Measurement of $\sigma(e^+e^- \to HZ) \times Br(H \to ZZ^*)$



Channels

 $Z_1(jj),$

Z(jj),

 $\mathrm{Z}^*(\ell\ell)$

 $Z_1(jj),$

 $Z(\ell\ell)$,

Z(jj),

 $Z^{\star}(\ell\ell)$

 $Z(\ell\ell)$,

 $Z^{\star}(jj)$

 $Z^{\star}(\ell\ell)$

 $Z(\ell\ell)$,

 $Z^{\star}(jj)$

 $Z_1(\nu_e\bar{\nu}_e)$

 $Z_1(\nu_{\mu,\tau}\bar{\nu}_{\mu,\tau}),$

 $\mathcal{P}_{e^-e^+}$ MC Lepton

eLpR 23989

eRpL 23845

eLpR 23261

eRpL 23132

eLpR 23059

eRpL 23096

eRpL 23225

eRpL

 $Z_1(\nu_{\mu,\tau}\bar{\nu}_{\mu,\tau}), \text{ eLpR} \quad 23189$

23862

events tagging,

events

16088

16027

eLpR $24044 17429 3.7 \cdot 10^{-3}$

eRpL 23910 $17259 7.9 \cdot 10^{-5}$

eLpR $23840 ext{ } 17103 ext{ } 4.1 \cdot 10^{-3}$

17168

21168



at the 250 GeV ILC published in Phys Rev D. DOI: 10.1103/PhysRevD.104.093007



E. Antonov (LPI/MEPhI) on behalf of the ILD Collaboration

Introduction

Width of the Higgs boson is difficult to measure at LHC in a model-independent approach (the uncertainty is expected to be ~20% after LHC luminosity upgrade) [DOI: 10.1093/ptep/ptaa104].

We propose to use the process $e^+e^- \rightarrow HZ$ with the subsequent decay $H \rightarrow ZZ^*$ to measure in ILC:

$$\sigma(e^+e^- \to HZ) \times Br(H \to ZZ^*) = C \cdot g_Z^4/\Gamma_H$$

Weighted

number

of events

338

79

1.7

71

2.7

Constant, Error < 1% expected arXiv:1403.7734

> Weight factors

 $2.1 \cdot 10^{-2}$

 $1.3 \cdot 10^{-3}$

 $20879 \quad 2.1 \cdot 10^{-2}$

 $20664 \quad 1.3 \cdot 10^{-3}$

 $21108 \quad 3.7 \cdot 10^{-3}$

 $21149 \quad 7.9 \cdot 10^{-5}$

 $21246 \quad 1.6 \cdot 10^{-4}$

 $1.6 \cdot 10^{-4}$

 $4.1 \cdot 10^{-3}$

Coupling HZZ Error < 0.5% expected arXiv:1903.01629

Higgs boson width

Full Geant-4-based detector simulation

One of secondary Z bosons (from Higgs boson decay) is reconstructed from two quarks:

$$Z_1 \rightarrow jj \ or \ \nu\nu, Z \rightarrow jj \ or \ ll, Z^* \rightarrow ll \ or \ jj$$
On-shell
Off-shell

 $e^+e^- \to Z_1(j_1j_2) H, \quad H \to Z(j_3j_4)Z^*(\ell_1\ell_2)$ **Channel 1:**

 $e^+e^- \to Z_1(j_1j_2) H, \quad H \to Z(\ell_1\ell_2)Z^*(j_3j_4)$ **Channel 2:**

 $e^+e^- \to Z_1(\nu\bar{\nu}) H, \quad H \to Z(j_1j_2)Z^*(\ell_1\ell_2)$ **Channel 3:**

 $e^+e^- \to Z_1(\nu\bar{\nu}) H, \quad H \to Z(\ell_1\ell_2)Z^*(j_1j_2)$ **Channel 4:**

Analysis steps

- 1. For detailed background studies we extract specific processes on generator level.
- 2. We identify two isolated lepton candidates.
- 3. ISR identification and removing procedure. DOI: 10.1140/epjc/s10052-020-08729-7
- 4. Jet reconstruction using clustering tools.
- 5. Applying weight factors to each event to get expected number of signal or background events.

