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Ion backflow in TPC

P. Colas

Motivation for the tracking at e+e- colliders as Higgs factories

- Need to reconstruct complex multi-track events (jets) in a noisy environment : calls for high segmentation
- Also need to reconstruct very accurately high energy tracks from Z recoil to Higgs. This translates into $O(10 \mu m)$ control of the systematics on sagitta and $\Delta(1/P)$ =10⁻⁴ GeV ⁻¹ for TPC only (2.10 -5 GeV -1 for the whole tracking system).
- Silicon detectors give point measurement accuracy, but also introduce multiple scattering, while a TPC provides a continuous 3D track reconstruction with minimal matter : useful for V0, kinks, connecting to vertex tracker, other silicon trackers, and to calorimeter.
- Also a TPC has dE/dx capability, for K/π separation, essential for Heavy Flavour physics at the Z pole.

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The Ion Problem

- Ions are copiously produced in the amplification device. A fraction of them will flow back in the drift space and add to the primary ions produced by the charged tracks and converted X-rays from the machine background.
- Ions drifting in the electric field of the TPC are very slow, of the order of a m/s
- As a consequence, this positive charge accumulates and gives rise to a space charge. The region where this space charge builds up being limited, and the density being non-uniform, this produces transverse E field components which produce distortions (enhanced by ExB effects)
- At the TeraZ, there is one Z produced in every drift time of the TPC. Even primary ionization produces enough space charge to distort the tracks.

• The ion backflow problem impacts TPCs at colliders (LEP, RHIC, LHC/Alice). It is usually escaped by stopping and neutralizing the ions with a gating grid of wires between collisions.

Parallel wires are set at alternating potentials to close the gate by producing a transverse field. The gate is open only during electron drift (50 µs at LEP)

C. Garabatos in workshop «New Horizons for TPCs», October 7, 2020

MWPC won't work for ALICE in Run 3

- In 2022 the LHC will deliver \bullet 50 kHz Pb-Pb collision rate
- At \sim 100 kHz/cm² the space \bullet charge near the anode wires would affect dE/dx resolution
- With a gating grid, only 3 kHz can be achieved
	- GG must stay closed while jons from the avalanche reach the wires, otherwise 10% of them escape and would produce \sim 1 m distortions in the drift volume

Need for gating

In TPCs, ions are produced and migrate very slowly (1 m/s) . They produce a charge density which can be one or two orders of magnitude above the primary ionization (IBF*Gain). Present-day electronics allows running at a gas gain of 2000 or a bit less. 500 could be in reach. Thus backflow fractions of 0.5% to 2% would not add to the primary ionization. The resulting electric field can be the origin of distortions.

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Charge density for 100 bunches

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One gating solution: the gating GEM

86 % transparency for electrons demonstrated for Delta V = 4 V over 50 μ m Preservation of the resolution demonstrated in a beam test Ion blocking expected for DeltaV =-40V

Other solution : equipotential wire grids (was not possible at LEP as there was no time to switch it on and off)

A. Sugiyama et al.

Summary on gating

- Gating is a solution for ILC @ nominal luminosity, for instance with a gating grid, thanks to the beam time structure at 5 or 10 Hz.
- However mechanically very difficult.
- In many other applications (sPhenix@EIC, ALICE@LHC, FCCee, CEPC), gating is not possible because of continuous beam.

ONE MUST FIND WAYS OF MITIGATING THE BACK-FLOW WHILE KEEPING GOOD ELECTRON EFFICIENCY

Note that at the TeraZ, even primary ionization rules out the use of TPC

Hadronic Z decays

The higher the magnetic field, the more curlers, increasing ionization

This estimate is probably a lower limit, if not an underestimate. No background is included. Assumes 20 kHz Z rate (should be up to 40 kHz) and 40 ions per cm (should be 100) Assumes IBF*gain=1

Possible handles : lower magnetic field, inner radius increased to 50cm?

