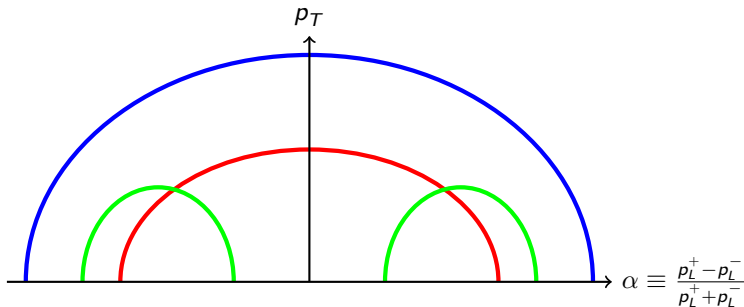


Investigating High Precision Tracker p-Scale Calibration

with K_S^0 , Λ , D^0 , J/ψ

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- ILC can and should make precision measurements of the **masses** of known fundamental particles (M_H , M_t , M_W , M_Z), and Γ_Z . Measure new ones M_X .
- A primary issue for most determinations is the measurement of the **absolute center-of-mass energy scale**. A method, the so-called \sqrt{s}_p method, has been proposed using only the momenta of muons in di-muon events.
- Critical issue for \sqrt{s}_p method: calibrating the **tracker momentum scale**.
- Up to now, foresaw $J/\psi \rightarrow \mu^+ \mu^-$ with mass known to 1.9 ppm as the gold standard method. Statistically limited though in e^+e^- collisions.

More details in my older studies of \sqrt{s}_p from [DESY](#), [ECFA LC2013](#), and of momentum-scale from [Fermilab](#), [AWLC 2014](#).

Today,

- 1 Explore method based on the Armenteros-Podolanski construction (inspired by [Rodriguez et al.](#)) using mainly K_S^0 , Λ .
- 2 Bonus: potential to also improve masses of parent and daughter particles.

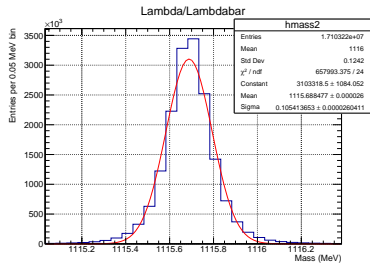
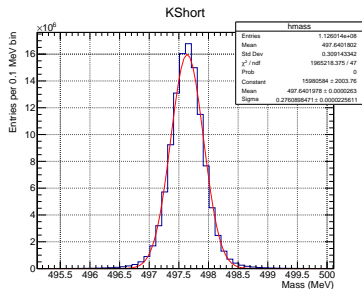
Particles of Interest

- Indeed J/ψ is the best measured parent particle, but other particles much more prevalent
- Daughter masses are known very well (except K^\pm)
- ppm = parts per million, ppb = parts per billion (10^9)

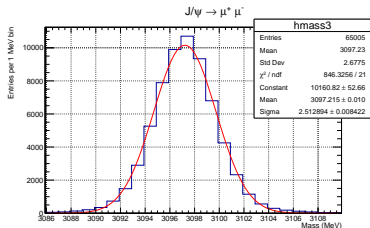
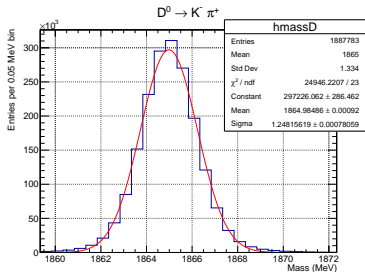
Particle	Mode	$\Delta M/M$ (PDG-2020)	$n_Z^{\text{had}} B$
K_S^0	$\pi^+ \pi^-$	26 ppm	0.71
Λ	$p \pi^-$	5.4 ppm	0.25
D^0	$K^- \pi^+$	27 ppm	0.018
J/ψ	$\mu^+ \mu^-$	1.9 ppm	0.00031
Z	$\mu^+ \mu^-$	23 ppm	(0.047)
μ	–	23 ppb	
π^\pm	–	1.3 ppm	
K^\pm	–	32 ppm	
p	–	6.4 ppb	

Following slide has the expected mass peaks from the main decay modes for a sample of 250 M hadronic Z's (91 GeV). Uses ILD momentum resolution parametrization. Hierarchy: $K_S^0, \Lambda, D^0, J/\psi$.

