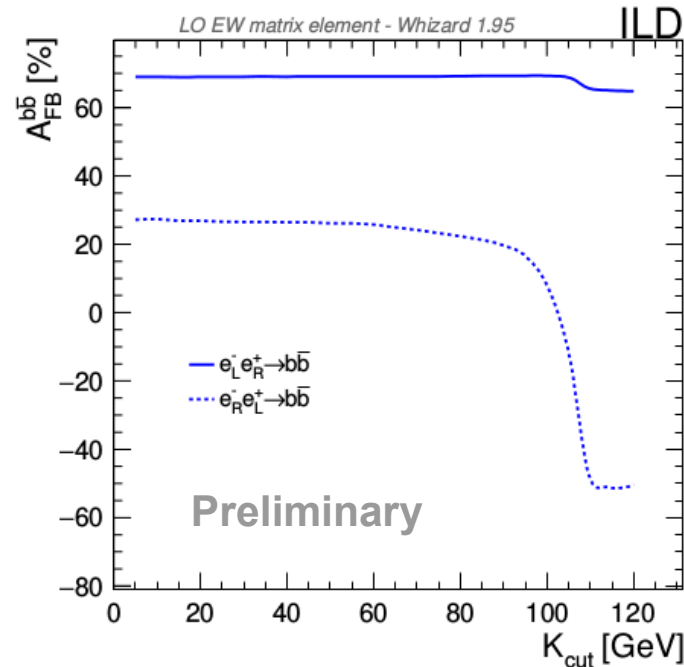
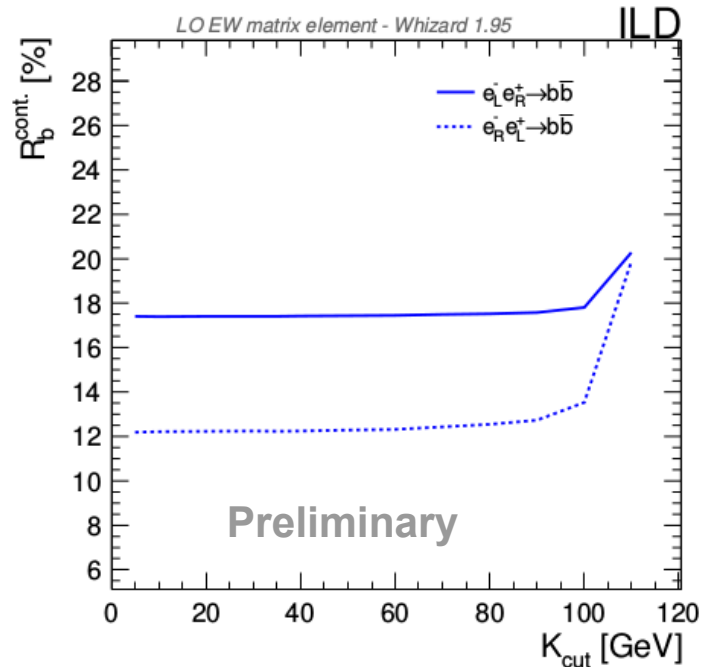
• Parton level



• Parton level

Predictions (as a function of the ISR)

$$\frac{d\sigma_{e^-e^+ \rightarrow q\bar{q}}^{cont.}}{d\cos\theta_q}(E_\gamma < K_{cut})$$



- ▶ The observables remain basically flat for a large range of the Kcut
- ▶ Drastic change when the photn ISR is large enough to produce a return to the Z-pole
 - We need to avoid that region of the phase space.

- ▶ Alternatives to $m(2\text{jets})$?
- ▶ Estimator of the energy of the photon ISR using only the two reconstructed jets.
 - From momentum conservation (if the photon/s are emitted parallel to the beam pipe):

$$|\vec{k}| \approx K_{reco} = \frac{250 \text{ GeV}}{\sin \Psi_{acol} + \sin \theta_1 + \sin \theta_2}$$