Still OK for CEPC parameters (1 to 2 orders of magnitude less in luminosity) (Huirong Qi assuming Gain*IBF=5 with a Micromegas + GEM)

Ion backflow mitigation/suppression

- Several ideas have been proposed to suppress ion backflow
- 4-GEMS, 2 or 3 GEMs + Micromegas, double or triple meshes, staggered, crossed…
- The situation is rather confused : impressive results like 10⁻⁴ or even 10⁻⁵ ion backflows are suspected to be due to screening by space charge . The measurement is extremely difficult.
- The suppression is often obtained at the cost of resolution

(in energy/dE/dx and position)

Ion backflow natural suppression in Micromegas

Electrons from the avalanche undergo transverse diffusion. The ions created follow field lines with negligible diffusion. Thus most of them (outside the funnel) go back to the mesh where they are neutralized.

The maximum backflow is thus proportional to the field ratio (0.3% in the ILC TPC case).

The inverse field ratio law is well verified experimentally

P. Colas *et al.*, NIMA535(2004)226

More detailed measurements have been carried out with Ingrid detectors. (M. Chefdeville's thesis)

J. Martin (2003) When the avalanche extends to neighbouring holes, the minimum backflow is attained

4 GEMs

ALICE with 4 GEMs Also adopted by SPhenix

Klaus Dehmelt

Piotr Gasik

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Double Micromegas

DMM (Jianbei Liu et al., USTC) Claim that misaligning the meshes reduces the backflow.

But resolution also degrades.

Also studied in F. Jeanneau et al. NIMA623(2010)94, but very difficult measurement

Caveats : Moiré modulation of the gain, effect of magnetic field

It is also possible to make use of a magnetic field in a high $\omega\tau$ gas (E. Shulga, Weizmann)

Also measurements have been done with Micromegas + 1, 2 or 3 GEMs At the end marginally better than Micromegas alone

One direction of research is to use Graphene, but this looks very difficult

Graphene Layers in Gas Detectors

- Goal: create a device fully transparent to electrons and fully opaque to ions
- Graphene is narrowest and thinnest possible conductive mesh with pore size $<$ 1 Å
- Study of charge transfer \bullet through graphene layer suspended on Cu meshes

AUCLI

31/10/2014

WIEN

Possible solution : shape the field so that they converge to a small region, from which electrons escape by diffusion. Necessitates to shape the anode at the hole scale level.

This can be done using micro-additive fabrication at the sub-micronic level (example from Tiago Silva, Sao Paulo)

Non-exhaustive list of potential issues with indicative expectations:

Peter Kluit, Nikhef, CEPC workshop, Shanghai, Oct. 2020

Conclusions

- Ion backflow has to be suppressed if one is to select Higgs recoils.
- If gating is not possible (case of FCC and CEPC) one needs to find a new structure to suppress the ion backflow without degrading the space resolution or dE/dx resolution.
- So far many trials with multiple GEMs or Micromegas or combinations have been done, but still IBF limited to 1 permil or a few.
- Micro-additive fabrication of the anode might provide a technical solution.

Peter Kluit : 'there is no need for a gating device (a double grid suffices)

Reducing the Ion back flow in a Pixel TPC

The Ion back flow can be reduced by adding a second grid to the device.

It is important that the holes of the grids are aligned. The lon back flow is a function of the geometry and electric fields. Detailed simulations - validated by data - have been presented in LCTPC WP #326.

With a hole size of 25 µm an IBF of 3 10⁻⁴ can be achieved and the value for IBF*Gain would be 0.6. Well below the specifications.

A TPC for ILC

Options : GEM and Micromegas with pads or pixels

R&D lead to a design of a Micromegas module with fully integrated readout and 2-phase CO2 cooling. The ERAM concept (Encapsultated Resistive Anode Micromegas) proved to 'do the job' : 200 points with 60 µm resolution

This ERAM concept is adopted by T2K/ND280 for its upgrade. If successful, it might be used at DUNE/SAND