Mass Peaks in 250 M hadronic Z's (91 GeV)



(this is just for Λ - multiply by two).



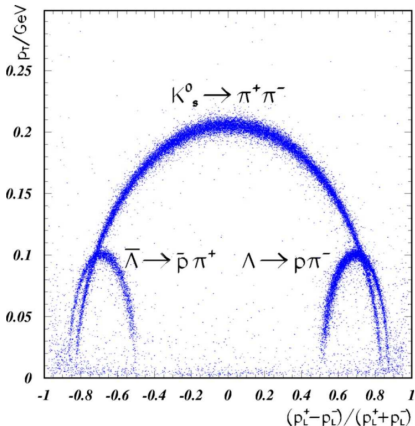
113M, 34M, 3.8M, 65k

- Study uses lots of events. Sufficient that many of the known particle masses could be improved by factors of up to 75 for 250M hadronic Zs.
- Likely a real reach in terms of eventual systematics. Discussed more later.
- Study so far is like an estimate of the average B -field absolute scale statistical uncertainty.
- For now I decided to be somewhat cavalier for several reasons.
 - ① High event counts are needed to uncover residual systematics in the fit procedure and figure out how to correct them.
 - ② Some channels need very high statistics to be viable especially in the context of the analysis method chosen.
 - ③ A high statistics run at the Z for EW precision measurements with potentially 4000 M Z's would have 16 times this statistics
 - ④ This study should give an idea of how well point-to-point systematics might be controllable for scan observables.

Armenteros-Podolanski Method for 2-body V Decays

J. Podolanski and R. Armenteros, "Analysis of V-events", Phil. Mag. 1954.
For a "V-decay", $M^0 \rightarrow m_1^+ m_2^-$, decompose the daughter momenta in the lab into components transverse and parallel to the parent momentum.

The resulting distributions of (daughter p_T , $\alpha \equiv \frac{p_L^+ - p_L^-}{p_L^+ + p_L^-}$) are related by the CM decay angle, θ^* , the underlying masses, (M, m_1, m_2) , that determine, p^* , and β .



❖ K^0 -decay

$$M = 0.498 \text{ GeV}$$

$$m_{1,2} = 0.140 \text{ GeV}$$

$$p_{cm} = 0.206 \text{ GeV}$$

$$\alpha_0 = 0$$

$$r_\alpha = 0.827$$

❖ Λ -decay

$$M = 1.116 \text{ GeV}$$

$$m_1 = 0.938 \text{ GeV}$$

$$m_2 = 0.140 \text{ GeV}$$

$$p_{cm} = 0.101 \text{ GeV}$$

$$\alpha_0 = \pm 0.691$$

$$r_\alpha = 0.181$$

plot from talk
by
M. Schmelling

Armenteros-Podolanski II

CM frame, so $p_1^* = p_2^* = p^* = \frac{\sqrt{[M^2 - (m_1 + m_2)^2][M^2 - (m_1 - m_2)^2]}}{2M}$.
 Transverse momentum, $p_T^* = p^* \sin \theta^* = p_T$ in CM and lab.



One can derive, that $\alpha = \frac{2p^*}{\beta M} \cos \theta^* + \frac{m_+^2 - m_-^2}{M^2}$ and can rewrite as

$$\alpha = \alpha_0 + \frac{r_\alpha}{\beta} \cos \theta^* \quad \text{where } \alpha_0 \equiv \frac{m_+^2 - m_-^2}{M^2}, \quad r_\alpha \equiv \frac{2p^*}{M}$$

$$\text{AP Ellipse : } \left(\frac{\alpha - \alpha_0}{(r_\alpha/\beta)} \right)^2 + \left(\frac{p_T}{p^*} \right)^2 = 1 \quad (1)$$