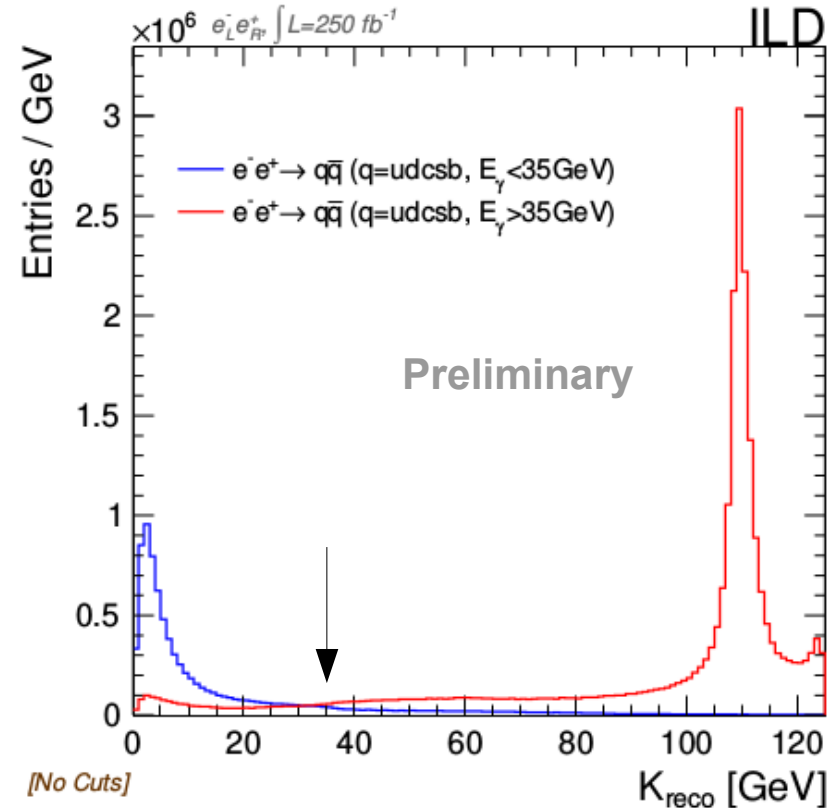
Two jet acolinearity

$$\sin \Psi_{acol} = \frac{|\vec{p}_{j_1} \times \vec{p}_{j_2}|}{|\vec{p}_{j_1}| \cdot |\vec{p}_{j_2}|}$$

Jet angular variables (w.r.t. detector frame)

Preselection : Kreco

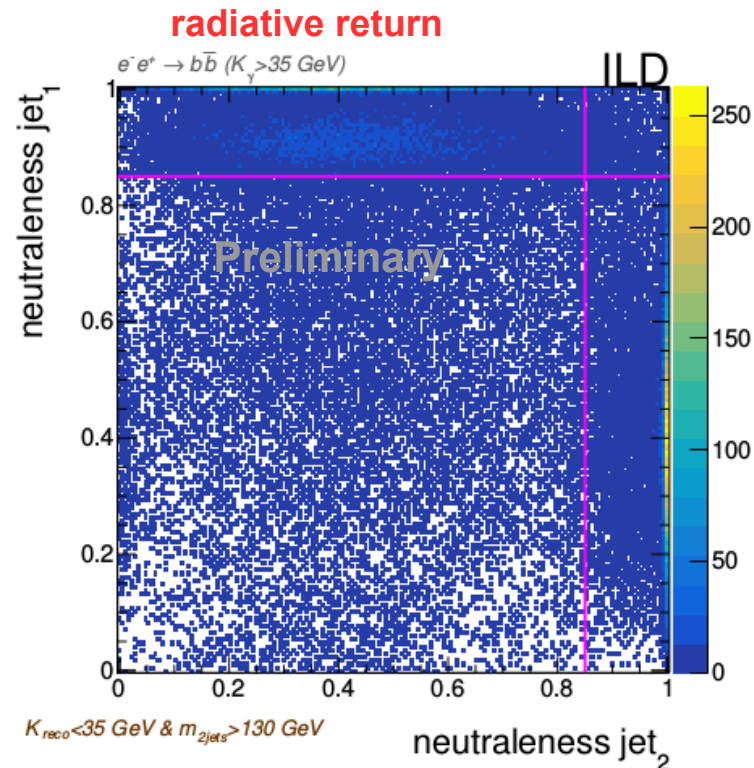
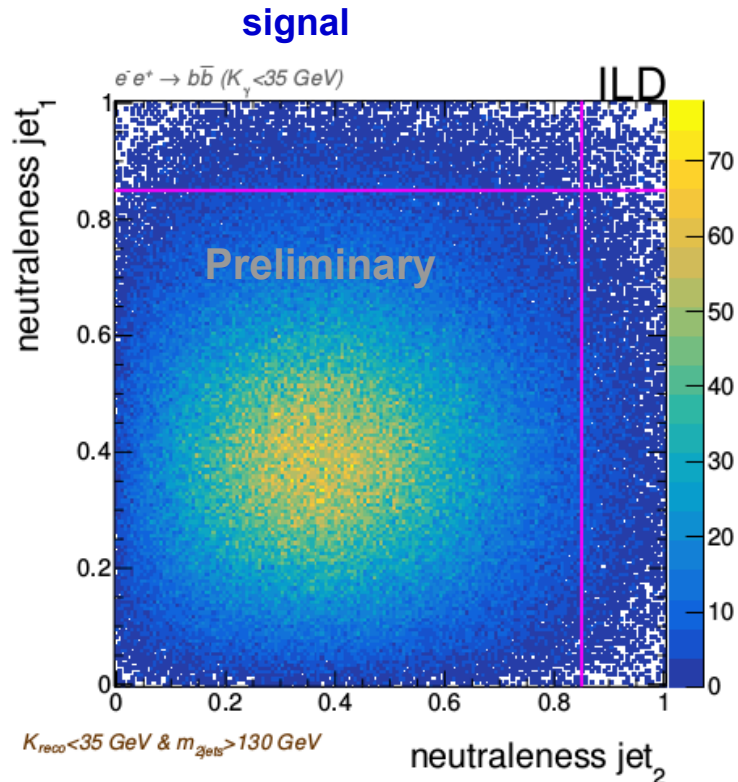
- ▶ Estimator of the energy of the photon ISR
- ▶ We apply a cut of $K_{reco} < 35$ GeV
- ▶ Some signal events have larger K_{reco} (~15%)
 - Because of detector resolution and double photon ISR
- ▶ Some radiative return events have $K_{reco} < 35$ GeV (~7%)
 - Because the photon(s) has not escaped through the beam pipe
- ▶ Can we identify the photon clustered in one or both jets and veto these events?