 $\left[\frac{1\pm0.8}{2}\cdot\frac{1\pm0.3}{2}\right]\cdot\frac{2\ ab^{-1}}{\mathcal{L}}$ \mathcal{L} - the sample integrated luminosities

Valencia algorithm is used to force the remaining particles into 2 or 4 jets.

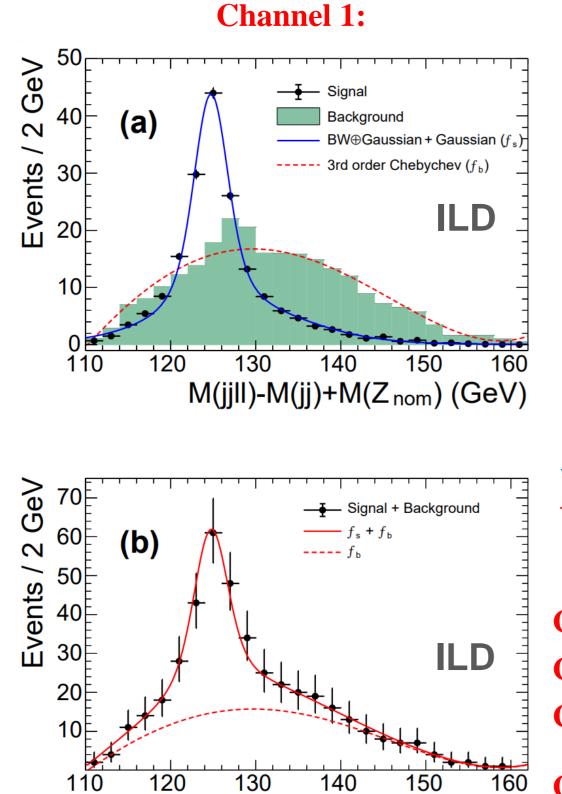
Optimized Valencia algorithm parameters chosen for the jet reconstruction in different channels:

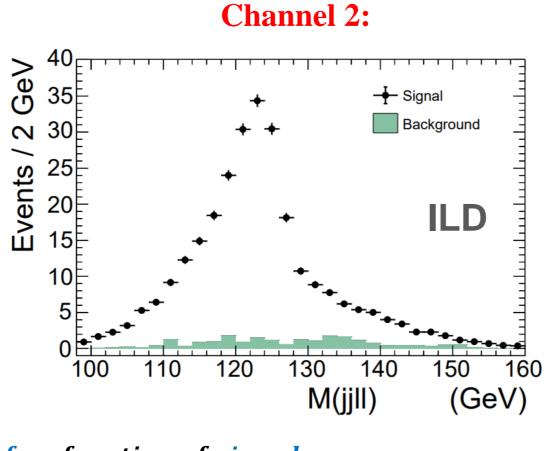
Valencia	$Z_1(jj),$	$Z_1(jj),$	$Z_1(\nu\bar{\nu}),$	$Z_1(\nu\bar{\nu}),$
parameters	Z(jj),	$Z(\ell\ell)$,	Z(jj),	$\mathrm{Z}(\ell\ell),$
	$\mathrm{Z}^*(\ell\ell)$	$Z^*(jj)$	$\operatorname{Z}^*(\ell\ell)$	$\mathrm{Z}^{\star}(jj)$
β	1.0	1.0	1.0	1.0
γ	0.4	0.4	0.6	0.3
R	1.6	0.7	1.4	1.4
	·			·

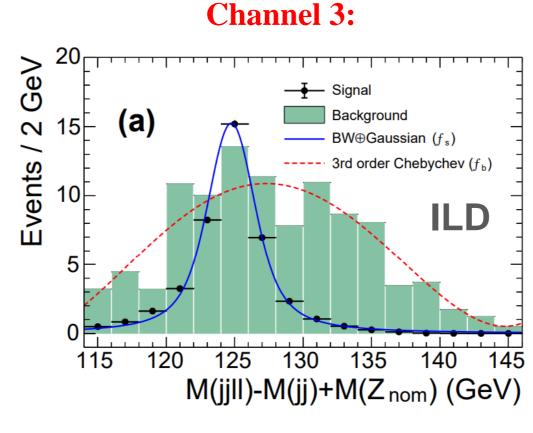
To cancel jet reconstruction uncertainties:

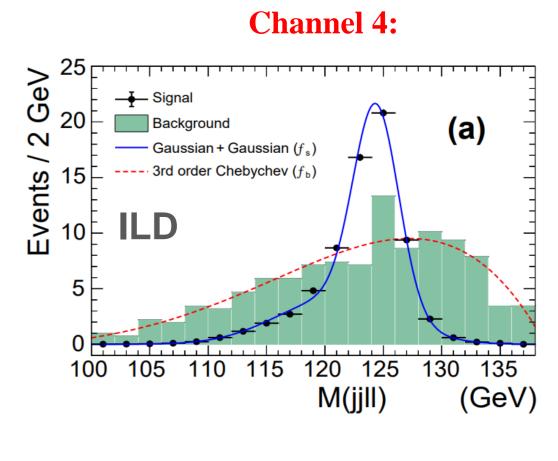
$$M_{\Delta} = M(jj\ell\ell) - M(jj) + M(Z_{nom})$$

Analysis results for four channels









 f_s – fraction of signal

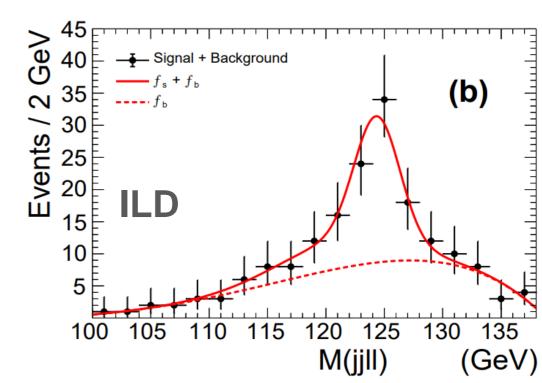
f_b – fraction of background

Dominant backgrounds:

 $W^+W^-\gamma^*$ and $ZZ\gamma^*$, where $\gamma^* \rightarrow ll$ No significant background

 $Z(2j)Z(2\tau)$ with τ leptonic decays, $W(2j)W(l\nu)$, $b\overline{b}$ semileptonic decays Ch 4: Similar to the channel 3 except the $b\overline{b}$

35 Ge/ $----f_s + f_b$ (b) **ILD** 125 130 120 $M(jjII)-M(jj)+M(Z_{nom})$ (GeV)



Combined signal significance estimate

We calculate the combined statistical uncertainty from statistical significance of individual channels:

 $M(jjII)-M(jj)+M(Z_{nom})$ (GeV)

$$S_{\text{comb}} = 1/\sqrt{\sum_{i=1}^{4} S_i^{-2}}$$

These results indicate that the Higgs width can be measured at ILC with an accuracy of about (5-6)% in the model-independent approach.

Number of signal events and uncertainties for each channel

	$Z_1(jj),$	$Z_1(jj),$	$Z_1(\nu\bar{\nu}),$	$Z_1(\nu\bar{\nu}),$	Sum
	Z(jj),	$\mathrm{Z}(\ell\ell),$	Z(jj),	$\mathrm{Z}(\ell\ell),$	
	$\mathrm{Z}^*(\ell\ell)$	$\mathrm{Z}^*(jj)$	$\mathrm{Z}^*(\ell\ell)$	$Z^{\star}(jj)$	
		2 ab^{-1}	eLpR		
Number	192.4	275.3	51.9	73.3	-
of events	± 24.9	±17.2	± 13.0	± 14.2	-
Statistical	12.9%	6.3%	25.1%	19.3%	5.29%
uncertainty					
($0.9~\mathrm{ab}^-$	1 eLpR +	-0.9 ab^{-1}	eRpL	1
Number	135.2	202.2	30.9	67.3	-
of events	± 20.4	±14.7	±10.7	± 14.3	-
Statistical	15.1%	7.3%	34.6%	21.2%	6.15%
uncertainty					

Conclusions

At 250 GeV the accuracy of this method is similar to one obtained in arXiv:1310.0763, arXiv:1403.7734 using the combination of four channels measurements. The results of both methods can be combined to further improve the accuracy.

Our measurement can be used to test the Higgs width value obtained within the SM, as well as within the EFT approach. The theoretical accuracy of the Higgs width is expected to be about 2% [DOI: 10.1103/PhysRevD.97.053003].

The Higgs boson width can be measured in ILC experimentally in a model-independent approach with accuracy is about (5-6)%.