Elliptical Transformation a la Rodriguez

Building on the recent preprint, arXiv:2012.03620 (Rodriguez et al), one can “flatten” the measured AP ellipse into (r, ϕ) variables defined using

$$\begin{cases} \alpha(r, \phi) = \tilde{\alpha}_0 & + & \frac{\tilde{r}_\alpha}{\beta} r \cos \phi \\ p'_T(r, \phi) = S_p p_T(r, \phi) = \tilde{p}^* r \sin \phi \end{cases}, \quad (2)$$

where the tilde based values assume some reference masses, p'_T is the measured p_T biased by a momentum scale factor, S_p . Straightforward solution is:

$$\begin{cases} r^2(\alpha, p'_T, \beta) = \left(\frac{\alpha - \tilde{\alpha}_0}{(\tilde{r}_\alpha/\beta)} \right)^2 + \left(\frac{p'_T}{\tilde{p}^*} \right)^2 \\ \phi(\alpha, p'_T, \beta) = \text{atan2}\left(\frac{p'_T}{\tilde{p}^*}, \frac{\alpha - \tilde{\alpha}_0}{(\tilde{r}_\alpha/\beta)} \right) \end{cases} \quad (3)$$

Equation 1 results in a quadratic in r that depends on ϕ with coefficients defined by combinations of true masses, reference masses, S_p , and β .

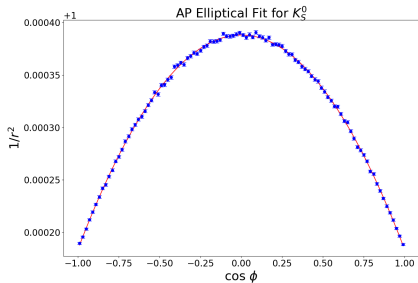
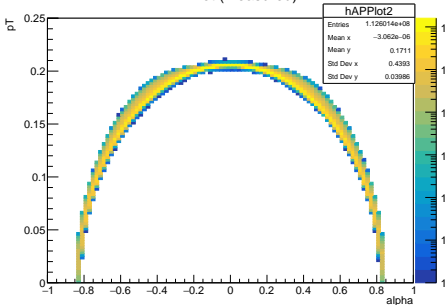
$$\left(\frac{(\tilde{\alpha}_0 - \alpha_0) + (\tilde{r}_\alpha/\beta) r \cos \phi}{(r_\alpha/\beta)} \right)^2 + \left(\frac{\tilde{p}^* r \sin \phi}{S_p p^*} \right)^2 = 1, \quad (4)$$

If all masses are as assumed, and $S_p = 1$, expect $r(\phi) = 1$ for $m_1 = m_2$.

Illustrate with K_S^0 (113M accepted decays)

Simulate measurement of (p_T, α) . For each decay, measure corresponding (r, ϕ) using Eqn. 3 with reference masses. For each bin in $\cos \phi$ (approximately the same as $\cos \theta^*$), find the mean value of $1/r^2$. For $S_p = 1.0001$.

AP Plot (measured)



Fit $1/r^2$ vs $\cos \phi$ for underlying parameters ($m_{K_S^0}$ and S_p) with m_π assumed known perfectly. Uses reference (tilde) value of the PDG K_S^0 mass + 0.1 MeV. Results: $\chi^2/\text{dof} = 108/98$ and relative uncertainties on $(m_{K_S^0}, S_p)$ of (0.3, 0.5) ppm with correlation of -97.9% . Deviations are $(+2.6, -3.9)\sigma$ (residual systematics? fix...)