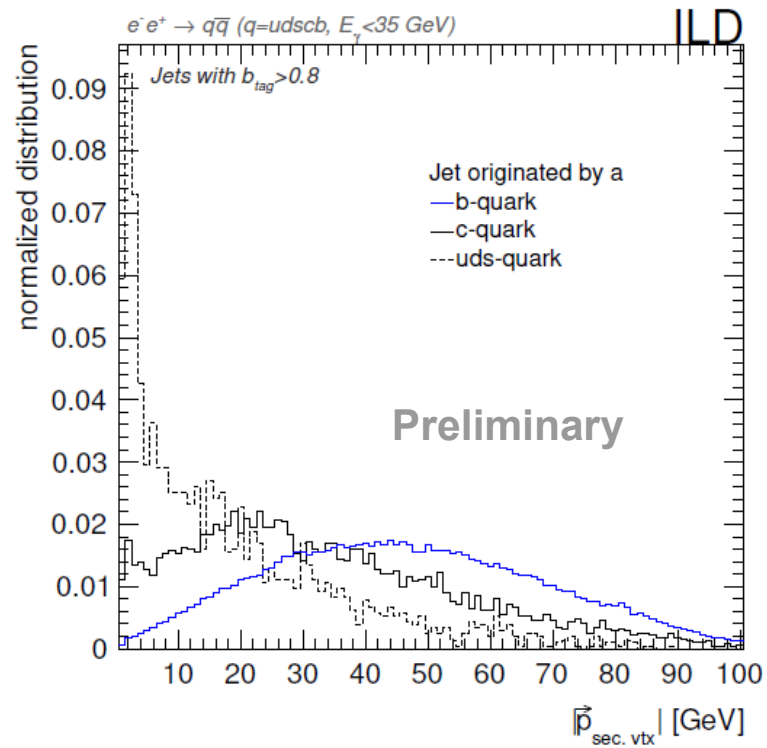
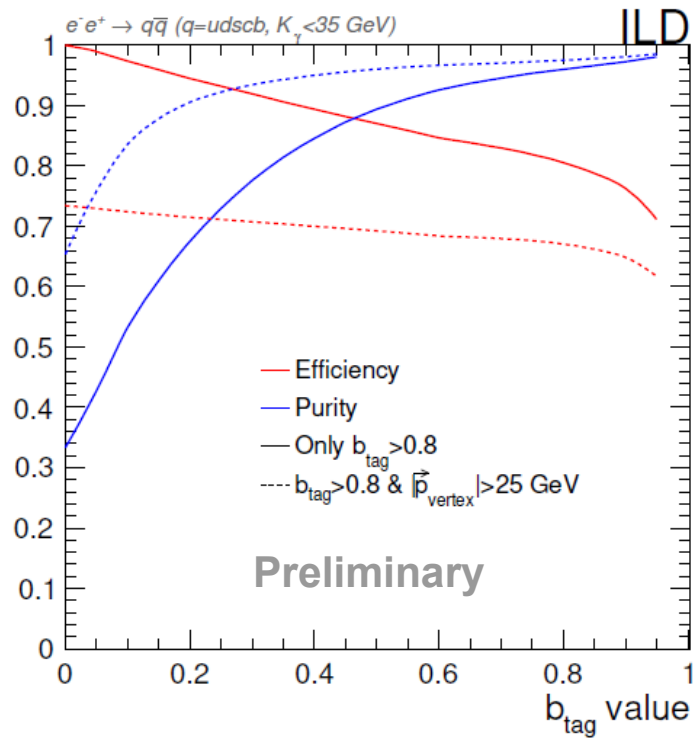