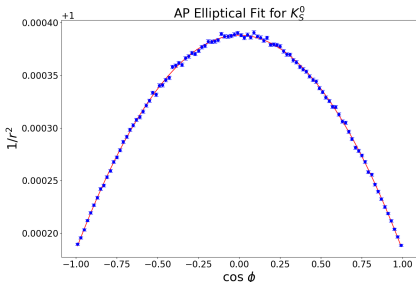
Parametric Form (for $m_1 = m_2 = m$)

For this simpler case, the solutions of the quadratic equation (with $B = 0$) lead to

$$\frac{1}{r^2} = -\frac{A}{C} = \left(\frac{\tilde{r}_\alpha}{r_\alpha}\right)^2 \cos^2 \phi + \left(\frac{\tilde{p}^*}{S_p p^*}\right)^2 \sin^2 \phi. \quad (5)$$

Note: no dependence on β . Define, $c_r = \frac{\tilde{r}_\alpha}{r_\alpha}$, $c_p = \frac{\tilde{p}^*}{S_p p^*}$. Eqn. 5 becomes

$$\frac{1}{r^2} = c_p^2 + (c_r^2 - c_p^2) \cos^2 \phi \quad (6)$$



- Can measure both r_α and the product, $S_p p^*$.
- These depend on M , m , and S_p .
- When m is well known. Can measure M and S_p .
- $r^{-2}(\cos \phi = 0, \pm 1) = (c_p^2, c_r^2)$
- p_T' depends on S_p . α does not.

Experimental Methods Used in Current Study*

Caveat: short-term goal to establish whether this method can work and get first estimates of precisions especially at the Z. Not meant to be an end-to-end study.

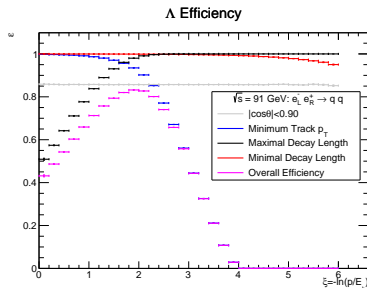
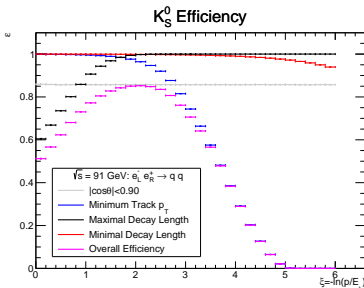
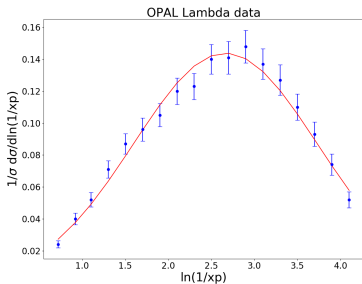
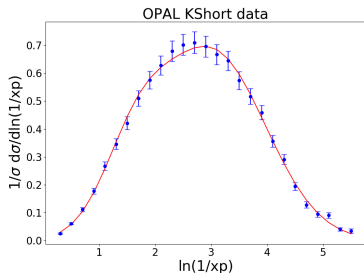
Methods

- Use toy MC generator with momentum spectrum from PYTHIA 6.
- Use parametrization of “normal” ILD momentum resolution (from fits to DBD curves) - checked with SGV. (trackresmodel.py code attached)
- Scale a and b parameters for reduced lever arm (decays). Scale b by $1/\beta$.
- $\frac{d\sigma}{d\cos\theta} = 1 + \cos^2\theta$. Rates from PDG.
- Neglect angular resolution and backgrounds.

Cuts

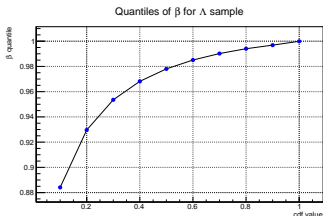
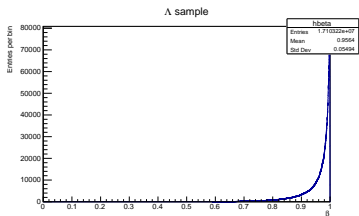
- Angular acceptance: $|\cos\theta| < 0.90$
- Minimum track detector p_T of 0.25 GeV
- Fiducial cut (decay vertex within 20 cm of outer edge of TPC in r and z)
- Require decay radius exceeds 250 μm (only for K_S^0, Λ)
- Require momentum resolution of each track $< 1\%$ (for late decays ...)

ξ Distributions and Efficiency



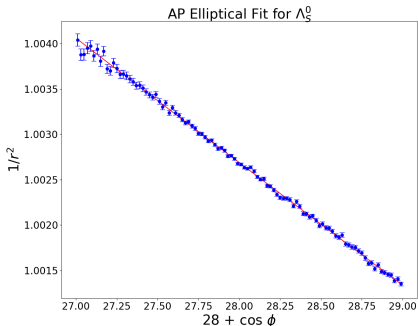
Note: efficiency plots neglect the 1% resolution cut (recent addition)

Lambda Fits (Eqn 4 with $\tilde{\alpha}_0 \neq \alpha_0$)



Lowest decile ($\beta < 0.885$) is excluded from all fits for now. (suspect β -bite is currently too big.)

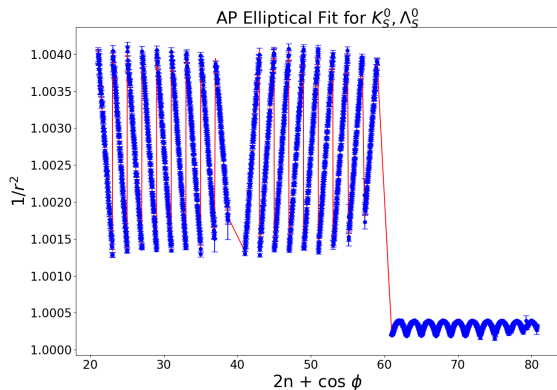
r vs ϕ dependence depends on β . So divide and conquer. Divide β range in tenths. Use $\langle \beta \rangle$ in fit for now. Example with $\beta \in [0.9853, 0.9904]$.



Again for $S_p = 1.0$ and $M + 0.1$ MeV. Fit $\chi^2/\text{dof} = 109/98$. Find (M, S_p) to $(0.3, 5)$ ppm, $\rho = -0.962$. Uses 1.7M Λ 's. Overall fit has $\times 18$ (for 250M Z).

Mega Fit for K_S^0 , Λ , $\bar{\Lambda}$ with 250M Z statistics

Fit deciles for each particle type (necessary for $\Lambda/\bar{\Lambda}$).
 9 deciles of Λ , 9 deciles of $\bar{\Lambda}$, 10 deciles of K_S^0 .



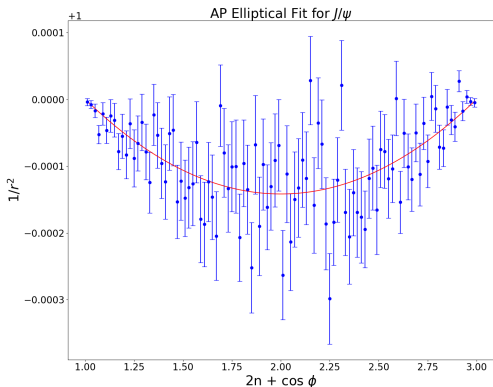
$$\chi^2 / \nu = 3117/2714 \text{ (NYGE)}$$

- ① $m_{K_S^0}$: 0.48 ppm
- ② m_Λ : 0.072 ppm
- ③ m_π : 0.46 ppm
- ④ S_p : 0.57 ppm

	$M(K_S^0)$	$M(\Lambda)$	m_π	S_p
$M(K_S^0)$	1.0	0.934	0.814	-0.942
$M(\Lambda)$	-	1.0	0.914	-0.818
m_π	-	-	1.0	-0.601
S_p	-	-	-	1.0

Here m_p is fixed. Technically, fits are set up with χ^2 penalty terms that constrain the particles to the PDG masses within known uncertainties. So in essence the fitted masses are new world averages.

Fix m_μ .



- ① $m_{J/\psi}$: 1.9 ppm (no improvement wrt PDG)
- ② S_p : 4.4 ppm
- ③ Correlation: -0.44

- Consistent with prior estimate of 1.0 ppm statistical uncertainty on S_p from $J/\psi \rightarrow \mu^+ \mu^-$ in 4 GZ hadronic (find 1.03 ppm here).
- Fit with $m(J/\psi)$ fixed gives 4.0 ppm uncertainty on S_p .