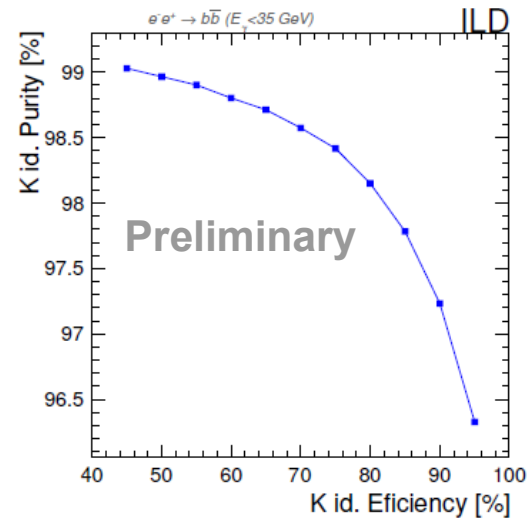
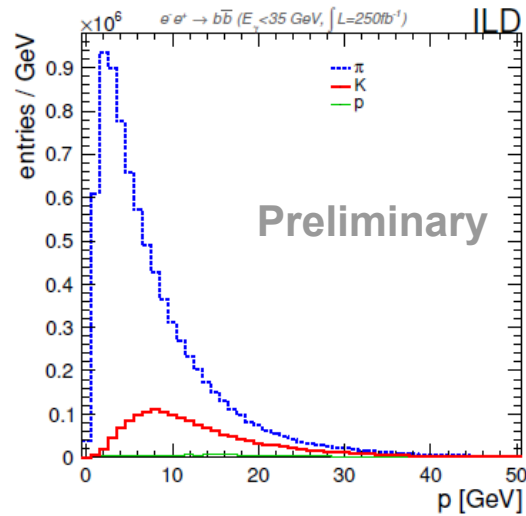
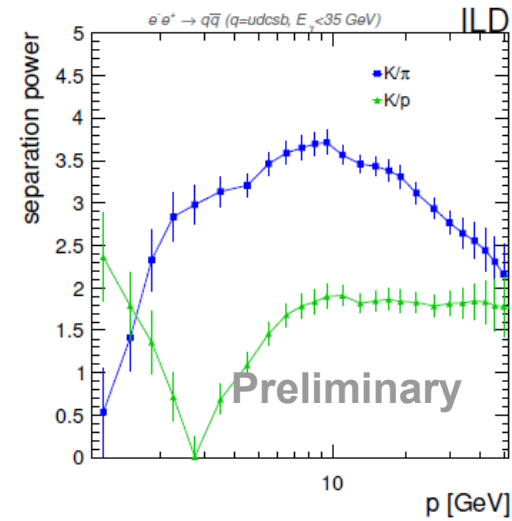
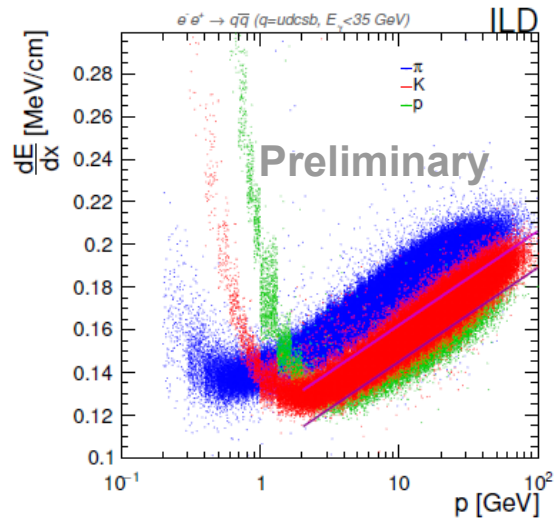
Preselection : Photon Veto