Intrinsic S_p sensitivity (mostly for reference)

For each particle decay separately, carry out fit with all particles set to correct masses, and fit only for S_p (again using 250M hadronic Zs).

Results:

Particle	Mode	S_p uncertainty	$n_Z^{\text{had}} B$	S
K_S^0	$\pi^+ \pi^-$	0.104 ppm	0.71	0.00139
Λ	$p \pi^-$	0.297 ppm	0.25	0.00235
D^0	$K^- \pi^+$	0.538 ppm	0.018	0.00114
J/ψ	$\mu^+ \mu^-$	3.98 ppm	0.00031	0.00111

Sensitivity, S , defined as S_p relative uncertainty per event, ie,

$$S = [\Delta(S_p)/S_p] \sqrt{250 \times 10^6 (n_Z^{\text{had}} B)}$$

Note that the sensitivity to S_p differs quite a lot (due to different Q values).

I am also looking into adding other decay modes into the mix. For now, adding D^0 with only the $K^- \pi^+$ decay mode would add uncertainty from both the D^0 mass and the charged kaon mass. $\phi(1020)$ is not so interesting given the width.

For example,

① $D^+ \rightarrow K_S^0 \pi^+$ (1.56%)

② $D^+ \rightarrow K^- \pi^+ \pi^+$ (9.38%)

③ $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ (2.80%)

So far fits to D^0 with only $K^- \pi^+$ (3.95%) with two free parameters, S_ρ and $m(K)$, (ie. $m(D^0)$ fixed) give S_ρ to 0.69 ppm and $m(K)$ to 2.0 ppm for 250M Zs. ($\rho = +0.63$).

More realistic and improved methods / systematics?

Full Reconstruction

- Need to develop performant V0 finder and fitter for large IP.
- With nano-beams - much potential.
- Looper reconstruction
- Angular resolution
- Backgrounds
- Expect some degradation in resolution for Si-poor tracks.

Systematics

- Field map precision
- Tracker alignment/survey
- Material distribution
- Energy loss corrections
- Radiative tails
- Variations with p , $\cos \theta$, decay point
- Smaller bins / better β treatment
- Understand r , $\cos \phi$ resolution

- Current measurements are based on the sample average from TProfile plots. Scope for improved measurement uncertainties simply from fits to the r distributions in each $\cos \phi$ bin and better use of errors.
- Note β estimate uses the measured P and measured mass of the two tracks.

$$M^2 = m_1^2 + m_2^2 + 2p_1 p_2 \left(\frac{1}{\beta_1 \beta_2} - \cos \psi_{12} \right), \quad \beta = P / \sqrt{P^2 + M^2}$$

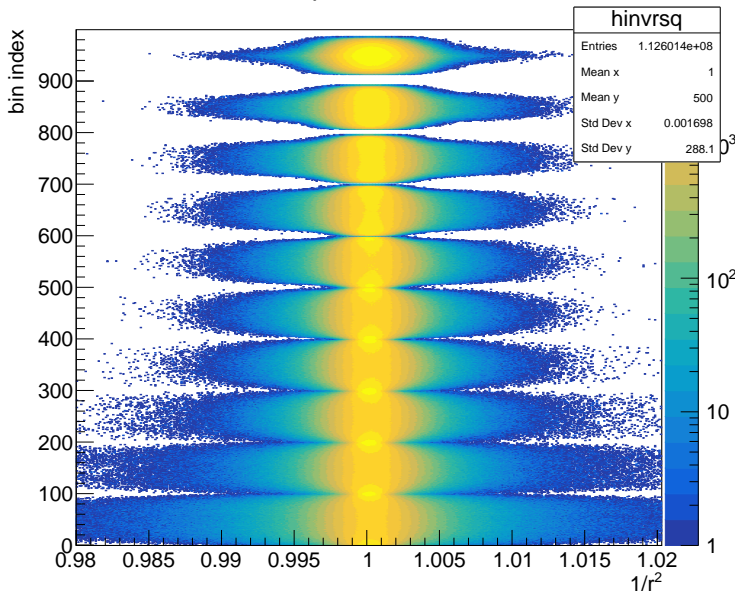
- Tremendous opportunity to target ppm type uncertainties on the momentum scale factor, S_p , at the Z.
- Would open up precision measurements of the masses of lots of known particles at the ppm level. In particular: K_S^0 , Λ , π^\pm , K^\pm .
- Would guarantee similar precision in the center-of-mass energy scale which would open a high precision Z program.
- When I started working on this a few years ago, 10 ppm seemed a sensible but very challenging goal. Maybe the bar should be set a bit higher still.
- Convincing people that this is realistic when typical experiments are at best at the 100 ppm level needs a lot of more realistic work, and work on designing this kind of functionality in from the start.
- (Now need to evaluate better limiting non- S_p systematics for \sqrt{s}_p method for \sqrt{s}).