- We look at the neutralness of the jets

$$\text{neutralness}_j = \frac{\sum_{i=nPFO} E_i}{E_j}$$

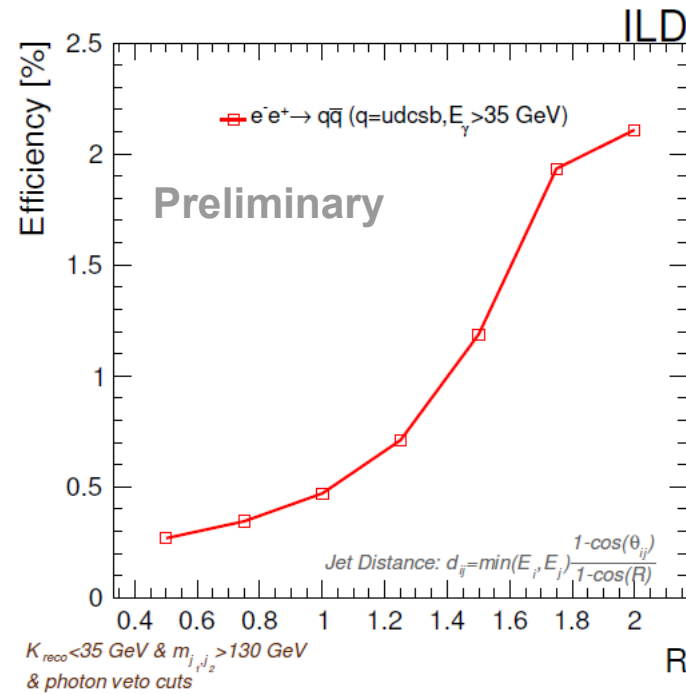
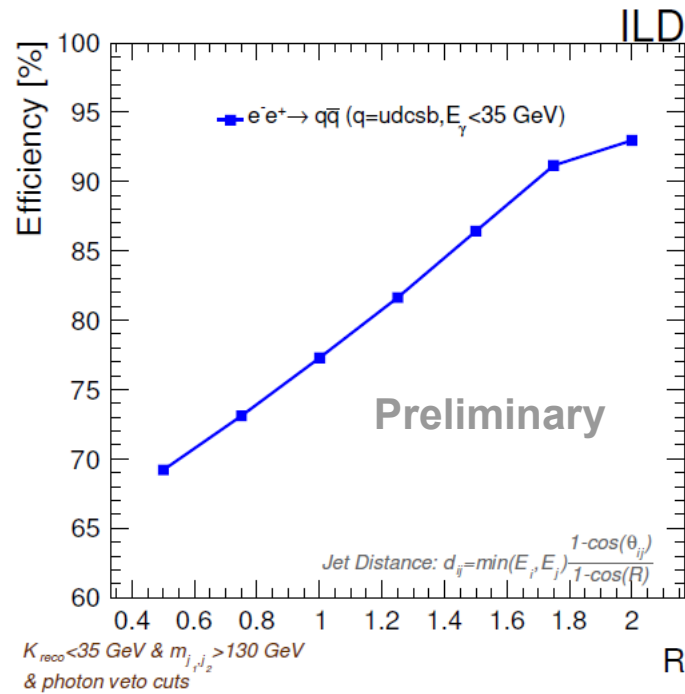