Some simplifications:

- Each of the decayed tracks is assumed to be in the same direction as the parent particle in terms of detector, θ (should fix). This affects the simulated detector acceptance and the momentum resolution formula.
- Each track's momentum magnitude along its momentum direction is smeared as detailed on p11. So no angular smearing.
- The AP variables are calculated using the components of the smeared decay particle momentum perpendicular and parallel to the parent particle's true flight direction.
- In particular the AP p_T variable is calculated as the average of the $|p_T|$ of the two tracks.
- Note that once the tracks are fitted to a common vertex and this neutral vertex is constrained to the nano beam spot, it is expected that the above assumptions are not unreasonable, especially in $r - \phi$.

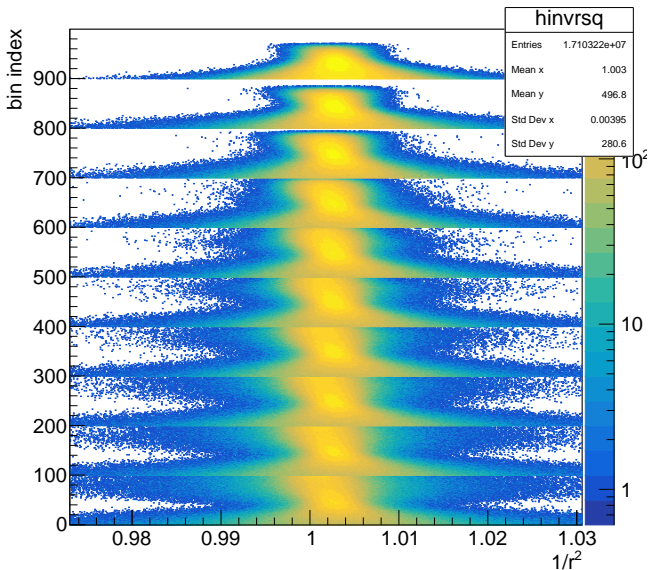
Backup: K_S^0 plot (β deciles and $\cos \phi$ bins)

Unrolled plot with resolution

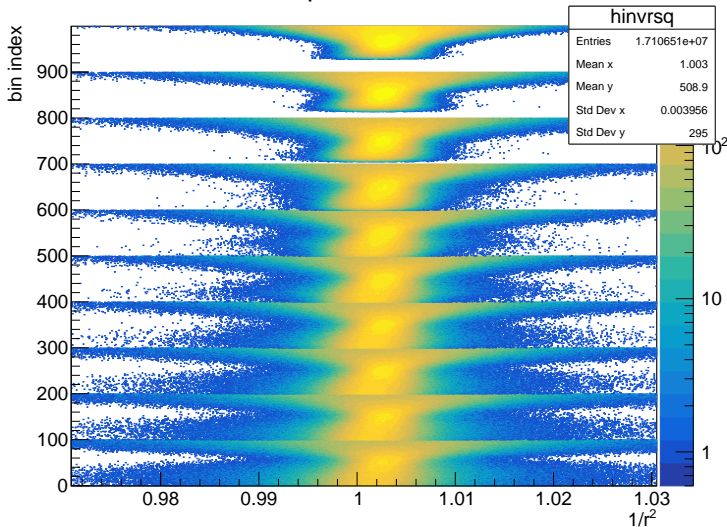


Backup: Λ plot (β deciles and $\cos \phi$ bins)