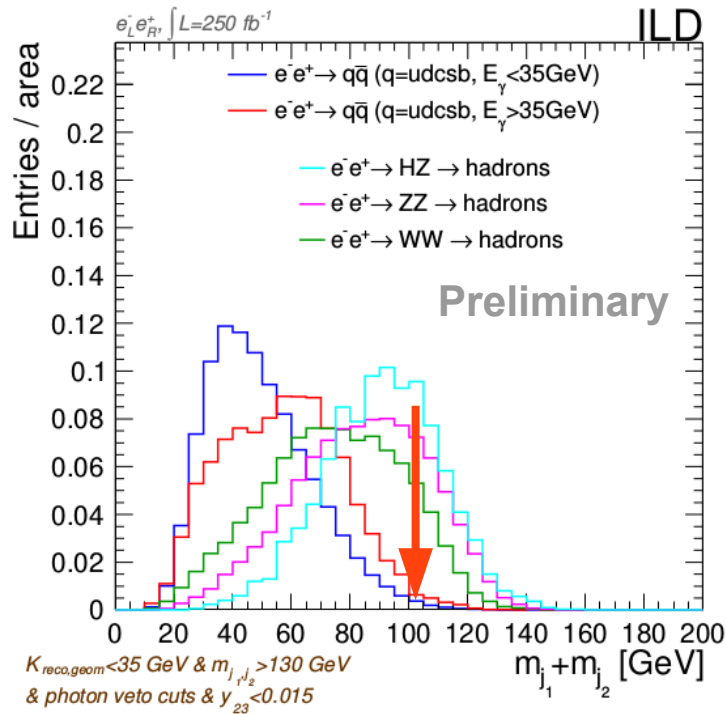
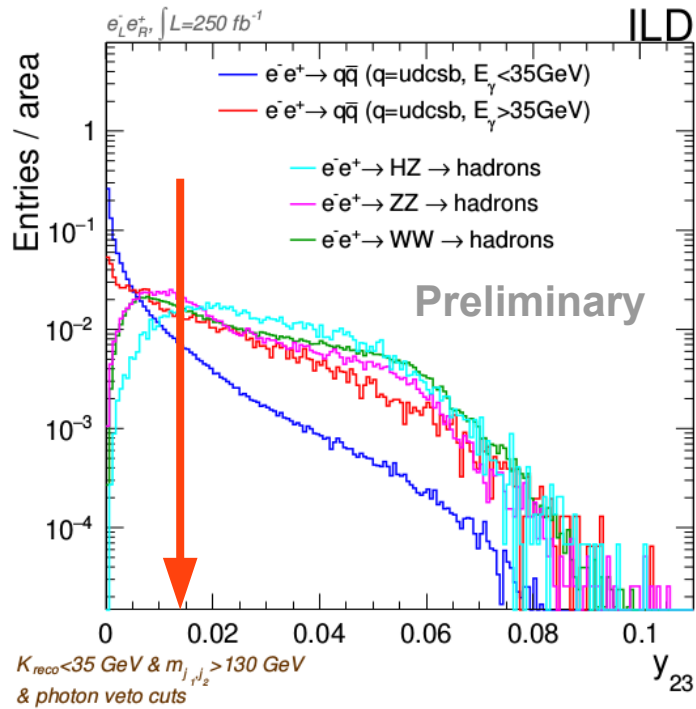
$$d_{ij} = \min(E_i^{2p}, E_j^{2p}) \frac{1 - \cos(\theta_{ij})}{1 - \cos(R)}$$

$$d_{iB} = E_i^{2p}$$



Final steps of the preselection

- ▶ Cut on $y_{23} < 0.015$ (jet distance at which the 2 jet event would be clustered in 3 jets)
- ▶ Cut on $m_{j1} + m_{j2} < 100$ GeV



Preselection summary

- ▶ Cut 1:
 - $K_{reco} < \text{GeV}$ & $m(2\text{jets}) > 130 \text{ GeV}$
- ▶ Cut 2:
 - Photon veto cuts
- ▶ Cut 3:
 - $y_{23} < 0.015$
- ▶ Cut 4:
 - $m_{j1} + m_{j2} < 100 \text{ GeV}$

- ▶ What is the preselection efficiency $\epsilon_{q\bar{q}}$ for each flavour?
 - It is flat in almost all the detector
 - Almost equal for all flavours

$\epsilon_{q\bar{q}}$