Unrolled plot with resolution

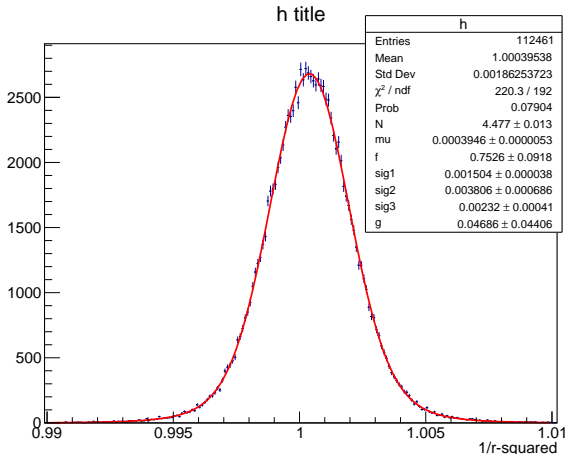


Unrolled plot with resolution



Backup: Example $1/r^2$ Distribution

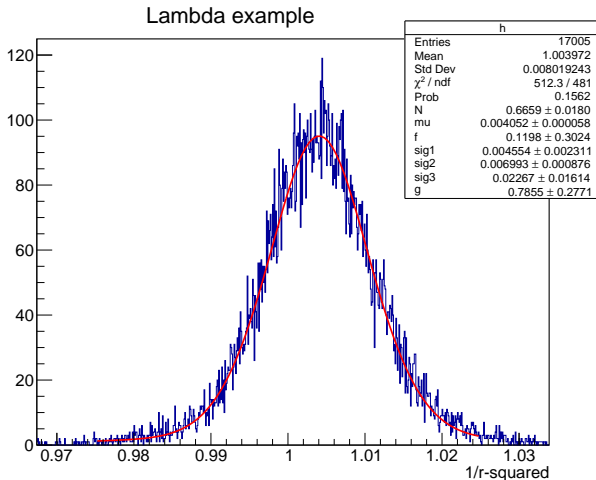
K_S^0 , $\beta \in [0.99037, 0.99418]$, $\cos \phi \in [-0.02, 0.0]$



Triple Gaussian fit with common mean of $1 + \mu$. (Note that histogram mean is currently used - overly sensitive to far tails ...)

Backup: Example $1/r^2$ Distribution

Λ , same β as slide 13, $\cos \phi \in [-1.0, -0.98]$



Triple Gaussian fit with common mean of $1 + \mu$. (Note that histogram mean is currently used)

Backup: Plots from Rodriguez Preprint

Note here ϕ definition differs by $\pi/2$.

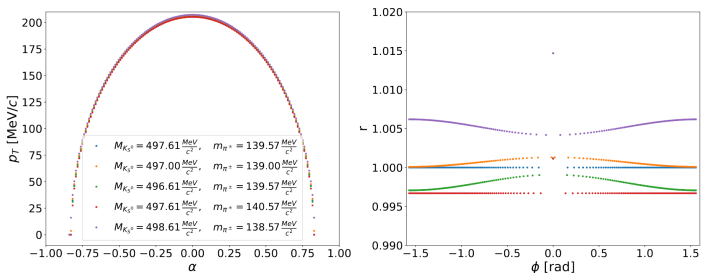


Figure 2: Functional form of a set of symmetric decays ($m_1 = m_2 = m$) in the Armenteros plot (left) and in the plot with elliptical coordinates (right).

In this case ($S_p = 1$). I included this mainly for illustration. It does not comport well with the plotting convention I used, nor the idea that there is a true AP ellipse, and a series of flattened elliptical coordinate plots for different reference values.

I think what is done is 5 different (M, m) assumptions are made, and the right plot is the analysis based on the PDG as reference (leading to the blue set at $r = 1$).